



# Observation of new baryons in the $\Xi_b^- \pi^+ \pi^-$ and $\Xi_b^0 \pi^+ \pi^-$ systems

LHCb collaboration<sup>†</sup>

## Abstract

The first observation and study of two new baryonic structures in the final state  $\Xi_b^0 \pi^+ \pi^-$  and the confirmation of the  $\Xi_b(6100)^-$  state in the  $\Xi_b^- \pi^+ \pi^-$  decay mode are reported using proton-proton collision data collected by the LHCb experiment, corresponding to an integrated luminosity of  $9 \text{ fb}^{-1}$ . In addition, the properties of the known  $\Xi_b^{*0}$ ,  $\Xi_b'^-$  and  $\Xi_b^{*-}$  resonances are measured with improved precision. The new decay mode of the  $\Xi_b^0$  baryon to the  $\Xi_c^+ \pi^- \pi^+ \pi^-$  final state is observed and exploited for the first time in these measurements.

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The  $\Xi_b^{(-,0)}$  baryons form an isospin doublet and are made of a  $b$  quark, an  $s$  quark and a lighter  $q$  ( $u$  or  $d$ ) quark. Their ground states have angular momentum  $L = 0$  between the  $b$  quark and the light diquark. Three isospin doublets of such non-excited states are expected in the quark model [1–3] with different spin-parity  $J^P$  and angular momentum of the  $sq$  diquark  $J_{sq}$ . The  $\Xi_b^{(-,0)}$ ,  $\Xi_b^{(0,-)}$  and  $\Xi_b^{*(0,-)}$  states are characterized by  $(J_{sq}, J^P)$  values of  $(0, \frac{1}{2}^+)$ ,  $(1, \frac{1}{2}^+)$  and  $(1, \frac{3}{2}^+)$ , respectively. Although five of these states have been observed experimentally [4–7], the  $\Xi_b^0$  baryon remains unobserved. This may be because its mass lies below the threshold for the  $\Xi_b^0 \rightarrow \Xi_b^- \pi^+$  decay [8, 9], meaning that it only decays to either the  $\Xi_b^0 \pi^0$  or  $\Xi_b^0 \gamma$  final states making it experimentally challenging to observe. A number of excited states of higher mass is expected, with predictions for their properties available, *e.g.* in Refs. [9–17]. The CMS collaboration has reported the observation of the  $\Xi_b(6100)^-$  resonance in the  $\Xi_b^- \pi^+ \pi^-$  final state, using  $\Xi_b^-$  decays to final states containing  $J/\psi$  mesons [18].

In this Letter, both the  $\Xi_b^- \pi^+ \pi^-$  and  $\Xi_b^0 \pi^+ \pi^-$  final states and their intermediate  $\Xi_b^- \pi^+$  and  $\Xi_b^0 \pi^-$  states are investigated experimentally (the inclusion of charge conjugate processes and the use of natural units are implicit throughout this Letter), using  $pp$  collision data collected by the LHCb experiment at centre-of-mass energies of 7, 8, 13 TeV, corresponding to an integrated luminosity of  $9 \text{ fb}^{-1}$ . The observation of three narrow states is reported and their properties are measured.

The LHCb detector [19, 20] is a single-arm forward spectrometer covering the pseudorapidity range  $2 < \eta < 5$ . The detector includes a high-precision tracking system consisting of a silicon-strip vertex detector surrounding the  $pp$  interaction region [21], a large-area silicon-strip detector located upstream of a dipole magnet with a bending power of about 4 Tm, and three stations of silicon-strip detectors and straw drift tubes [22, 23] placed downstream of the magnet. Different types of charged hadrons are distinguished using information from two ring-imaging Cherenkov detectors [24]. Simulated data samples are produced with the software packages described in Refs. [25–29] and are used to model the detector resolution and optimize the selection criteria.

Samples of  $\Xi_b^-$  ( $\Xi_b^0$ ) candidates are formed from  $\Xi_c^0 \pi^-$  ( $\Xi_c^+ \pi^-$ ) and  $\Xi_c^0 \pi^- \pi^+ \pi^-$  ( $\Xi_c^+ \pi^- \pi^+ \pi^-$ ) combinations, where the  $\Xi_c^0$  ( $\Xi_c^+$ ) baryon is reconstructed in the  $pK^-K^- \pi^+$  ( $pK^- \pi^+$ ) final state. To suppress background coming from promptly produced particles, all  $\Xi_b^{(-,0)}$  decay products are required to be displaced significantly from all primary  $pp$  collision vertices (PVs) in the event. The reconstructed  $\Xi_c^0$  ( $\Xi_c^+$ ) candidates are required to have a mass within 20 MeV (25 MeV) of the respective world-average mass values [30]. Displaced pion tracks are combined with  $\Xi_c$  candidates to form  $\Xi_b^{(-,0)}$  candidates, requiring good vertex-fit quality and significant displacement of the  $\Xi_b^{(-,0)}$  decay point from any PV in the event. All charged particles are required to have particle-identification (PID) information consistent with their respective mass hypotheses. PID variables are based on neural network algorithms [31] and their distributions in simulation are calibrated using data [32]. If multiple candidates per collision event pass the selection requirements, all of them are preserved in the data sample. The topological, kinematic and PID variables are used as inputs to a Boosted Decision Tree (BDT) classifier [33] that discriminates  $\Xi_b^{(-,0)}$  signal candidates from random track combinations. The classifier is trained using simulated  $\Xi_b^{(-,0)}$  decays as a signal proxy and  $\Xi_b^{(-,0)}$  data candidates in the sideband  $5900 < m(\Xi_c^{(0,+)} \pi^-, \Xi_c^{(0,+)} \pi^- \pi^+ \pi^-) < 6000 \text{ MeV}$  as a background proxy. The mass distributions of the selected  $\Xi_b^{(-,0)} \rightarrow \Xi_c^{(0,+)} \pi^-$  and  $\Xi_b^{(-,0)} \rightarrow \Xi_c^{(0,+)} \pi^- \pi^+ \pi^-$  candidates are

shown in the supplemental material [34]. The decay mode  $\Xi_b^0 \rightarrow \Xi_c^+ \pi^- \pi^+ \pi^-$  is observed for the first time experimentally. The selection requirement on the BDT classifier output is optimized for the observation of new states and retains 96% of the  $\Xi_b^-$  and 92% of the  $\Xi_b^0$  signal candidates. Additional vetoes are imposed, as described in Ref. [35], to suppress other abundant processes with displaced vertices *e.g.* those coming from  $D^0$ ,  $D^+$ ,  $D_s^+$  and  $\Lambda_c^+$  decays with a misidentified particle.

The  $\Xi_b^{(-,0)}$  candidates within a  $\pm 75$  MeV window around the peak position are combined with one charged pion (two pions) to investigate the  $\Xi_b^- \pi^+$  and  $\Xi_b^0 \pi^-$  ( $\Xi_b^- \pi^+ \pi^-$  and  $\Xi_b^0 \pi^- \pi^+$ ) mass spectra. In order to improve the mass resolution, the obtained candidates are refitted with the masses of the  $\Xi_b^{(-,0)}$  and  $\Xi_c^{(0,+)}$  baryon candidates constrained to their known values [30] and the  $\Xi_b^{(-,0)}$  flight direction to originate from the PV [36]. Additional requirements are applied to the  $\Xi_b \pi^- \pi^+$  candidates, where the  $\Xi_b^{*0}$ ,  $\Xi_b'^-$  and  $\Xi_b^{*-}$  intermediate states are selected according to their observed widths and known mass values [30]. Signal mass windows for the  $\Xi_b \pi$  intermediate resonances are defined as  $|m(\Xi_b \pi) - m_{\Xi_b^{*0}}| < 3$  MeV,  $|m(\Xi_b \pi) - m_{\Xi_b'^-}| < 1.25$  MeV and  $|m(\Xi_b \pi) - m_{\Xi_b^{*-}}| < 5$  MeV, each corresponding to  $2.5\sigma$  of the observed experimental peak.

The signal yields and lineshape parameters of the signal resonances are determined with an extended unbinned maximum-likelihood fit to the  $Q$ -value distributions defined as  $m_{\Xi_b \pi} - m_{\Xi_b} - m_\pi$  and  $m_{\Xi_b \pi \pi} - m_{\Xi_b} - 2m_\pi$  for  $\Xi_b \pi$  and  $\Xi_b \pi \pi$  decays, respectively. The mass distributions of the  $\Xi_b^- \pi^+$  and  $\Xi_b^0 \pi^-$  ( $\Xi_b^- \pi^+ \pi^-$  and  $\Xi_b^0 \pi^- \pi^+$ ) samples are shown in Fig. 1 (Fig. 2), together with the results of the fit. All signal components are modeled using relativistic Breit–Wigner distributions [37] including Blatt–Weisskopf form factors [38] with a radius of  $3 \text{ GeV}^{-1}$ . The orbital angular momentum between the  $\Xi_b^{(-,0)}$  baryons and the pions is assumed according to the expected spin assignment. The relativistic Breit–Wigner distributions are convolved with functions parameterizing the detector resolution. These resolution models are determined from simulation samples and are consistent with a resolution that scales as  $\sqrt{Q}$ . Simulation shows that for each resonance, the resolution is comparable to or smaller than the measured natural widths of the peaks, with the exception of the  $\Xi_b'^-$  baryon. The background contribution is parameterized as  $(Q - d)^n$ , where  $d$  and  $n$  parameters vary freely in the fit. This function, which is validated using wrong-charge  $\Xi_b^- \pi^-$  and  $\Xi_b^0 \pi^+$  candidates, is sufficient to describe the smooth background coming from random track combinations. Additional components are included in the fit to the  $\Xi_b^0 \pi^-$  spectrum to describe partially reconstructed candidates coming from higher mass resonances. These components are referred to as *reflections* in the rest of this Letter. The reflections of the newly observed states in the  $\Xi_b \pi \pi$  mode to the  $\Xi_b \pi$  spectrum are studied with simulation and dedicated components are included in the fit, modeled as Gaussian functions with power-law tails [39]. The means of the reflection components vary freely in the fits to data and their fitted values are consistent with expectations from simulated backgrounds and cross-checks in data. The fit confirms the presence of partially reconstructed  $\Xi_b(6100)^- \rightarrow \Xi_b^{*0}(\Xi_b^0 \pi^0) \pi^-$  decays and shows hints of a contribution from the decay chain  $\Xi_b(1P, 1/2)^- \rightarrow \Xi_b'^0(\Xi_b^0 \pi^0) \pi^-$ , where neither the  $\Xi_b(1P, 1/2)^-$  state, the expected lighter resonance equivalent to that found in the neutral system, nor the  $\Xi_b'^0$  state has been observed experimentally to date (Fig. 2). This component has been validated simulating different mass and width hypotheses for the two particles involved, taking into account expected isospin splittings given the masses of their charged partners. However, a precise estimation of the  $\Xi_b(1P, 1/2)^-$  and  $\Xi_b'^0$  state

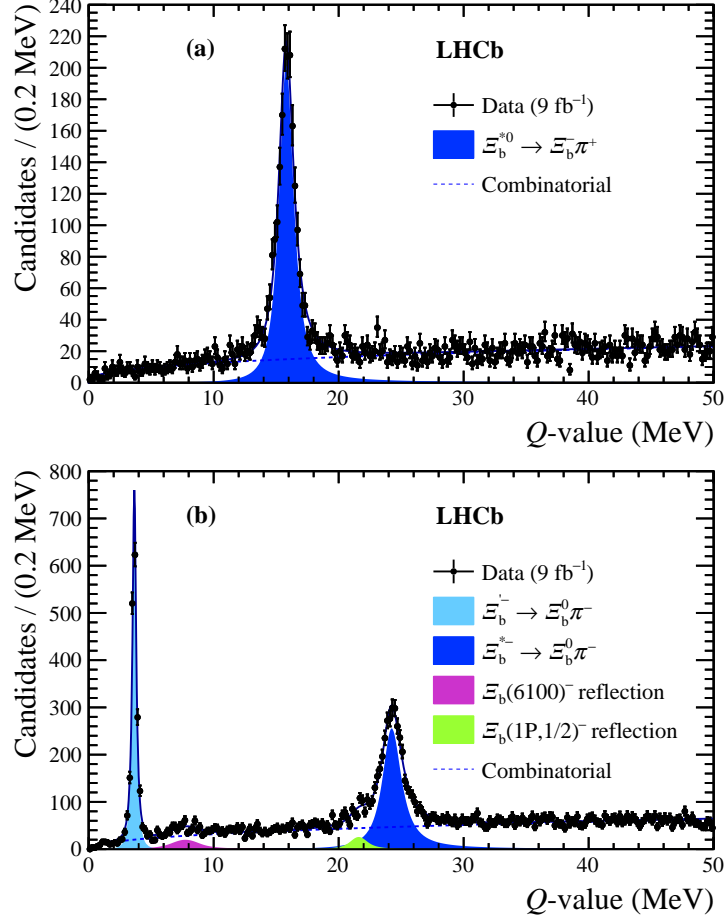


Figure 1: Distributions of the mass difference  $Q = m_{\Xi_b \pi} - m_{\Xi_b} - m_{\pi}$  for selected (a)  $\Xi_b^{*0} \pi^+$  and (b)  $\Xi_b^0 \pi^-$  candidates. The fit results are superimposed.

properties is not possible due to the limited yield and the presence of two unknown mass values. The resolutions of the signal components are fixed to the values obtained from simulation. The fit models are validated with pseudoexperiments and no significant bias is found on any of the parameters of interest.

The fitted yields in the  $\Xi_b \pi$  mass spectra are  $2019 \pm 58$  for the  $\Xi_b^{*0}$  baryon,  $1750 \pm 50$  for the  $\Xi_b^0$  baryon and  $3380 \pm 110$  for the  $\Xi_b^{*-}$  baryon. The  $\Xi_b(6100)^-$  state [18] is confirmed in the  $\Xi_b^{*0} \pi^-$  mass distribution (Fig. 2a), while two new peaks are observed in the  $\Xi_b^{*-} \pi^+$  (Fig. 2b) and  $\Xi_b^{*0} \pi^+$  (Fig. 2c) mass distributions, referred to as  $\Xi_b(6087)^0$  and  $\Xi_b(6095)^0$  in this Letter. For the newly observed states, the signal yields are  $136 \pm 17$  for the  $\Xi_b(6100)^-$  resonance,  $147 \pm 19$  for the  $\Xi_b(6087)^0$  resonance and  $69 \pm 14$  for the  $\Xi_b(6095)^0$  resonance. The three peaks are observed with local significances of  $18\sigma$ ,  $15\sigma$  and  $9\sigma$ , based on the differences in log-likelihood between a fit with the signal yield fixed to zero and the default fit. These significances are reduced to  $12\sigma$ ,  $10\sigma$  and  $8\sigma$ , once systematic uncertainties on the yields are taken into account.

Different sources of systematic uncertainties are considered in the determination of the resonance parameters. All these systematic uncertainties are summarized in Table 1. One of the most important contributions to the mass measurements comes from the knowledge of the momentum scale at LHCb. The associated systematic uncertainty is assigned as the

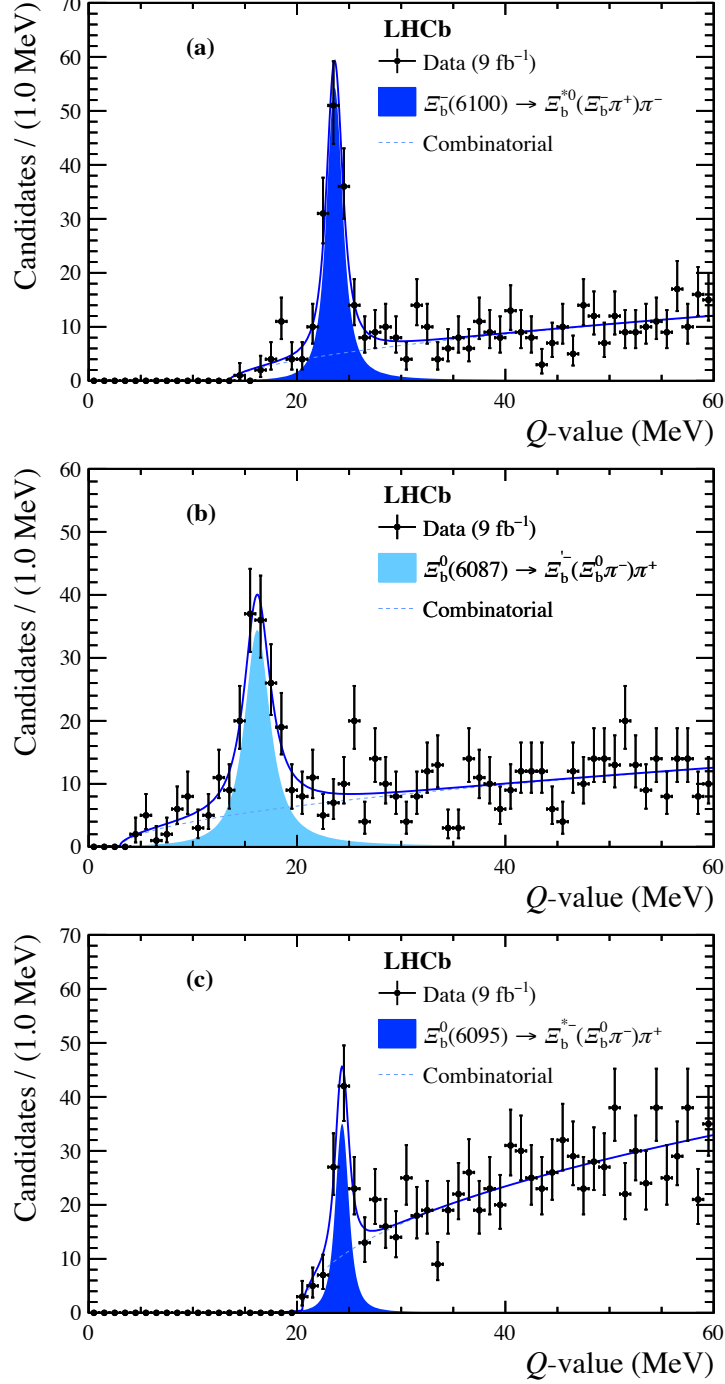


Figure 2: Distributions of the mass difference  $Q = m_{\Xi_b \pi \pi} - m_{\Xi_b} - 2m_\pi$  for selected (a)  $\Xi_b^{*-0} \pi^+ \pi^-$  candidates in the  $\Xi_b^{*0}$  mass window, (b)  $\Xi_b^{*0} \pi^- \pi^+$  candidates in the  $\Xi_b^{*0}$  mass window and (c)  $\Xi_b^{*0} \pi^- \pi^+$  candidates in the  $\Xi_b^{*0}$  mass window. The fit results are superimposed.

larger of the changes in the measured parameters when the momentum scale is changed by its uncertainty, which is estimated to be  $3 \times 10^{-4}$  [40]. An additional uncertainty arises from the empirical description of the background shapes and is estimated by modeling them with alternative functions. A third uncertainty is assigned by varying the description of the reflection components and their properties, obtained either from simulation or

Table 1: Systematic uncertainties (MeV) on the measured physical properties. The parameters  $Q_0$  and  $\Gamma$  are the mean and the width of the Breit–Wigner distribution, respectively.

Source	$\Xi_b^{*0}$		$\Xi_b^{\prime-}$		$\Xi_b^{*-}$	
	$Q_0$	$\Gamma$	$Q_0$	$\Gamma$	$Q_0$	$\Gamma$
Momentum scale	0.006	0.001	0.001	0.001	0.008	0.001
Background	0.003	0.029	0.000	0.006	0.004	0.073
Reflections			0.000	0.000	0.002	0.007
Resolution	0.001	0.038	0.002	0.027	0.000	0.033
BW param.	0.001	0.001	0.000	0.000	0.001	0.002
Total	0.007	0.048	0.002	0.028	0.010	0.081
Source	$\Xi_b(6100)^-$		$\Xi_b(6087)^0$		$\Xi_b(6095)^0$	
	$Q_0$	$\Gamma$	$Q_0$	$\Gamma$	$Q_0$	$\Gamma$
Momentum scale	0.008	0.002	0.007	0.001	0.009	0.006
Background	0.004	0.035	0.022	0.089	0.023	0.025
Resolution	0.004	0.054	0.001	0.035	0.006	0.073
BW param.	0.016	0.050	0.056	0.007	0.001	0.079
Total	0.019	0.081	0.060	0.096	0.026	0.111

from data, where relevant. Further sources of uncertainty on the measurement of the natural widths are included to describe the known differences in resolution between data and simulated events. Differences are expected to be within 5%, based on previous studies [41–44], therefore uncertainties are estimated by varying the resolution function width and the parametrization of the mass resolution function by that amount. Possible uncertainties can arise from the assumed relativistic Breit–Wigner distributions and their parameters. Lower mass states are assumed to decay with angular momentum  $l = 0$ , while higher mass states with  $l = 1$ . For the newly observed states, the hypotheses assuming  $l = 0, 2, 3$  are tested and the largest shifts of the fitted parameters with respect to the default fit are assigned as systematic uncertainties. A further uncertainty on the baryon mass  $m_0$  is assigned due to the limited knowledge of the  $\Xi_b^-$  and  $\Xi_b^0$  baryon masses [30].

The numerical results are summarized in Table 2. The properties of the  $\Xi_b^{*0}$ ,  $\Xi_b^{\prime-}$  and  $\Xi_b^{*-}$  baryons are measured with world-leading precision. For the narrow  $\Xi_b^{\prime-}$  state, its natural width is compatible with zero once the systematic uncertainties are considered, and an upper limit  $< 0.05$  MeV is estimated at 90% confidence level.

In summary, the first observation of two new baryons  $\Xi_b(6087)^0$  and  $\Xi_b(6095)^0$ , with quark content  $bsu$ , is reported in the  $\Xi_b^0\pi^+\pi^-$  final state. Additionally, this Letter confirms the observation of the  $\Xi_b(6100)^-$  charged state by the CMS collaboration [18], with improved significance and sensitivity on its physical parameters. This measurement uses final states with up to nine tracks, most of which are pions, showing excellent performance of the LHCb tracking, reconstruction and PID systems. Finally, the decay mode  $\Xi_b^0 \rightarrow \Xi_c^+\pi^-\pi^+\pi^-$  is observed for the first time. The properties of the  $\Xi_b^{*0}$ ,  $\Xi_b^{\prime-}$  and  $\Xi_b^{*-}$  baryons are measured with high precision. Determination of the spin and parity for the new baryons is not possible given the low signal yields. However, data indicate that the  $\Xi_b(6100)^-$  baryon decays mainly through the  $\Xi_b^{*0}\pi^-$  state, the  $\Xi_b(6087)^0$  baryon mainly through the  $\Xi_b^{\prime-}\pi^+$  state, and the  $\Xi_b(6095)^0$  baryon mainly through the  $\Xi_b^{*-}\pi^+$  state, with no significant contributions to the signals from events outside their

Table 2: Masses and widths of the states considered in this Letter. The first uncertainty is statistical, the second systematic. The third uncertainty on  $m_0$  is due to limited knowledge of the  $\Xi_b^-$  and  $\Xi_b^0$  baryon masses [30].

State	Observ.	Value (MeV)
$\Xi_b(6100)^-$	$Q_0$	$23.6 \pm 0.11 \pm 0.02$
	$\Gamma$	$0.94 \pm 0.30 \pm 0.08$
	$m_0$	$6099.74 \pm 0.11 \pm 0.02 \pm 0.6 (\Xi_b^-)$
$\Xi_b(6087)^0$	$Q_0$	$16.20 \pm 0.20 \pm 0.06$
	$\Gamma$	$2.43 \pm 0.51 \pm 0.10$
	$m_0$	$6087.24 \pm 0.20 \pm 0.06 \pm 0.5 (\Xi_b^0)$
$\Xi_b(6095)^0$	$Q_0$	$24.32 \pm 0.15 \pm 0.03$
	$\Gamma$	$0.50 \pm 0.33 \pm 0.11$
	$m_0$	$6095.36 \pm 0.15 \pm 0.03 \pm 0.5 (\Xi_b^0)$
$\Xi_b^{*0}$	$Q_0$	$15.80 \pm 0.02 \pm 0.01$
	$\Gamma$	$0.87 \pm 0.06 \pm 0.05$
	$m_0$	$5952.37 \pm 0.02 \pm 0.01 \pm 0.6 (\Xi_b^-)$
$\Xi_b'^-$	$Q_0$	$3.66 \pm 0.01 \pm 0.00$
	$\Gamma$	$0.03 \pm 0.01 \pm 0.03$
	$m_0$	$5935.13 \pm 0.01 \pm 0.00 \pm 0.5 (\Xi_b^0)$
$\Xi_b^{*-}$	$Q_0$	$24.27 \pm 0.03 \pm 0.01$
	$\Gamma$	$1.43 \pm 0.08 \pm 0.08$
	$m_0$	$5955.74 \pm 0.03 \pm 0.01 \pm 0.5 (\Xi_b^0)$

respective  $m_{\Xi_b\pi}$  mass windows [34]. These patterns closely resemble those observed in the  $\Xi_c^0$  and  $\Xi_c^+$  baryon systems [30]. An interpretation would be that the new states are  $P$ -wave states ( $l = 1$  between the  $b$  quark and the  $qs$  diquark) coupling to the  $b$  quark to give a pair of states with  $J^P = \frac{1}{2}^-$  and  $\frac{3}{2}^-$ . One might expect the dominant decay mode of the lighter one to be  $\Xi_b'^{(0,-)}\pi$  and for the heavier one  $\Xi_b^{*(0,-)}\pi$ . The lighter  $\Xi_b(1P, 1/2)^-$  state could not be observed as it would likely decay primarily through the intermediate  $\Xi_b'^0$  resonance which is below threshold to decay to  $\Xi_b^-\pi^+$ . However, hints of such  $\Xi_b(1P, 1/2)^- \rightarrow \Xi_b'^0(\Xi_b^0\pi^0)\pi^-$  decay are observed in the  $\Xi_b^0\pi^-$  spectrum as a partially reconstructed feed-down component.

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## Supplemental material

Figure 3 shows the mass distributions of the selected  $\Xi_b^-$  and  $\Xi_b^0$  signal candidates. These distributions are fitted with an asymmetric Gaussian-like function with power-law tails for the signal component, an exponential function for the combinatorial background and an empirical function for partially reconstructed backgrounds. The fitted signal yields are  $12\,020 \pm 140$  for  $\Xi_b^- \rightarrow \Xi_c^0 \pi^-$ ,  $27\,300 \pm 210$  for  $\Xi_b^0 \rightarrow \Xi_c^+ \pi^-$ ,  $5\,680 \pm 100$  for  $\Xi_b^- \rightarrow \Xi_c^0 \pi^- \pi^+ \pi^-$ ,  $11\,690 \pm 180$  for  $\Xi_b^0 \rightarrow \Xi_c^+ \pi^- \pi^+ \pi^-$  decays.

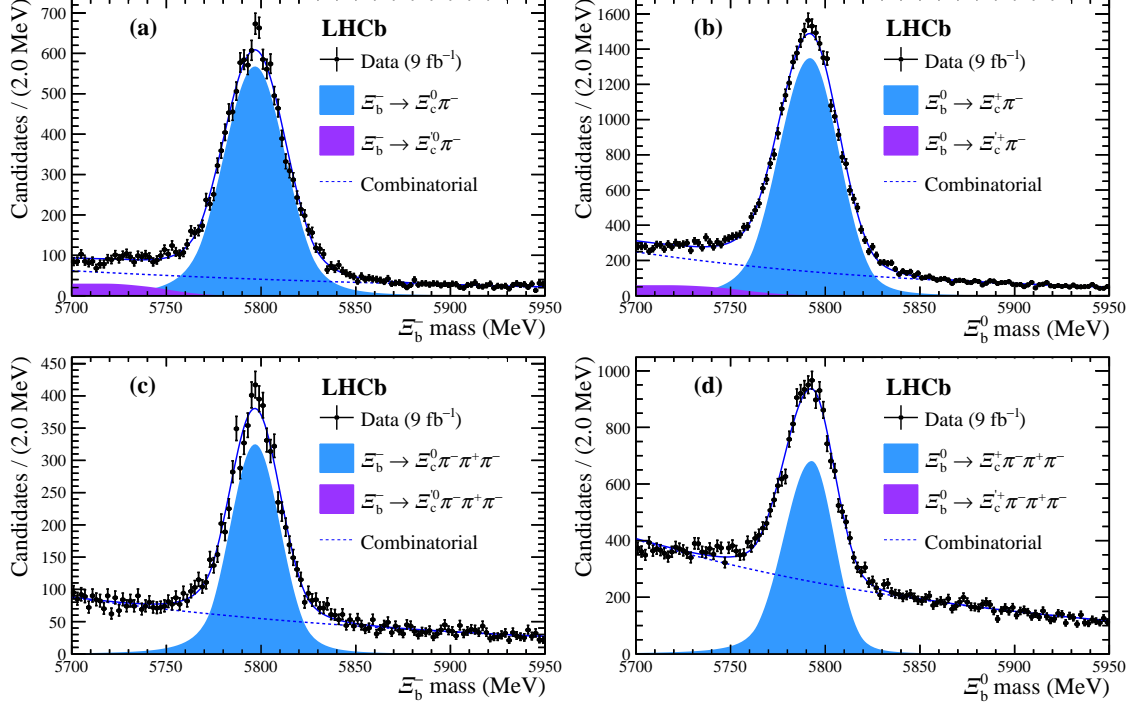


Figure 3: Mass distributions of the (a)  $\Xi_b^- \rightarrow \Xi_c^0 \pi^-$ , (b)  $\Xi_b^0 \rightarrow \Xi_c^+ \pi^-$ , (c)  $\Xi_b^- \rightarrow \Xi_c^0 \pi^- \pi^+ \pi^-$  and (d)  $\Xi_b^0 \rightarrow \Xi_c^+ \pi^- \pi^+ \pi^-$  selected signal candidates.

Figures 4 and 5 show the mass-difference ( $Q = m_{\Xi_b \pi \pi} - m_{\Xi_b} - 2m_\pi$ ) distributions inside and outside the mass windows corresponding to the known  $\Xi_b \pi$  intermediate resonances. Simultaneous fits are performed using the fit models presented in the Letter, with the resonance parameters fixed to their nominal values. Yields are allowed to vary in the fit. The observed suppression of decays not proceeding through the respective intermediate resonances confirms the decay pattern described in the text.

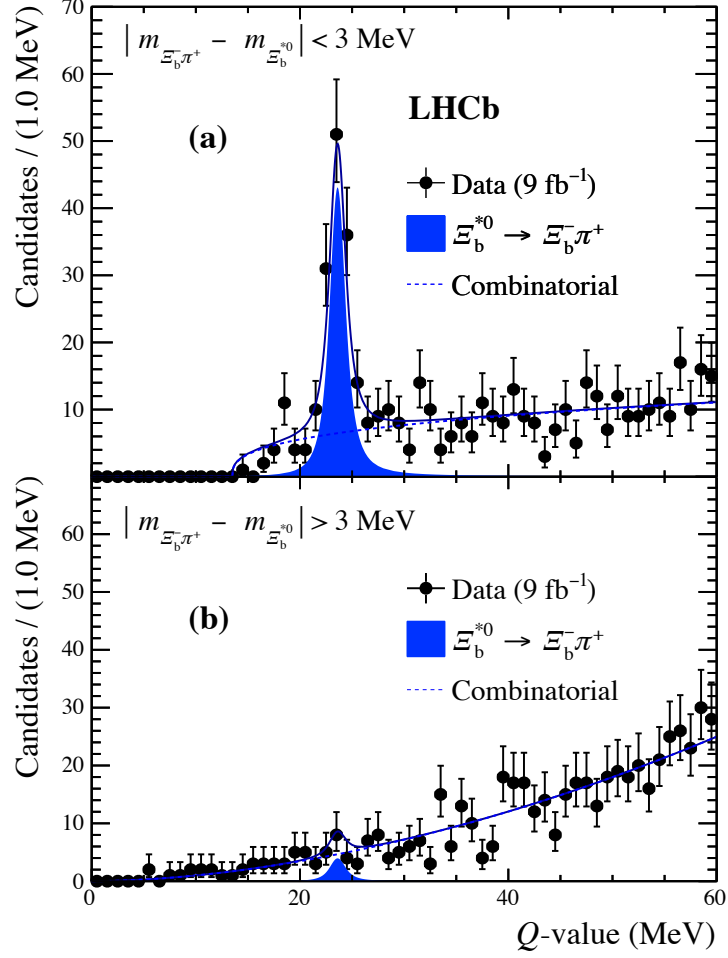


Figure 4: Distributions of the mass difference  $Q = m_{\Xi_b \pi \pi} - m_{\Xi_b} - 2m_\pi$  for selected  $\Xi_b^- \pi^+ \pi^-$  candidates, (a) in the  $\Xi_b^{*0}$  mass window and (b) outside the  $\Xi_b^{*0}$  mass window. The results of the simultaneous fit are superimposed.

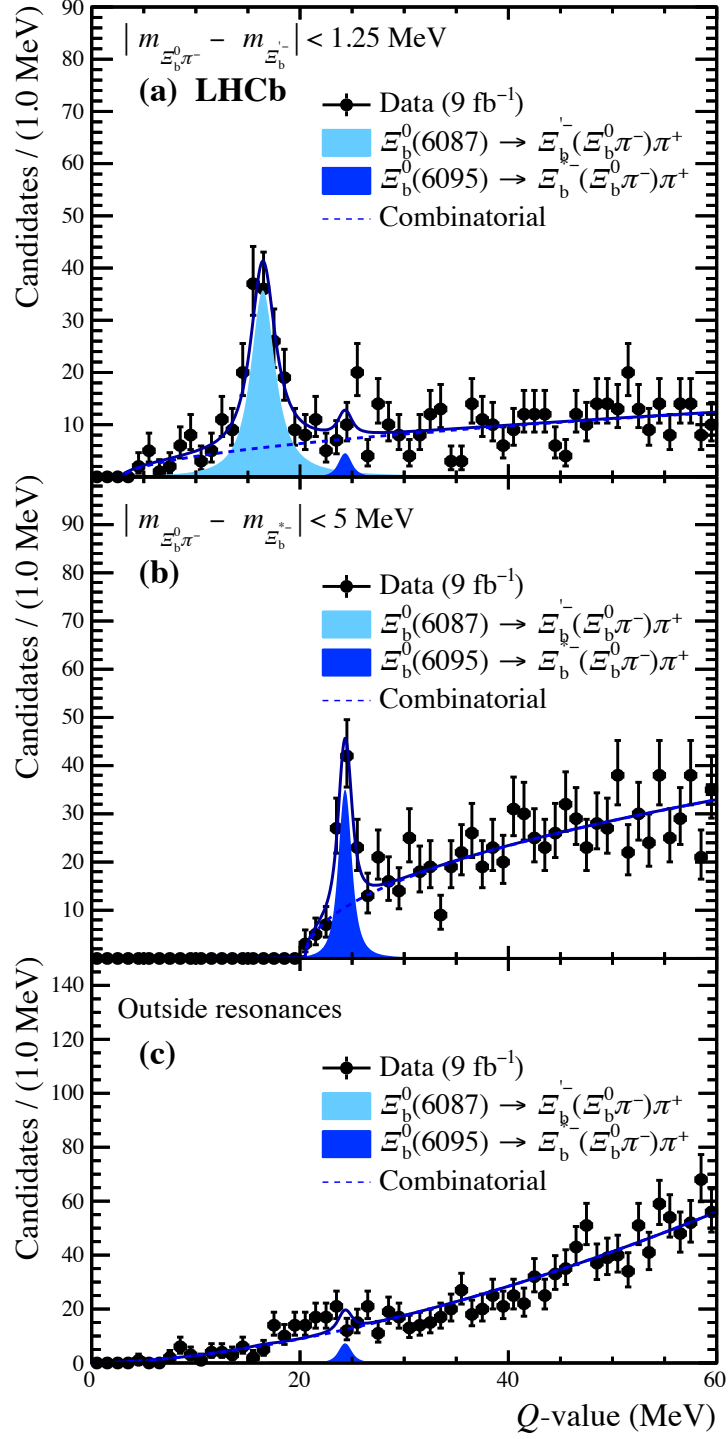


Figure 5: Distributions of the mass difference  $Q = m_{\Xi_b^0 \pi^+ \pi^-} - m_{\Xi_b} - 2m_\pi$  for selected  $\Xi_b^0 \pi^+ \pi^-$  candidates, (a) in the  $\Xi_b'^-$  mass window, (b) in the  $\Xi_b^{*-}$  mass window and (c) outside the mass windows of the intermediate resonances. The results of the simultaneous fit are superimposed.

## LHCb collaboration





































































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