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To cite this article: H. Weisen, P. Blanchard, M. Vallar, A. N. Karpushov, J. Dubray, A. Merle, B. P. Duval, J. Cazabonne, D. Testa, H. Hamac Elaian, and the TCV Team, A. Žohar, L. Snoj, B. Kos, M. Fortuna, A. Čufar, F. Tesse, F. Fontana, C. Gloor, R. Iannarelli, H. Palacios, C. Tille & M. Molteni (2023): TCV Tokamak Neutron Shielding Upgrade for Dual NBI Operation, Fusion Science and Technology, DOI: [10.1080/15361055.2023.2209490](https://doi.org/10.1080/15361055.2023.2209490)

To link to this article: <https://doi.org/10.1080/15361055.2023.2209490>



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Published online: 02 Jun 2023.



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# TCV Tokamak Neutron Shielding Upgrade for Dual NBI Operation

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Received December 15, 2022

Accepted for Publication April 22, 2023

**Abstract** — The Tokamak à Configuration variable (TCV) is equipped with two neutral beam injection (NBI) systems delivering up to 1.2 MW each for pulse durations of up to 2 s. The first system (NB1), designed for an injection energy in the range of 25 to 30 keV has been operational since 2016. The existing concrete neutron shielding of the experimental hall proved insufficient for fully protecting human accessible areas, limiting the number of daily plasma pulses using NB1. The recently commissioned second system (NB2) is designed for injection synergies in the range 50 to 60 keV. Both systems are tangentially oriented in opposite directions in order to permit experiments with low or no net torque.

Calculations with the TRANSP and ORBIS heating codes show that neutron rates from deuterium-deuterium fusion reactions may be as high as  $10^{14}$  n/s, up to 10 times higher than with the lower energy beam only. This is due both to the ~five times larger beam-plasma neutron rates from the higher energy beam and to an exceptionally high contribution from beam-beam reactions between the opposing beams. The radiation protection policy at the Swiss Plasma Center is that all staff members be considered as members of the general public, limiting the daily personal dose to 4 μSv. This is also the maximum admissible daily dose in any publicly accessible zone, whether occupied or not.

Currently, with only the lower energy beam, this limit can be attained in the control room adjacent to the device hall after only five NBI pulses out of a possible 30 daily pulses. To allow for exploitation of the two beams at full specifications, the source side of the existing barite concrete walls of the 15 × 20 × 8 m large TCV hall will be covered with 20-cm-thick polythene (PE) cladding and a ceiling made of 35-cm-thick PE will be added. The total mass of PE will be 200 tons. The usage of PE at this scale for neutron shielding is unprecedented at any fusion research facility.

**Keywords** — Radiation protection, neutron dose, tokamak, safety.

**Note** — Some figures may be in color only in the electronic version.

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## I. INTRODUCTION

The scientific missions of the Tokamak à Configuration (TCV) tokamak<sup>[1]</sup> include researching plasma confinement, in particular concerning the impact of plasma shaping and the physics and control of plasma-wall interactions, as well as the physics of energetic particles and energetic particle modes,<sup>[2]</sup> which will dominate in future fusion reactors because of the high-energy alpha particles produced by fusion reactions.

At the time of construction (1991), TCV was only intended to have electron cyclotron heating (ECH) and electron cyclotron current drive.<sup>[3]</sup> Consequently, most of the radiation exposure was expected to be from X-rays and gamma rays. The 8-m-high walls of the TCV experimental hall are open to the building roof (Fig. 1) and consist of 0.5-m-thick barite concrete. This shielding proved unsatisfactory when, in 2016, the existing ECH system on TCV

was complemented with a neutral beam injector (NBI1) with power of 1.2 MW and 19- to 30-keV injection energy for pulse durations up to 2 s.<sup>[4,5]</sup> This injector has allowed for achieving ion temperatures up to near 3.5 keV, which would have been impossible with ECH alone.<sup>[5,6]</sup> For the purpose of investigating energetic particle modes<sup>[2]</sup> and confinement in plasmas with near-zero momentum input, and hence, low toroidal rotation (as expected for ITER), a second 1-MW neutral beam injection (NBI) system (NBI2) with 50- to 60-keV acceleration energy was commissioned in 2021.<sup>[7]</sup> In order to achieve near-zero NBI torque, the injection direction of NBI2 is opposite to that of NBI1. The higher injection energy and the beam-beam collisions between ions from the opposing sources lead to far higher neutron rates than with NBI1 alone, making a very substantial upgrade of the shielding an essential requirement for the operation of the NBI system.

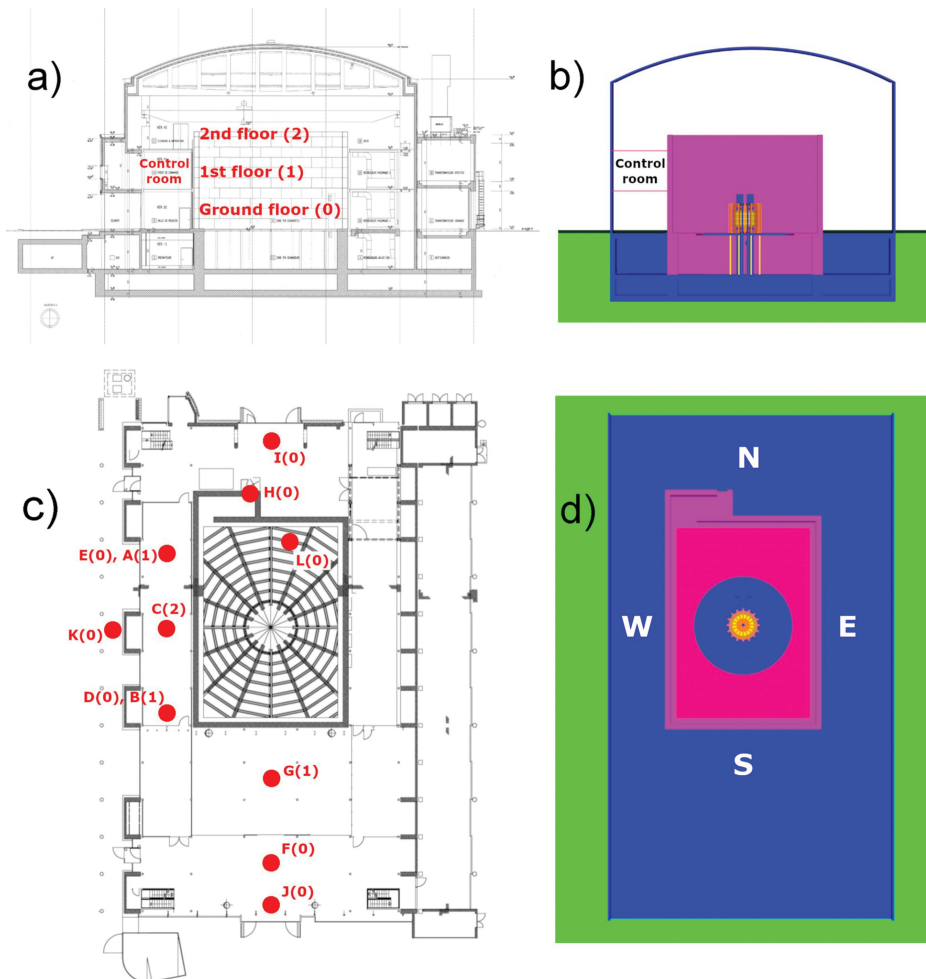


Fig. 1. (a) East-west engineering cross section of the TCV building in its current conditions. (b) East-West cross section of the TCV building in MCNP model. (c) Horizontal cross section of the TCV building with dose measurement positions (identical with MCNP tally locations). The numbers in brackets indicate the floor levels of the tally locations. The TCV device is at the center of the TCV hall on the web-like segmented dismantable ground floor. (d) Horizontal cross section of the MCNP model.

At the relatively modest temperatures of most TCV plasmas heated with deuterium beams, the neutron production of 2.45-MeV D-D fusion neutrons is dominated by beam-plasma reactions. Neutron rate calculations were performed using the heating codes TRANSP-NUBEAM<sup>[8]</sup> and a version of ORBIS<sup>[4]</sup> modified to include beam-beam reactions. They show that with experimentally observed plasma parameters, neutron source rates up to  $10^{13}$  n/s can be obtained with NBI1 only. For NBI1 pulse durations of 1 s and an injection power 1 MW, this can lead to attaining the daily allowed control room dose limit of 4  $\mu$ Sv in as little as 5 discharges out of up to 30 possible discharges per day. As a result of the sharp increase of the cross section for the D-D reaction with energy, the beam-plasma neutron rate for NBI2 is approximately five times higher than for NBI1 for the same power and plasma parameters. A further increase was expected from beam-beam interactions between fast ions from the two injectors and moving in opposite directions.

As the TCV hall currently has no ceiling, a large fraction of neutrons escaping vertically are backscattered from the concrete roof above back down to all locations within the building. In order to allow for the operation of both NBI sources simultaneously with full specifications, the publicly accessible rooms in the TCV building, such as the control room, the neutron dose needs to be reduced by at least two orders of magnitude. A classical shielding solution using a concrete ceiling and thicker concrete walls was excluded because the weight of the concrete would exceed the design specifications for the foundations of the TCV shielding walls. Instead, a solution based on polythene (PE), which is 8.7 times more effective by weight than TCV concrete for slowing down fusion neutrons, was developed (i.e., the PE mass required for equivalent shielding is 8.7 times lower than for TCV concrete). The additional shielding requirements are satisfied by a 35-cm-thick PE ceiling and a 20-cm-thick PE cladding inside the walls of the TCV experimental hall. The specifications for this novel solution for protection from neutron radiation have been established with the help of neutron transport calculations using a state-of-the-art hybrid (deterministic/stochastic) particle transport methodology. It combines the ADVANTG<sup>[9]</sup> code to determine efficient variance reduction parameters based on a rough deterministic transport simulation followed by a high-fidelity continuous-energy stochastic particle transport simulation using MCNP.<sup>[10]</sup>

## II. RADIATION PROTECTION AT THE TCV FACILITY

The maximum daily radiation dose for staff and visitors adopted for the TCV facility is 4  $\mu$ Sv based on Article 22 of the Swiss federal ordinance on radioprotection,<sup>[11]</sup> which specifies a maximum dose of 1 mSv per calendar year for members of the public. The Swiss Plasma Center has made the choice that all staff members should be regarded as members of the public insofar as radiation safety is concerned, obviating the need for individual dosimetry. In order to simplify access management, the 4- $\mu$ Sv daily limit is also applicable to any areas accessible by staff or members or the public and excludes only areas where any human access is prohibited due to electrical hazard, i.e., the high-voltage power supply rooms at the eastern side of the building and the TCV hall itself. This limit applies whether these areas are occupied or not. Ambient dose levels are routinely measured in the control room (position A in Fig. 1) and at the platform supporting the gyrotrons (position G).

Figure 1 shows four cross sections of the TCV building comparing the construction drawings to the simplified model for Monte Carlo simulations. The TCV hall walls are mostly constituted of  $2 \times 1 \times 0.5$ -m sized blocks made from barite concrete. A horizontal cross section is shown in Figs. 1c and 1d. The TCV device is positioned at the center of the hall at a height 1.4 m above the ground level (level 0). Neutron dose measurements are made using three available LUPINE 5401 BF3-NP PSI dosimeters<sup>[12]</sup> at 12 locations using reproducible TCV pulses with NBI1 operation. Gamma doses are measured using a NAUSICAA probe from the same manufacturer.<sup>[13]</sup> These dosimeters are designed to have a response approximating the International Commission on Radiological Protection H\*10 ambient dose equivalent,<sup>[14]</sup> i.e., to be representative of the effect of neutron and gamma radiation on biological tissue.

Radiation levels were measured in 2021 at multiple positions in the building in reproducible discharges with NBI heating using the available NBI1 injector. The positions of the dose measurement and MCNP tally locations are indicated by red dots in Fig. 1c. They are all at a height of 1 m above the respective floor levels. Table I shows the neutron and the gamma-ray doses obtained at the measurement locations in columns 3 and 4, denoted L213 and N227. One of the neutron dosimeters, denoted L214 in Table I, was used as a fixed monitor at position A in the control room. Although these pulses were designed to be reproducible, the pulse-to-pulse variations of the neutron doses were more than 20%, reflecting variations of the TCV neutron source. In order to compensate for these pulse-to-pulse



TABLE I

Neutron and Gamma Dose Measurements in a Series of Reproducible TCV Pulses at Different Tally Locations in the TCV Building

TCV Shot Number	Tally Point Number	L213 (Neutron) ( $\mu\text{Sv}$ ) Mobile	N227 (Gamma) ( $\mu\text{Sv}$ ) Mobile	L214 (Neutron) ( $\mu\text{Sv}$ ) Fixed at A	L213* ( $\mu\text{Sv}$ ) Normalized
70181	A	0.35	0.04	0.35	0.41
70180	B	0.245	0.03	0.32	0.32
70070	C	3.8	0.14	0.458	3.41
70188	D	0.07	0.017	0.39	0.074
70186	E	0.09	0.02	0.395	0.094
70190	F	0.15	0.012	0.39	0.16
70183	G	0.67	0.038	0.385	0.72
70197	H	0.95	0.048	0.44	0.89
70074	I	0.39	0.02	0.47	0.34
70192	J	0.162	0.011	0.4	0.17
70195	K	0.017	0.044	0.42	0.017
70588	L	54.5	0.79	0.52	43.2

variations, the raw measurements using the mobile neutron probe (L213) were normalized to ones obtained with the reference probe such that  $L213^* = L213 / (L214 / \langle L214 \rangle)$  where  $\langle L214 \rangle$  is the mean value of all L214 measurements. The normalized measurements are shown in the column denoted L213\* in Table I.

Outside the TCV hall, the highest neutron dose rates in the building were obtained at level 2 (position C, 3.8  $\mu\text{Sv/s}$ ), at the level 0 entrance labyrinth (position H, 0.95  $\mu\text{Sv/s}$ ), and at the platform supporting the gyrotron of the ECH system (position G, 0.67  $\mu\text{Sv/s}$ ). The lowest dose was measured at the ground floor below the control room (position D, 0.07  $\mu\text{Sv/s}$ ). The control room dose rates (positions A and B) were in the range 0.32 to 0.52  $\mu\text{Sv/s}$ . Prompt gamma doses were typically 4 to 50 times lower than the neutron doses. Because of the high doses at level 2 (position C) and on the gyrotron platform (position G), these areas are currently off limits during NBI operation. NBI operation is terminated for the day when a total dose of 4  $\mu\text{Sv}$  is obtained in the control room, sometimes limiting NBI operation to as few as five pulses.

Exposure to delayed gammas inside the TCV hall due to the activation of machine parts, equipment, and the walls following TCV pulses has so far not been a concern. Measurements of the post-pulse gamma dose show that for every 1  $\mu\text{Sv}$  of total dose measured in the control room, 0.06  $\mu\text{Sv}$  of gamma dose is measured in the immediate vicinity of the tokamak. Half of the gamma dose is measured in the first 2 min after a TCV pulse. Since it takes 2 min to electrically isolate the TCV device from the high-voltage power supplies before entry is permitted, a person spending the whole time between pulses after

those first 2 min would be exposed to a gamma dose of 0.24  $\mu\text{Sv}$  until NBI operation must be discontinued for the day. The higher post-pulse gamma doses expected in the TCV hall when both beams will operate together, however, will lead to access restrictions, which will be defined based on gamma dose rate measurements.

### III. MODELING OF THE TCV NEUTRON SOURCE

A series of experiments using NBI1 was performed in 2020 as a basis for modeling calculations aimed at estimating the future neutron rates using both NBI sources, and hence, the shielding requirements. They included a density scan with volume-averaged electron densities in the range 1.1 to  $5.5 \times 10^{19} \text{ m}^{-3}$ , two power levels, 0.47 and 0.97 MW, and two plasma currents, 220 and 350 kA.

#### III.A. Calculations for Neutron Rates with a Single Neutral Beam

Calculations using the codes TRANSP<sup>[8]</sup> and the much simpler slowing-down orbit code ORBIS<sup>[4]</sup> were performed in 2020 using NBI1 operating at energies in the range 19 to 25 keV depending on the power requested. A NUCSAFE monitor<sup>[15]</sup> was used for time-resolved, albeit uncalibrated, neutron rate measurements. In the implementation used, ORBIS takes account only of shine-through and first-orbit losses, but includes a calculation of the beam-beam neutron rates. The electron temperatures and densities for these calculations

were obtained from laser Thomson scattering.<sup>[16]</sup> There was no available ion temperature measurement. For use in TRANSP, the ion temperatures were estimated using the ASTRA code<sup>[17]</sup> and an ad hoc transport model with  $T_i/T_e$  ranging from 0.8 to 2.4 depending on conditions.

For ORBIS,  $T_i/T_e = 2$  was assumed in all cases, as this temperature ratio was the highest previously measured by charge exchange spectroscopy.<sup>[5]</sup> The ion temperature had a modest effect, as thermal-thermal reactions were calculated to be always below 11% of the beam-thermal neutron rate. Plasma rotation in the direction of the injection slightly reduces the beam-thermal neutron rates. As no rotation measurement was available, zero rotation was assumed. TRANSP includes orbit losses during slowing down, and most importantly, charge exchange losses calculated on the basis of the measured vessel pressure. As a result of the neglect of these losses in ORBIS, beam-thermal neutron rates by ORBIS were on average some 20% higher for NBI1 than those calculated by TRANSP.

The largest differences, up to 45% of the TRANSP prediction, were seen at the lowest densities where neutrals penetrate farthest into the plasma. At the highest densities, where charge exchange losses are least important, ORBIS and TRANSP agree within 10%. The highest neutron rates, near  $10^{13}$  n/s, were predicted by TRANSP at the highest power (0.97 MW) and plasma current (350 kA) at medium volume average density ( $\sim 3 \times 10^{19} \text{ m}^{-3}$ ). Figure 2 shows that there is a high degree of proportionality between the neutron rates from the uncalibrated NUCSAFE neutron detector<sup>[15]</sup> and the TRANSP predictions. The higher neutron rates obtained at the higher current level (350 kA) are explained by the better orbit confinement and the higher electron temperatures obtained at the higher current.

### III.B. Extrapolation of the Neutron Rates to Operation with Two Neutral Beams

The next step was to repeat the modeling, assuming that both NBI1 and NBI2 were operating simultaneously at a power of 1 MW each, with NBI1 counter-injecting (opposite to the direction of the plasma current) with 25-keV acceleration voltage and NBI2 co-injecting (same direction as the plasma current) at 50-keV acceleration voltage. This modeling was not self-consistent, as it was assumed that the densities and temperatures were the same as with NBI1 only. From basic confinement scaling ( $W_p \propto P_{\text{tot}}^{0.5}$ ) and neglecting the ohmic power, the electron temperature may rise by up to  $\sim 40\%$  when the NBI power is doubled provided the density remains constant. In our experience, however, the density rises and the temperature remains roughly constant when beam power is applied or increased.

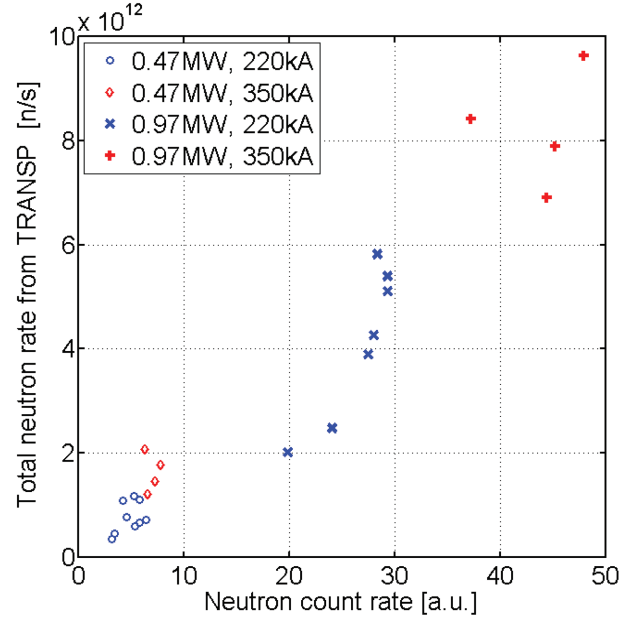


Fig. 2. Neutron rates calculated by TRANSP versus the uncalibrated neutron count rate measured using the NUCSAFE 1 detector. The symbols refer to different NBI1 power levels and plasma currents.

As a result, the dominant beam-plasma neutron rate, which depends mainly on electron temperature and not on density, is adequately predicted with our assumptions.

The predictions for the two codes are shown in Fig. 3. Only two cases at 1 MW, those which produced the highest neutron rates in the experiment, were modeled with TRANSP (triangles). Due to the absence of charge exchange neutrals, ORBIS predictions for the neutron rates in these two cases are higher by a factor of 2. The TRANSP and ORBIS calculations for the 350-kA case with  $\langle n_e \rangle \approx 3 \times 10^{19} \text{ m}^{-3}$  show that neutron rates with both beams are respectively 5 times and 10 times above those with NBI1 only. The neutron rates include a substantial contribution from beam-beam reactions from reactions between the ion populations from NBI1 and NBI2, which at birth have relative energies up to 145 keV (assuming injection energies of 25 and 50 keV). In the case of the ORBIS calculation, the beam-beam contribution was about half of the total neutron rate for  $\langle n_e \rangle \approx 3 \times 10^{19} \text{ m}^{-3}$ . Initial short pulse ( $\sim 50$ -ms) operation in 2021 has confirmed that the neutron rates with both beams can exceed those from NBI1 by an order of magnitude.

There is a clear inverse relation of ORBIS predictions with density (Fig. 3) as the fast ion densities scale as  $n_f \propto \tau_{\text{slow}} \propto n_e T_e^{3/2}/n_e$ , and hence, for the beam-beam neutron rate  $R_{\text{nbb}} \propto n_f^2 \propto P_{\text{NBI1}} P_{\text{NBI2}} T_e^3/n_e^2$ .  $T_e$  is seen to scale weakly inversely with density in these discharges,

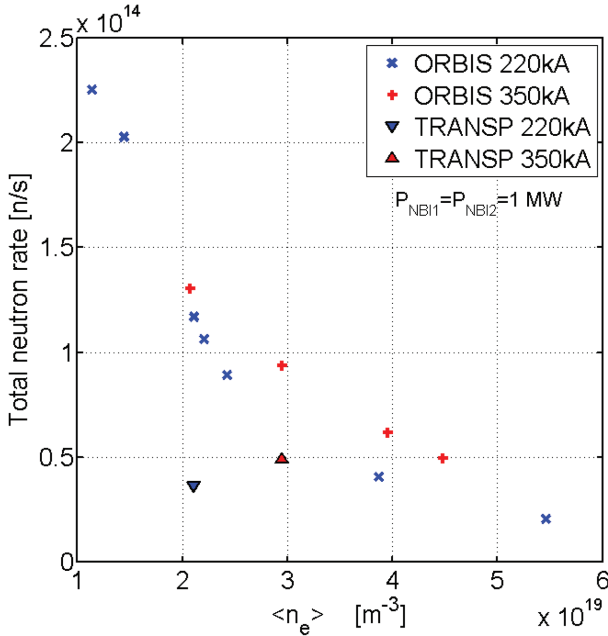


Fig. 3. Predictions for neutron rates with NBI1 and NBI2 together for  $P_{NBI1} = P_{NBI2} = 1$  MW.

$T_e \propto n_e^{-1/3}$ , hence  $R_{nbb} \propto P_{NBI1} P_{NBI2} n_e^{-3}$ . This simple scaling does not take into account the fact that the beam deposition profile depends on plasma density. The ORBIS predictions at low to medium density ( $\langle n_e \rangle < 3 \times 10^{19} \text{ m}^{-3}$ ) are clearly overestimates for current conditions in TCV, but may be relevant if neutral densities in the plasmas, and hence, charge exchange losses, can be significantly lowered, e.g., by wall conditioning.

### III.C. Assessment of Shielding Requirements

We conservatively assume, based on the ORBIS predictions for  $\langle n_e \rangle \approx 3 \times 10^{19} \text{ m}^{-3}$  and  $I_p = 350 \text{ kA}$ , shown in Fig. 3, that the neutron rates can be up to  $10^{14} \text{ n/s}$  when both NBIs are operated together, i.e., 10 times higher than with NBI1 alone. We also assume that the pulse duration, now typically 1 s, can be as long as 2 s. We use the current worst conditions for NBI1 operation, which limit the daily operation to 5 pulses (out of a technically possible 30 pulses) as a scaling basis. In order to remain within the daily radiation dose limit for 30 consecutive pulses with the highest expected neutron rates, for the longest possible duration, and for 30 pulses per day, the shielding needs to reduce the radiation dose by a factor 120. We note that it is very unlikely that 30 worst-case pulses would actually be performed on a single experimental day and much less so that such worst-case experimental days would be performed on every day of the year.

## IV. SHIELDING MATERIAL STUDIES

First, several types of concrete as well as borated and unborated PE were assessed for their suitability as materials for additional shielding. The types of concrete included the barite concrete used for the existing side walls of the TCV hall, ordinary concrete, and a handful of promising concrete types selected from the Pacific Northwest National Laboratory (PNNL) Compendium of Material Composition Data for Radiation Transport Modeling<sup>[18]</sup> for their high hydrogen and boron content. The TCV concrete has a density of  $3.4 \times 10^3 \text{ kg/m}^3$  and contains 50.8% barium by weight. By comparison, ordinary concrete has a density of  $2.4 \times 10^3 \text{ kg/m}^3$ . Two of the concrete types in the PNNL Compendium<sup>[18]</sup> have a modestly higher neutron moderating power than the TCV concrete for a similar mass density (numbers 87 and 88 in the Compendium). However, a concrete based solution, which would have implied doubling the wall thicknesses and building a 70- to 90-cm-thick concrete ceiling had to be abandoned because the foundations of the TCV hall walls were only specified for barite concrete walls up to a 1-m thickness but no ceiling. An additional constraint was the capacity of the overhead crane in the TCV building, which is limited to 10 tons and is therefore too low for a demountable ceiling made of 15-m-long heavy concrete beams of practical cross section.

### IV.A. Spherical Model for Material Comparisons

In order to compare the dose attenuation properties of different materials, a simple spherical model with the source inside was devised. The simulation was performed with a hollow sphere with different wall thicknesses and a doughnut-shaped neutron source similar to a TCV plasma. For each of the studied wall configurations, a MCNP simulation using MCNP5 v1.60<sup>[10]</sup> together with the FENDL 3.1d<sup>[19,20]</sup> nuclear data library was performed using  $10^8$  source particles to ensure the statistical uncertainty of the results was below 1%. Figure 4 shows the relative flux averaged over the surface of the sphere per 2.45-MeV source neutron penetrating shieldings of various thicknesses of some of the materials that were under consideration. We can see that for neutron shielding, the PE-based solutions performed vastly better than concrete. PE is one of the most effective solid materials for neutron moderation by weight that is available in bulk. A 20-cm thickness of PE has a neutron

moderating power equivalent to 50 cm of TCV barite concrete for an 8.7 times lower weight (materials only, without iron reinforcements or support structures).

Two of the PE based shields calculated were laminates with the last 2.5, respectively 5 cm of ordinary PE replaced by PE doped with 5% of boron (labeled SWX 201 in Fig. 4).<sup>[21]</sup> Boron has an exceptionally high neutron capture cross section at energies below 1 eV. Figure 4 shows that the two laminates perform better than pure PE and doped PE. It also shows that 20 cm of PE is equivalent in neutron attenuation to 50 cm of barite concrete.

Figure 5 shows the dose rates corresponding to the fluxes in Fig. 4. The ambient dose equivalent is defined as  $H^*(10) = \int \Phi h^*(10) dE$ , where the integral extends over the entire spectrum.<sup>[14]</sup> The spectral weighting factor,  $h^*(10)$ , represents the energy-dependent detrimental biological effects of neutron radiation. The presence of boron in PE makes virtually no difference to the neutron dose rate, as the neutrons removed by capture by the boron are all of low energy ( $<1$  eV), and hence contribute little to the  $H^*(10)$  ambient dose equivalent. This can be understood by inspecting the neutron spectra transmitted through the shields shown in Fig. 6. The vertical axis is the flux per unit lethargy  $f(u) = I(E)$ , where  $\Phi$  is the flux in neutrons/s/cm and  $u = \ln(E_0/E)$  is called the neutron lethargy, and

$E_0 = 2.45$  MeV is the neutron energy at the source. The unit lethargy spectral interval is an interval of energy  $[E_1 E_2]$  such that  $E_2/E_1 = e$ . We see that above a few electron-volts, all spectra are similar and the spectra of the three PE-based shields are essentially identical. The two

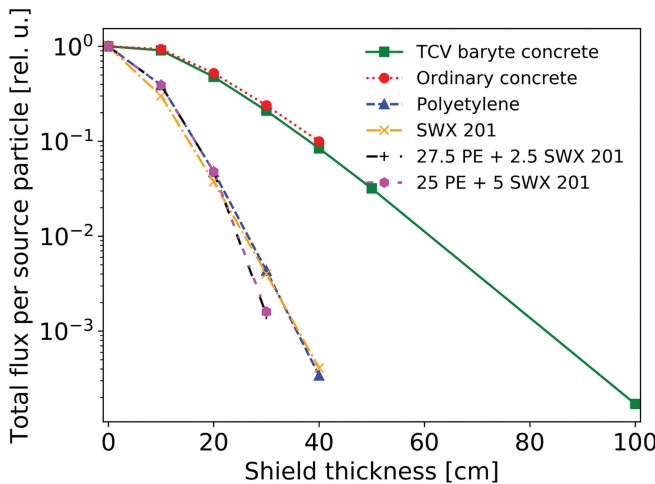


Fig. 4. Neutron flux per source particle in dependence of the position in the shield for different PE and SWX laminates and two types of concrete, calculated using the simple spherical model.

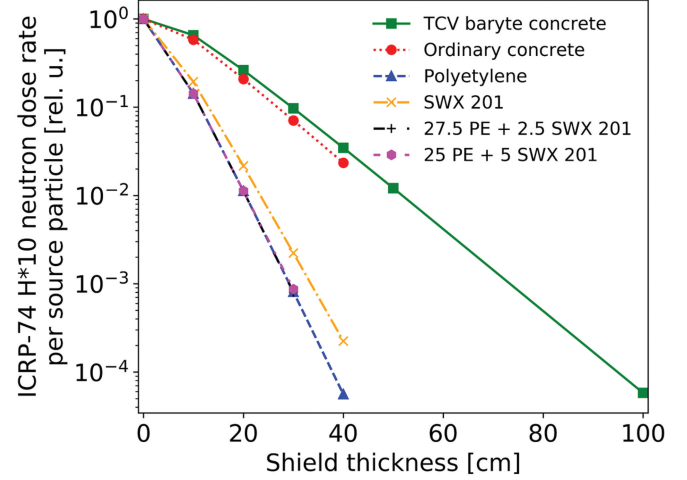


Fig. 5. Neutron  $H^*(10)$  ambient dose equivalent per source particle in dependence of the position for different PE and SWX laminates and two types of concrete calculated using the simple spherical model.

boron laminates remove the large spectral peak in the thermal region (from  $\sim 10^{-2}$  to  $\sim 10^{-1}$  eV). This narrow region has very little weight in comparison with the wide high-energy part of the spectrum. In addition to the much greater spectral width of the high-energy part, the spectral weighting factor  $h^*(10)$  rises sharply for  $E > 1$  keV and is  $\sim 30$  times larger for  $0.7 \text{ MeV} < E < 5 \text{ MeV}$  than in the thermal range.

This means that even after substantial attenuation, the ambient dose from a fusion neutron source remains dominated by energetic neutrons and the shielding needs to be optimized for fast, rather than for thermal neutrons.

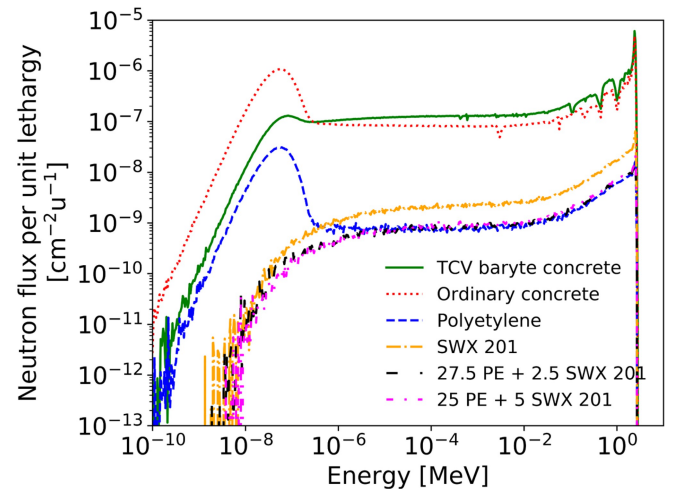


Fig. 6. Neutron lethargy spectra in 640 energy groups after 30 cm for different PE and SWX laminates and two types of concrete calculated using the simple spherical model. The spectra are given in units per source particle.



However, when large dose attenuation factors are required, as is the case with TCV, gamma-ray transmission and gamma-ray emission following neutron capture in the shielding materials have to be accounted for too. The advantage of boron-doped PE is that most of the captured neutrons lead to the emission of a gamma photon with 0.5-MeV energy, whereas capture by hydrogen, the main capture process in pure PE, leads to the emission of a 2.1-MeV photon, which is more difficult to stop and is associated with a higher dose. The gamma-ray contribution was evaluated in the full building model described in [Sec. VI](#).

#### IV.B. Experimental Characterization of PE as a Shielding Material

Before committing several million Swiss francs to this project, the neutron dose attenuation by PE was experimentally tested for thicknesses up to 30 cm using a set of six nestable PE boxes made from 5-cm high-density PE 500 sheets.<sup>[22]</sup> The four innermost boxes are depicted in [Fig. 7](#) without the 5-cm-thick PE lids. A mobile dosimeter was placed inside the innermost box, and the boxes were closed with the lids for the measurements. The boxes were placed in the TCV hall with one side facing the TCV device at a distance of 7 m. Measurements were made for different total thicknesses of PE in a series of reproducible TCV discharges with NB11 and normalized to measurements by a reference detector. [Figure 8](#) shows the measured attenuation together, for comparison, with the MCNP predictions



Fig. 7. Photograph of the four innermost PE boxes (without their lids) used for experimental measurements of neutron and gamma-ray dose attenuation.

for the spherical model, demonstrating that a dose reduction by a factor  $>3000$  was obtained with 30 cm of PE.

This result validated the choice of PE as a neutron shielding material and exceeded the expectations based on the spherical model. The nearly exponential decrease of the dose with thickness may be expressed by an effective dose attenuation coefficient of PE, which amounts to  $29 \text{ m}^{-1}$  for the experimental result and  $24 \text{ m}^{-1}$  for the calculations with the spherical model. The former is likely an overestimate due to the geometry of the boxes because detected neutrons incident onto the sides of the boxes experience a higher effective thickness.

The experiment was repeated for a PE thickness of 20 cm using the NAUSICA gamma dosimeter. Prompt gamma emission was detected without the PE box, originating from neutron capture in the walls and equipment in the TCV hall. The gamma dose measured inside the box was only a factor 1.77 lower than without the box. We conjecture that the gammas detected consisted in part of attenuated gammas incident onto the assembly and of gammas produced by neutron capture in the PE. Future MCNP calculations will determine what fraction of the gamma dose inside the box is due to neutron capture in PE.

#### V. DESIGN AND MODELING OF THE TCV RADIATION SHIELDING

The MCNP shielding models and the engineering design progressed hand in hand, with the latter being subjected to many supplementary constraints. A detailed

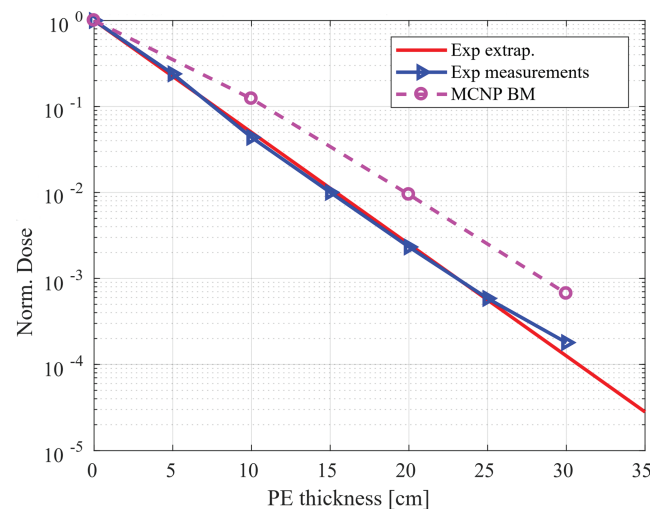


Fig. 8. Measured normalized dose versus shielding thickness (blue triangles), lin-log fit, and expectations from the spherical MCNP model (broken lines).



description of the final MCNP model is available in Fortuna et al.<sup>[23]</sup> These include the weight limit (10 tons) of the available overhead crane, the requirement to use nonmagnetic steel support beams for supporting the PE ceiling, room for a light (1-ton) overhead crane below the PE ceiling, accessibility, fire safety, ventilation, and safety requirements for the installation and removal of the ceiling elements. A basic previous neutronics model of the building, provided in 2016 by Culham Centre for Fusion Energy on a commercial basis, was upgraded in steps to include a model of the TCV device with a steel vessel and copper coils to provide a better representation of the neutron and gamma spectra produced by the device. Some of the major structures in the TCV hall, the TCV hall floor (partly made of wood and partly of concrete) and of the 10-cm-thick wooden floors between levels 0, 1 (control room), and 2 (storage area) were included for better realism. More recently, far more details of the building, including the equipment in the vicinity of the TCV device and the structure of the roof, were included in the model. This provided a substantial improvement between the predicted and measured doses throughout the building.<sup>[23]</sup>

Initial shielding designs had all or part of the wall PE shielding on the outside of the wall for simplicity and accessibility (requiring less equipment to be removed from the inside walls). The MCNP<sup>[10]</sup> calculation showed that this is satisfactory for neutron dose attenuation, but not for the gamma doses. The current design, therefore, has all of the PE wall shielding adjacent to zones accessible during TCV pulses on

the inside, allowing for most of the gammas to be stopped by the barite concrete walls. The height of the concrete wall had to be reduced by 1.5 m to allow access with the overhead crane because of the extra thickness of the ceiling and supporting structure. The concrete blocks recovered will be reused as additional shielding at the lower part of the hall, up to a height of 3 m in parts of the hall, as their neutron shielding is equivalent to that of 20 cm of PE (PE will be used above the concrete blocks). The hall is topped by a 35-cm-thick PE ceiling and no further gamma shielding. The PE-covered areas will total 800 m<sup>2</sup> for a total weight of 200 tons. Construction started in February 2023 and is expected to take several months.

Figure 9 shows a north-south cross section of the junction of the ceiling and the supporting concrete walls at the south of the hall. The PE slabs rest on eleven 50-cm-high beams spaced 2 m apart, which themselves rest on a pair of 30-cm-high beams placed on 20-cm-thick PE above the east and west walls. Radiation shielding in the horizontal direction is provided by a total of 30 cm of PE and a parapet made of 20 cm of barite concrete. The latter is essential to stop the majority of the gammas produced by neutron capture inside the PE. An equivalent arrangement, from the point of view of neutron and gamma shielding, exists at the junctions of all four walls with the ceiling.

The hall entrance at the northwest corner will be widened to 3 m and equipped with a custom-made sliding door consisting of two 17.5-cm-thick boron-doped PE

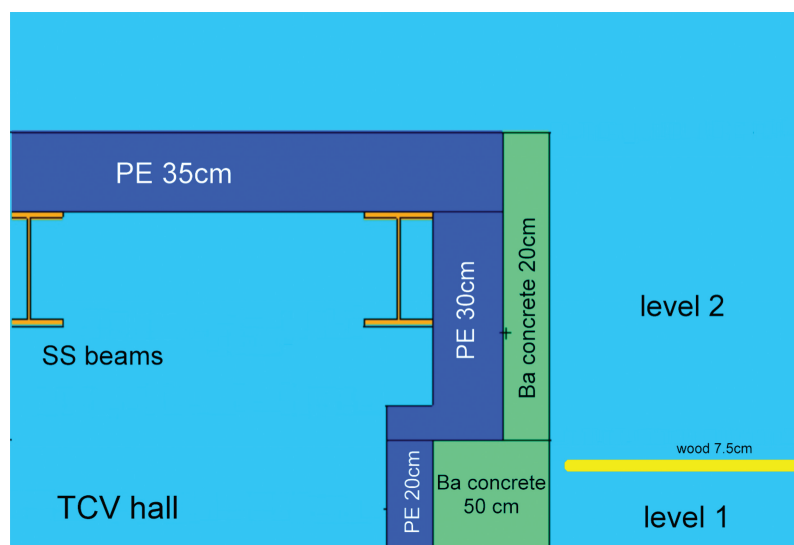


Fig. 9. Junction of ceiling and the 50-cm-thick south wall of the TCV experimental hall (north-south cut through the middle of the hall). The 20-cm-thick barite concrete side wall (parapet) at level 2 is necessary for the reduction of gamma doses at levels 1 and 2 produced by neutron capture in the PE neutron shielding at level 2.

sheets. The remainder of the entrance above a height of 3 m and up to the ceiling will be made of vertically stacked boron-doped PE slabs totaling a thickness of 35 cm that must be light in order to be easily removable for transferring large pieces of equipment. Boron-doped PE was chosen for these locations in order to reduce gamma-ray doses, as, because of weight restrictions, there is no additional gamma-ray shielding by high-atomic-mass materials. At level -1, the barite walls are 1 m thick and no PE cladding is needed. However, there are numerous passages in the east and west walls, approximately  $50 \times 50$  cm wide for power and signal cables. All of these will be plugged with a thickness of at least 40 cm of PE or filled with PE beads. Access to the TCV hall at level -1 is granted by a sliding boron-doped PE door similar to the one at level 0. The ceiling is made from two overlapping layers of PE slabs, 17.5 cm thick, 6 m long, and 2 m wide and will be supported by 15-m-long steel beams spaced by 2 m and oriented in the east-west direction. As nonmagnetic steel beams are not standard products, they will be made by welding austenitic type 1.4307 steel sheets together with cross sections equivalent to that of standard beams.<sup>[24]</sup>

## VI. MCNP VALIDATION OF THE ENGINEERING DESIGN

The MCNP engineering model includes, in addition to the detailed features of the building and the contents of the TCV hall, all essential features of the engineering design, i.e., the structural elements such as the beams and realistic models of the sliding doors. All simulations were performed with MCNP5 v1.60<sup>[10]</sup> together with the FENDL 3.1d<sup>[19,20]</sup> nuclear data library and ADVANTG 3.2.1<sup>[9]</sup> for production of variance reduction parameters for tally position located behind different walls of the TCV building. The neutron source intensity used in the simulations was  $4 \times 10^{13}$  neutrons per second. The east-west cuts of the calculated neutron dose rate fields for unshielded and shielded conditions are shown in Figs. 10 and 11. The corresponding gamma radiation fields are shown in Figs. 12 and 13.

An initial calculation without the concrete parapet shown in Fig. 9 established that neutron doses were at or below the design targets throughout the TCV building. However, gamma doses at level 2 (tally point C in Fig. 1), at the gyrotron platform (G), and in the control room (A and B) substantially exceeded the design objectives. At tally point C, the gamma dose exceeded the design target by a factor 11.

In a subsequent set of calculations, the origin of the gammas was determined by defining particular parts of

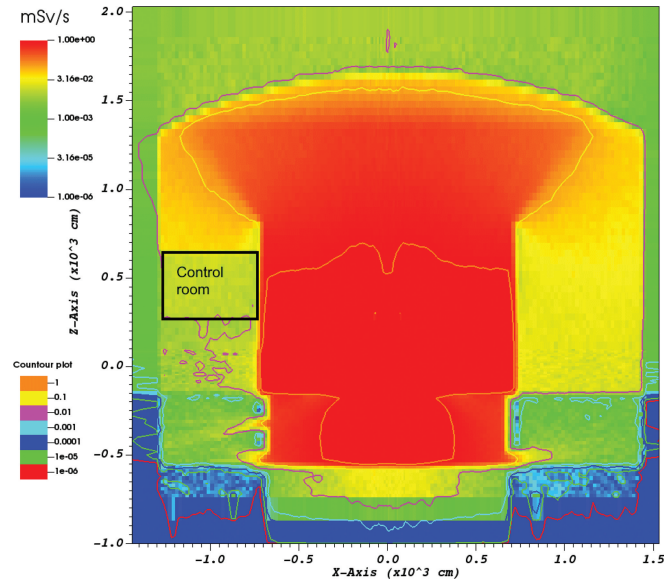


Fig. 10. Neutron dose rate distribution for an east-west cut through the TCV building for a MCNP model reflecting current conditions.

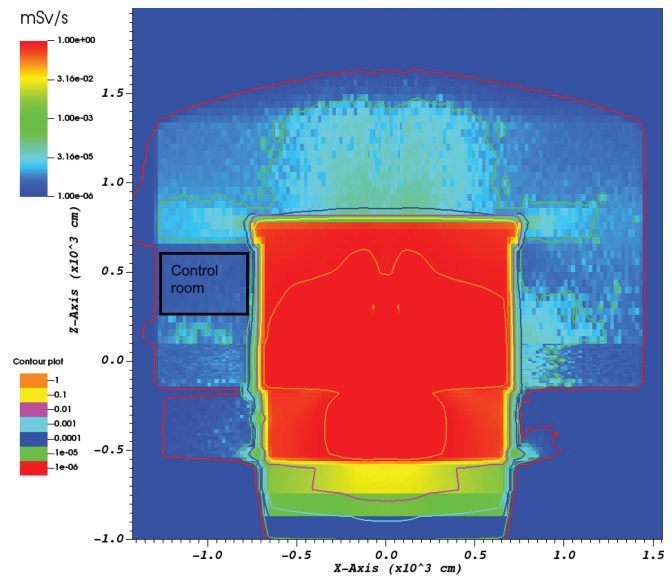


Fig. 11. Neutron dose rate distribution for an east-west cut through the TCV building for a MCNP model with proposed shielding.

the PE shielding suspected of being major contributors to the gamma-ray doses in the inhabited zones as “cells.” The gamma-ray “importance” of the cells was set to 0, meaning that all gamma rays entering or generated in the cell would be absorbed. In this way, combinations of calculations using different cells with importance set to 0 were used to determine the

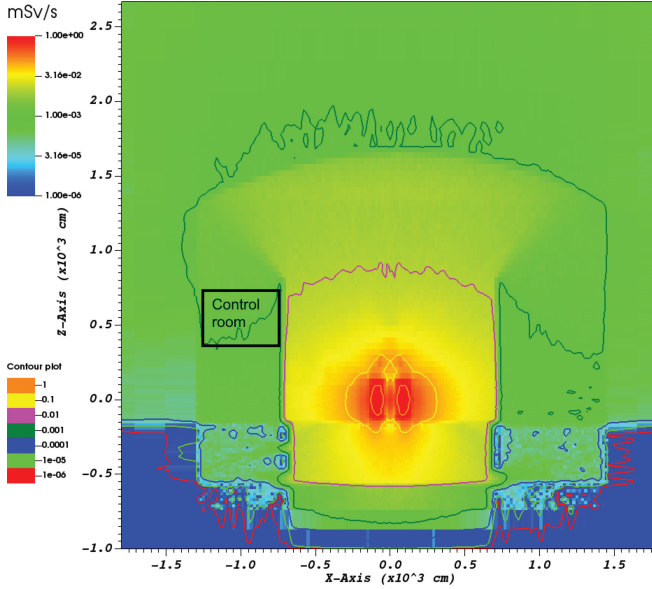


Fig. 12. Gamma dose rate distribution for an east-west cut through the TCV building for a MCNP model reflecting current conditions.

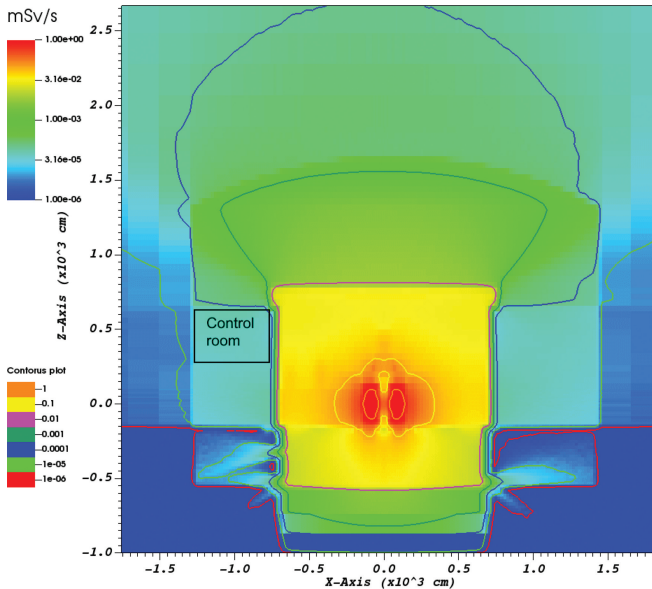


Fig. 13. Gamma dose rate distribution for an east-west cut through the TCV building for a MCNP model with proposed shielding.

contributions of the various parts of PE shielding to the dose rates at the tally points. The results established that the majority of the gammas contributing to the doses at the tally locations were indeed emitted by the PE side walls at level 2. This prompted the addition of a 20-cm-thick barite concrete parapet at level 2, as shown in

Fig. 9. The calculations also showed that the contribution of gammas emitted by or reflected from the ceiling is low, obviating the need for gamma-ray shielding on top of the ceiling.

Tables II and III show the expected dose reductions for neutrons and gammas, respectively, at the various tally points for a source rate of  $4 \times 10^{13}$  n/s. The columns denoted “Reference Model” refer to the model of the TCV building without additional shielding. The results of the calculations were subjected to a normalization to the experiment dose measurements for evaluating the highest expected daily dose expected for 30 pulses of a duration of 2 s and a neutron yield 10 times higher than with NBI1 only. The magnitude of the source rate therefore has no influence on the calculation of the expected highest daily dose.

The relative uncertainties for the doses were determined by the Monte Carlo statistics. We see that in the control room (A and B) and the gyrotron platform (G), the target dose reduction is met, with a total highest expected daily dose near or below  $4 \mu\text{Sv}$ . At level 2 (tally point C), however, the highest expected total daily dose (neutrons + gammas) is  $11.2 \mu\text{Sv}$ , i.e., almost a factor of 3 above the maximum allowed daily dose, implying the necessity of access management. This is deemed tolerable, however, as level 2 is a storage area requiring only occasional access and the achievement of highest expected daily doses is considered to be exceptional.

## VII. CONCLUSIONS

The current work has established that a retrofit PE shielding of the TCV hall satisfies the stringent requirements of a total neutron and gamma dose reduction by a factor 120. We emphasize that this is a conservative figure and we regard it as very unlikely that TCV would regularly produce daily neutron doses requiring such a high level of additional protection. While the use of PE allows for weight savings of a factor of 8.7 compared to barite concrete for neutron shielding, it cannot be used on its own if it is in direct view of human occupied zones because neutron capture by hydrogen in the material leads to significant gamma emission. Gamma emission must be stopped by a high-atomic-charge material such as barium. This required the addition of a 20-cm-thick parapet made from barite concrete on top of the side walls where naked PE would otherwise be in direct

TABLE II

Calculated Neutron H\*10 Dose Rates for the Engineering Design for a Source Rate of  $4 \times 10^{13}$  n/s

Location		Reference Model		Reference + PE Shield Model		Relative Dose Reduction	Highest Daily Dose ( $\mu$ Sv)
		Dose (mSv/s)	Relative Uncertainty (%)	Dose (mSv/s)	Relative Uncertainty (%)		
Control room, north	A	1.03E-02	3.0	5.86E-06	4.1	1757.82	0.35
Control room, south	B	9.94E-03	2.9	3.26E-06	5.3	3046.84	0.20
Second floor	C	5.37E-02	1.3	2.42E-05	1.8	2219.83	1.45
Workshop	D	5.18E-03	4.0	4.69E-06	4.6	1105.37	0.28
Diagnostic lab	E	5.66E-03	4.3	1.27E-05	3.0	446.48	0.76
South extension	F	5.30E-03	4.0	1.68E-06	7.5	3146.45	0.10
Gyrotron platform	G	1.62E-02	2.3	3.11E-06	5.6	5213.46	0.19
TCV entrance turnstile	H	1.07E-02	3.8	1.26E-05	3.8	849.57	0.76
North door	I	8.98E-03	3.2	3.61E-06	5.7	2487.87	0.22
South door	J	5.04E-03	4.1	1.46E-06	8.6	3461.35	0.09
West side outside	K	5.91E-04	5.6	1.47E-06	9.2	401.43	0.09
TCV hall	K	6.42E-01	0.4	8.61E-01	0.1	0.75	5.16E+4

TABLE III

Calculated Gamma H\*10 Dose Rates for the Engineering Design for a Source Rate of  $4 \times 10^{13}$  n/s

Location		Reference Model		Reference + PE Shield Model		Relative Dose Reduction	Highest Daily Dose ( $\mu$ Sv)
		Dose (mSv/s)	Relative Uncertainty (%)	Dose (mSv/s)	Relative Uncertainty (%)		
Control room, north	A	6.23E-04	2.4	3.33E-05	0.3	18.71	2.00
Control room, south	B	6.37E-04	2.5	2.10E-05	0.3	30.35	1.26
Second floor	C	1.65E-03	1.5	1.62E-04	0.2	10.16	9.75
Workshop	D	4.66E-04	3.0	1.11E-05	0.5	42.00	0.67
Diagnostic lab	E	3.70E-04	3.3	2.28E-05	0.4	16.22	1.37
South extension	F	3.57E-04	3.4	7.95E-06	0.4	44.90	0.48
Gyrotron platform	G	8.65E-04	2.2	3.82E-05	0.3	22.62	2.29
TCV entrance turnstile	H	2.41E-04	5.1	6.16E-05	0.3	3.91	3.69
North door	I	3.76E-04	3.2	2.23E-05	0.3	16.85	1.34
South door	J	3.24E-04	3.5	8.82E-06	0.5	36.75	0.53
West side outside	K	1.06E-04	6.2	6.32E-06	0.9	16.76	0.38
TCV hall	L	1.02E-02	0.67	2.38E-02	0.01	0.43	1.42E+3

view. However, we found that it is acceptable not to cover the PE ceiling of the TCV hall with a gamma-ray-stopping material, as the top of the ceiling is not in direct view from any human accessible zone and the roof above the ceiling is a weak contributor to gamma emission. Following implementation of the radiation shield in 2023, a series of experiments will be undertaken to experimentally validate the design and the protection level afforded.

## Acknowledgments

This work has been carried out in part within the framework of the EUROfusion Consortium, funded by the European Union via the Euratom Research and Training Programme (grant agreement no. 101052200—EUROfusion). The views and opinions expressed, however, are those of the authors only and do not necessarily reflect those of the European Union or the European Commission. Neither the European Union nor the European Commission can be held responsible for them.



The JSI team acknowledges the financial support from the Slovenian Research Agency research core funding no. P2-0073 and P2-0405.

## Disclosure Statement

No potential conflict of interest was reported by the authors.

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## References

1. H. REIMERDES et al., “Overview of the TCV Tokamak Experimental Programme,” *Nucl. Fusion*, **62**, 4, 2018 (2022); <https://doi.org/10.1088/1741-4326/ac369b>.
2. B. GEIGER et al., “Observation of Alfvén Eigenmodes Driven by Off-Axis Neutral Beam Injection in the TCV Tokamak,” *Plasma Phys. Control. Fusion*, **62**, 095017 (2020); <https://doi.org/10.1088/1361-6587/aba19e>.
3. T. P. GOODMAN et al., “An Overview of Results from the TCV Tokamak,” *Nucl. Fusion*, **43**, 1619 (2003); <https://doi.org/10.1088/0029-5515/43/12/008>.
4. A. N. KARPUSHOV et al., “Feasibility Studies of the Neutral Beam Heating System for the TCV Tokamak,” *Proc. 36th EPS Conf. on Plasma Physics*, Vol. 33E, p. P-2.140, Sofia, Bulgaria, June 29–July 3, 2009 (2009); [http://ocs.ciemat.es/EPS2009/pdf/P2\\_140.pdf](http://ocs.ciemat.es/EPS2009/pdf/P2_140.pdf).
5. A. N. KARPUSHOV et al., “Neutral Beam Heating on the TCV Tokamak,” *Fusion Eng. Des.*, **123**, 468 (Nov 2017); <https://doi.org/10.1016/j.fusengdes.2017.02.076>.
6. M. VALLAR et al., “Status, Scientific Results and Technical Improvements of the NBH on TCV Tokamak,” *Fusion Eng. Des.*, **146**, Part A, 773 (Sep., 2019); <https://doi.org/10.1016/j.fusengdes.2019.01.077>.
7. A. N. KARPUSHOV et al., “Upgrade of the Neutral Beam Heating System on the TCV Tokamak—Second High Energy Neutral Beam,” *Fusion Eng. Des.*, **187**, 113384 (2023); <https://doi.org/10.1016/j.fusengdes.2022.113384>.
8. A. PANKIN et al., “The Tokamak Monte Carlo Fast Ion Module NUBEAM in the National Transport Code Collaboration Library,” *Comput. Phys. Commun.*, **159**, 157 (2004); <https://doi.org/10.1016/j.cpc.2003.11.002>.
9. S. W. MOSHER et al., “ADVANTG—An Automated Variance Reduction Parameter Generator,” ORNL/TM-2013/416, Oak Ridge National Laboratory (2013).
10. B. KIEDROWSKI et al., “MCNP5-1.6, Feature Enhancements and Manual Clarifications,” LA-UR-10-06217, Los Alamos National Laboratory (2010).
11. “Le Conseil fédéral suisse, Ordonnance sur la radioprotection (ORaP) 26 April (2017),” <https://fedlex.data.admin.ch/filestore/fedlex.data.admin.ch/eli/cc/2017/502/20210101/fr/pdf-a/fedlex-data-admin-ch-eli-cc-2017-502-20210101-fr-pdf-a.pdf>.
12. “LUPIN BF3-NP, Neutron REM Counter for Pulsed Fields,” ELSE Nuclear,; <http://www.elsenuclear.com/en/lupin-series>.
13. “NAUSICCAA IC-T, ICP-T Ion Chamber–Based Gamma Monitoring Units,” ELSE Nuclear; [http://www.elsenuclear.com/media/k2/attachments/DAT\\_IONCHAMBERDET\\_EN\\_1.10.pdf](http://www.elsenuclear.com/media/k2/attachments/DAT_IONCHAMBERDET_EN_1.10.pdf).
14. “The 2007 Recommendations of the International Commission on Radiological Protection,” *Ann. ICRP*, **37**, 2–4, 294 (2007).
15. Rapiscan Systems, data sheet, <https://www.rapiscan-ase.com/radiation-detection-products/detectors/guardian-glass>.
16. H. ARNICHAND et al., “New Capabilities of the Incoherent Thomson Scattering Diagnostics in the TCV Tokamak: Divertor and Real-Time Measurements,” *JINST*, **14**, C09013 (2019); <https://doi.org/10.1088/1748-0221/14/09/C09013>.
17. G. V. PEREVERZEV, “ASTRA. An Automatic System for Transport Analysis in a Tokamak,” MPG. PuRe Publication Repository (1991); <http://hdl.handle.net/11858/00-001M-0000-0027-646C-F>.
18. R. J. MCCONN et al., “Compendium of Material Composition Data for Radiation Transport Modelling,” Pacific Northwest National Laboratory; [https://www.pnnl.gov/main/publications/external/technical\\_reports/PNNL-15870Rev1.pdf](https://www.pnnl.gov/main/publications/external/technical_reports/PNNL-15870Rev1.pdf) (current as of 2011).
19. R. FORREST et al., “FENDL-3 Library—Summary Documentation,” INDC(NDS)-0628, International Atomic Energy Agency (2012).
20. “FENDL-3.1d: Fusion Evaluated Nuclear Data Library Ver.3.1d,” International Atomic Energy Agency (Jan. 24, 2018); <https://www-nds.iaea.org/fendl31d/>.
21. Radiation Safety & Protection, LCC, data sheets, <https://www.radsafetypro.com/shieldwerx-products.html>.
22. “Technical Sheet HMW-PE/PE-500, ESD 90, black,” ERIKS; [https://eriksdigitalcdn.azureedge.net/hlr-system/maagtechnic/datasheets/technische%20datenblätter/kunststoffe/hmw-pe-pe-500\\_eng.pdf](https://eriksdigitalcdn.azureedge.net/hlr-system/maagtechnic/datasheets/technische%20datenblätter/kunststoffe/hmw-pe-pe-500_eng.pdf).
23. M. FORTUNA et al., “Evaluation of Neutron Dose Rates at the TCV Tokamak Facility,” *Fusion Eng. Des.*, **191**, 113562 (2023); <https://doi.org/10.1016/j.fusengdes.2023.113562>.
24. “Dimensions of Steel Beams Type HEB,” PIPING-WORLD; <https://www.piping-world.com/dimensions-of-steel-beams-type-heb>.