



Computerised Tomography

Mathematical Foundations of Signal Processing

Dr. Matthieu Simeoni

November 23, 2020

EPFL

Table of contents

1 Background Concepts

- Physics of Positron Emission Tomography
- Radon Transform
- Projection-Slice Theorem
- Sinogram
- From Line Integrals to Tube Integrals

2 Radon-based Tomographic Reconstruction

- Filtered Back-Projection
- Discrete FBP Algorithm

3 Sampling & Interpolation in PET

- Sampling Operator
- Ideally-Matched Interpolation
- Computing the Gram Matrix

4 Regularisation

- Spectral truncation, Tikhonov regularisation
- Sketching

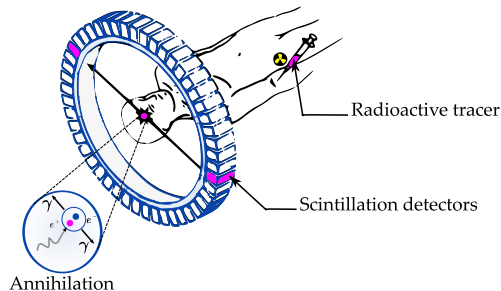
Background Concepts

Positron Emission Tomography (PET)

Definition: (Positron Emission Tomography)

Positron Emission Tomography (PET) is a medical diagnostic technique that enables a physician to study blood flow in and metabolic activity of an organ in a visual way.

- A biochemical metabolite labeled with a positron emitting radioactive material is introduced into the organ.
- The biochemical (typically sugar for the brain) concentrates in regions of high metabolic activity.
- Positron emissions occur randomly are counted by a PET scanner (ring of scintillation detectors).

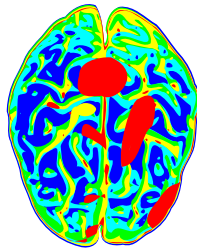
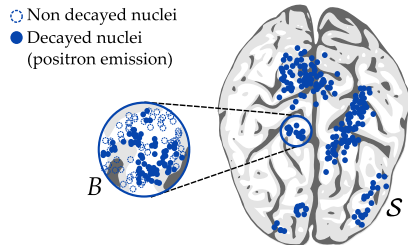


Positron Emissions as Poisson Process

- Positron emissions occur **randomly** in \mathcal{S} according to a **Poisson process**.
- Rate of occurrence is characterised by an **intensity function** $\lambda: \mathcal{S} \rightarrow \mathbb{R}$.
- For a given region $B \subset \mathcal{S}$ of the organ, the expected **number of positron emissions** $N(B)$ is given by

$$\mathbb{E}[N(B)] = \int_B \lambda(\mathbf{x}) d\mathbf{x}.$$

- The intensity function is assumed to be **proportional to the metabolic activity of interest**.

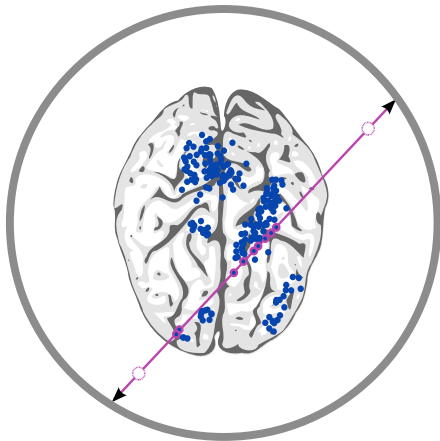


Indirectly Observed Poisson Process

- In practice, we **cannot directly observe** positron emissions.
- Instead, we observe **gamma rays** induced by annihilation with neighbouring electrons.
- Coincident gamma rays are recorded on a **detector ring**.
- We speak of an **indirectly observed Poisson process**. The number n_d of gamma rays recorded by each detector pair is **Poisson distributed**, with rate

$$n_d \sim \mathcal{P}(\check{\lambda}_d), \quad \check{\lambda}_d = \mathbb{E}[N(L_d)] = \int_{L_d} \lambda(\mathbf{x}) d\mathbf{x}, \quad d = 1, \dots, D,$$

where D denotes the **total number of detector pairs** on the detector ring. $\check{\lambda}_d$ corresponds to the **line integral** of λ along the line L_d linking the d -th pair of detectors on the detector ring.



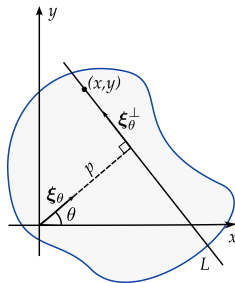
The Radon Transform

- Line integrals can be seen as samples of the so-called **Radon transform**.
- The **chord** linking two points on a ring can be parametrised as:

$$L = \{x \in \mathbb{R}^2 : \langle x, \xi_\theta \rangle = p\},$$

where $p \in \mathbb{R}$, $\theta \in [0, \pi)$ and $\xi_\theta = [\cos(\theta), \sin(\theta)] \in \mathbb{S}^1$.

- The **Radon transform** maps a function λ onto its **line integrals**:



Definition (Radon Transform)

The **Radon transform** $\check{\lambda} : [0, \pi[\times \mathbb{R} \rightarrow \mathbb{R}$ of a function $\lambda : \mathbb{R}^2 \rightarrow \mathbb{R} \in \mathcal{L}^2(\mathbb{R}^2)$ is

$$\check{\lambda}(\theta, p) = (\mathcal{R}\lambda)(\theta, p) := \int_{\mathbb{R}^2} \lambda(x) \delta(p - \langle x, \xi_\theta \rangle) dx.$$

Example: Radon Transform of a Gaussian

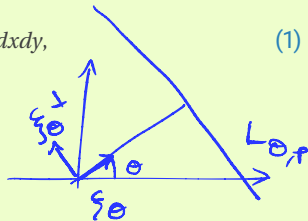
Example: Radon Transform of a Gaussian

Let $f(x, y) = e^{-\pi(x^2+y^2)}$, $\forall (x, y) \in \mathbb{R}^2$. Then,

$$\check{f}(\theta, p) = \int_{-\infty}^{+\infty} \int_{-\infty}^{+\infty} e^{-\pi(x^2+y^2)} \delta(p - \cos(\theta)x - \sin(\theta)y) dx dy, \quad (1)$$

We perform the following orthogonal transformation:

$$\begin{pmatrix} u \\ v \end{pmatrix} = \begin{pmatrix} \cos \theta & \sin \theta \\ -\sin \theta & \cos \theta \end{pmatrix} \begin{pmatrix} x \\ y \end{pmatrix}.$$

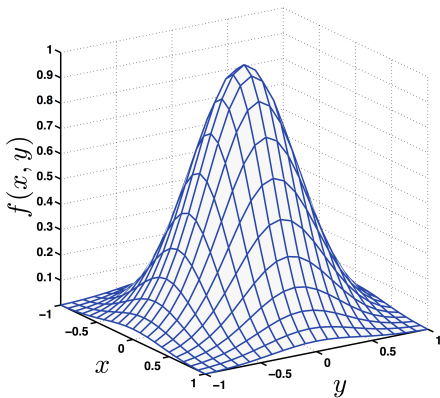


The transformation is **orthogonal**, and thus $u^2 + v^2 = x^2 + y^2$. Moreover, in this new basis, the equation of the line $L = \{(x, y) \in \mathbb{R}^2 : \cos(\theta)x + \sin(\theta)y = p\}$ becomes simply: $L = \{(u, v) \in \mathbb{R}^2 : u = p\}$. Eq. (1) hence becomes:

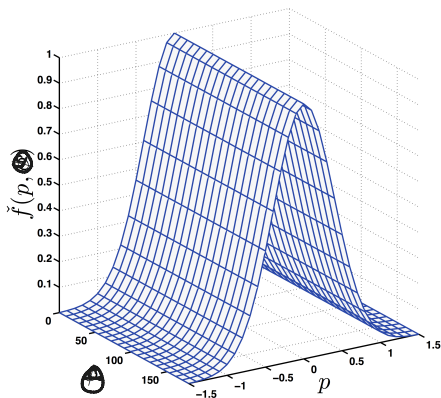
$$\check{f}(\theta, p) = \int_{-\infty}^{+\infty} \int_{-\infty}^{+\infty} e^{-\pi(u^2+v^2)} \delta(p - u) du dv = e^{-\pi p^2} \int_{-\infty}^{+\infty} e^{-\pi v^2} dv = e^{-\pi p^2},$$

where we have used the well known result $\int_{-\infty}^{+\infty} e^{-t^2} dt = \sqrt{\pi}$. Hence, we have: $\mathcal{R}\{e^{-\pi(x^2+y^2)}\} = e^{-\pi p^2}$.

Example: Radon Transform of a Gaussian



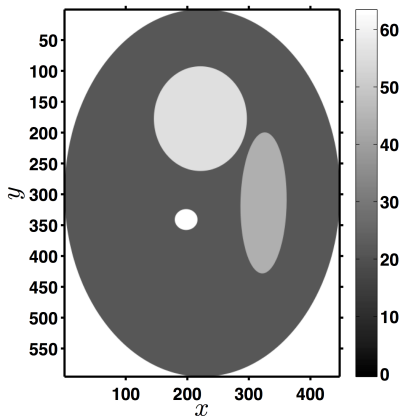
(a) The Gaussian distribution.



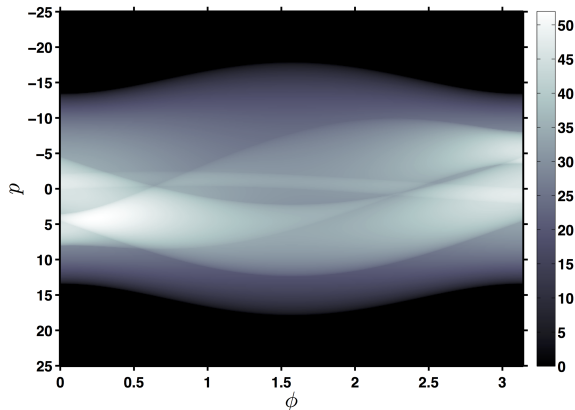
(b) Radon transform of the Gaussian distribution.

Figure: The Gaussian distribution and its Radon transform.

Example: Radon Transform of Ellipses [1, Example 2.2]



(a) Sum of ellipses' characteristic functions with heights given by the color of each ellipse.



(b) Radon transform.

Figure: The Radon transform of a sum of ellipses' characteristic functions.

Basic Properties of the Radon Transform

Proposition: (Properties of the Radon Transform)

- **Linearity:** Let f and g be two functions and $\alpha, \beta \in \mathbb{R}$. Then,

$$\mathcal{R}\{\alpha f + \beta g\} = \alpha \mathcal{R}f + \beta \mathcal{R}g.$$

- **Shifting Property:** Let $g(\mathbf{x}) = f(\mathbf{x} - \mathbf{a})$ for some $\mathbf{a} \in \mathbb{R}^2$. Then we have:

$$\mathcal{R}g(\theta, p) = \mathcal{R}f(\theta, p - \langle \xi_\theta, \mathbf{a} \rangle), \quad \forall (\theta, p) \in [0, \pi) \times \mathbb{R}.$$

- **Scaling Property:** Let $g(\mathbf{x}) = f(\alpha \mathbf{x})$ for some $\alpha \neq 0$. Then we have:

$$\mathcal{R}g(\theta, p) = \frac{1}{\alpha^2} \mathcal{R}f(\theta, \alpha p), \quad \forall (\theta, p) \in [0, \pi) \times \mathbb{R}.$$

Proof:

- **Linearity:** $\mathcal{R}\{\alpha f + \beta g\} = \int (\alpha f(\mathbf{x}) + \beta g(\mathbf{x})) \delta(p - \xi \cdot \mathbf{x}) d\mathbf{x} = \alpha \check{f} + \beta \check{g}.$
- **Shifting Property:** $\mathcal{R}\{f(\mathbf{x} - \mathbf{a})\} = \int f(\mathbf{x} - \mathbf{a}) \delta(p - \xi \cdot \mathbf{x}) d\mathbf{x} = \int f(\mathbf{y}) \delta(p - \xi \cdot \mathbf{a} - \xi \cdot \mathbf{x}) d\mathbf{x}.$
- **Scaling Property:** $\mathcal{R}\{f(\alpha \mathbf{x})\} = \int f(\alpha \mathbf{x}) \delta(p - \xi \cdot \mathbf{x}) d\mathbf{x} = \frac{1}{\alpha} \int f(\mathbf{y}) \delta(p - \frac{1}{\alpha} \xi \cdot \mathbf{y}) d\mathbf{y} = \frac{1}{\alpha^2} \int f(\mathbf{y}) \delta(\alpha p - \xi \cdot \mathbf{y}) d\mathbf{y}.$

Link with 2D Fourier Transform

The Radon and Fourier transforms are linked by the **projection-slice theorem**:

Lemma: (Projection-Slice Theorem)

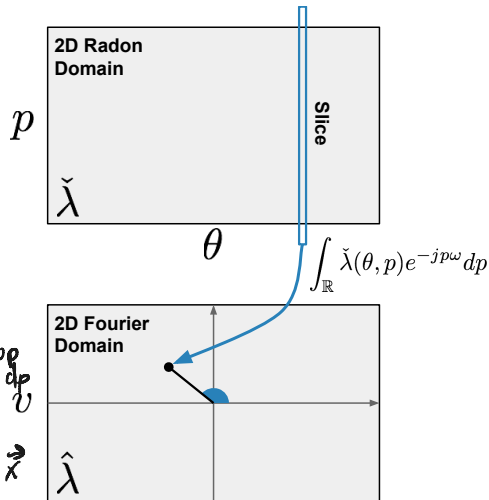
For any $\theta \in [0, \pi)$ we have:

$$\int_{\mathbb{R}} \mathcal{R}\lambda(\theta, p) e^{-j\omega p} dp = \hat{\lambda}(\omega \cos \theta, \omega \sin \theta),$$

where $\hat{\lambda}: \mathbb{R}^2 \rightarrow \mathbb{C}$ is the **2D Fourier transform** of λ .

Proof  :

$$\begin{aligned} \int_{\mathbb{R}} \check{\lambda}(\theta, p) e^{-j\omega p} dp &= \int_{\mathbb{R}^2} \left(\int_{\mathbb{R}} \lambda(\vec{x}) \delta(p - \langle \vec{x}, \vec{\xi}_{\theta} \rangle) d\vec{x} \right) e^{-j\omega p} dp \\ &= \int_{\mathbb{R}^2} \lambda(\vec{x}) \left(\int_{\mathbb{R}} e^{-j\omega p} \delta(p - \langle \vec{x}, \vec{\xi}_{\theta} \rangle) dp \right) d\vec{x} \\ &= \int_{\mathbb{R}^2} \lambda(\vec{x}) e^{-j\omega \langle \vec{x}, \vec{\xi}_{\theta} \rangle} d\vec{x} = \hat{\lambda}(\omega \vec{\xi}_{\theta}) = \hat{\lambda}(\omega \cos \theta, \omega \sin \theta) \end{aligned}$$



Back to PET

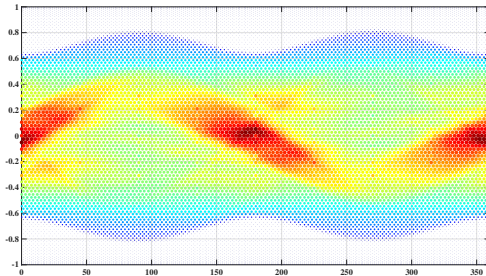
- A PET scanner **samples** the Radon transform of the metabolic activity.
- Samples are **polluted by Poisson noise**, resulting in the data:

$$n(\theta_d, p_d) \sim \mathcal{P}(\check{\lambda}(\theta_d, p_d)), \quad d = 1, \dots, D.$$

- It is customary to represent the data in the $(\theta, p) \in [0, \pi) \times \mathbb{R}$ plane, yielding a **sinogram**.



(a) Metabolic activity



(b) PET data as a Sinogram

Complications with Real-life Scanners

$$\chi_d(p) := \begin{cases} 1 & \text{if } |p_d - p| \leq \epsilon_d \\ 0 & \text{otherwise} \end{cases}$$

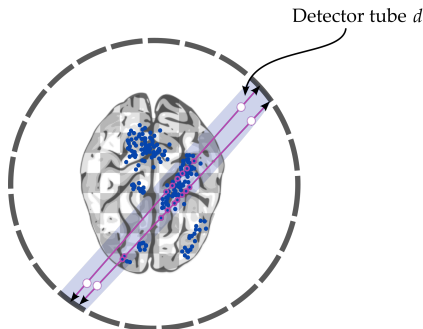
- In practice, the detectors have a certain *width*. We have hence **detector tubes** instead of lines,

$$\begin{aligned} \mathcal{T}_d &= \{ \mathbf{x} \in \mathbb{R}^2 : |\langle \mathbf{x}, \boldsymbol{\xi}_d \rangle - p_d| \leq \epsilon_d \} \\ &= \{ \mathbf{x} \in \mathbb{R}^2 : \underline{\chi}_d(p_d - \langle \mathbf{x}, \boldsymbol{\xi}_d \rangle) = 1 \}. \end{aligned}$$

- Actual data hence consist in **tube integrals** and not line integrals:

$$\check{\lambda}(\theta_d, p_d) = \mathbb{E}[N(\mathcal{T}_d)] = \int_{\mathbb{R}^2} \lambda(\mathbf{x}) \chi_d(p_d - \langle \mathbf{x}, \boldsymbol{\xi}_d \rangle) d\mathbf{x}.$$

- Formulating the data model in terms of the Radon transform is hence only valid in the limit for **infinitely thin detectors**.



Radon-based Tomographic Reconstruction

Inverting the Radon Transform

The Radon transform is **invertible**. Inversion formula in 2D is given by the **filtered back-projection (FBP)** formula [2]:

Theorem: (Filtered Back-Projection)

Let $\lambda: \mathbb{R}^2 \rightarrow \mathbb{R}$ be **sufficiently smooth**. Then we have

$$\lambda(x) = \frac{1}{(2\pi)^2} \int_0^\pi [\check{\lambda}(\theta, \cdot) * h](\langle x, \xi_\theta \rangle) d\theta,$$

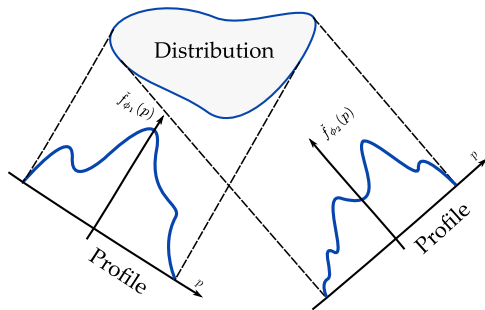
where $h: \mathbb{R} \rightarrow \mathbb{R}$ – called the **Ramp filter** – is defined in terms of its Fourier transform $\hat{h}(\omega) := |\omega|, \forall \omega \in \mathbb{R}$.

$$(f * h)(p) \underset{\text{FT}}{\longleftrightarrow} |\omega| \hat{f}(\omega)$$

(multiplication theorem)

Inverting the Radon Transform

- **Holographic result:** Reconstruct 2D object from many 1D projections (profiles).
- The **Ramp filter**, with Fourier transform $\hat{h}(\omega) = |\omega|$, is a *roughening filter*, acting as a **derivative**.^a It makes objects more singular, with **sharper edges**.
- The function λ **must hence be sufficiently smooth** for the inversion formula to be well-defined!



^aRecall that in the Fourier domain **differentiating accounts to multiplying** by $j\omega$: $\mathcal{F}\{g\}(\omega) = j\omega \hat{g}(\omega)$.

Interpretation of FBP

- The **adjoint** of the Radon transform \mathcal{R} is given, for all $\mu: [0, \pi) \times \mathbb{R} \rightarrow \mathbb{R}$, by

$$\mathcal{R}^* \mu = \int_0^\pi \mu(\theta, \langle \mathbf{x}, \boldsymbol{\xi}_\theta \rangle) d\theta.$$

Indeed,

$$\begin{aligned} \langle \mathcal{R}\lambda, \mu \rangle &= \int_0^\pi \int_{\mathbb{R}} \left[\int_{\mathbb{R}^2} \lambda(\mathbf{x}) \delta(p - \langle \mathbf{x}, \boldsymbol{\xi}_\theta \rangle) d\mathbf{x} \right] \mu(\theta, p) dp d\theta \\ &= \int_{\mathbb{R}^2} \lambda(\mathbf{x}) \left[\int_0^\pi \mu(\theta, \langle \mathbf{x}, \boldsymbol{\xi}_\theta \rangle) d\theta \right] d\mathbf{x} = \langle \lambda, \mathcal{R}^* \mu \rangle. \end{aligned}$$

- \mathcal{R}^* performs a **back-projection**: all sinogram points to which \mathbf{x} contributed to are summed together.
- The inverse Radon transform decomposes as:

$$\mathcal{R}^{-1} = \frac{1}{(2\pi)^2} \mathcal{R}^* \tilde{\lambda}, \quad \text{where} \quad \tilde{\lambda}(\theta, p) = \int_{\mathbb{R}} \check{\lambda}(\theta, p-t) h(t) dt = \frac{1}{2\pi} \int_{\mathbb{R}} |\omega| \hat{\lambda}(\omega \boldsymbol{\xi}_\theta) e^{j\omega p} d\omega,$$

where the last equality results from the **projection-slice** and **convolution/multiplication** theorems.

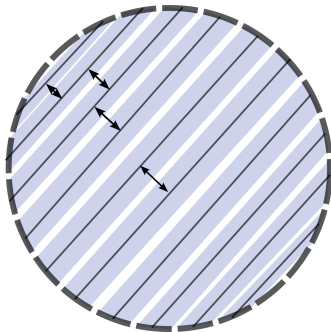
- In summary: first apply 1D **filter** h to slices $\check{\lambda}(\theta, \cdot)$ of the Radon transform, then **back-project** with \mathcal{R}^* . Hence the name: **filtered back-projection**.

Filtered Backprojection in Practice

- Tubes are assumed *narrow* and **approximated** by lines with coordinates (θ_d, p_d) .
- Samples are typically **non-uniform** and **noisy**.
- Discrete FBP for **uniform samples** in the p -direction is given by

$$\lambda_{FBP}(\mathbf{x}) = \frac{\Delta\theta}{(2\pi)^2} \sum_{n=1}^{N_\theta} \left(\lambda[\theta_n, \cdot] \circledast \mathbf{h} \right) \left[\lfloor \langle \mathbf{x}, \boldsymbol{\xi}_{\theta_n} \rangle / \Delta p \rfloor \right], \quad \mathbf{x} \in \mathbb{R}^2.$$

- Data must be **gridded**.
- Convolution is approximated by a **circular discrete convolution**, efficiently implemented via **FFT/iFFT**.
- The routine `skimage.transform.iradon` implements the discrete FBP in Python.

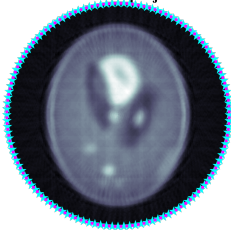


Example

Ground Truth



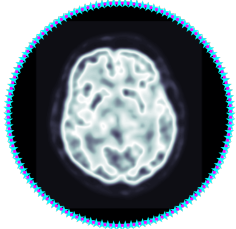
Filtered Backprojection



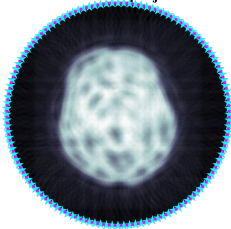
Error filtBackprojection



Ground Truth



Filtered Backprojection

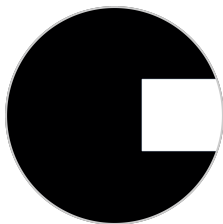


Error filtBackprojection

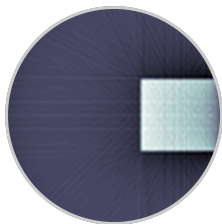


Issues with FBP

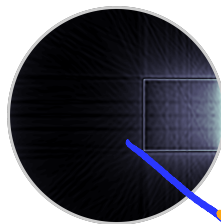
- Gridding is **expensive** and ad-hoc.
- Discrete formula makes a **lot of approximations** (circular convolution, interpolation).
- **Border effects** due to circular convolution.
- Ramp filter boosts high frequencies, generally **polluted by noise**. FBP is **unstable** and must be regularised by **truncating** or more generally **windowing** the Ramp filter (optimal window?).



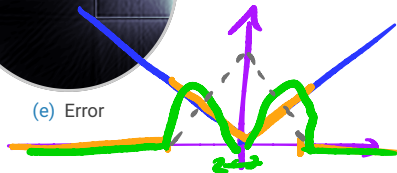
(c) Metabolic activity



(d) FBP Estimate



(e) Error



Sampling & Interpolation in PET

22

Sampling Operator

- Define $\mathcal{H} = \mathcal{L}^2(\mathbb{B}_1)$ with \mathbb{B}_1 the **unit disk** (the brain space).
- We can link the measurements (n_1, \dots, n_D) to the Poisson process of interest N through a **sampling operator** $\Phi^* : \mathcal{H} \rightarrow \mathbb{R}^D$:

$$\begin{bmatrix} n_1 \\ \vdots \\ n_D \end{bmatrix} = \Phi^* N = \begin{bmatrix} \langle N, \phi_1 \rangle_{\mathcal{H}} \\ \vdots \\ \langle N, \phi_D \rangle_{\mathcal{H}} \end{bmatrix},$$

where ϕ_d are the **indicator functions of the detector tubes**:

$$\phi_d(x) = \chi_d(x) = \begin{cases} 1 & \text{if } x \in \mathcal{T}_d, \\ 0 & \text{otherwise.} \end{cases}$$

Sampling Operator

- Since Φ^* is **linear**, on expectation we have

$$\begin{bmatrix} \lambda_1^* \\ \vdots \\ \lambda_D^* \end{bmatrix} = \Phi^* \lambda = \begin{bmatrix} \langle \lambda, \phi_1 \rangle_{\mathcal{H}} \\ \vdots \\ \langle \lambda, \phi_D \rangle_{\mathcal{H}} \end{bmatrix}.$$

- Measurements give us *evidence about the components* of λ in $\mathcal{R}(\Phi) = \text{span}\{\phi_1, \dots, \phi_D\}$.
- **Goal**: to solve this **inverse problem** and find an estimate λ .
- This is an **ill-posed problem**: any component in $\mathcal{N}(\Phi^*)$ is **unaccessible to us**

$$\Phi^* \lambda = \Phi^* \lambda_1 + \underbrace{\Phi^* \lambda_2}_{=0} = \Phi^* \lambda_1,$$

where $\lambda = \lambda_1 + \lambda_2 \in \mathcal{R}(\Phi) \oplus \mathcal{N}(\Phi^*) = \mathcal{H}$.

Least-Squares Estimate

- An estimate of λ can be obtained by solving

$$\lambda^{\star} \in \operatorname{argmin}_{\lambda \in \mathcal{H}} \|\mathbf{n} - \Phi^* \lambda\|_2^2$$

- No unique solution!
- Impose minimal \mathcal{L}_2 norm and $\lambda \in \mathcal{R}(\Phi)$ for uniqueness.
- Leads to the generalised Moore-Penrose pseudo-inverse solution:

$$\lambda^{\star} = (\Phi^*)^{\dagger} \mathbf{n}.$$

Ideally-Matched Interpolation

- Generalised pseudoinverse is given by

$$(\Phi^*)^\dagger = \Phi(\Phi^* \Phi)^{-1},$$

and the **least-squares estimate** is hence given by

$$\lambda^*(x) = (\Phi(\Phi^* \Phi)^{-1} \mathbf{n})(x) = \sum_{d=1}^D \tilde{\mathbf{n}} \phi_d(x),$$

where $\tilde{\mathbf{n}} = (\Phi^* \Phi)^{-1} \mathbf{n} \in \mathbb{R}^D$.

- Recovery in two steps:** apply *Gram correction* to the data, and interpolate using the **synthesis operator** Φ , adjoint of Φ^*

$$\Phi: \begin{cases} \mathbb{R}^D \rightarrow \mathcal{R}(\Phi), \\ \mathbf{y} \mapsto (\Phi \mathbf{y})(x) = \sum_{d=1}^D y_d \phi_d(x). \end{cases}$$

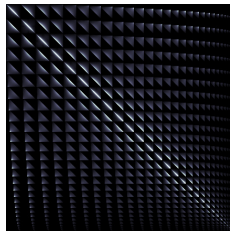
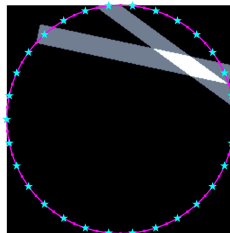
- We have **consistency** $\Phi^*(\Phi^*)^\dagger = I_D$ and hence $(\Phi^*)^\dagger \Phi^*$ is an **orthogonal projection**.

About the Gram Matrix

- The quantity $\Phi^* \Phi \in \mathbb{R}^{D \times D}$ is the **Gram matrix**.
- An element of the Gram matrix is given by

$$(\Phi^* \Phi)_{ij} = \langle \phi_i, \phi_j \rangle.$$

- Need to compute *areas of parallelograms*! Can be **efficiently computed analytically**.
- **Dense** matrix! Basis elements are **not localised**...
- **Dense = ill-conditioned** (often)



Regularisation

Eigenfunctions of the Integral Operator

- We have

$$\Phi(\Phi^* \Phi)^{-1} \Phi^* \Phi \alpha = \Phi \alpha.$$

- Hence any element of $\mathcal{R}(\Phi)$ is an **eigenfunction with eigenvalue 1**. To get orthogonal eigenspaces, we need to find $\{\alpha_1, \dots, \alpha_D\} \subset \mathbb{R}^D$ s.t.

$$\langle \Phi \alpha_i, \Phi \alpha_j \rangle = \alpha_j^T (\Phi^* \Phi) \alpha_i = 0.$$

- Choosing $\{\alpha_1, \dots, \alpha_D\}$ **eigenvectors of $\Phi^* \Phi$** yields the spectral decomposition:

$$\Phi(\Phi^* \Phi)^{-1} \Phi^* = \sum_{d=1}^D \frac{1}{\eta_d} (\Phi \alpha_d) (\Phi \alpha_d)^*,$$

where $\eta_d = \|\Phi \alpha_d\|_2^2 = \alpha_d^T (\Phi^* \Phi) \alpha_d$.

Noisy Measurements

- We can hence re-write the pseudo-inverse estimate as

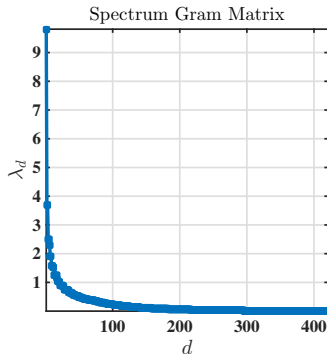
$$\lambda^{\star} = \sum_{d=1}^D \frac{\langle \lambda, \Phi \alpha_d \rangle}{\eta_d} (\Phi \alpha_d) = \sum_{d=1}^D \frac{\alpha_d^T \mathbf{n}}{\eta_d} (\Phi \alpha_d).$$

- In practice, measurements are **noisy** (Poisson noise). For **high rates** we can approximate

$$\mathbf{n} = \Phi^* \lambda + \epsilon,$$

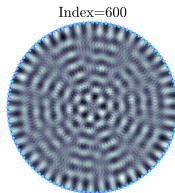
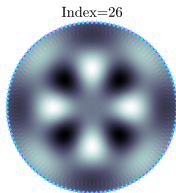
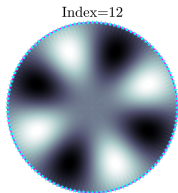
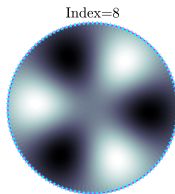
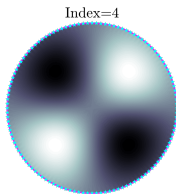
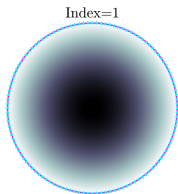
with $\epsilon \sim \mathcal{N}(0, \Sigma)$.

- Small η_d may lead to **numerical instability**! Need **regularisation**...



Back to the Eigenfunctions

- Small η_d correspond to *high-frequency eigenfunctions*! Without regularisation the estimate will be dominated by those eigenfunctions and hence **very wiggly**...



Regularisation

- Two avenues: Spectral truncation vs. Tikhonov regularization.

- Spectral truncation:

$$\lambda_{ST}^* = \sum_{d=1}^{\tau} \frac{\alpha_d^T \mathbf{n}}{\eta_d} (\Phi \alpha_d),$$

with $\tau \leq D$ some integer, truncation parameter.

- Tikhonov regularization:

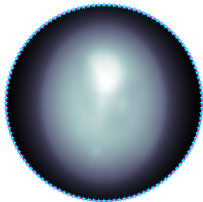
$$\lambda_{\rho}^* = \sum_{d=1}^D \frac{\alpha_d^T \mathbf{n}}{\eta_d + \rho} (\Phi \alpha_d),$$

with $\rho > 0$, called the ridge parameter.

- Act as smoothers. Tikhonov performs better (in terms of consistency), but spectral truncation is more economic computationally (less terms in the summation).

Spectral Truncation

No Gram Correction



Truncation= $1e-02 \times \lambda_{max}$



Truncation= $5e-03 \times \lambda_{max}$



Truncation= $1e-03 \times \lambda_{max}$



Truncation= $1e-07 \times \lambda_{max}$

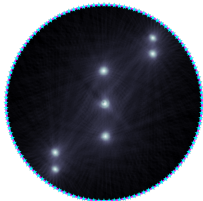


Truncation= $1e-10 \times \lambda_{max}$

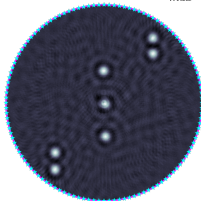


Spectral Truncation (Point Spread Function)

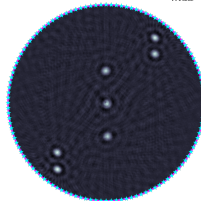
No Gram Correction



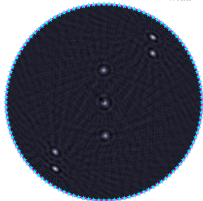
Truncation= $1e-02 \times \lambda_{max}$



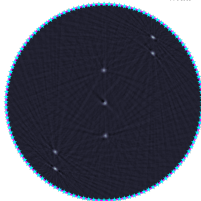
Truncation= $5e-03 \times \lambda_{max}$



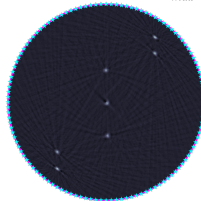
Truncation= $1e-03 \times \lambda_{max}$



Truncation= $1e-07 \times \lambda_{max}$

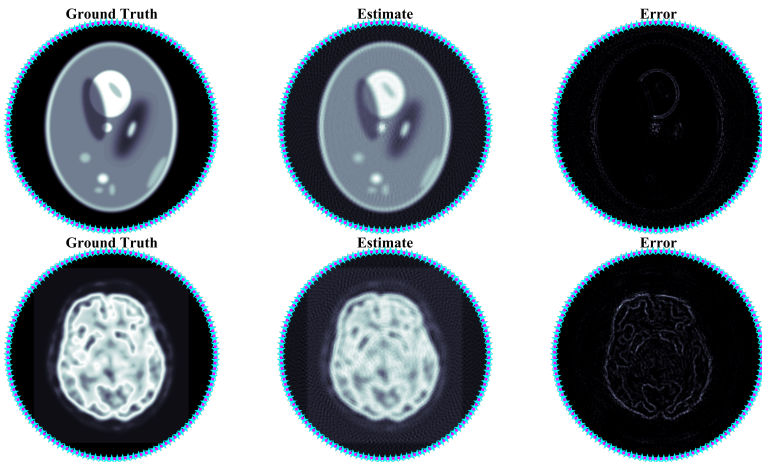


Truncation= $1e-10 \times \lambda_{max}$

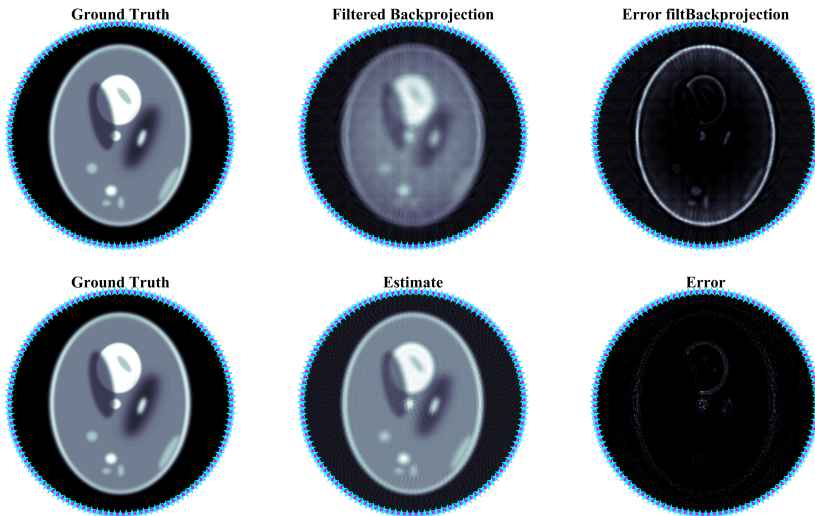


Final Estimate

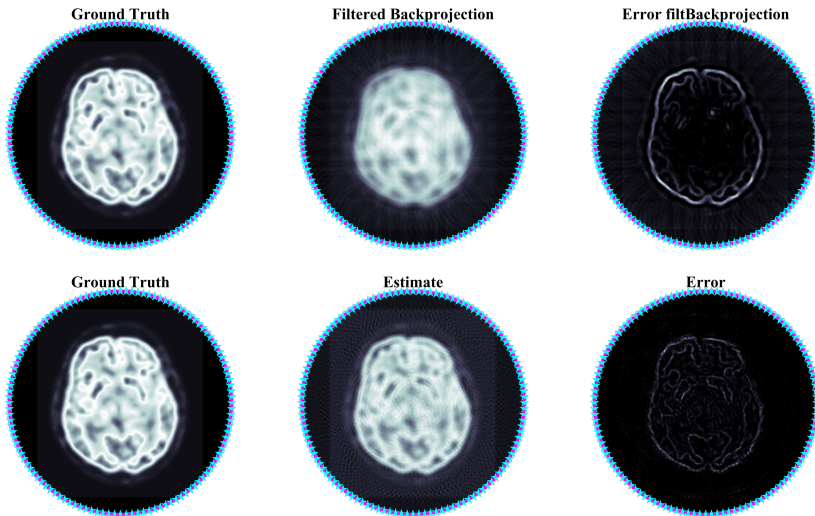
- Optimal truncation parameter chosen according to the width of point spread function main lobe.



Comparison: FBP vs. Interpolation



Comparison: FBP vs. Interpolation



Comparison: FBP vs. Interpolation

- Interpolation produces **more accurate images** than state of the art method.
- Estimate can be sampled and displayed at **any resolution (continuous estimate)**.
- Filtered backprojection scales however better with the number of detectors.
- Indeed the Gram matrix is **expensive** to compute and invert.
- Need to **investigate dimensionality reduction** via **sketching** (random projections):

$$\Psi^* = W^H \Phi^*,$$

with $W^H : \mathbb{R}^D \rightarrow \mathbb{R}^L, L \ll D$.

Gaussian Sketching

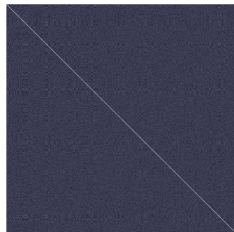
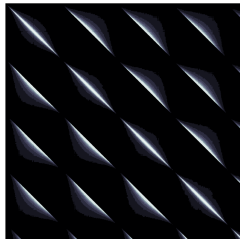
- Choose $\psi_m(x) = \sum_{d=1}^D W_{dm} \phi_d(x)$, with

$$W_{dm} \stackrel{i.i.d}{\sim} \mathcal{N}(0, \sigma).$$

- New Gram is $G_\Psi = W^H \Phi^* \Phi W$.
- If $\sigma = 1/\text{trace}(\Phi^* \Phi)$ we can show that

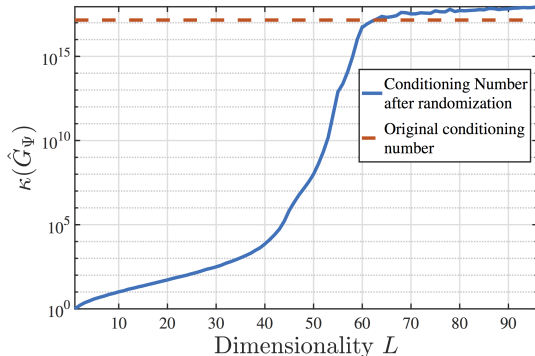
$$\mathbb{E} \left[W^H \Phi^* \Phi W \right] = I_{L \times L}.$$

- Sketching acts as a **preconditioner** (improves conditioning). In expectation, the Gram is identity...
- Basis functions are **less coherent**.



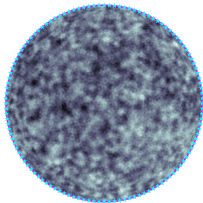
Compression Factor

- In practice we only have **one** random realization of $W^H \Phi^* \Phi W$.
- No guarantee it would fall **near the mean**!
- For **small dimensions**, this is more **likely** to be the case.
- Example: for $D = 7140$ and $L = 3160$, we go from $\kappa(G_\Phi) = 1.9458 \times 10^{21}$ to $\kappa(G_\Psi) = 7476$.
- **Huge** improvement (18 orders of magnitude)!

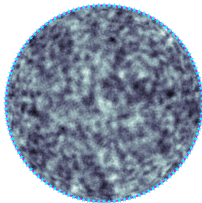


Results ($D = 7140$, $M = 3160$)

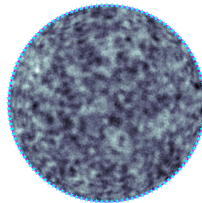
Basis Function 1



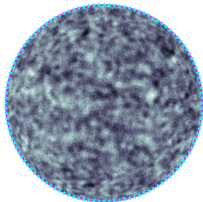
Basis Function 632



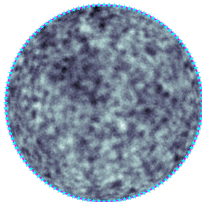
Basis Function 1264



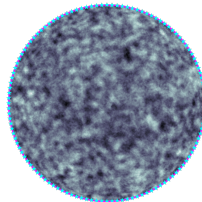
Basis Function 1896



Basis Function 2528

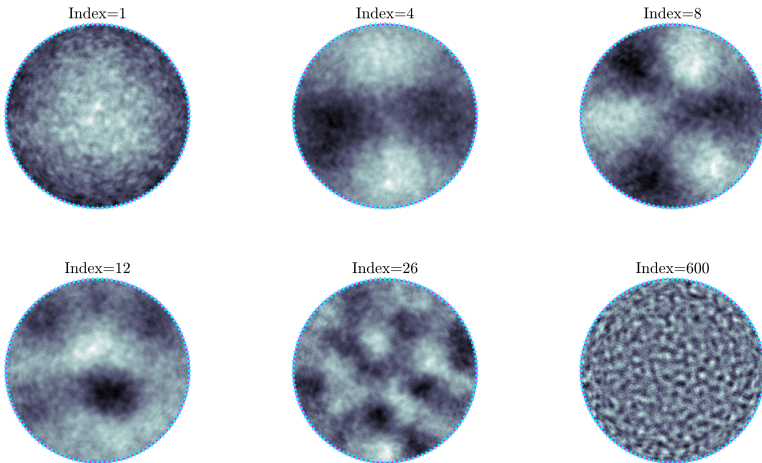


Basis Function 3160

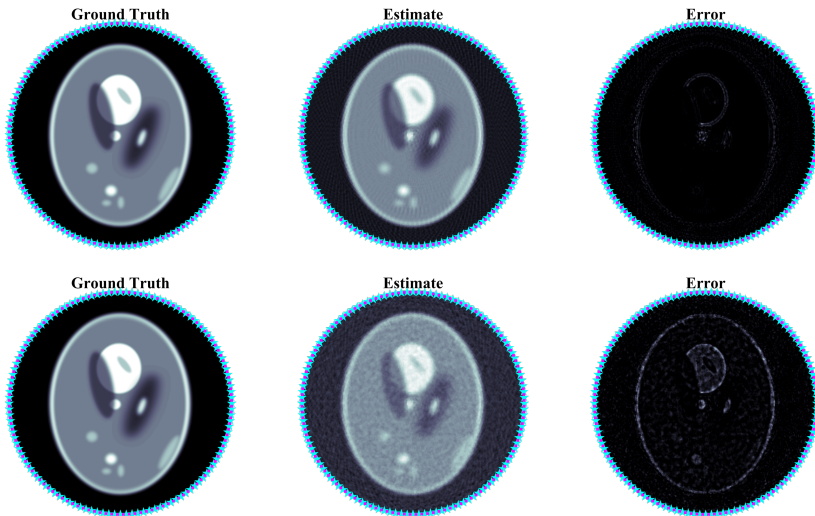


Results ($D = 7140$, $M = 3160$)

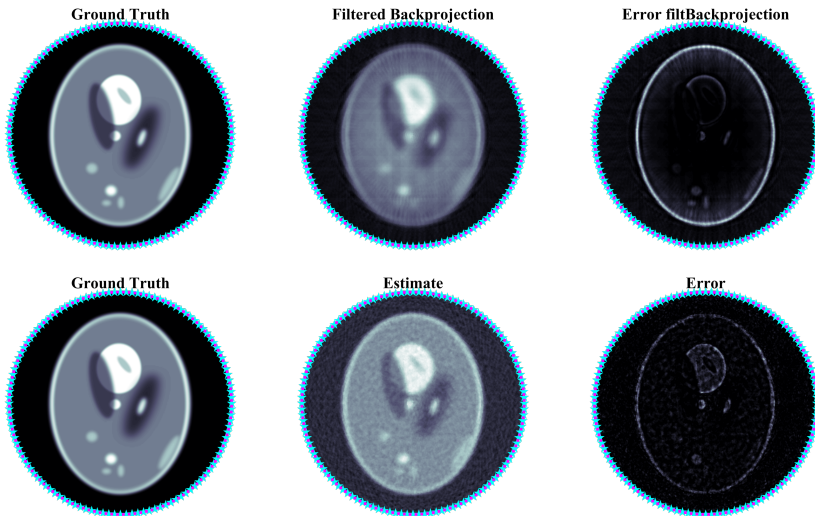
- Eigenfunctions of the sketched sampling operator Ψ^* :



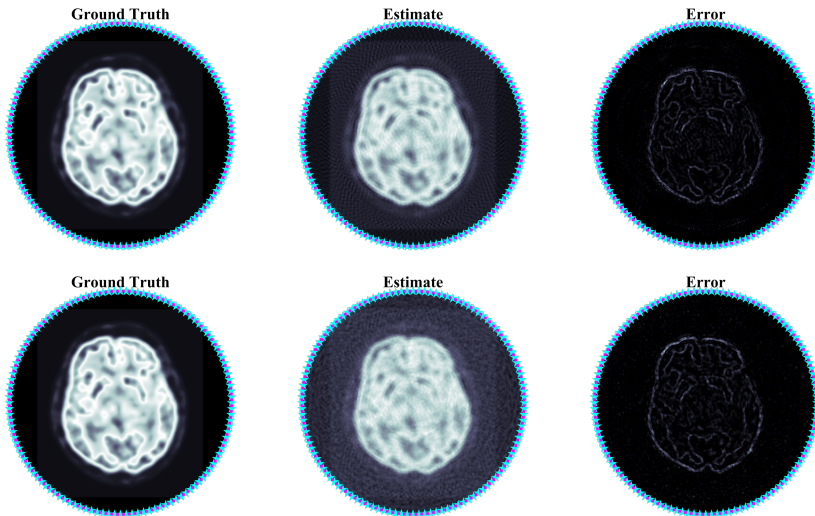
Comparison without/with Sketching



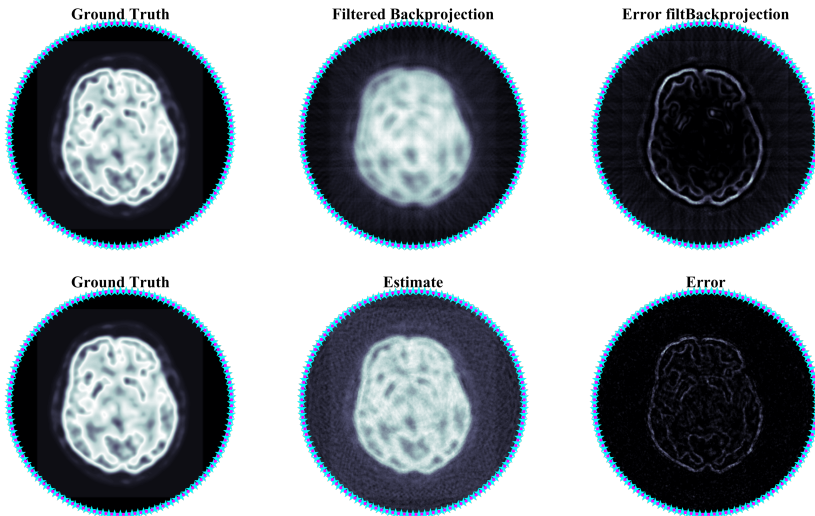
Comparison Sketching vs. FB



Comparison without/with Sketching



Comparison Sketching vs. FB



References I

- [1] Matthieu Martin Jean-Andre Simeoni.
Statistical inference in positron emission tomography.
Technical report, 2014.
- [2] Stanley R Deans.
The Radon transform and some of its applications.
Courier Corporation, 2007.