

**Measurement of time-dependent CP violation
in $B^0 \rightarrow D^{\mp} \pi^{\pm}$ decays and optimisation
of flavour tagging algorithms at LHCb**

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Beauty is the only thing that time cannot harm. Philosophies fall away like sand, creeds follow one another, but what is beautiful is a joy for all seasons, a possession for all eternity
O. Wilde

The ability to quote is a serviceable substitute for wit
W. S. Maugham

Abstract

This thesis presents the results of a time-dependent analysis of $B^0 \rightarrow D^\mp \pi^\pm$ decays using 3 fb^{-1} of proton-proton collision data collected with the LHCb detector at CERN's Large Hadron Collider during Run 1 with a centre-of-mass energy of 7 (2011) and 8 (2012) TeV. The LHCb experiment is dedicated to the study of the properties of b -flavoured hadrons, in particular CP violation in the B meson system. The Standard Model of Particle Physics describes very precisely the mechanism and the amount of CP violation expected in the Universe. However, the observed matter-antimatter asymmetry is larger by several order of magnitude compared to the predictions. This could be explained by the existence of a new source of CP violation, originating in New Physics beyond the Standard Model. The time-dependent analysis of $B^0 \rightarrow D^\mp \pi^\pm$ decays provides constraints on the angle γ of the Unitarity Triangle, one of the fundamental parameters of the Standard Model related to CP violation. Since no sizeable high-order Standard Model processes are expected to contribute, any deviation from the predictions would be an unambiguous signature of New Physics. The current experimental precision on γ is significantly lower than that of theoretical predictions. This motivates the effort for new experimental determinations of γ in order to reduce its uncertainty. The analysis of $B^0 \rightarrow D^\mp \pi^\pm$ decays, although not as sensitive as that obtained from decays of charged B mesons into $D^{(*)0} K^{(*)+}$ final states, represents an independent and uncorrelated estimation of γ that contributes to the global combination of all γ measurements. The result obtained in this thesis is more precise than previous determinations from other experiments (BaBar, Belle) using $B^0 \rightarrow D^\mp \pi^\pm$ decays. Although based on a very large sample of about half a million signal events, it is still dominated by statistical uncertainties, indicating good prospects for future improvements in precision with more data from Run 2 and beyond.

In addition to the $B^0 \rightarrow D^\mp \pi^\pm$ analysis, this thesis also summarizes the studies to improve the performances of the flavour tagging algorithms used by the LHCb collaboration to infer the flavour of neutral B mesons in time-dependent analyses. The performance of these algorithms, being correlated with the kinematics of the reconstructed particles as well as the complexity of the event (number of tracks and primary vertices), showed a significant decrease on Run 2 data (2015–2018), which were collected at a centre-of-mass energy of 13 TeV. Thanks to new implementations, these algorithms now have a performance similar to that obtained with Run 1 data.

Keywords: B physics, CKM angle γ , CP violation, flavour tagging, mixing, LHCb, LHC.

Résumé

Cette thèse présente les résultats d'une analyse des désintégrations $B^0 \rightarrow D^\mp \pi^\pm$ dans un échantillon de collisions proton-proton de 3 fb^{-1} enregistré par le détecteur LHCb durant le Run 1 du Grand Collisionneur de Hadrons (LHC) du CERN, à des énergies dans le centre de masse de 7 (2011) et 8 (2012) TeV. L'expérience LHCb a pour objet l'étude des propriétés des hadrons lourds, et en particulier de la violation de CP (CPV) dans les mésons B . Le Modèle Standard (MS) de la physique des particules décrit très précisément la CPV attendue dans l'Univers. Cependant, l'asymétrie observée entre matière et antimatière dépasse les prédictions de plusieurs ordres de grandeur. Cela pourrait s'expliquer par l'existence d'une source supplémentaire de CPV au-delà du MS. L'analyse des désintégrations $B^0 \rightarrow D^\mp \pi^\pm$ permet de contraindre l'un des paramètres fondamentaux de la CPV dans le MS : l'angle γ du Triangle d'Unitarité (UT). Une déviation par rapport aux prédictions théoriques serait un signe de Nouvelle Physique. La détermination expérimentale du paramètre γ est actuellement bien moins précise que les prédictions théoriques, ce qui justifie de nouvelles mesures de γ . Bien que l'analyse des désintégrations $B^0 \rightarrow D^\mp \pi^\pm$ soit moins sensible que celle des désintégrations $B^+ \rightarrow D^{(*)0} K^{(*)+}$, elle permet d'estimer γ de façon indépendante et ainsi de contribuer à la combinaison globale de toutes les mesures de ce paramètre. Le résultat présenté dans cette thèse est plus précis que toutes les mesures déjà obtenues par d'autres expériences (BaBar et Belle) en utilisant les désintégrations $B^0 \rightarrow D^\mp \pi^\pm$. Bien que basée sur un demi-million d'événements de signal, la mesure demeure dominée par les incertitudes statistiques, ce qui ouvre de belles perspectives pour de futures améliorations de la précision en utilisant les données du Run 2 et au-delà.

De plus, cette thèse donne un aperçu des études réalisées pour améliorer les performances des algorithmes utilisés par la collaboration LHCb pour déterminer la saveur des mésons B neutres étudiés dans les analyses temporelles. La performance de ces algorithmes, étant tributaire de la cinématique des particules reconstruites et de la complexité des événements (nombre de traces et de vertex primaires), a significativement chuté sur les données du Run 2 (2015–2018) à une énergie dans le centre de masse de 13 TeV. Grâce à de nouvelles implémentations, ces algorithmes atteignent de nouveau un niveau de performance similaire à celui sur les données du Run 1.

Mots clés : physique du meson B , angle CKM γ , violation de CP , étiquetage de saveur, oscillations, LHCb, LHC.

Sommario

Questa tesi riporta i risultati di un'analisi dei decadimenti $B^0 \rightarrow D^\mp \pi^\pm$ effettuata su 3 fb^{-1} di dati ottenuti da collisioni protone-protone, i quali sono stati collezionati dal rivelatore LHCb al Grande Collisionatore di Adroni (LHC) durante il Run 1 con un'energia nel centro di massa di 7 (2011) e 8 (2012) TeV. LHCb è dedicato allo studio delle proprietà di adroni contenenti quark b , in particolare la violazione di CP nei mesoni B . Il Modello Standard (MS) della Fisica delle Particelle descrive in modo preciso la violazione di CP (CPV) prevista nell'Universo, ma l'asimmetria osservata fra materia ed antimateria è più grande di diversi ordini di grandezza rispetto alle attese. Ciò può essere spiegato tramite nuovi meccanismi di CPV non contemplati dal MS. L'analisi dei decadimenti $B^0 \rightarrow D^\mp \pi^\pm$ consente di ottenere vincoli sull'angolo γ del Triangolo di Unitarietà (UT), un parametro fondamentale del MS legato alla CPV. Poiché nessun contributo apprezzabile è atteso dal MS, ogni deviazione dalle predizioni di quest'ultimo sarebbe un segnale di Nuova Fisica. La precisione sperimentale attuale su γ è inferiore a quella teorica; ciò motiva gli sforzi per ottenere nuove misure di γ per ridurre l'incertezza. L'analisi dei decadimenti $B^0 \rightarrow D^\mp \pi^\pm$, anche se non consente una precisione paragonabile a quella ottenuta da decadimenti $B^+ \rightarrow D^{(*)0} K^{(*)+}$, rappresenta una stima indipendente e complementare che contribuisce alla combinazione globale di tutte le misure di γ . Il risultato ottenuto in questa tesi migliora la precisione di misure precedenti realizzate da altri esperimenti (BaBar, Belle) nello stesso decadimento. Nonostante il grande numero di decadimenti analizzati (circa mezzo milione di eventi di segnale), l'incertezza statistica rappresenta il contributo dominante all'incertezza totale; ciò implica che tale incertezza potrà essere migliorata utilizzando nuovi dati a partire dal Run 2.

Oltre all'analisi fin qui descritta, questa tesi riassume gli studi effettuati per migliorare le prestazioni degli algoritmi di etichettatura utilizzati a LHCb per dedurre il sapore dei mesoni B neutri nel contesto di analisi del tempo di decadimento. Queste prestazioni, correlate alla cinematica delle particelle ed alla complessità degli eventi (numero di tracce e vertici primari), hanno mostrato un peggioramento significativo sui dati dal Run 2 (2015–2018), i quali sono stati raccolti con un'energia nel centro di massa pari a 13 TeV. Grazie a nuove implementazioni di questi algoritmi, le prestazioni sono ora paragonabili a quelle ottenute sui dati dal Run 1.

Parole chiave: fisica del mesone B , angolo CKM γ , violazione di CP , etichettatura del sapore, oscillazioni, LHCb, LHC.

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1 Introduction

1.1 The Standard Model of particle physics

The *Standard Model* (SM) of particle physics [1, 2, 3] is a *non-abelian, Yang-Mills quantum field theory* based on the $SU(3) \times SU(2) \times U(1)$ gauge symmetry group. This model provides a coherent, unified and experimentally-established picture of electromagnetic, weak and strong interactions, as well as a description of the known elementary particles (quarks, leptons, gauge bosons and Higgs boson, Fig. 1.1).

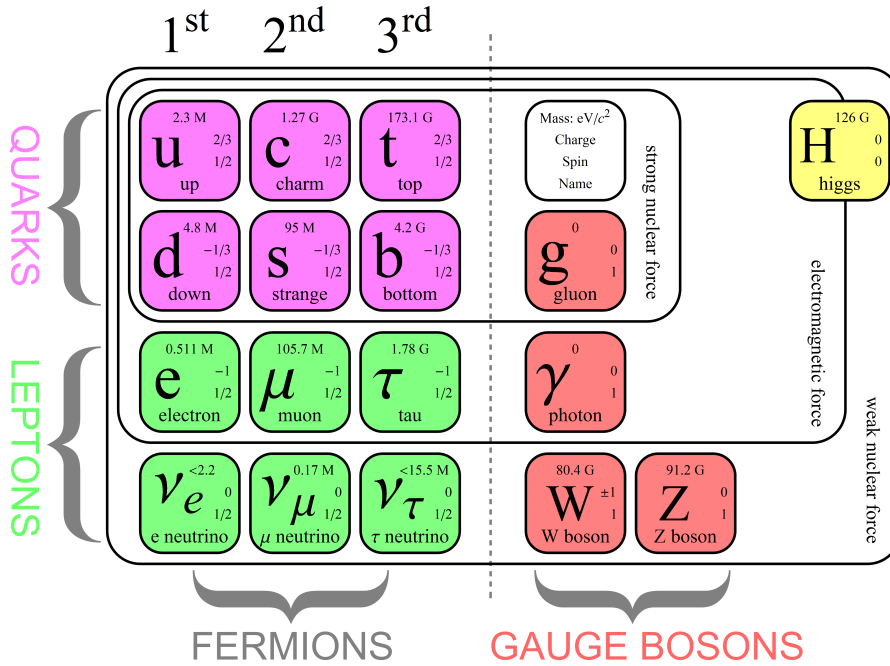


Figure 1.1 – Elementary particles described by the SM [4].

All particles are either *fermions* or *bosons*. The fermions (leptons, quarks) have half-integer spin and follow the Fermi-Dirac statistics [5, 6], whereas bosons (gauge bosons, Higgs boson) have integer spin and follow the Bose-Einstein statistics [7].

Leptons (spin- $\frac{1}{2}$) include three charged, massive particles (electron e^- , muon μ^- and tau τ^-), which interact via the electromagnetic and weak interactions, and three neutral, (nearly) massless particles, called *neutrinos* (ν_e , ν_μ and ν_τ), which only experience weak interactions.

Six different flavours of quarks (spin- $\frac{1}{2}$) exist: the *up-type* quarks up (u), charm (c) and top/truth (t), having charge¹ $+\frac{2}{3}$, and the *down-type* quarks down (d), strange (s), and bottom/beauty (b), which have charge $-\frac{1}{3}$. They can interact via electromagnetic, weak and strong interactions, and they are all massive.

¹Electric charge is always quoted in units of the fundamental charge, defined as minus the charge of the electron.

1.2. The Cabibbo-Kobayashi-Maskawa matrix

The fundamental interactions are *mediated* by gauge bosons (spin-1). The photon (γ) is responsible for the electromagnetic interaction, whereas the Z^0 and W^\pm bosons are the mediators for the weak interaction. These two forces are considered to be different manifestations of a single *electroweak* interaction, which is responsible for all electric and magnetic phenomena as well as some radioactive decays. The strong interaction among quarks is mediated by the gluons g . Photon and gluons are massless, whereas the weak force gauge bosons have a non-zero mass.

Each particle has an *antiparticle* partner, which has the same mass as the corresponding particle, but opposite quantum numbers (electric charge, lepton numbers, etc. . .).

Quarks do not exist in a free state: they can only be bound inside *hadrons* via the *confinement* mechanism, a feature of the strong interaction. A hadron can be composed by a quark-antiquark pair (*meson*), or by three quarks or antiquarks (*baryons*). Examples of mesons include the B^0 ($\bar{b}d$) and D^+ ($c\bar{d}$) mesons, whereas the proton (uud) and the neutron (udd) are examples of baryons. Recently more complex states (tetraquarks [8], pentaquarks [9]) have been evidenced.

The non-zero mass of leptons, quarks and weak force gauge bosons would require a gauge symmetry breaking term in the SM Lagrangian density. The *Brout-Englert-Higgs mechanism* [10, 11, 12] introduces a scalar (spin-0) field, called Higgs field, and a potential that allows the Higgs field to have a non-zero vacuum expectation value. This implies that the gauge symmetry is broken *dynamically*, and that the masses of the particles arise from the resulting interaction with the Higgs field. The quantum of the Higgs field is known as Higgs boson, the last SM particle discovered experimentally [13, 14].

The fourth fundamental interaction, the gravitational force, is described by another field theory, the General Relativity (GR), currently not unified with the SM.

Any experimental signature that is not described by the SM would be a hint for *new physics* (NP). Although the SM is known to be an incomplete theory because of different unsolved problems, such as dark matter, *naturalness*, matter-antimatter asymmetry, lack of SM-GR unification, etc. . . , no evidence for NP has been found so far.

1.2 The Cabibbo-Kobayashi-Maskawa matrix

The Lagrangian density describing the weak interactions between quarks and W^\pm (*charged current interaction*) can be written as

$$\mathcal{L}_{cc} = \frac{g}{\sqrt{2}} \left(\bar{u} \quad \bar{c} \quad \bar{t} \right) V_{CKM} \gamma^\mu \frac{(1 - \gamma^5)}{2} \begin{pmatrix} d \\ s \\ b \end{pmatrix} W_\mu^+ + h.c., \quad (1.1)$$

Chapter 1. Introduction

where g is a coupling constant, γ^μ are Dirac matrices and V_{CKM} , known as the Cabibbo-Kobayashi-Maskawa (CKM) matrix [15, 16], couples the *flavour* eigenstates d , s and b to the *mass* (or *physical*) eigenstates d' , s' and b' :

$$V_{\text{CKM}} = \begin{pmatrix} V_{ud} & V_{us} & V_{ub} \\ V_{cd} & V_{cs} & V_{cb} \\ V_{td} & V_{ts} & V_{tb} \end{pmatrix}. \quad (1.2)$$

The CKM matrix is unitary ($V_{\text{CKM}}^\dagger V_{\text{CKM}} \equiv 1$), so it can be written in terms of four independent parameters, namely three angles and a *complex phase* δ . The latter is the source of all *CP-violating* phenomena in the SM, i.e. asymmetries between particles and antiparticles; in fact, the *complexity* of V_{CKM} implies that the SM Lagrangian density is non *CP*-invariant, in agreement with the experimentally observed *CP* violation.

A first, standard parameterisation of the CKM matrix [17] gives

$$V_{\text{CKM}} = \begin{pmatrix} c_{12}c_{13} & s_{12}c_{13} & s_{13}e^{-i\delta} \\ -s_{12}c_{23} - c_{12}s_{23}s_{13}e^{i\delta} & c_{12}c_{23} - s_{12}s_{23}s_{13}e^{i\delta} & s_{23}c_{13} \\ s_{12}s_{23} - c_{12}c_{23}s_{13}e^{i\delta} & -c_{12}s_{23} - s_{12}c_{23}s_{13}e^{i\delta} & c_{23}c_{13} \end{pmatrix}, \quad (1.3)$$

where $s_{ij} = \sin \theta_{ij}$ and $c_{ij} = \cos \theta_{ij}$.

Another, useful parameterisation is given by *Wolfenstein* [18] and points out the order of magnitude of each matrix element. By defining the quantities λ , A , ρ and η with

$$s_{12} = \lambda = \frac{|V_{us}|}{\sqrt{|V_{ud}|^2 + |V_{us}|^2}}, \quad s_{23} = A\lambda^2 = A \left| \frac{V_{cb}}{V_{us}} \right|, \quad (1.4)$$

$$s_{13}e^{i\delta} = V_{ub}^* = A\lambda^3(\rho + i\eta), \quad (1.5)$$

the V_{CKM} matrix can be rewritten as a series expansion in powers of λ , given that λ is a small number:

$$V_{\text{CKM}} = \begin{pmatrix} 1 - \lambda^2/2 & \lambda & A\lambda^3(\rho - i\eta) \\ -\lambda & 1 - \lambda^2/2 & A\lambda^2 \\ A\lambda^3(1 - \rho - i\eta) & -A\lambda^2 & 1 \end{pmatrix} + \mathcal{O}(\lambda^4). \quad (1.6)$$

From Eq. 1.6, one can see that quark transitions within the same family (*e.g.* $u \rightarrow d$) are more probable, whereas transitions between different families (*e.g.* $b \rightarrow c$) are more suppressed. *CP* violation is a consequence of $\eta \neq 0$ and $\eta \neq \pi$.

The unitarity condition $V_{\text{CKM}}^\dagger V_{\text{CKM}} \equiv 1$ can be rewritten in terms of six scalar equations.

1.3. Physics of the neutral B mesons

Two of them are particularly relevant for the b -hadron phenomenology:

$$V_{ud}V_{ub}^* + V_{cd}V_{cb}^* + V_{td}V_{tb}^* = 0, \quad (1.7)$$

$$V_{us}V_{ub}^* + V_{cs}V_{cb}^* + V_{ts}V_{tb}^* = 0. \quad (1.8)$$

These two equations can be graphically represented as *triangles* in the $(\bar{\rho}, \bar{\eta})$ complex plane, where $\bar{\rho}$ and $\bar{\eta}$ are defined in terms of the series expansions $\bar{\rho} = \rho(1 - \lambda^2/2 + \dots)$ and $\bar{\eta} = \eta(1 - \lambda^2/2 + \dots)$, respectively. Having introduced the following angles,

$$\alpha = \phi_2 = \arg \left[-\frac{V_{td}V_{tb}^*}{V_{ud}V_{ub}^*} \right], \quad \beta = \phi_1 = \arg \left[-\frac{V_{cd}V_{cb}^*}{V_{td}V_{tb}^*} \right], \quad (1.9)$$

$$\gamma = \phi_3 = \arg \left[-\frac{V_{ud}V_{ub}^*}{V_{cd}V_{cb}^*} \right], \quad \beta_s = \chi = \arg \left[-\frac{V_{cb}V_{cs}^*}{V_{tb}V_{ts}^*} \right], \quad (1.10)$$

the triangles given by Eqs. 1.7 and 1.8 can be depicted as shown in Fig. 1.2. The first triangle, defined by Eq. 1.7, is known as the *Unitarity Triangle* (UT) and its elements can be measured from analyses of B^0 , B_s^0 and B^\pm decays. The other triangle (Eq. 1.8) can be studied from decays of B_s^0 mesons.

The amount of CP violation in the SM is given by the Jarlskog invariant J [19], which satisfies

$$\Im [V_{ij}V_{kl}V_{il}^*V_{kj}^*] = J \sum_{m,n} \varepsilon_{ikm}\varepsilon_{jln}, \quad (1.11)$$

where ε denotes the fully-antisymmetric tensor. The value of J is too small by several orders of magnitude to explain the observed matter-antimatter asymmetry in the universe, according to the baryogenesis model [20]. So, new sources of CP violation not foreseen by the SM have to exist, and thus measuring the UT with the highest possible precision is crucial to constrain these new physics scenarios.

1.3 Physics of the neutral B mesons

The theory of neutral B meson oscillation, decays and CP violation presented here is derived from Refs. [22, 23].

1.3.1 Oscillation of neutral mesons

Neutral B meson states are characterised by the following quark content:

$$|B^0\rangle = |d\bar{b}\rangle, \quad |\bar{B}^0\rangle = |\bar{d}b\rangle, \quad (1.12)$$

$$|B_s^0\rangle = |s\bar{b}\rangle, \quad |\bar{B}_s^0\rangle = |\bar{s}b\rangle. \quad (1.13)$$

1.3. Physics of the neutral B mesons

All neutral mesons will be denoted as P^0 or \bar{P}^0 hereafter. The \bar{P}^0 state is obtained from P^0 via the CP operator up to an arbitrary phase factor $e^{i\phi_{CP}}$. Since charged currents do not conserve flavour quantum numbers (e.g. strangeness, beauty etc...), a neutral meson can transform itself into its own anti-meson, and vice versa. So, the time evolution of a neutral B meson system can be generally written as

$$|\Psi(t)\rangle = a(t)|P^0\rangle + b(t)|\bar{P}^0\rangle + \sum_i c_i(t)|f_i\rangle, \quad (1.14)$$

where $|f_i\rangle$ are all the possible final states with $c_i(0) = 0$ as initial condition.

Since the typical timescale of weak interactions is much longer than the strong interaction timescale, we can neglect all weak interactions among final states (*Weisskopf-Wigner approximation*). So, we can write the Schrodinger equation for $|\Psi(t)\rangle$ in terms of an effective, non-hermitian hamiltonian \mathcal{H} :

$$i\partial_t \begin{pmatrix} a(t) \\ b(t) \end{pmatrix} = \mathcal{H} \begin{pmatrix} a(t) \\ b(t) \end{pmatrix} = \begin{pmatrix} H_{11} & H_{12} \\ H_{21} & H_{22} \end{pmatrix} \begin{pmatrix} a(t) \\ b(t) \end{pmatrix}. \quad (1.15)$$

The \mathcal{H} matrix can be rewritten as the sum of two hermitian matrices M and Γ :

$$\mathcal{H} = M - \frac{i}{2}\Gamma = \begin{pmatrix} M_{11} & M_{12} \\ M_{21} & M_{22} \end{pmatrix} - \frac{i}{2} \begin{pmatrix} \Gamma_{11} & \Gamma_{12} \\ \Gamma_{21} & \Gamma_{22} \end{pmatrix}. \quad (1.16)$$

Assuming CPT invariance ($H_{11} = H_{22} = H_0$, $M_{11} = M_{22} = M_0$, $\Gamma_{11} = \Gamma_{22} = \Gamma_0$), the eigenvalues of \mathcal{H} are

$$\lambda_L = m_L - \frac{i}{2}\Gamma_L = H_0 + \sqrt{H_{12}H_{21}} = H_0 + \sqrt{\left(M_{12} - \frac{i}{2}\Gamma_{12}\right) \left(M_{12}^* - \frac{i}{2}\Gamma_{12}^*\right)}, \quad (1.17)$$

$$\lambda_H = m_H - \frac{i}{2}\Gamma_H = H_0 - \sqrt{H_{12}H_{21}} = H_0 - \sqrt{\left(M_{12} - \frac{i}{2}\Gamma_{12}\right) \left(M_{12}^* - \frac{i}{2}\Gamma_{12}^*\right)}, \quad (1.18)$$

where L (“light”) and H (“heavy”) refer to the value of the mass for each eigenstate. The corresponding eigenvectors are

$$|P_L\rangle = p|P^0\rangle + q|\bar{P}^0\rangle, \quad |P_H\rangle = p|P^0\rangle - q|\bar{P}^0\rangle, \quad (1.19)$$

where p and q satisfy $|p|^2 + |q|^2 = 1$ and are given by

$$\frac{q}{p} = \sqrt{\frac{H_{21}}{H_{12}}} = \sqrt{\frac{M_{21} - \frac{i}{2}\Gamma_{21}}{M_{12} - \frac{i}{2}\Gamma_{12}}} = \sqrt{\frac{M_{12}^* - \frac{i}{2}\Gamma_{12}^*}{M_{21}^* - \frac{i}{2}\Gamma_{21}^*}}. \quad (1.20)$$

Chapter 1. Introduction

The relative phase ϕ_M between M_{12} and Γ_{12} is an observable quantity describing indirect CP violation (Sec. 1.3.3):

$$M_{12} = M_{12}^* e^{i\phi_{CP}}, \quad \Gamma_{12} = \Gamma_{12}^* e^{i\phi_{CP}} e^{i\phi_M}. \quad (1.21)$$

For the neutral B meson system, the ratio $|\Gamma_{12}/M_{12}|$ is expected to be small in the SM; as a consequence, it can be shown that

$$\frac{q}{p} = -e^{-i\phi_M} \left[1 - \frac{1}{2} \left| \frac{\Gamma_{12}}{M_{12}} \right| \sin \phi_M + \mathcal{O} \left(\left| \frac{\Gamma_{12}}{M_{12}} \right|^2 \right) \right], \quad (1.22)$$

which implies $|q/p| \sim 1$.

The difference and the average of the masses and widths of the two mass eigenstates are defined as

$$\Delta m = m_H - m_L = \Re(\lambda_H - \lambda_L), \quad m = \frac{m_L + m_H}{2} = \frac{\Re(\lambda_H + \lambda_L)}{2}, \quad (1.23)$$

$$\Delta\Gamma = \Gamma_L - \Gamma_H = -2\Im(\lambda_L - \lambda_H), \quad \Gamma = \frac{\Gamma_L + \Gamma_H}{2} = -\frac{\Re(\lambda_H + \lambda_L)}{4}. \quad (1.24)$$

The sign convention for $\Delta\Gamma$ is chosen to have a positive experimental value for the B_s^0 system (for B^0 , experiments give a result compatible with zero, in agreement with the SM).

The time evolution of the states $|P^0(t)\rangle$ and $|\bar{P}^0(t)\rangle$ when they are initially produced as $|P^0(0)\rangle$ and $|\bar{P}^0(0)\rangle$ can be obtained from the effective hamiltonian:

$$|P^0(t)\rangle = g_+(t)|P^0(t)\rangle + \frac{q}{p}g_-(t)|\bar{P}^0(t)\rangle, \quad (1.25)$$

$$|\bar{P}^0(t)\rangle = g_+(t)|\bar{P}^0(t)\rangle + \frac{p}{q}g_-(t)|P^0(t)\rangle. \quad (1.26)$$

The functions $g_{\pm}(t)$ are built in terms of the eigenvalues:

$$g_{\pm}(t) = \frac{1}{2} \left(e^{-im_H t - \frac{1}{2}\Gamma_H t} \pm e^{-im_L t - \frac{1}{2}\Gamma_L t} \right). \quad (1.27)$$

The probabilities that a state initially produced as P^0 or \bar{P}^0 becomes a P^0 or \bar{P}^0

at time t are

$$|\langle P^0(0)|P^0(t)\rangle|^2 = |g_+(t)|^2 = \frac{e^{-\Gamma t}}{2} \left(\cosh \frac{\Delta\Gamma t}{2} + \cos \Delta mt \right), \quad (1.28)$$

$$|\langle \bar{P}^0(0)|P^0(t)\rangle|^2 = \left| \frac{q}{p} \right|^2 |g_-(t)|^2 = \left| \frac{q}{p} \right|^2 \frac{e^{-\Gamma t}}{2} \left(\cosh \frac{\Delta\Gamma t}{2} - \cos \Delta mt \right), \quad (1.29)$$

$$|\langle P^0(0)|\bar{P}^0(t)\rangle|^2 = \left| \frac{p}{q} \right|^2 |g_-(t)|^2 = \left| \frac{p}{q} \right|^2 \frac{e^{-\Gamma t}}{2} \left(\cosh \frac{\Delta\Gamma t}{2} - \cos \Delta mt \right), \quad (1.30)$$

$$|\langle \bar{P}^0(0)|\bar{P}^0(t)\rangle|^2 = |g_+(t)|^2 = \frac{e^{-\Gamma t}}{2} \left(\cosh \frac{\Delta\Gamma t}{2} + \cos \Delta mt \right). \quad (1.31)$$

The equations above describe the *oscillation* of the B^0 and B_s^0 mesons.

1.3.2 Decay of neutral mesons

The amplitude for the decay of a neutral meson into a final state f can be obtained from the effective hamiltonian \mathcal{H} :

$$A_f = \langle f|\mathcal{H}|P^0\rangle, \quad \bar{A}_f = \langle f|\mathcal{H}|\bar{P}^0\rangle, \quad (1.32)$$

$$A_{\bar{f}} = \langle \bar{f}|\mathcal{H}|P^0\rangle, \quad \bar{A}_{\bar{f}} = \langle \bar{f}|\mathcal{H}|\bar{P}^0\rangle. \quad (1.33)$$

After defining the parameters

$$\lambda_f = \frac{q \bar{A}_f}{p A_f}, \quad \frac{1}{\lambda_{\bar{f}}} = \lambda_{\bar{f}} = \frac{q \bar{A}_{\bar{f}}}{p A_{\bar{f}}}, \quad (1.34)$$

it is possible to write the *decay rates* for neutral mesons decaying into f or \bar{f} :

$$\begin{aligned} \frac{d\Gamma(P^0 \rightarrow f)}{dt}(t) = N_f |A_f|^2 \frac{1 + |\lambda_f|^2}{2} e^{-\Gamma t} \\ \left[\cosh \left(\frac{\Delta\Gamma t}{2} \right) + D_f \sinh \left(\frac{\Delta\Gamma t}{2} \right) + C_f \cos(\Delta mt) - S_f \sin(\Delta mt) \right], \end{aligned} \quad (1.35)$$

$$\begin{aligned} \frac{d\Gamma(\bar{P}^0 \rightarrow f)}{dt}(t) = N_f |A_f|^2 \left| \frac{p}{q} \right|^2 \frac{1 + |\lambda_f|^2}{2} e^{-\Gamma t} \\ \left[\cosh \left(\frac{\Delta\Gamma t}{2} \right) + D_f \sinh \left(\frac{\Delta\Gamma t}{2} \right) - C_f \cos(\Delta mt) + S_f \sin(\Delta mt) \right], \end{aligned} \quad (1.36)$$

$$\frac{d\Gamma(P^0 \rightarrow \bar{f})}{dt}(t) = N_f |\bar{A}_{\bar{f}}|^2 \left| \frac{q}{p} \right|^2 \frac{1 + |\bar{\lambda}_{\bar{f}}|^2}{2} e^{-\Gamma t} \left[\cosh\left(\frac{\Delta\Gamma t}{2}\right) + D_{\bar{f}} \sinh\left(\frac{\Delta\Gamma t}{2}\right) + C_{\bar{f}} \cos(\Delta m t) - S_{\bar{f}} \sin(\Delta m t) \right], \quad (1.37)$$

$$\frac{d\Gamma(\bar{P}^0 \rightarrow \bar{f})}{dt}(t) = N_f |\bar{A}_{\bar{f}}|^2 \frac{1 + |\bar{\lambda}_{\bar{f}}|^2}{2} e^{-\Gamma t} \left[\cosh\left(\frac{\Delta\Gamma t}{2}\right) + D_{\bar{f}} \sinh\left(\frac{\Delta\Gamma t}{2}\right) - C_{\bar{f}} \cos(\Delta m t) + S_{\bar{f}} \sin(\Delta m t) \right], \quad (1.38)$$

where N_f is a time-independent normalisation factor and

$$D_f = -\frac{2\Re\lambda_f}{1 + |\lambda_f|^2}, \quad C_f = \frac{1 - |\lambda_f|^2}{1 + |\lambda_f|^2}, \quad S_f = \frac{2\Im\lambda_f}{1 + |\lambda_f|^2}, \quad (1.39)$$

$$D_{\bar{f}} = -\frac{2\Re\bar{\lambda}_{\bar{f}}}{1 + |\bar{\lambda}_{\bar{f}}|^2}, \quad C_{\bar{f}} = -\frac{1 - |\bar{\lambda}_{\bar{f}}|^2}{1 + |\bar{\lambda}_{\bar{f}}|^2}, \quad S_{\bar{f}} = -\frac{2\Im\bar{\lambda}_{\bar{f}}}{1 + |\bar{\lambda}_{\bar{f}}|^2} \quad (1.40)$$

are known as *CP coefficients*.

1.3.3 CP violation in neutral meson systems

Three types of *CP* violation can occur. They are briefly sketched in Fig. 1.3 and described in the following paragraphs.

CP violation in decays

CP violation in decays, also known as *direct CP* violation, happens when the decay rate for $P \rightarrow f$ is different from that of the *CP*-conjugated process $\bar{P} \rightarrow \bar{f}$:

$$\left| \frac{\bar{A}_{\bar{f}}}{A_f} \right| \neq 1. \quad (1.41)$$

This kind of *CP* violation occurs if, for each decay, at least two amplitudes with different weak (ϕ_j) and strong (δ_j) phases contribute:

$$A_f = \sum_j |A_j| e^{i(\delta_j + \phi_j)}, \quad \bar{A}_{\bar{f}} = \sum_j |\bar{A}_j| e^{i(\delta_j - \phi_j)}. \quad (1.42)$$

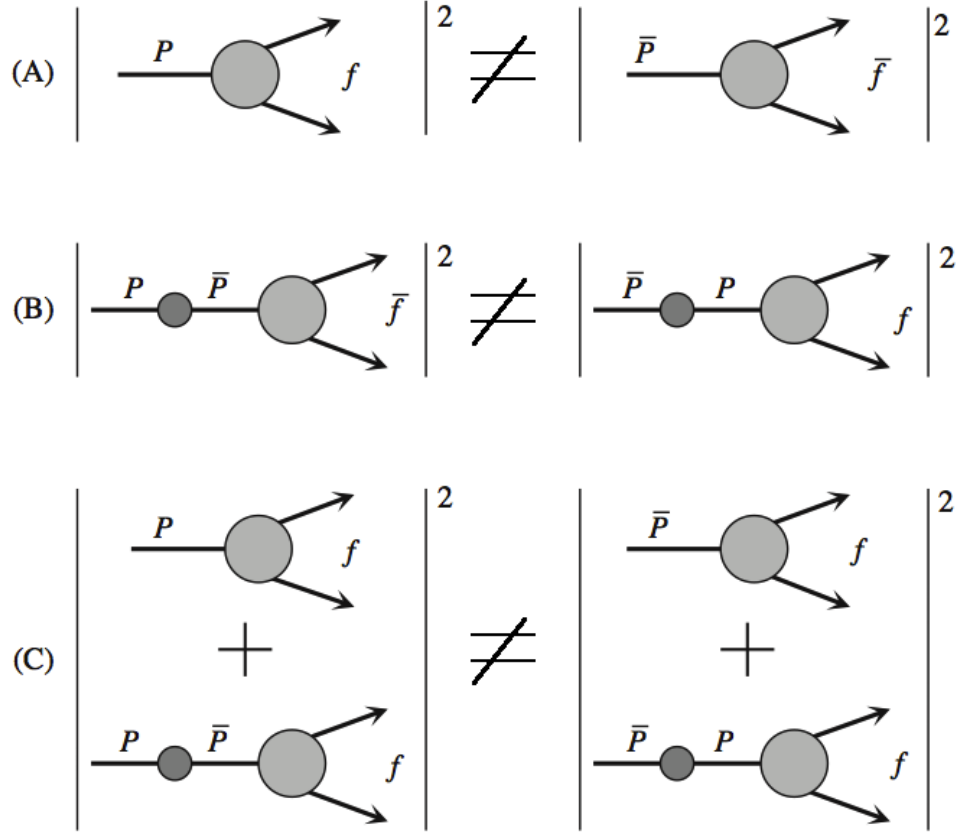


Figure 1.3 – Graphical representation of CP violation in decay (A), mixing (B) and interference between mixing and decay (C) [23].

In fact, the strong phases are invariant under CP conjugation, whereas the weak phases change sign. The following asymmetry between final states can be measured to determine direct CP violation experimentally for *charged mesons*, where mixing effects are absent:

$$\mathcal{A}_{f^\pm} = \frac{\Gamma(P^- \rightarrow f^-) - \Gamma(P^+ \rightarrow f^+)}{\Gamma(P^- \rightarrow f^-) + \Gamma(P^+ \rightarrow f^+)} = \frac{\left| \frac{\bar{A}_f}{A_f} \right|^2 - 1}{\left| \frac{\bar{A}_f}{A_f} \right|^2 + 1} \quad (1.43)$$

CP violation in mixing

CP violation in mixing, also called *indirect CP violation*, occurs when the oscillation rate for $\bar{P}^0 \rightarrow P^0$ is different from that of the CP -conjugated process $P^0 \rightarrow \bar{P}^0$. These two oscillation probabilities are given by Eqs. 1.29 and 1.30. It turns out that they are

identical unless

$$\left| \frac{q}{p} \right| \neq 1. \quad (1.44)$$

From Eq. 1.22, it can be seen that CP violation in mixing occurs when the relative phase ϕ_M is different from any multiple of π . It is possible to measure the $|q/p|$ ratio by comparing the oscillation rates in flavour-specific, semileptonic decays of neutral mesons $P^0 \rightarrow l^+ X$ and $\bar{P}^0 \rightarrow l^- X$, where no direct CP violation occurs. The decays where oscillation occurred are identified by reconstructing “wrong sign” leptons. The so-called semileptonic asymmetry

$$\mathcal{A}_{\text{SL}} = \frac{\frac{d\Gamma(\bar{P}^0 \rightarrow l^+ X)}{dt} - \frac{d\Gamma(P^0 \rightarrow l^- X)}{dt}}{\frac{d\Gamma(\bar{P}^0 \rightarrow l^+ X)}{dt} + \frac{d\Gamma(P^0 \rightarrow l^- X)}{dt}} = \frac{1 - |q/p|^4}{1 + |q/p|^4} \quad (1.45)$$

is independent of time.

CP violation in the interference between mixing and decay

This type of decay occurs when a neutral meson can decay directly to a given final state, $P^0 \rightarrow f$, or via mixing, $P^0 \rightarrow \bar{P}^0 \rightarrow f$. This can happen only if the final state f is common to both P^0 and \bar{P}^0 . This type of CP violation can occur also if other sources of CP violation (mixing or decay) are absent. In general, the interference between mixing and decay can be accessed by studying the following asymmetries:

$$\mathcal{A}_f(t) = \frac{\frac{d\Gamma(P^0 \rightarrow f)}{dt} - \frac{d\Gamma(\bar{P}^0 \rightarrow f)}{dt}}{\frac{d\Gamma(P^0 \rightarrow f)}{dt} + \frac{d\Gamma(\bar{P}^0 \rightarrow f)}{dt}}, \quad \mathcal{A}_{\bar{f}}(t) = \frac{\frac{d\Gamma(P^0 \rightarrow \bar{f})}{dt} - \frac{d\Gamma(\bar{P}^0 \rightarrow \bar{f})}{dt}}{\frac{d\Gamma(P^0 \rightarrow \bar{f})}{dt} + \frac{d\Gamma(\bar{P}^0 \rightarrow \bar{f})}{dt}}. \quad (1.46)$$

A relevant example is the case of neutral B mesons, where $|q/p| = 1$. Using Eqs. 1.35–1.38, the asymmetries of Eq. 1.46 take the following forms:

$$\mathcal{A}_f(t) = \frac{C_f \cos(\Delta mt) - S_f \sin(\Delta mt)}{\cosh\left(\frac{\Delta\Gamma t}{2}\right) + D_f \sinh\left(\frac{\Delta\Gamma t}{2}\right)}, \quad \mathcal{A}_{\bar{f}}(t) = \frac{-C_{\bar{f}} \cos(\Delta mt) + S_{\bar{f}} \sin(\Delta mt)}{\cosh\left(\frac{\Delta\Gamma t}{2}\right) + D_{\bar{f}} \sinh\left(\frac{\Delta\Gamma t}{2}\right)}. \quad (1.47)$$

The CP coefficients can be directly measured from a time-dependent analysis of certain B decays.

1.4 $B^0 \rightarrow D^{\mp} \pi^{\pm}$ analysis motivation and phenomenology

In this thesis, a decay-time dependent analysis of the decay $B^0 \rightarrow D^{\mp} \pi^{\pm}$ is presented, where the D^{\pm} meson is reconstructed as $D^{\pm} \rightarrow K^{\mp} \pi^{\pm} \pi^{\pm}$. The pion produced together with the D^{\pm} meson will be named *bachelor* or *companion* hereafter. The objective of this

1.4. $B^0 \rightarrow D^\mp \pi^\pm$ analysis motivation and phenomenology

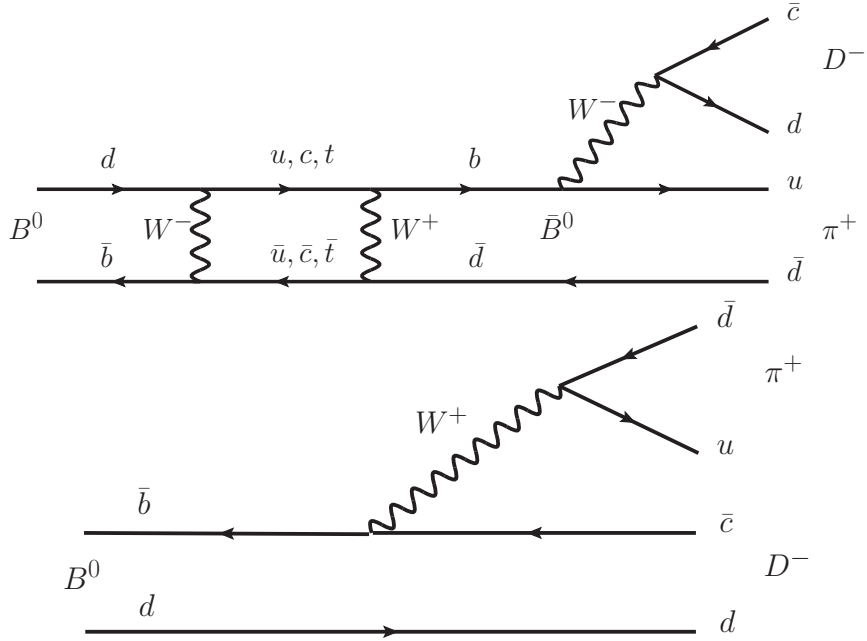


Figure 1.4 – Feynman diagrams contributing to $B^0 \rightarrow D^- \pi^+$, with (top) and without (bottom) mixing.

study is to perform a measurement of CP asymmetries, in order to constrain the CKM angle γ [15, 16]. The γ angle is particularly important to test the CKM picture of the CP violation. In fact, γ is the least-known parameter of the UT, whereas the theoretical predictions for its value are very clean and free from hadronic uncertainties [24, 25, 26]. So, improving the experimental precision on γ is a milestone of the current Flavour Physics programme. CP violation appears in the interference between the Cabibbo favoured $b \rightarrow c$ amplitude without mixing, $A(B^0 \rightarrow D^- \pi^+)$, and the Cabibbo suppressed $b \rightarrow u$ amplitude with mixing, $A(B^0 \rightarrow \bar{B}^0 \rightarrow D^- \pi^+)$. Two of the corresponding Feynman diagrams for these amplitudes are depicted in Fig. 1.4.

The measurement is performed by analysing the four time-dependent decay rates $\frac{d\Gamma(B^0 \rightarrow D^- \pi^+)}{dt}$, $\frac{d\Gamma(B^0 \rightarrow D^+ \pi^-)}{dt}$, $\frac{d\Gamma(\bar{B}^0 \rightarrow D^- \pi^+)}{dt}$ and $\frac{d\Gamma(\bar{B}^0 \rightarrow D^+ \pi^-)}{dt}$. Identifying the final state as $f = D^- \pi^+$ or $\bar{f} = D^+ \pi^-$, and assuming CPT symmetry, no CP violation in mixing ($|q/p| = 1$) and decay ($|A_f|^2 = |\bar{A}_{\bar{f}}|^2$, $|A_{\bar{f}}|^2 = |\bar{A}_f|^2$), and $\Delta\Gamma = 0$, the time-dependent

decay rates for B mesons initially produced as B^0 or \bar{B}^0 can be written as

$$\frac{d\Gamma(B^0 \rightarrow f)}{dt}(t) = \frac{1}{4\tau} e^{-t/\tau} [1 + C_f \cos(\Delta mt) - S_f \sin(\Delta mt)], \quad (1.48)$$

$$\frac{d\Gamma(\bar{B}^0 \rightarrow f)}{dt}(t) = \frac{1}{4\tau} e^{-t/\tau} [1 - C_f \cos(\Delta mt) + S_f \sin(\Delta mt)], \quad (1.49)$$

$$\frac{d\Gamma(B^0 \rightarrow \bar{f})}{dt}(t) = \frac{1}{4\tau} e^{-t/\tau} [1 + C_{\bar{f}} \cos(\Delta mt) - S_{\bar{f}} \sin(\Delta mt)], \quad (1.50)$$

$$\frac{d\Gamma(\bar{B}^0 \rightarrow \bar{f})}{dt}(t) = \frac{1}{4\tau} e^{-t/\tau} [1 - C_{\bar{f}} \cos(\Delta mt) + S_{\bar{f}} \sin(\Delta mt)], \quad (1.51)$$

where Δm and $\tau = 1/\Gamma$ are given by Eqs. 1.23 and 1.24, respectively. The CP coefficients defined in Eqs. 1.39 and 1.40, can be expressed as

$$S_f = -\frac{2r_{D\pi} \sin[\delta - (\gamma + 2\beta)]}{1 + r_{D\pi}^2}, \quad S_{\bar{f}} = \frac{2r_{D\pi} \sin[\delta + (\gamma + 2\beta)]}{1 + r_{D\pi}^2}, \quad (1.52)$$

$$C_f = -C_{\bar{f}} = C = \frac{1 - r_{D\pi}^2}{1 + r_{D\pi}^2}, \quad (1.53)$$

where β (Eq. 1.9) is related to the B^0 mixing phase,

$$r_{D\pi} = \frac{|\bar{A}_f|}{|A_f|} = \frac{|A_{\bar{f}}|}{|\bar{A}_{\bar{f}}|} \quad (1.54)$$

is the magnitude of the ratio between the doubly Cabibbo suppressed and favoured amplitudes, and δ is the strong phase difference between these amplitudes.

A measurement of γ can be obtained by measuring the CP coefficients and taking external measurements of β and $r_{D\pi}$ as input. The angle β is known with very high precision, both theoretically and experimentally [27]. An estimation of $r_{D\pi}$ was performed by the BaBar and Belle collaborations [28, 29], by measuring the branching fraction of $B^0 \rightarrow D_s^{(*)+} \pi^-$ decays and assuming SU(3) symmetry, yielding an average of about 1.7% with a relative error around 15%, mainly due to SU(3) symmetry breaking. Because of the very small value of $r_{D\pi}$, this analysis is not sensitive to C_f and $C_{\bar{f}}$; for this reason, these coefficients are simply fixed to +1 and -1, respectively.

The small value for the $r_{D\pi}$ parameter, which reduces the sensitivity on $S_{f/\bar{f}}$, makes this measurement challenging as compared to similar analyses like $B_s^0 \rightarrow D_s^\mp K^\pm$. However, the B_s^0 production rate is significantly smaller than the B^0 production fraction ($f_s/f_d = 0.259 \pm 0.015$ [30]). Moreover, the $B_s^0 \rightarrow D_s^\mp K^\pm$ branching ratio, $(2.27 \pm 0.19) \times 10^{-4}$ [31], is much smaller than the $B^0 \rightarrow D^- \pi^+$ branching ratio, $(2.52 \pm 0.13) \times 10^{-3}$ [32]. So, the resulting, larger number of $B^0 \rightarrow D^\mp \pi^\pm$ candidates compensates for this reduced sensitivity.

Measurements of $\sin(2\beta + \gamma)$ in $B^0 \rightarrow D^{(*)\mp} \pi^\pm$ were performed previously by the BaBar

and Belle collaborations [33, 34, 35, 36]. I’m one of the authors of a measurement of $2\beta_s - \gamma$ in $B_s^0 \rightarrow D_s^\mp K^\pm$ decays with 3 fb^{-1} of data [37].

The measurement presented in this thesis is performed in terms of a flavour-tagged, decay-time dependent analysis of the Run 1 dataset (Sec. 2.2.5). The dataset includes two sub-samples recorded with opposite directions of the magnetic field (“up” and “down”) in the spectrometer dipole. The selection of the decays, which is explained in detail in Sec. 4.1, includes the use of vetoes to reduce the number of components that must be modelled in the sample, and a boosted decision tree (BDT) to reduce the amount of *combinatorial* background. The expected sample composition after the selection is discussed in Sec. 4.2 based on studies with simulated samples. A fit to the invariant mass distribution of the resulting dataset is performed to extract *sWeights* [38] for the signal component. The fit is described in detail in Sec. 4.3. The training and calibration of the flavour-tagging algorithms, which infer the initial flavour of the reconstructed B^0 candidates, is summarised in Sec. 3.1. Finally, the estimation of the *CP* coefficients is the result of an unbinned, *sWeighted* likelihood fit to the distributions of the decay time and the flavour tagging observables.

1.5 Personal contribution

The analysis work presented in this thesis, published in Ref. [39] has been carried out in close cooperation with LHCb members of the Technische Universität of Dortmund. My main responsibilities related to the $B^0 \rightarrow D^\mp \pi^\pm$ analysis during my PhD time were:

- opposite-side flavour tagging calibration (Sec. 3.2.1);
- correction of particle identification (Sec. 4.2.1);
- mass fit for *sWeights* calculation (Sec. 4.3);
- time-dependent analysis, in particular acceptance parameterisation (Sec. 5.2), decay-time fit (Sec. 5.3), and estimation of systematic uncertainties (Sec. 5.5).

Moreover, I was responsible for:

- calibration and performance studies of the LHCb Silicon Tracker (Sec. 2.2.1). The results of these studies are reported in Ref. [40];
- monitoring and maintenance of the Silicon Tracker as on-call expert during data-taking operations (*piquet*);
- opposite-side electron tagger optimisation (Sec. 3.3).

2 Experimental setup

2.1 The Large Hadron Collider

The *Large Hadron Collider* (LHC) is a circular collider with a circumference of 26.66 km. It is located at CERN, near Geneva, between Switzerland and France. The LHC is designed to produce proton-proton (pp) collisions with a *luminosity* of 10^{34} $\text{cm}^{-2}\text{s}^{-1}$ and a centre-of-mass energy of 14 TeV. In the first data-taking period before the first long shutdown, called Run 1 (2010–2012), the centre-of-mass energy reached 7 TeV (2010–2011) and 8 TeV (2012).

The proton bunches, produced from hydrogen gas, pass through different intermediate accelerating stages (Fig. 2.1):

- LINAC 2 (50 MeV);
- Proton Synchrotron Booster (1.4 GeV);
- Proton Synchrotron (25 GeV);
- Super Proton Synchrotron (450 GeV).

Finally, they are injected clockwise and counter-clockwise into the LHC and accelerated to their final energy. Each bunch contains $\sim 10^{11}$ protons, and the nominal number of bunches per beam is 2808. At LHC, in addition to LHCb, there are two general-purpose detectors (ATLAS and CMS), a detector dedicated to quark matter and quark-gluon plasma physics (ALICE) and other smaller experiments (TOTEM, LHCf, MoEDAL) dedicated to different topics.

The LHC can also accelerate particles other than protons, such as lead or xenon nuclei, in order to collect data samples for specific studies.

2.2 The LHCb experiment

The *Large Hadron Collider beauty* (LHCb) experiment [41] is a single-arm forward spectrometer (see Fig. 2.2) that exploits the forward production of the b - and c - quark pairs in pp collisions (Fig. 2.3). The LHCb angular coverage is comprised between 15 mrad and 250 (300) mrad in the vertical (horizontal) plane. The LHCb coordinate system consists of a right-handed set of axes, x , y , z , where the positive z direction extends into the LHCb detector, y is perpendicular to the LHCb cavern ground and x is orthogonal to the other two.

The LHCb experiment is composed of different sub-detectors. The tracking system includes a vertex and tracking detector called *VErtex LOcator* (VELO), the *Tracker Turicensis* (TT), located upstream a magnetic dipole with an integrated field of 4 Tm, the *Inner Tracker* (IT), situated downstream the magnet in three separated stations around the beryllium beam pipe, and the *Outer Tracker* (OT), installed in the same stations as the IT. The *Particle IDentification* (PID) system comprises two *Ring Imaging CHerenkov*

2.2. The LHCb experiment

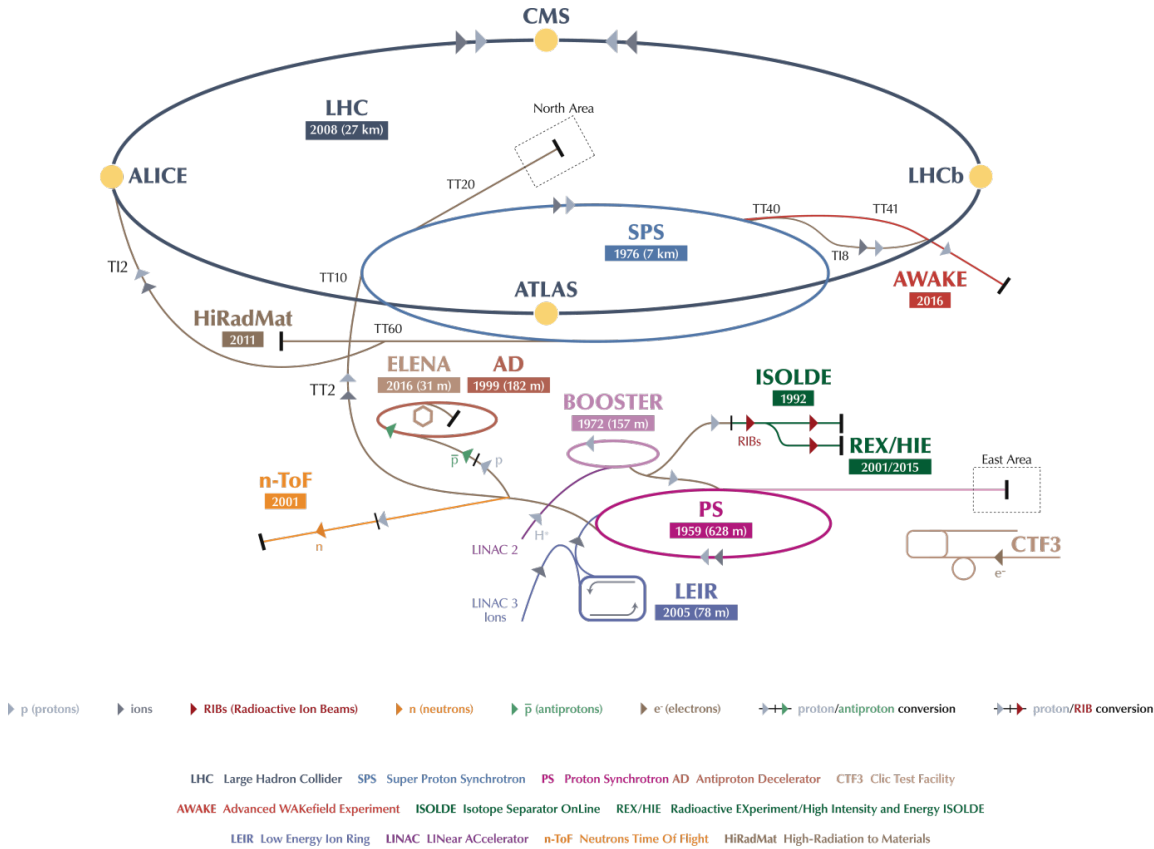


Figure 2.1 – Overview of the CERN accelerators complex.

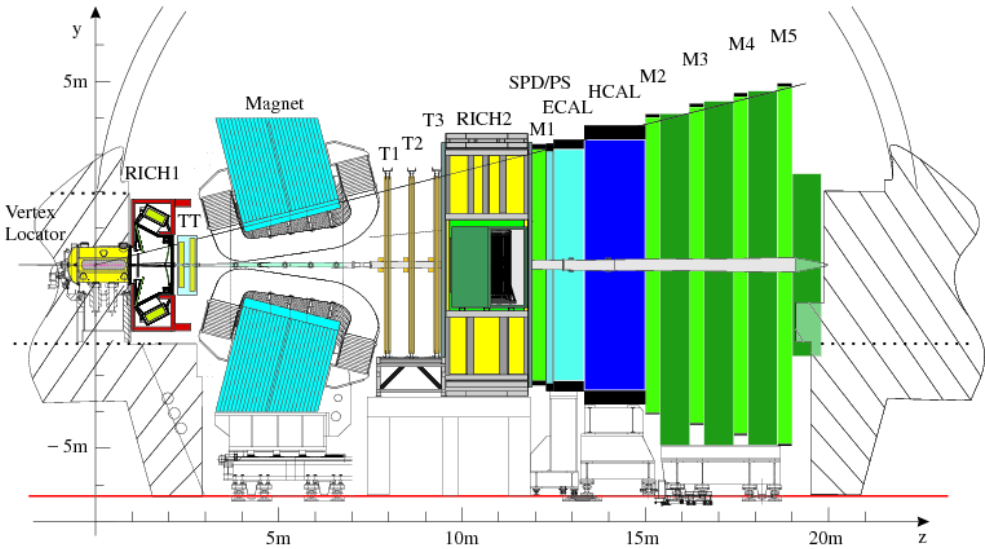


Figure 2.2 – Side view of the LHCb detector.

Chapter 2. Experimental setup

detectors (RICH), an *Electromagnetic CALorimeter* (ECAL), which also includes a *Pre-Shower* (PS) and *Scintillator Pad Detector* (SPD), a *Hadronic CALorimeter* (HCAL) and five *muon stations* (M1–M5).

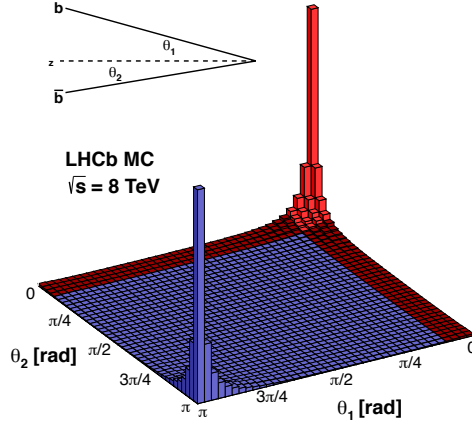


Figure 2.3 – Joint distribution of the b and \bar{b} production angles with respect to the beam direction at $\sqrt{s} = 8$ TeV, as obtained from simulation. The red part shows the LHCb acceptance.

2.2.1 Tracking system

The tracking system plays a crucial role for the time-dependent analysis of $B^0 \rightarrow D^{\mp} \pi^{\pm}$ decays at LHCb. The excellent vertex and impact-parameter resolution allows to separate true, long-lived B^0 mesons from combinations of other random tracks, and to improve the decay-time resolution. Moreover, the optimal momentum resolution implies an excellent resolution on the reconstructed invariant mass, which is crucial to separate true $B^0 \rightarrow D^{\mp} \pi^{\pm}$ decays from other physical background.

The tracking system is also essential for the flavour tagging algorithms, which rely on the quality of the reconstructed tracks and vertices to discriminate correctly tagged neutral B mesons from wrongly tagged ones.

A summary of the performances of the tracking system (momentum, impact-parameter, and decay-time resolution) is shown in Fig. 2.4.

The VERtEx LOCator (VELO)

The VELO [42] (Fig. 2.5) is a silicon micro-strip detector surrounding the interaction point, which detects charged particles, performs the first track reconstruction step and

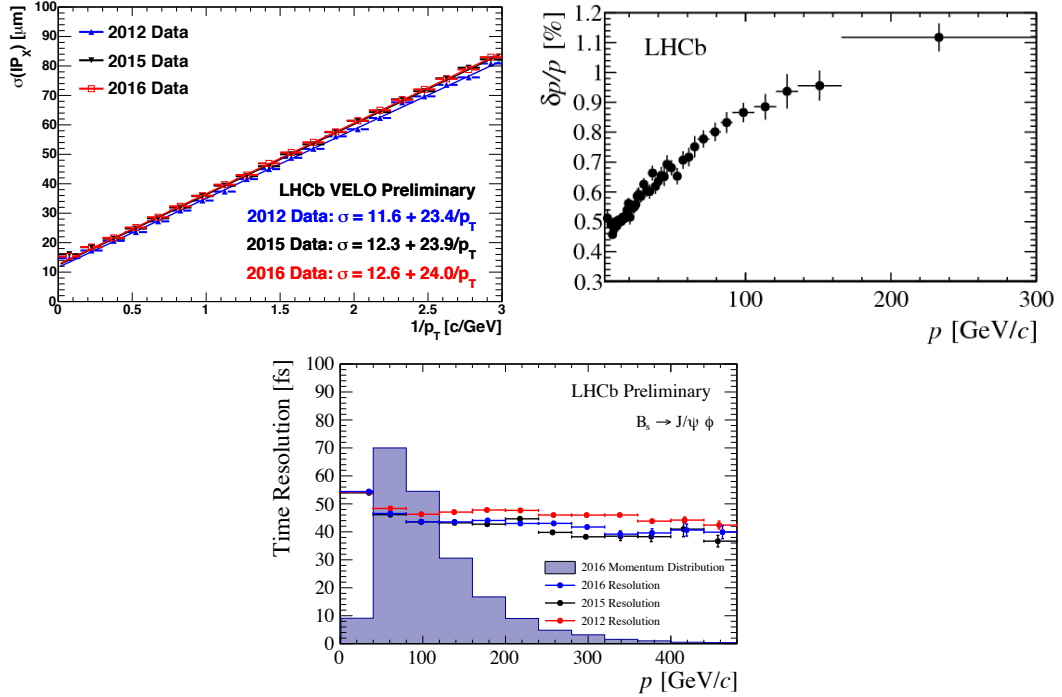


Figure 2.4 – Top left: resolution on the x coordinate of the impact parameter as a function of the inverse of the reconstructed transverse momentum in 2012 (blue), 2015 (black) and 2016 (red) data. The result of a linear fit for each dataset is superimposed. Top right: relative momentum resolution as a function of the reconstructed momentum. Bottom: decay-time resolution as a function of momentum for reconstructed $B_s^0 \rightarrow J/\psi\phi$ decays in 2012 (red), 2015 (black) and 2016 (blue) data. The momentum distribution is superimposed as the solid, purple histogram.

identifies decay vertices. The sensitive region of the VELO is composed of n-on-n silicon micro-strip half-disk sensors with two different read-out strip geometries, called r -type and ϕ -type, which measure the radial (r) and azimuthal (ϕ) position in polar coordinates. The silicon sensors are 8.4 cm in diameter and have an inner hole with radius 0.8 cm. The strip pitch ranges from 38 to 108 μm (38 to 97 μm) for r (ϕ) sensors, while the sensor thickness is 300 μm . The VELO consists of 21 stations placed perpendicular to the beam axis. Each station has two independent halves that can be moved apart during beam injection and then closed again when the beam orbit is stabilised. Each half-station is composed by one r -type and one ϕ -type sensor. The total length of the VELO detector is about 1 m. The impact parameter (IP) resolution of a track is measured to be $\sigma_{\text{IP}} = 11.6 \pm 23.4/p_{\text{T}} \mu\text{m}$ in x and $\sigma_{\text{IP}} = 11.2 \pm 23.2/p_{\text{T}} \mu\text{m}$ in y , where p_{T} is the *transverse momentum* (in GeV/c) of the particle with respect to the beam axis.

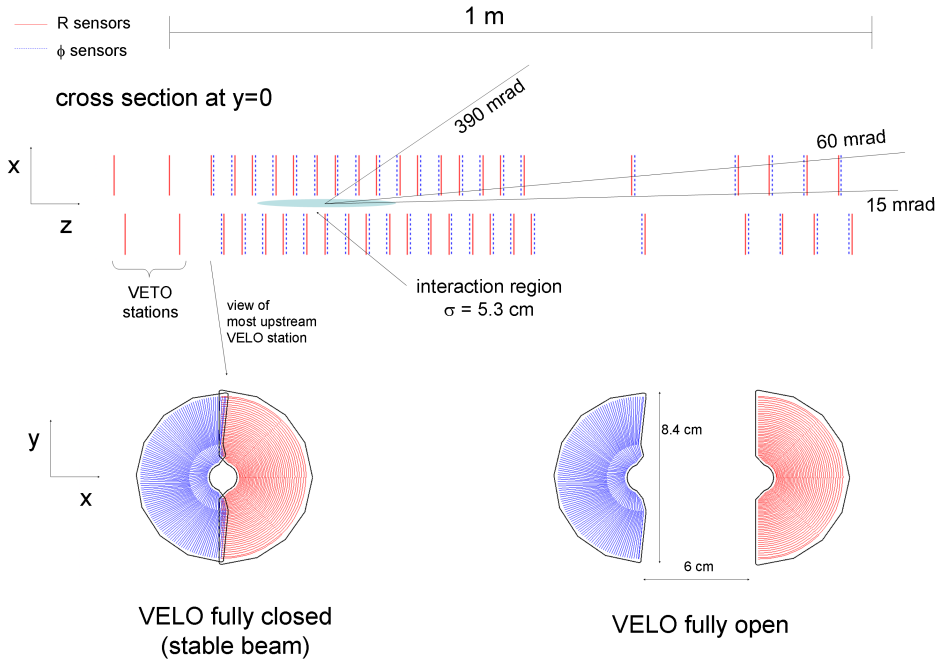


Figure 2.5 – Schematic view of the VELO detector (top) and its sensors (bottom).

The Tracker Turicensis (TT)

The TT [43] (Fig. 2.6) is a silicon micro-strip detector covering a total area of about 7.9 m^2 upstream the magnet and divided into two separate stations (TTa, TTb). Each station has two layers. The TT helps in improving the track momentum resolution and detecting long-lived particles that decay outside the VELO acceptance. TTa is composed of a X and an U layer, while TTb includes a V and an X layer. The X layers have readout strips aligned vertically, whereas the U and V *stereo* layers are rotated by $+5^\circ$ and -5° with respect to the vertical in the xy plane.

The TT active area is made of p-on-n silicon micro-strip sensors. Since the sensors are exposed to a significant radiation dose due to the high track multiplicity, they are cooled to 0°C in order to minimise the effect of reverse annealing due to radiation damage.

A TT readout module contains from one to four sensors connected in series, resulting in read-out strips up to 37 cm long. The strip pitch is $183 \mu\text{m}$ and the sensor thickness is $500 \mu\text{m}$. The hit resolution is about $50 \mu\text{m}$.

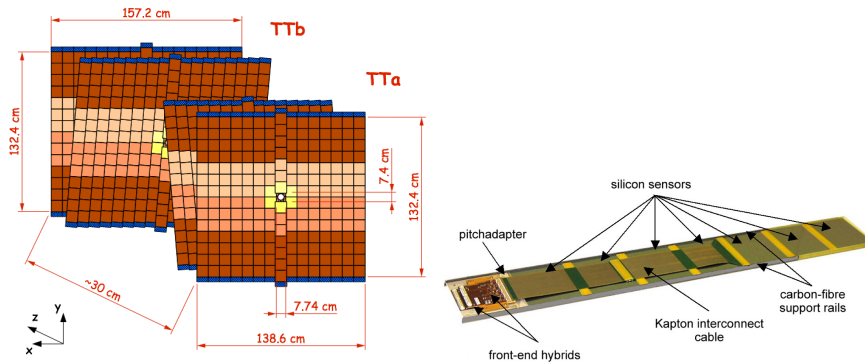


Figure 2.6 – Schematic view of the TT stations/layers (left) and one of the TT readout modules (right).

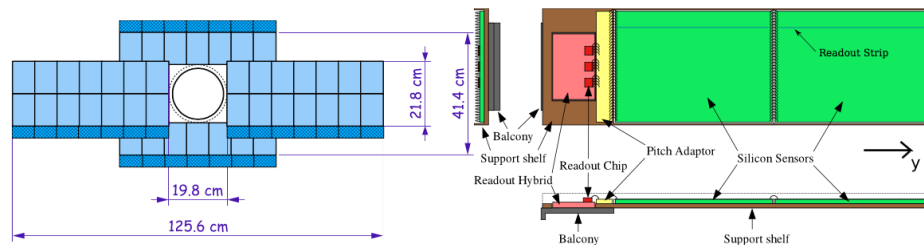


Figure 2.7 – Schematic view of an IT station (left) and one of the long IT readout modules (right).

The Inner Tracker (IT)

The IT [44] (Fig. 2.7) is also a silicon micro-strip detector. Together with the TT, it forms the *Silicon Tracker* (ST). It is dedicated to detect charged particles in the high track-density region around the beam pipe downstream the magnet. It is separated into three stations, where each station consists of four boxes. Each box has four layers made of seven read out modules arranged in a X-U-V-X layout similar to that of the TT. The total coverage of the IT is about 4.2 m². The boxes directly above and below the beam pipe are made of single-sensor modules, called *short modules*, whereas the side boxes are made of two bonded silicon sensor modules, called *long modules*. The IT strip pitch is 198 μm , while the p-on-n sensor thickness is 320 (410) μm for the short (long) modules. The hit resolution is about 50 μm .

The Outer Tracker (OT)

The OT [45] (Fig. 2.8) is a gaseous straw-tube detector filled with an Ar/CO₂/O₂ (70%/28.5%/1.5%) gas mixture. It is dedicated to the detection of charged particles in

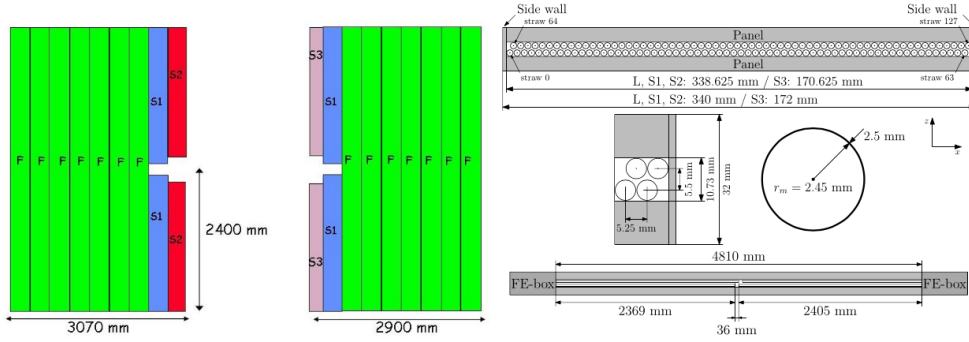


Figure 2.8 – Schematic view of an OT layer (left) and an OT module layout (right).

the low track density region outside the IT acceptance and covers a large area of about 340 m^2 . The OT is composed of three stations, where each station has four layers in a X-U-V-X configuration. Each station is separated physically in left and right sides with respect to the beam pipe mounted in support structures called C frames because of their shape. Each layer is divided into two mono-layers. The OT has different types of modules, the long F modules and the short S1, S2, S3 modules that are cut in two pieces to leave space for the IT. The straw tube and anode wire diameters are 5 mm and $25 \mu\text{m}$ respectively. The hit resolution is about $200 \mu\text{m}$.

Spillover noise in the Silicon Tracker

Starting from 2015, the time spacing between the LHC proton *bunches* has been 25 ns, half of the value adopted before. This had a direct impact on the front-end electronics of the silicon detectors (VELO, TT, IT), because the width in time of the analogue signal produced by the front-end electronics is of the same order of magnitude. This means that it is possible to still have a non negligible amount of signal in the subsequent collision, which can be misidentified as coming from particles produced in that event. This source of noise is called *spillover*. The starting seeds of tracking algorithms, called *clusters*, can be polluted by spillover clusters which may increase the number of fake (or *ghost*) reconstructed tracks.

In the first part of my PhD activity, I studied the effect of spillover clusters in the ST using both simulated events and real collision data. This study shows that this time spacing has little impact on the detector *occupancy*, and that the increase of ghost tracks is negligible. Moreover, it is shown that the charge deposited by particles in the detector can be exploited as a *feature* in multivariate analyses in order to further reduce the ghost track contamination. These results are documented in a note [40] internal to the LHCb collaboration.

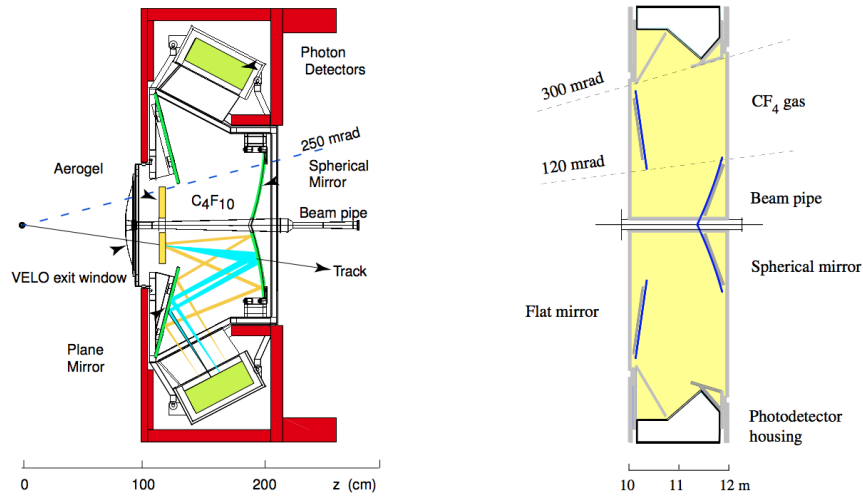


Figure 2.9 – Side view of RICH1 (left) and top view of RICH2 (right).

2.2.2 Particle identification (PID)

The Ring Imaging CHerenkov (RICH) detectors

When a charged particle is travelling faster than the speed of light in a medium, Cherenkov light is produced at an angle that depends on the velocity of the particle and the refractive index of the medium (*radiator*). By knowing the momentum from the tracker and the velocity from the RICH detectors, the mass can be determined and therefore provide particle identification. Two RICH detectors [46] (Fig. 2.9) are used in order to provide PID in different momentum ranges.

RICH1 is responsible for providing PID in the momentum range from 1 to 60 GeV/c . The angular acceptance ranges from 25 to 50 (300) mrad in the vertical (horizontal) plane. The adopted radiator is fluorobutane (C_4F_{10}). RICH1 is located between the VELO and the TT. The Cherenkov photons are guided to Hybrid Photon Detectors (HPD) via dedicated mirrors.

The average kaon identification efficiency in the momentum range from 2 to 100 GeV/c is $\sim 95\%$. The average probability that pions are wrongly identified as kaons is $\sim 5\%$ in the same momentum range.

RICH2 is optimised for the momentum range from 15 to 100 GeV/c . The angular acceptance ranges from 15 to 100 (120) mrad in the vertical (horizontal) plane where most of the high-momentum tracks are produced. RICH2 uses tetrafluoromethane (CF_4) as radiator.

The Electromagnetic CALorimeter (ECAL), Pre-Shower (PS) and Scintillator Pad Detector (SPD)

The ECAL [47] is used for the detection and measurement of the energy of electrons and photons. The ECAL is built as a sandwich of alternating scintillators and lead layers in the xy plane. Scintillation light produced by the shower of particles generated by the lead plates is read out by Wave-Length Shifter (WLS) fibres coupled to PhotoMultiplier Tubes (PMTs). The SPD is installed upstream the ECAL to separate neutral pions from photons. The PS is installed between the SPD and the ECAL. Both SPD and PS use scintillator pads read out by WLS fibres coupled to MultiAnode PhotoMultiplier Tubes (MAPMT). The acceptance range of the ECAL is from 25 up to 300 (250) mrad in the horizontal (vertical) plane. The relative energy resolution of the ECAL is given by $\sigma_E/E = 10\%/\sqrt{E} \oplus 1\%$, where E is given in GeV.

The Hadronic CALorimeter (HCAL)

The HCAL [47] is used for the detection and measurement of the energy of hadrons for the first level trigger. A HCAL cell is a sampling device made of alternating iron and scintillator tiles, where the latter are located along the beam direction. The HCAL has the same acceptance coverage of the ECAL. The relative energy resolution of the HCAL is given by $\sigma_E/E = (69 \pm 5)\%/\sqrt{E} \oplus (9 \pm 2)\%$, where E is given in GeV.

Muon detectors

The muon system [48] (Fig. 2.10) is a gaseous detector composed of five stations (M1 to M5) interleaved by 80 cm thick iron filters. The gaseous detectors are Multi-Wire Proportional Chambers (MWPC), except for the innermost part of M1, where triple Gas Electron Multipliers (GEM) detectors are used to cope with the higher track density. The angular acceptance ranges from 20 (16) to 308 (256) mrad in the horizontal (vertical) plane. The muon detector has 1380 chambers and covers a total area of 435 m². Each muon chamber is composed of four layers of MPWC, except for M1, where two layers are used. The hit efficiency of the chambers is higher than 99% enabling a trigger efficiency greater than 95% for muons. The adopted gas mixture (40% Ar, 55% CO₂, 5% CF₄) allows a fast triggering on muons (40 MHz).

2.2.3 Trigger system

The bunch crossing rate of LHC is very high (40 MHz) because more than 99% of the pp collisions do not produce interesting events. It is not possible to record events with such a high rate: therefore, a trigger system [49] is required to reduce the rate from 40 MHz down to a few kHz. The rate reduction is achieved via selection criteria which

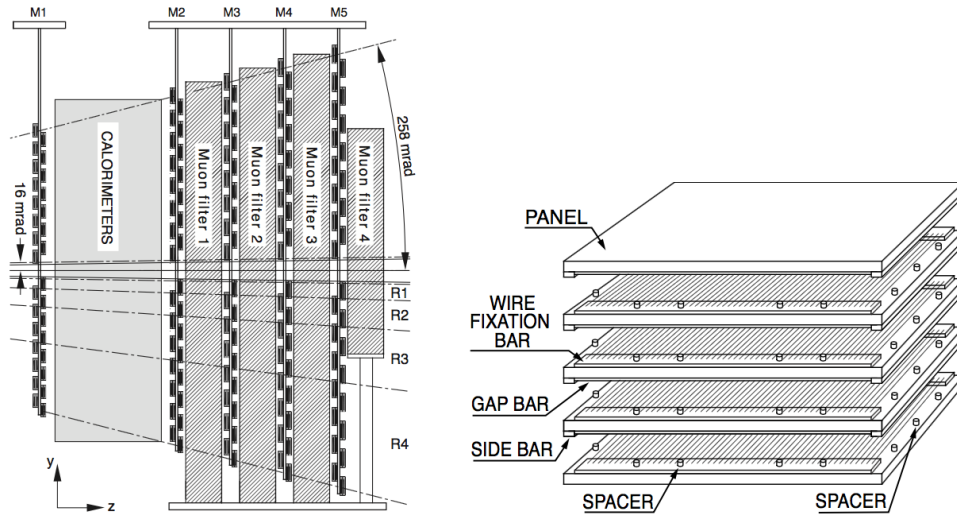


Figure 2.10 – Schematic view of the muon system (left) and a MWPC (right).

ensure that events containing heavy flavour decays are stored. The signatures of these interesting decays include high transverse momentum (p_T) and transverse energy (E_T) of the decay products, as well as displaced decay vertices or tracks with large impact parameter (IP) with respect to the pp collision point due to the relative long lifetimes of b and c hadrons. The trigger is divided into two sequential stages: a hardware stage called Level-0 (L0) trigger, and a software stage called High Level Trigger (HLT). Different trigger decisions are separated into various *lines*, each of which provides information on different physics processes (*e.g.* decay topology, presence of muons etc. . .). All the trigger steps are summarised in Fig. 2.11.

Two types of trigger response are assigned offline, when some physics channel is analysed. The TOS (*Trigger On Signal*) trigger occurs when the presence of the signal is sufficient to have a positive trigger decision. The TIS (*Trigger Independent of Signal*) trigger occurs when, after removing signal tracks and hits associated to them, another signature in the event is sufficient to have a positive trigger decision.

After the trigger stage, the data go through further offline selection steps, where exclusive (*e.g.* $B^0 \rightarrow D^\mp \pi^\pm$, $B^\pm \rightarrow J/\psi K^\pm$) and inclusive (*e.g.* b hadron $\rightarrow J/\psi X$, $J/\psi \rightarrow \mu^+ \mu^-$) decays are reconstructed at higher quality than was possible in the strict timing requirements of the trigger, and further selections are applied. This offline selection step is known as *stripping*, and each set of selection requirements is called *stripping line*.

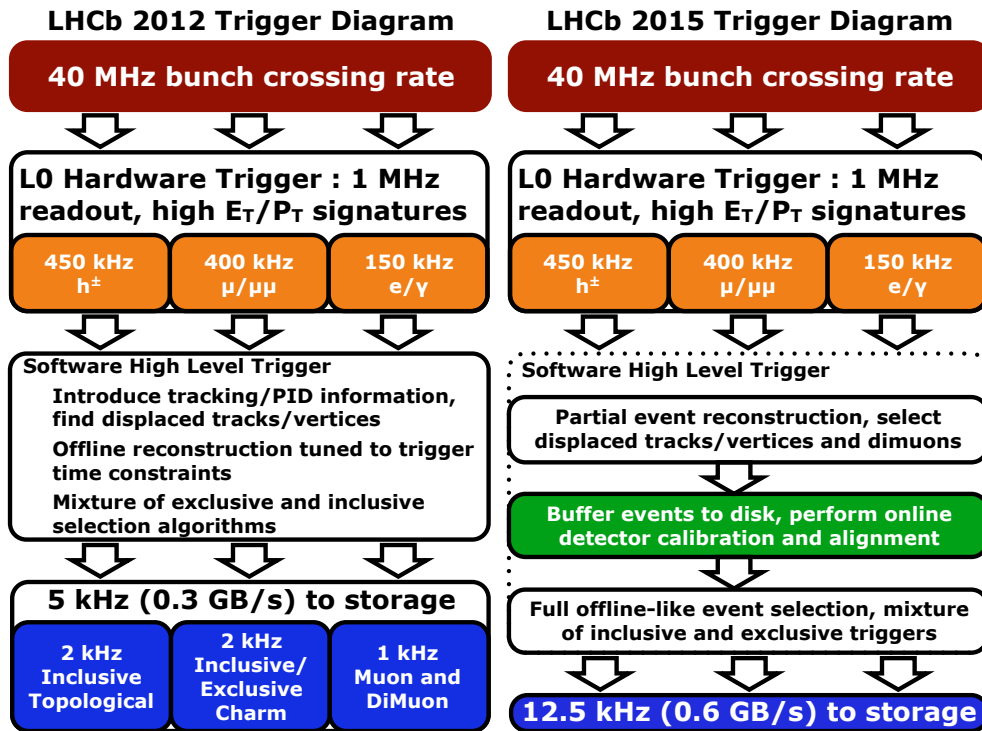


Figure 2.11 – Summary of the trigger strategies followed during the Run 1 (2011–2012, left) and Run 2 (2015–2017, right) data-taking periods. During Run 2, an online detector and calibration alignment was introduced, as well as full event selections (both inclusive and exclusive).

Level-0 trigger (L0)

The L0 mainly exploits the calorimeters and muon chambers. The idea behind the L0 is to select events that contain high p_T muons and high E_T hadrons, electrons and photons, which very likely come from b - and c -hadron decays, by using simple signatures that do not require the reconstruction of the event. The L0 reduces the data rate from 40 MHz down to 1 MHz, which is the rate limit of the LHCb readout electronics.

High Level Trigger (HLT)

The HLT is separated into two stages, HLT1 and HLT2, and runs on about 29 000 commercial CPU cores.

At the HLT1 level, the full detector information is read out, and then track/vertex reconstruction and PID are performed. The exploited signatures are mainly the presence of high p_T tracks, high transverse energy calorimeter clusters (photons and π^0), high

invariant mass of muon pairs, and tracks with large IP. All the HLT1 trigger lines are *inclusive*, meaning that only decay products common to various decay processes are selected rather than specific decays. After the HLT1, the rate goes down to about 70 kHz.

The HLT2 is a combination of inclusive selections and algorithms that reconstruct entirely (*exclusively*) some decay processes. The main lines are topological lines using Multi-Variate Analysis (MVA) techniques with different sets of kinematic and position features as input, exclusive charm lines and high mass displaced di-hadron/lepton lines. After the HLT2, the events are stored on tape for further offline analysis.

2.2.4 Event reconstruction and simulation

Track and vertex reconstruction

Starting from the *hits* in the tracking detectors, tracks and vertices are reconstructed via dedicated algorithms. Different track types are reconstructed, as shown in Fig. 2.12. Each track is characterised by hits collected in different sub-detectors. For example, downstream tracks, with no hits in the VELO, are typically associated to long-lived particles such as Λ and K_S^0 . Because of the vertical magnetic field, tracks are bend in the xz plane. By knowing the reconstructed particle trajectory and the magnetic field map, it is possible to measure the momentum of the particle.

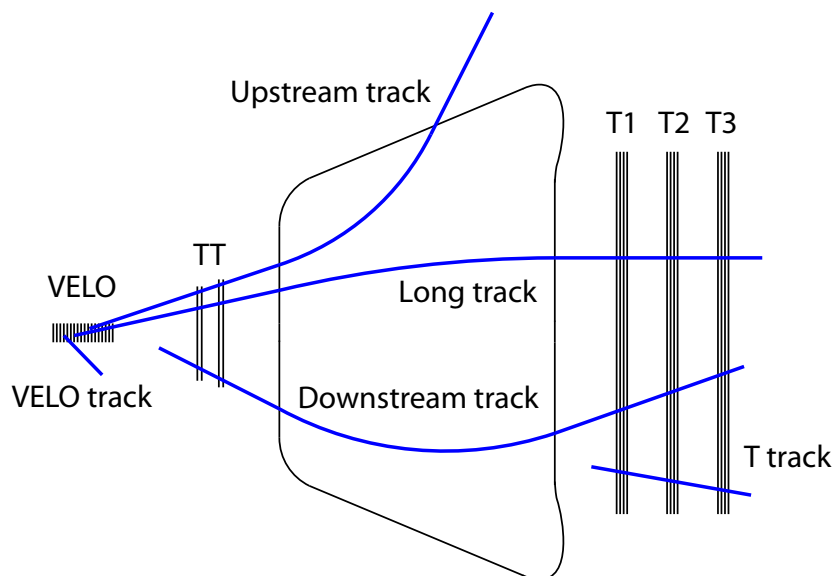


Figure 2.12 – Schematic description of the different track types reconstructed in LHCb.

PID

The Cherenkov radiation emitted by charged particles in the RICH radiators produces *rings* in the RICH acceptance, which are reconstructed via pattern recognition algorithms. For each reconstructed pattern, the *likelihood* \mathcal{L}_π for the ring to be produced by a pion (the most common particle in the LHCb environment) is computed. The momentum of the particle is also used in the likelihood computation. Then, likelihood functions for other hypotheses (kaon, proton, electron, muon) are computed and compared with the pion likelihood. For a given particle X ($X=e, \mu, p, K$), the PIDX observable is defined as

$$\text{PIDX} = \ln \mathcal{L}_X - \ln \mathcal{L}_\pi. \quad (2.1)$$

In typical LHCb analyses, requirements on the PID observables are applied to suppress physical backgrounds due to wrongly-identified particles. The PID observables are adopted, together with other kinematic and detector-related observables, as input feature for neural networks which predict the probability for a particle to be either an electron, a muon, a proton, a kaon, or a pion. This probability is denoted as PROBNNX ($X=e, \mu, p, K, \pi$).

Simulation

The Monte Carlo (MC) simulation of pp collisions, particle decays, and interactions with the detector are crucial in the validation of physics analyses. The parton-parton collision and hadronisation simulation is performed by `PYTHIA` [50], interfaced to `EvtGen` [51] for the decay of the hadrons and leptons for standard productions. The QED corrections to the decay (*i.e.* the emission of radiation photons) is generated by the `PHOTOS` package [52]. The interactions of particles with the detector material and their tracking in the magnetic field are simulated by `GEANT4` [53, 54].

2.2.5 Data collected by LHCb

The collision rate R [s^{-1}] in LHC can be expressed in terms of the *cross-section* σ [cm^2] and the *luminosity* \mathcal{L} [$\text{s}^{-1}\text{cm}^{-2}$] as

$$R = \mathcal{L}\sigma. \quad (2.2)$$

For a given data-taking period, the *integrated luminosity* L [cm^{-2}] is a measure of the amount of recorded data. The typical unit for luminosity is the inverse *barn*, which corresponds to 10^{24} cm^{-2} . The LHCb integrated luminosity is of the order of an inverse *femtobarn* (fb^{-1}); one inverse femtobarn corresponds to the production of about $10^{11} b\bar{b}$ quark pairs.

The LHCb detector collected data produced mainly from pp collisions in the 2010–2017 period, so far. During the 2011–2012 data-taking period, called *Run 1*, about 3 fb^{-1} of

2.2. The LHCb experiment

data were collected. The centre-of-mass energy \sqrt{s} of the pp interactions was 7 TeV and 8 TeV in 2011 and 2012, respectively, and the time spacing between *bunches* of protons in the LHC was 50 ns. The 2013–2014 period, known as *Long Shutdown 1* (LS1), was dedicated to work on the LHC machine to bring the energy up. The *Run 2* data-taking period started in 2015, and is planned to last until the end of 2018. The centre-of-mass energy of the pp collisions during Run 2 is 13 TeV, whereas the time spacing between proton bunches is 25 ns. In the 2015–2017 data taking period, about 3.7 fb^{-1} of data were collected. The data collected during each year is summarised in Fig. 2.13.

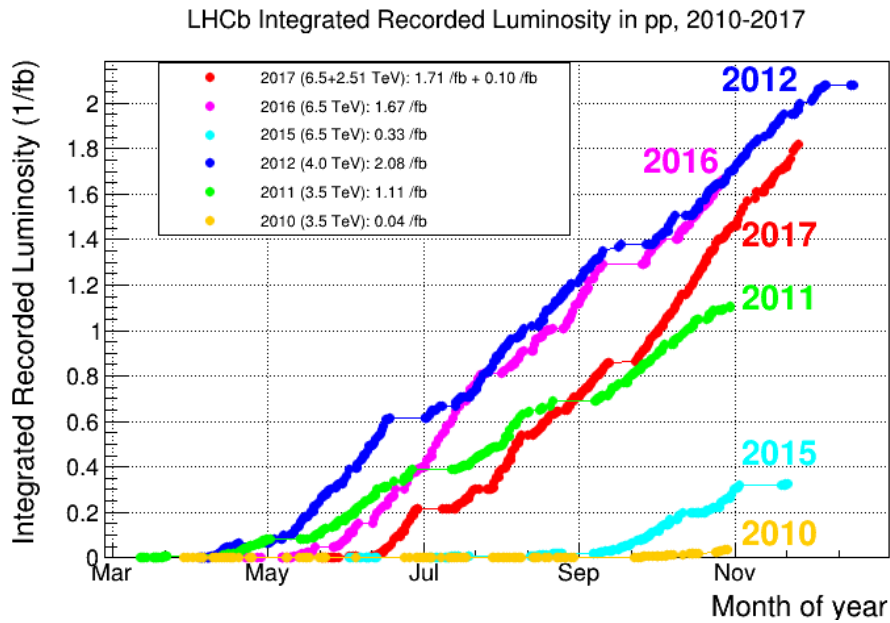


Figure 2.13 – Integrated luminosity of pp collisions collected each year by LHCb.

3 Flavour tagging

3.1 Flavour tagging algorithms

In this chapter, a description of the *flavour tagging* techniques at LHCb is reported. After a brief introduction to the methods, the calibration of the *Opposite Side* (OS) and *Same Side* (SS) algorithms for the $B^0 \rightarrow D^\mp \pi^\pm$ time-dependent analysis are described. Finally, a reoptimisation of the OS *electron* (OSe) tagger on both Run 1 and Run 2 data is reported. This work was made in collaboration with the University of Dortmund. During my PhD work activity, I focused mainly on the OS calibration and OSe reoptimisation.

The identification of the flavour at production time of a neutral B meson, *i.e.* whether it contained a b or a \bar{b} quark at production, is a key element for any time-dependent analysis aiming at the measurement of oscillations and CP asymmetries, and in particular for the $B^0 \rightarrow D^\pm \pi^\mp$ analysis reported in this thesis. In fact, this information is needed in order to measure the decay rates or asymmetries introduced in Sec. 1.3. Complications arise from two facts:

- neutral B meson oscillate, so the flavour at the production time might differ from the flavour at the decay time;
- there are many final states, such as $D^\mp \pi^\pm$, that can be produced by the decay of both B and \bar{B} mesons; so, the flavour cannot be obtained from the charge of the final state, even if there were no oscillations.

For these reasons, the flavour has to be reconstructed by exploiting particles not produced in the decay of the neutral B meson being analysed.

Techniques to infer the initial flavour of a reconstructed candidate are called flavour tagging algorithms. Several flavour tagging algorithms exist in LHCb; they can be categorised into same-side taggers (SS taggers) and opposite-side taggers (OS taggers). A schematic representation of the taggers that can be used for tagging B^0 mesons is shown in Fig. 3.1.

The SS taggers infer the production flavour of the signal B meson by selecting charged particle candidates that have a high chance of being remnants of the hadronisation process of the B candidate [55]. For B^0 mesons, the same-side pion tagger (SS π), which exploits charged π mesons produced in the hadronisation of the B^0 meson, and the same-side proton tagger (SS p), which looks for co-produced protons, have been developed. For both taggers, the charge of the pion or proton is correlated with the production flavour of the signal B^0 meson. The response of the two taggers is combined into a common SS response.

In contrast, the OS taggers exploit the predominant production process of B mesons via $b\bar{b}$ quark pair production [56]. They partially reconstruct the decay of the *other* b hadron produced along with each reconstructed signal B meson, and infer its initial flavour. In fact, the flavour of the signal B meson and the other b hadron produced

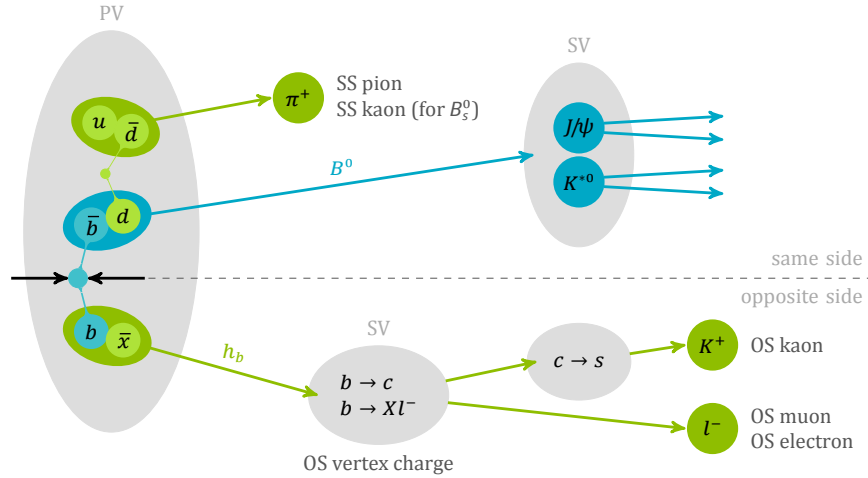


Figure 3.1 – Flavour tagging algorithms used in LHCb. In this cartoon, the signal channel is considered to be $B^0 \rightarrow J/\psi K^{*0}$.

in the same collision are opposite. Several OS taggers have been developed in LHCb, where the combination of the OS kaon (OSK), muon (OS μ), electron (OSE), and vertex charge (OSVtx) tagging algorithms represents the current standard OS combination. An additional OS tagger, the OS charm tagger (OSCharm) [57], can be exploited, and can be combined with the OS standard combination.

Given a reconstructed candidate, each flavour tagging algorithm provides a flavour tag d and a prediction η for the probability of the tag to be wrong. This mistag probability η is defined in the range $[0, 0.5]$ and is based on the output of multivariate classifiers, which are trained on datasets of flavour-specific decays, and combine several kinematic and geometric information on the tagging particle(s) and the event. The flavour tag takes the values $d = +1$ for an initial B^0 , $d = -1$ for an initial \bar{B}^0 , and $d = 0$ when no tag could be assigned; this happens, for example, if the tagging particle fails the selection criteria of a given tagging algorithm, or if its trajectory lies outside the detector acceptance.

More details on flavour tagging at LHCb can be found in Refs. [58, 59, 60].

Performance characteristics

The performance of flavour tagging algorithms can be characterised by different quantities. If N_U is the number of untagged candidates and N_W (N_R) is the number of wrongly (rightly) tagged candidates, the *tagging efficiency* (*i.e.* the fraction of tagged candidates) can be defined as

$$\epsilon_{\text{tag}} = \frac{N_R + N_W}{N_R + N_W + N_U}. \quad (3.1)$$

Chapter 3. Flavour tagging

The fraction of wrongly tagged candidates, or *mistag fraction*, is given by

$$\omega = \frac{N_W}{N_R + N_W}. \quad (3.2)$$

A non-zero mistag fraction dilutes the time-dependent asymmetries, reducing the experimental sensitivity to them. For instance, the measured decay rates for a $B \rightarrow f$ decay and its CP -conjugate decay are

$$\frac{d\Gamma^{\text{meas}}}{dt} = (1 - \omega) \frac{d\Gamma}{dt} + \omega \frac{d\bar{\Gamma}}{dt}, \quad (3.3)$$

$$\frac{d\bar{\Gamma}^{\text{meas}}}{dt} = \omega \frac{d\Gamma}{dt} + (1 - \omega) \frac{d\bar{\Gamma}}{dt}. \quad (3.4)$$

As a consequence, the measured CP asymmetry is

$$A^{\text{meas}}(t) = \frac{\frac{d\bar{\Gamma}^{\text{meas}}}{dt} - \frac{d\Gamma^{\text{meas}}}{dt}}{\frac{d\bar{\Gamma}^{\text{meas}}}{dt} + \frac{d\Gamma^{\text{meas}}}{dt}} = (1 - 2\omega) \frac{\frac{d\bar{\Gamma}}{dt} - \frac{d\Gamma}{dt}}{\frac{d\bar{\Gamma}}{dt} + \frac{d\Gamma}{dt}} = DA^{\text{phys}}(t), \quad (3.5)$$

where A^{phys} is the physical (true) CP asymmetry. The quantity $D = 1 - 2\omega$ is known as average *dilution*. If $\omega = 0$ (perfect tagger), then $D = 1$ and no asymmetry dilution occurs. If $\omega = 0.5$ (random tagger), then $D = 0$, and it is not possible to measure the asymmetry anymore.

The quantity that can be interpreted as the figure of merit to optimise a tagging algorithm is the *effective tagging efficiency*, also called *tagging power*:

$$\varepsilon_{\text{eff}} = \varepsilon_{\text{tag}} (1 - 2\omega)^2 = \varepsilon_{\text{tag}} D^2. \quad (3.6)$$

Assuming that ε_{eff} is known without uncertainty, it can be shown that the statistical uncertainty on the physical asymmetry is given by

$$\sigma_{A^{\text{phys}}} = \frac{\sqrt{1 - A^{\text{meas}^2}}}{\sqrt{N\varepsilon_{\text{tag}}(1 - 2\omega)}} = \frac{\sqrt{1 - A^{\text{meas}^2}}}{\sqrt{N\varepsilon_{\text{eff}}}}, \quad (3.7)$$

where N is the total number of candidates. So, according to Eq. 3.7, the greater the tagging power, the smaller the resulting statistical uncertainty on the CP asymmetry. Instead of using an average mistag fraction or *probability* ω , it is possible to exploit the mistag probability η estimated by the tagging algorithm. This probability η is evaluated for each B candidate individually, rather than being a global quantity. Usually, η needs to be *calibrated* via a function $\omega(\eta)$ in order to return the true mistag probability (details in Sec. 3.1). So, the tagging power can be rewritten as

$$\varepsilon_{\text{eff}} = \frac{1}{N} \sum_{i=1}^N D_i^2 = \frac{1}{N} \sum_{i=1}^N (1 - 2\omega(\eta_i))^2, \quad (3.8)$$

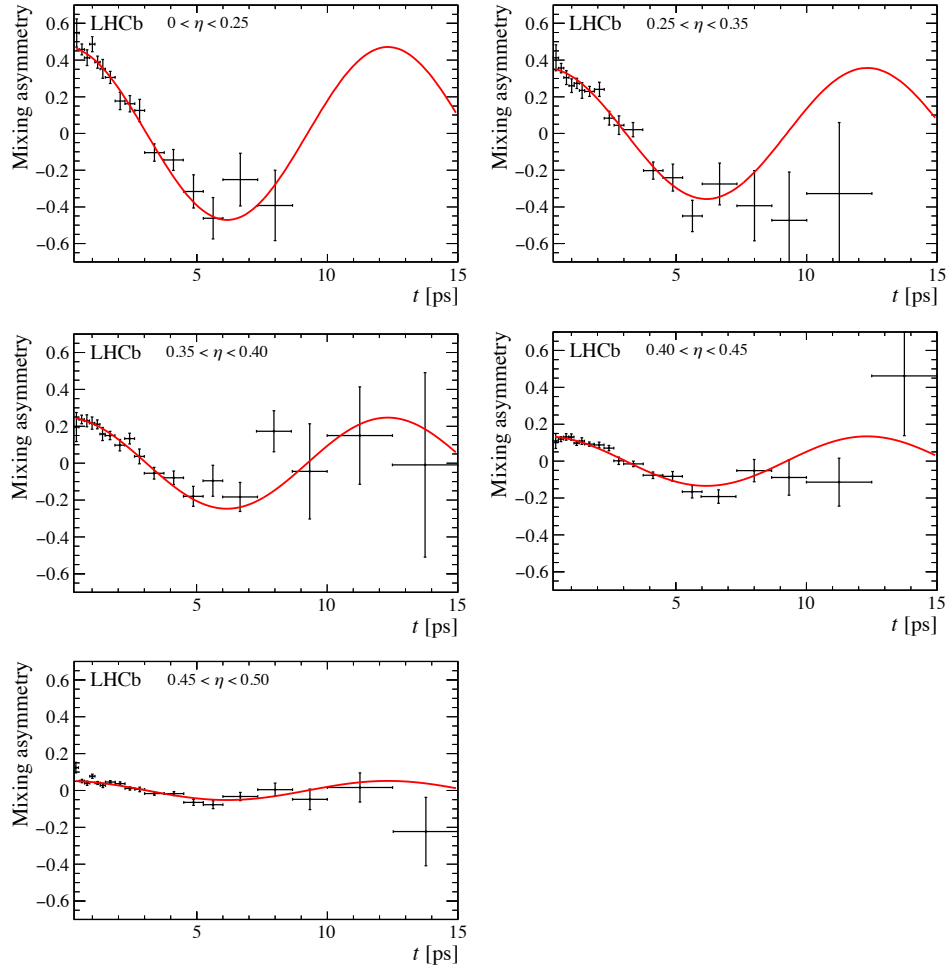


Figure 3.2 – Mixing asymmetry for SS-pion-tagged $B^0 \rightarrow D^- \pi^+$ candidates in bins of increasing estimated mistag η [55].

where $\omega(\eta_i) = 0.5$ ($D_i = 0$) for untagged candidates.

An example of dilution effect can be seen in Fig. 3.2, which shows how the measured amplitude of an asymmetry gets smaller for increasing values of η .

Calibration of the tagging output

The output of the flavour tagging algorithms is the result of training multivariate classifiers (MVA) using datasets of flavour-specific B decays, and transforming the classifier output into mistag estimates η through regression. However, as the training and validation samples are different from the signal sample used in the CP measurement (*e.g.* in terms of trigger and selection criteria that affect the distribution of the MVA input features),

Chapter 3. Flavour tagging

the output needs to be calibrated. Again, using control samples of flavour-specific decays, calibration functions $\omega(\eta)$ are obtained to transform the mistag estimate η of the algorithm to the mistag probability ω measured in the control sample.

A common choice for the calibration function is a linear function:

$$\omega(\eta) = p_0 + p_1 (\eta - \langle \eta \rangle). \quad (3.9)$$

The use of the arithmetic mean $\langle \eta \rangle$ of the η distribution aims at a decorrelation of p_0 and p_1 , hence a perfect calibration of the taggers would result in $p_0 = \langle \eta \rangle$ and $p_1 = 1$.

The performance of the flavour taggers is not necessarily independent of the initial flavour of the B^0 . The charged decay products, like the K^\pm mesons that are used by the OS kaon tagger, can have significantly different interaction rates with the detector material and therefore different reconstruction efficiencies. This can result in different tagging efficiencies ε_{tag} and mistag probabilities ω for B^0 and \bar{B}^0 . These tagging asymmetries can dilute or enhance the observed raw asymmetry and need to be corrected for. The asymmetries of the mistag probability, *i.e.* the difference of the tagging calibration parameters p_0 and p_1 for initial B^0 and \bar{B}^0 , can be parameterised with two independent calibration functions:

$$\begin{aligned} \omega^{B^0}(\eta) &= p_0^{B^0} + p_1^{B^0} (\eta - \langle \eta \rangle), \\ \omega^{\bar{B}^0}(\eta) &= p_0^{\bar{B}^0} + p_1^{\bar{B}^0} (\eta - \langle \eta \rangle). \end{aligned} \quad (3.10)$$

Equivalently, we can parameterise the calibration parameters p_i (with $i = 0, 1$) as

$$p_i^{B^0} = p_i + \frac{\Delta p_i}{2}, \quad p_i^{\bar{B}^0} = p_i - \frac{\Delta p_i}{2}. \quad (3.11)$$

The difference between the mistag of B^0 and \bar{B}^0 can be written as

$$\Delta\omega(\eta) = \omega^{B^0}(\eta) - \omega^{\bar{B}^0}(\eta) = \Delta p_0 + \Delta p_1 (\eta - \langle \eta \rangle). \quad (3.12)$$

In this thesis, new models for the calibration functions are adopted instead of the standard linear calibrations. These different parameterisations are called *Generalised Linear Models* (GLM), and are implemented in the EPM (*Espresso Performance Monitor*) package [61]. As will be explained in this section, these models allow a great flexibility to cope with non-linearities, and solve technical issues that may occur in fits that make use of flavour tagging. During my PhD, I worked to refine some of these models and to include them in the fitting routines used for decay-time fits.

In general, a GLM of order N that relates the predicted mistag probability η to the

calibrated probability ω can be written as follows:

$$\omega(\eta) = g(h(\eta)) = g\left(g^{-1}(\eta) + \sum_{i=1}^N \left(p_i + \frac{d\Delta p_i}{2}\right) f_i(\eta)\right). \quad (3.13)$$

The functions $f_i(\eta)$ are called *basis functions*, and they can be chosen as polynomials or spline functions. The set on basis functions is automatically orthogonalised by the EPM by using the Gram-Schmidt method [62]; this ensures that the corresponding calibration parameters p_i and Δp_i are correlated as little as possible.

The parameter d is the tagging decision, which is incorporated into the model in order to parameterise $\omega(\eta)$ for the two possible flavours.

The function g is known as *link function*. Usually, this is chosen as the inverse of a cumulative distribution function in order to map input values into the interval $[0, 1]$, such that the output can be naturally interpreted as a probability.

For the $B^0 \rightarrow D^\pm \pi^\mp$ analysis presented in this thesis, the adopted link function g is a *modified logistic function*, defined as

$$g(h) = \frac{1}{2(1 + e^h)}, \quad (3.14)$$

where h is defined in Eq. 3.13. This link function is built such that the calibrated mistag probability is defined in the interval $(0, 0.5)$. This choice solves a numerical issue that often occurs when standard link functions (*e.g.* identity or logistic) are adopted. In fact, if $\omega > 0.5$, then an arbitrary prescription has to be taken (*e.g.*, label the candidate as untagged, or flip the tagging decision and take $1 - \omega$ as new calibrated mistag). If the calibration parameters are free in a time-dependent fit, this choice has to be made during the minimisation process, according to the values ω takes at each iteration. This means that the relative number of B and \bar{B} , or the relative number of tagged and untagged candidates, may change during the fit, which leads to numerical instabilities due to discontinuous changes in the likelihood function.

The EPM estimates the calibration parameters p_i and Δp_i via an unbinned maximum likelihood fit called *binomial regression*; this is an improvement over traditional, binned least-squares fits, which are affected by a systematic uncertainty due to the binning choice.

Combination of multiple taggers

When more than one tagger is available per event, the tagging decisions and mistag probabilities provided by each tagger can be combined into a single decision and a single

probability using the equations

$$p(b) = \prod_i \left(\frac{1}{2} - d_i \left(\frac{1}{2} - \eta_i \right) \right), \quad p(\bar{b}) = \prod_i \left(\frac{1}{2} + d_i \left(\frac{1}{2} - \eta_i \right) \right), \quad (3.15)$$

where $p(\bar{b}/b)$ is the probability that the signal B^0 contains a \bar{b}/b , d_i is the decision taken by the i -th tagger and η_i is the predicted mistag probability of the i -th tagger. These probabilities are normalised as

$$P(\bar{b}) = \frac{p(\bar{b})}{p(\bar{b}) + p(b)}, \quad P(b) = 1 - P(\bar{b}). \quad (3.16)$$

If $P(\bar{b}) > P(b)$ the combined tagging decision is $d = +1$ and the final mistag probability is $\eta = 1 - P(\bar{b})$. Otherwise if $P(b) > P(\bar{b})$ the combined tagging decision and the mistag probability are $d = -1$ and $\eta = 1 - P(b)$.

Equation 3.15 is valid under the assumption that all taggers in the combination are independent. In the $B^0 \rightarrow D^\mp \pi^\pm$ analysis presented in this thesis, the OS taggers are combined in a single OS combination, and the same is done for the SS taggers. Effects due to correlations among taggers within a combination are corrected for by calibrating the combined predicted mistag.

3.2 Flavour tagging strategy for the $B^0 \rightarrow D^\mp \pi^\pm$ time-dependent analysis

In the $B^0 \rightarrow D^\pm \pi^\mp$ analysis presented in this thesis, the OS combination (including the OS charm tagger) and the SS combination are used. The implementation of the OS algorithms used in the combination are the same as described in Refs. [56, 57]; the OS algorithms other than the OS charm tagger were built as neural networks trained on $B^+ \rightarrow J/\psi K^+$ Run 1 data, whereas the OS charm tagger was implemented with a BDT trained on a cocktail of simulated $B^+ \rightarrow J/\psi K^+$, $B^0 \rightarrow J/\psi K^{*0}$ and $B_s^0 \rightarrow J/\psi \phi$ decays. The SS taggers have been reimplemented for this specific analysis by exploiting $B^0 \rightarrow J/\psi K^{*0}$ decays. The functional form of the tagging calibrations is studied in control samples of flavour-specific decays properly corrected to resemble the signal decay. The calibration parameters are determined directly in the decay-time-dependent fit of the signal described in Sec. 5.3; they are nuisance parameters of the likelihood function. Determining the calibration parameters from the data along with the CP observables is possible because the CP coefficients C_f and $C_{\bar{f}}$ of Eqs. 1.35–1.38 are fixed in this analysis (to 1 and -1 respectively). Hence, the cosine terms give sensitivity to the calibration parameters independently of the sine terms, which are proportional to the S_f and $S_{\bar{f}}$ coefficients. A heuristic explanation is presented in Fig. 3.3. This strategy avoids any assumption on the portability to the signal sample of the calibration parameters determined from the control

3.2. Flavour tagging strategy for the $B^0 \rightarrow D^\mp \pi^\pm$ time-dependent analysis

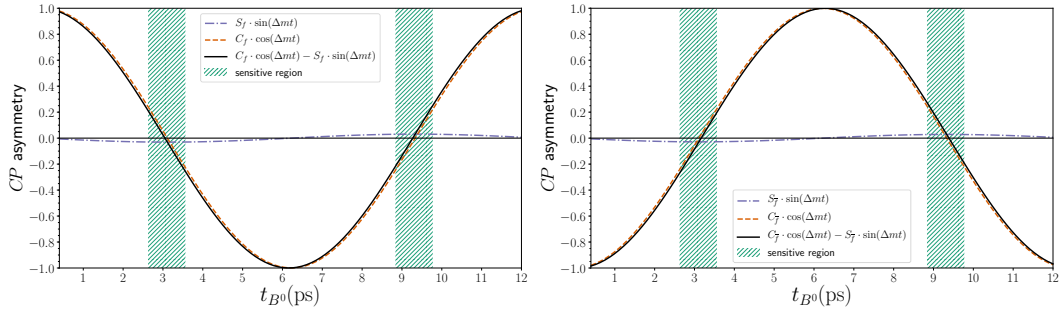


Figure 3.3 – \bar{B}^0 versus B^0 time-dependent asymmetries for the $D^- \pi^+$ (left) and $D^+ \pi^-$ (right) final states. The values of C_f , $C_{\bar{f}}$, S_f and $S_{\bar{f}}$ are the ones used in simulation (see Appendix M). The sensitivity to S_f and $S_{\bar{f}}$ is maximised in the intervals called “sensitive regions”, since the $\sin(\Delta m)$ amplitude becomes of the same order of the $\cos(\Delta m)$ amplitude, which is close to zero. In the outer regions, since C_f ($C_{\bar{f}}$) is fixed to 1 (-1) in the fit, the mistag dilution (which depends on the flavour tagging calibration parameters) adapts to fit the $\cos(\Delta m)$ amplitude, giving sensitivity to the calibration parameters.

data. Such a strategy was studied extensively on simulation: the increase of the statistical uncertainty of the S_f and $S_{\bar{f}}$ coefficients given by the additional degrees of freedom of the calibration parameters is smaller than the systematic uncertainties associated with the calibration portability. Moreover, the use of the calibration parameters from the control sample causes biases on S_f and $S_{\bar{f}}$ of the order of their statistical uncertainty; when letting the calibration parameters float in the fit, such biases are suppressed or disappear, at the cost of a moderate increase of the statistical uncertainty. In addition, while the precision of the OS tagger calibration from the control sample is similar to the one from the signal sample, the calibration of the SS tagger derived from the signal sample (Tab. 5.4) is much more precise than that from the control sample (Tab. 3.4).

In what follows, the study of the tagging calibration from the control sample is presented. For all reasons discussed above, these studies are not meant for determining the calibration parameters to use in the time fit to the signal data (usual strategy adopted in all flavour-tagged time-dependent analyses), but they serve the purpose of: i) determining the best functional form of the calibration functions to be used in the fit to the signal; ii) having some reference values for the calibration parameters for a comparison with those extracted from the signal.

The calibration for the OS combination are determined using $B^+ \rightarrow \bar{D}^0 \pi^+$ decays, as described in Sec. 3.2.1. The SS pion and the SS proton taggers were developed using $B^0 \rightarrow D^\mp \pi^\pm$ data and assuming negligible CP violation. The use of these algorithms in this analysis could bias the measurement. Therefore, the SS taggers are retrained using $B^0 \rightarrow J/\psi K^{*0}$ decays. The calibration of the SS combination is described in Sec. 3.2.2.

3.2.1 Calibration of the opposite-side tagger combination

Data sample selection

The calibration parameters of the OS tagger combination (namely the combination of the OS electron, muon, kaon, vertex charge, and charm algorithms) are determined using $B^+ \rightarrow \bar{D}^0(\rightarrow K^+\pi^-)\pi^+$ candidates reconstructed in 3 fb^{-1} of data. Such a control decay mode provides very high statistics (more than 300k OS-tagged signal candidates) and is very similar to the signal decay $B^0 \rightarrow D^-\pi^+$.

Candidate $B^+ \rightarrow \bar{D}^0\pi^+$ decays are selected through the `B2D0PiD2HHBeauty2CharmLine` stripping line, versions `S21r1` (2011 data) and `S21` (2012 data), of the `BhadronCompleteEvent` stream. The B^+ candidates are required to be TOS, *i.e.* to trigger on `Hlt1TrackAllL0Decision` at the HLT1 stage, and at least one among `Hlt2Topo2BodyBBDTDecision`, `Hlt2Topo3BodyBBDTDecision`, and `Hlt2Topo4BodyBBDTDecision` at HLT2. The additional requirements listed in Table 3.1 are applied to further suppress backgrounds and enhance the signal purity.

A fit to the mass distribution of B^+ candidates is done to calculate $sWeights$, used in the subsequent steps of the analysis to subtract the backgrounds surviving the selection. This fit is described in details in Appendix A.1.

Event-by-event weights are calculated to equalise the $B^+ \rightarrow \bar{D}^0\pi^+$ and $B^0 \rightarrow D^\mp\pi^\pm$ distributions of the variables on which the tagging calibration can depend. The procedure and the results of this reweighting are reported in Appendix A.2. Additionally, the number of B^+ and B^- candidates are made equal in the sample to avoid any spoil of the calibration parameters due to a B^+/B^- production asymmetry or a detection asymmetry. All these weights, along with $sWeights$ [38], are applied during the calibration procedure.

Calibration

The calibration of the estimated mistag η is performed on the fully reweighted $B^+ \rightarrow \bar{D}^0\pi^+$ dataset. A GLM model with NSpline basis function [61] is adopted. The projection of the fitted calibration function over the $B^+ \rightarrow \bar{D}^0\pi^+$ dataset is shown in Fig. 3.4, whereas the fitted calibration parameters are listed in Table 3.2. The number of free parameters in the adopted GLM model (10) has been chosen in order to have satisfactory goodness-of-fit (GOF) metrics (details in Appendix A.3).

Calibration portability

The aim of the calibration is to return a mistag ω as close as possible to the *true* mistag, which would be given by a *true calibration*. The latter is not defined for $B^0 \rightarrow D^-\pi^+$ decays in data, but it is possible to estimate it for $B^0 \rightarrow D^-\pi^+$ decays on MC. In fact,

3.2. Flavour tagging strategy for the $B^0 \rightarrow D^\mp \pi^\pm$ time-dependent analysis

Table 3.1 – Selection requirements for the $B^+ \rightarrow \bar{D}^0 \pi^+$ candidates.

Description	Variable	Requirement
Bachelor track		
muon identification criteria	IsMuon	= 0
ghost probability	p_{ghost}	< 0.1
quality of track	$\chi_{\text{track}}^2/\text{ndof}$	< 2
D^0 daughter tracks		
$\ln L_K - \ln L_\pi$	PIDK	> -2 (kaon), < 8 (pion)
ghost probability	p_{ghost}	< 0.1
quality of track	$\chi_{\text{track}}^2/\text{ndof}$	< 2.5
D^0 candidate		
invariant mass	$m_{K\pi}$	$m_{K\pi} \in [1830, 1904]$ MeV/ c^2
B^+ candidate		
decay time	τ_{B^+}	$\tau_{B^+} \in [0.2, 15]$ ps
minimum IP χ^2 w.r.t. PV	$\text{MIN}\chi^2\text{IP}_{\text{PV}}$	< 15

Table 3.2 – Fitted OS calibration parameters on the $B^+ \rightarrow \bar{D}^0 \pi^+$ reweighted dataset.

Parameter	Fitted value
p_0	-0.136 ± 0.019
p_1	-0.006 ± 0.022
p_2	-0.0107 ± 0.0083
p_3	-0.5 ± 0.10
p_4	-0.85 ± 0.46
Δp_0	-0.129 ± 0.038
Δp_1	0.042 ± 0.045
Δp_2	-0.020 ± 0.017
Δp_3	0.42 ± 0.21
Δp_4	1.91 ± 0.92

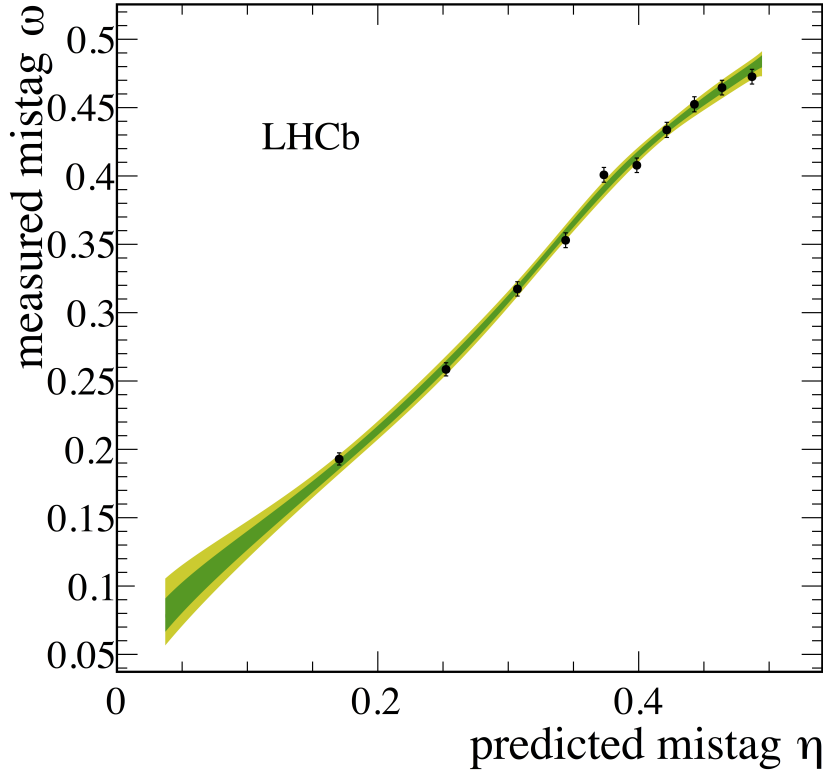


Figure 3.4 – Mistag ω measured in bins of predicted mistag η for reweighted $B^+ \rightarrow \bar{D}^0\pi^+$ candidates (data points) and fitted calibration function. The green (yellow) band indicates the 68% (95%) confidence interval on the calibration function.

since the true flavour of the B^0 meson is known in MC, this true MC calibration can be done in the same way as $B^+ \rightarrow \bar{D}^0\pi^+$, where the true flavour is given by the B charge.

This $B^0 \rightarrow D^-\pi^+$ calibration is performed after equalising the number of B^0 and \bar{B}^0 in the sample, in order to disentangle tagging asymmetries from CP violation and production asymmetries.

The $B^+ \rightarrow \bar{D}^0\pi^+$ MC calibration is performed in exactly the same way as described in Sec. 3.2.1, except that no *sWeights* are considered, since only true MC signal decays are used.

The two calibrations using the $B^0 \rightarrow D^-\pi^+$ and $B^+ \rightarrow \bar{D}^0\pi^+$ MC samples are shown in Fig. 3.5 and compared in Table 3.3. A more robust comparison is obtained from a χ^2 function describing the discrepancy between the two calibrations by taking the covariance matrices into account. The overall discrepancy (corresponding to the χ^2 minimum) is around 2σ .

3.2. Flavour tagging strategy for the $B^0 \rightarrow D^\mp \pi^\pm$ time-dependent analysis

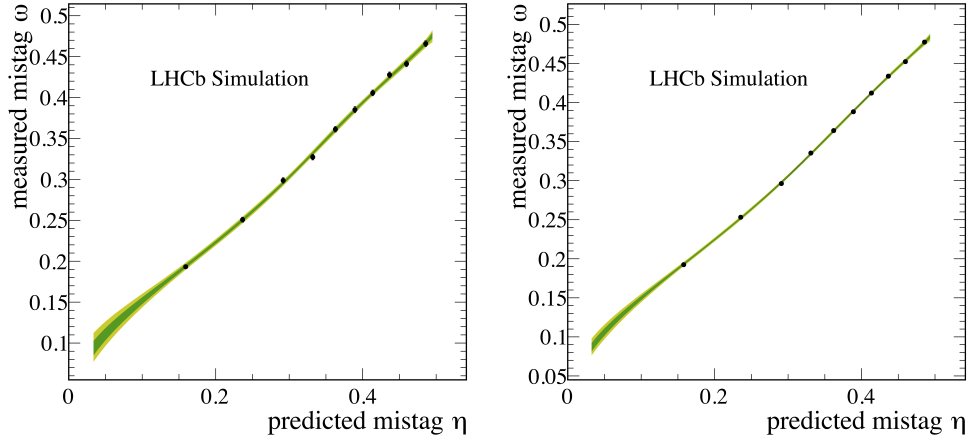


Figure 3.5 – Mistag ω measured in bins of predicted mistag η for reweighted $B^+ \rightarrow \bar{D}^0\pi^+$ (left) and $B^0 \rightarrow D^-\pi^+$ (right) candidates (data points) and fitted calibration functions. The green (yellow) band indicates the 68% (95%) confidence interval on the calibration functions.

Table 3.3 – Comparison between the fitted OS tagging calibration parameters using truth-matched $B^+ \rightarrow \bar{D}^0\pi^+$ and $B^0 \rightarrow D^-\pi^+$ MC decays. The discrepancy in each parameter is computed assuming independent datasets.

Parameter	$B^+ \rightarrow \bar{D}^0\pi^+$	$B^0 \rightarrow D^-\pi^+$	Discrepancy (σ)
p_0	-0.065 ± 0.011	-0.0996 ± 0.0066	2.70
p_1	-0.190 ± 0.012	-0.1492 ± 0.0077	-2.84
p_2	-0.0105 ± 0.0044	-0.0191 ± 0.0029	1.63
p_3	-0.295 ± 0.054	-0.234 ± 0.036	-0.93
p_4	-0.42 ± 0.26	-0.14 ± 0.20	-0.85
Δp_0	-0.059 ± 0.022	-0.058 ± 0.013	-0.03
Δp_1	0.044 ± 0.024	0.030 ± 0.015	0.46
Δp_2	-0.0012 ± 0.0088	-0.0126 ± 0.0058	1.08
Δp_3	-0.08 ± 0.11	-0.046 ± 0.073	-0.25
Δp_4	-0.34 ± 0.53	-0.29 ± 0.39	-0.08

3.2.2 Calibration of the same-side tagger combination

As described in Ref. [55], the SS pion and proton taggers were both trained on the 2012 data sample of $B^0 \rightarrow D^\mp \pi^\pm$ decays. As the effect of CP violation was neglected during the training the algorithms and the underlying MVAs cannot be blindly used when measuring CP violation in the same decay channel. Thus, $B^0 \rightarrow J/\psi K^{*0}$ decays are chosen instead, as they represent a flavour-specific B^0 decay with a large signal yield of about 350000 candidates in 2012 data.

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Table 3.4 – Fitted SS calibration parameters obtained on the $B^0 \rightarrow J/\psi K^{*0}$ data sample (calibration subsample).

p_0	p_1	Δp_0	Δp_1
-0.091 ± 0.059	-0.027 ± 0.065	0.034 ± 0.084	0.032 ± 0.094

Once the SS pion and proton taggers are implemented, they are combined into a single SS combination as described in Sec. 3.1.

Calibration

The calibration is performed on a $B^0 \rightarrow J/\psi K^{*0}$ data subsample that is not used for the SS pion and proton training. A GLM model having a first order polynomial is chosen as basis function and a modified logistic function (Eq. 3.14) is used as link. The number of free parameters in this model (4) is tuned in order to have satisfactory goodness-of-fit (GOF) metrics. Together with the *sWeights*, additional weights to correct the $B^0 \rightarrow J/\psi K^{*0}$ data to resemble the $B^0 \rightarrow D^\mp \pi^\pm$ data are applied during the calibration. The resulting calibration parameters are listed in Table 3.4 and a graphical representation of the calibration is presented in Fig. 3.6.

Calibration portability

In the same way as for the OS taggers (Sec. 3.2.1), the portability of the SS tagging calibration is checked on Monte Carlo. For $B^0 \rightarrow D^- \pi^+$ the calibration is performed using the true flavour of the B^0 meson after equalising the number of B^0 and \bar{B}^0 in the sample, in order to disentangle tagging asymmetries from *CP* violation and production asymmetries. Also on $B^0 \rightarrow J/\psi K^{*0}$ the true flavour of the B^0 meson is used for the calibration, and no *sWeights* are needed, since only the true MC signal decays are used.

The two calibrations using the $B^0 \rightarrow D^- \pi^+$ and $B^0 \rightarrow J/\psi K^{*0}$ Monte Carlo samples are shown in Fig. 3.7 and compared in Table 3.5. A full comparison that takes into account the correlation between the parameters is obtained from a χ^2 test similar to the one described in Sec. 3.2.1. The agreement is around 0.1σ . Even though this test doesn't hint to issues of portability between the decay modes, the same strategy used for the OS calibrations is followed, i.e fitting the parameter directly in data with the *CP* asymmetries. This is motivated by the fact that the $B^0 \rightarrow D^- \pi^+$ signal sample has much more sensitivity to determine the parameters than the $B^0 \rightarrow J/\psi K^{*0}$ sample. In addition, with this approach no systematic related to calibration portability is necessary, consistent with the OS tagger treatment.

3.2. Flavour tagging strategy for the $B^0 \rightarrow D^\mp \pi^\pm$ time-dependent analysis

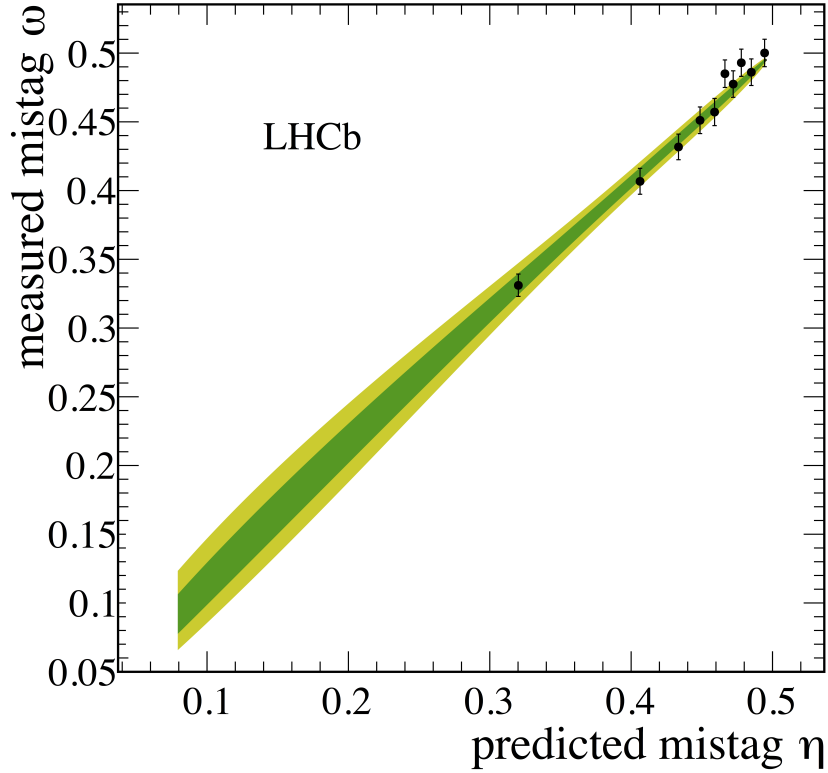


Figure 3.6 – Mistag ω measured in bins of predicted mistag η for reweighted $B^0 \rightarrow J/\psi K^{*0}$ candidates (data points) and fitted calibration function. The green (yellow) band indicates the 68% (95%) confidence interval on the calibration function.

Table 3.5 – Comparison between the fitted SS tagging calibration parameters using truth-matched $B^0 \rightarrow J/\psi K^{*0}$ and $B^0 \rightarrow D^- \pi^+$ MC decays. The discrepancy in each parameter is computed assuming independent datasets.

Parameter	$B^0 \rightarrow J/\psi K^{*0}$	$B^0 \rightarrow D^- \pi^+$	Discrepancy (σ)
p_0	-0.016 ± 0.017	-0.019 ± 0.008	-0.19
p_1	0.063 ± 0.021	0.060 ± 0.010	-0.14
Δp_0	-0.029 ± 0.033	-0.027 ± 0.015	0.04
Δp_1	-0.026 ± 0.041	0.015 ± 0.019	0.90

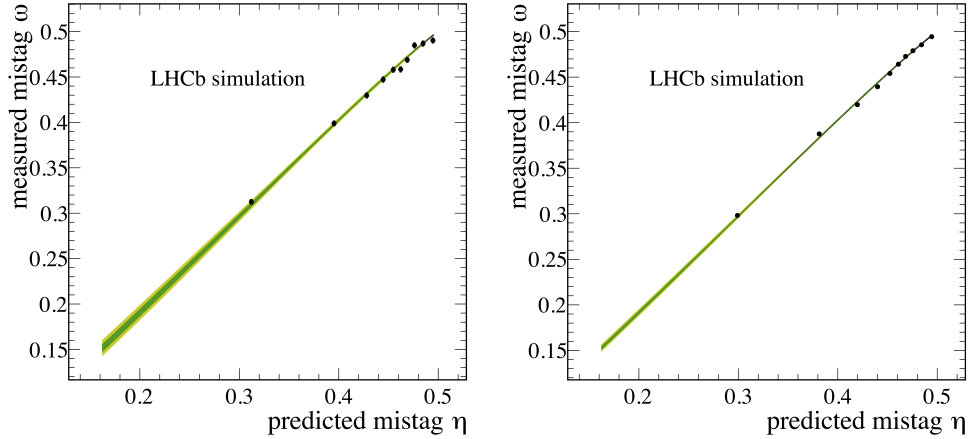


Figure 3.7 – Mistag ω measured in bins of predicted mistag η for reweighted $B^0 \rightarrow J/\psi K^{*0}$ (left) and $B^0 \rightarrow D^- \pi^+$ (right) candidates (data points) and fitted calibration functions. The green (yellow) band indicates the 68% (95%) confidence interval on the calibration functions.

3.3 Optimisation of the opposite-side electron tagger

The performance of the flavour tagging algorithms depends on the data taking conditions, in particular the centre-of-mass energy of the pp collision.

On one hand, the tagging power of the SS taggers shows an increase on Run 2 data as compared to Run 1, thanks to either a higher tagging efficiency (SS π and SS p) or a lower mistag rate (SS K). This is due to the higher boost of the $b\bar{b}$ quark pair at 13 TeV, which makes the momentum spectrum of B mesons and fragmentation tracks harder, and increases the acceptance of the fragmentation tracks.

On the other hand, the tagging power of the existing OS e , OS μ , and OS K taggers decreases on Run 2 data. The reason for this degradation is mainly due to the higher track multiplicity, which increases the probability to have a wrong tag decision. Moreover, because of the different Run 2 kinematics, the criteria to select the tagging particles are no longer optimal, thus giving a lower tagging efficiency.

The performance of the OSCharm and OSVtx algorithms is, on average, compatible or better on Run 2 as compared to Run 1.

In this section, the reoptimisation of the OS e tagger is presented. This reoptimisation is performed both on Run 2 data, in order to recover the observed loss in tagging power, and on Run 1 data, to further improve the already existing algorithm. This reoptimisation consists of two main steps. First, selection criteria are applied to select electron-like

3.3. Optimisation of the opposite-side electron tagger

particles, yielding a sample of B signal candidates with a low average tagging power. Then, a BDT classifier is applied to discriminate between B candidates with right and wrong tag decisions for each selected track. Finally, for each B candidate, the BDT output for the track with the highest transverse momentum is converted into a predicted mistag probability.

A similar approach is followed for the development of the OSK and $OS\mu$ taggers. In this case, the reoptimisation on Run 1 data does not show any gain in performance, whereas some significant gain is found on Run 2 data.

3.3.1 Sample definition

The OSe algorithm is developed in a data-driven fashion by using *sWeighted* samples of $B^+ \rightarrow J/\psi K^+$ decays. The full Run 1 dataset (2011+2012) is used to optimise the algorithm on Run 1 conditions, whereas the 2016 dataset is exploited to optimise the tagger on Run 2 conditions. An alternative optimisation on *sWeighted* 2016 $B^+ \rightarrow \bar{D}^0 \pi^+$ data is performed in parallel. The motivation for this is to cross-check the Run 2 implementation with an independent decay mode, which is characterised by different kinematics than that of $B^+ \rightarrow J/\psi K^+$. Hereafter, the OSe tagger optimised on Run 1 $B^+ \rightarrow J/\psi K^+$ data will be indicated as “Run 1 new” version, in order to distinguish it from the previous “Run 1 old” version introduced in Ref. [56], which was based on simple selection criteria and a neural network for the mistag estimation. The OSe tagger optimised on Run 2 $B^+ \rightarrow J/\psi K^+$ and $B^+ \rightarrow \bar{D}^0 \pi^+$ data will be denoted as “Run 2 B2CC” and “Run 2 B2OC” versions, respectively. Moreover, it is understood that the tunings of all the PROBNN features mentioned in this section are `MC12TuneV2` for Run 1 data and `MC15TuneV1` for Run 2 data.

Each dataset is divided in four subsamples:

- the first subset, including $\sim 25\%$ of the total data, is used for the optimisation of the electron preselection (Sec. 3.3.2);
- the second subset (*training sample*), including $\sim 50\%$ of the total data, is adopted as training set for the BDT classifier used for the predicted mistag estimation (Sec. 3.3.3). This sample is also used for tuning some *hyperparameters* (number of trees, maximum depth) which define the BDT classifier.
- the third and the fourth subsets (*evaluation set 1* and *2*), each including $\sim 12.5\%$ of the total data, are adopted together as test sets to check for overtraining (Sec. 3.3.4). The evaluation set 1 is also used to calibrate the obtained tagger, which is then applied to the second evaluation set in order to measure the performance; the procedure is then repeated by swapping the two samples (*two-fold validation*).

3.3.2 Preselection optimisation

Electron-like particles are selected by means of a set of requirements. The reconstructed tracks must not be associated to the B signal decay tree, and must not have hits in the muon detector in order to exclude muons (already exploited by the $\text{OS}\mu$ algorithm). Moreover, these tracks have to be of type long, lie in the ECAL acceptance, and have sufficient reconstruction quality ($\chi^2/\text{ndof} < 3$). Also, the inverse of the rigidity e/p has to be comprised between 0.85 and 2, and the charge deposited in the VELO detector must be smaller than 1.4 normalised Analogic-to-Digital Converter (ADC) counts. Finally, tracks where the fit for the track IP with respect to the primary vertex did not converge are excluded.

A further selection is applied in order to enhance the average tagging power of the resulting sample, which is defined as

$$\langle \varepsilon_{\text{eff}} \rangle = \varepsilon_{\text{tag}} \left[1 - 2 \frac{\sum_{i=1}^N w_i f_{\text{R}}^i}{\sum_{i=1}^N w_i (f_{\text{R}}^i + f_{\text{W}}^i)} \right]^2, \quad (3.17)$$

where w_i are the *sWeights*, while f_{R}^i (f_{W}^i) is the fraction of particles giving the right (wrong) flavour for the i th B candidate, with $f_{\text{R}}^i + f_{\text{W}}^i = 1$ for every candidate.

The expression of Eq. 3.17 is taken as the figure of merit to maximise during the selection optimisation. This maximisation is performed numerically by using gradient boosted regression trees to model $\langle \varepsilon_{\text{eff}} \rangle$ as a function of the applied cuts [63]. The cuts are optimised separately for the Run 1 new, Run 2 B2CC and Run 2 B2OC algorithms; in all cases, about 25% of the available data for each sample is used.

The resulting, optimised requirements are reported in Table 3.6, while the convergence plots of the minimisation are shown in Fig. 3.8.

After the optimisation, the performance of the selection (including the average tagging power defined in Eq. 3.17) is evaluated on the remaining 75% of data for each sample, yielding the results shown in Table 3.7.

3.3.3 BDT classifier implementation

The selection described in Sec. 3.3.2 is applied on the remaining part of the data ($\sim 75\%$) used by each OSe implementation. The BDT classifier is trained to identify B candidates as correctly or incorrectly tagged. The list of the features considered to build the BDTs are reported in Table 3.8. The distributions of the input features for the two possible values of the target are shown in Figs. 3.9, 3.10 and 3.11 for the Run 1 new, Run 2 B2CC

3.3. Optimisation of the opposite-side electron tagger

Table 3.6 – Optimised requirements for the preselection of tracks used by the OSe algorithm. p_{ghost} is the probability for the track to be a fake combination of hits. IPPU denotes the impact parameter with respect to the pile-up vertex in the event, which might be reconstructed in the event in addition to the nominal PV. $\Delta\phi$ is the difference in azimuthal angle between the track and the signal B candidate.

Requirement	Run 1 new	Run 2 B2CC	Run 2 B2OC
$p_{\text{ghost}} <$	0.861	0.843	0.348
PROBNN $\pi <$	0.934	0.983	0.980
PROBNN $p <$	0.719	0.271	0.732
PROBNN $K <$	0.765	0.695	0.954
PROBNN $e >$	0.061	0.243	0.040
PROBNN $\mu <$	0.938	0.158	0.263
PID $e >$	4.555	4.333	-0.691
$p_{\text{T}} >$ (MeV/ c)	1132	1403	1263
$p >$ (MeV)	3114	5035	2246
$\sigma_{\text{IP}}/\text{IP} >$	0.020	0.042	1.410
$\sigma_{\text{IPPU}}/\text{IPPU} >$	12.101	9.335	2.758
min $\Delta\phi >$	0.00803	0.0167	0.0299

Table 3.7 – Performance of the preselection (OSe algorithm) applied on the data not used for the preselection optimisation ($\sim 75\%$ of the total dataset for each sample). The average tagging power $\langle \varepsilon_{\text{eff}} \rangle$ is defined in Eq. 3.17.

Algorithm	ε_{tag} (%)	$\langle \omega \rangle$ (%)	$\langle \varepsilon_{\text{eff}} \rangle$ (%)
Run 1 new	3.440 ± 0.019	33.31 ± 0.27	0.383 ± 0.007
Run 2 B2CC	2.514 ± 0.017	33.50 ± 0.32	0.274 ± 0.006
Run 2 B2OC	3.664 ± 0.024	34.32 ± 0.32	0.360 ± 0.006

and Run 2 B2OC samples, respectively. The Pearson correlation coefficients between the input features are reported in Fig. 3.12.

The BDT classifier consists of an ensemble of 300 gradient-boosted decision trees [64], where each tree can have a maximum depth of 3. The objective of the classifier is a binary logistic loss function plus a quadratic regularisation term to control model complexity (with regularisation parameter $\lambda = 1$). Some hyperparameters were tested by means of a cross-validation+bootstrapping method on the training set, as described in Appendix B. The importance (or F score) of each feature, defined as the total number of times a feature is chosen as split node by any tree in the BDT ensemble, is presented in Fig. 3.13, while the *partial dependence* of the predicted mistag η (on the training set) as a function of each input feature is shown in Appendix B. The receiver operating characteristic (ROC) curves, which report the *true positive rate* as a function of the *false positive rate*, are shown in Fig. 3.14. The true (false) positive rate is the fraction of true, correctly (incorrectly)

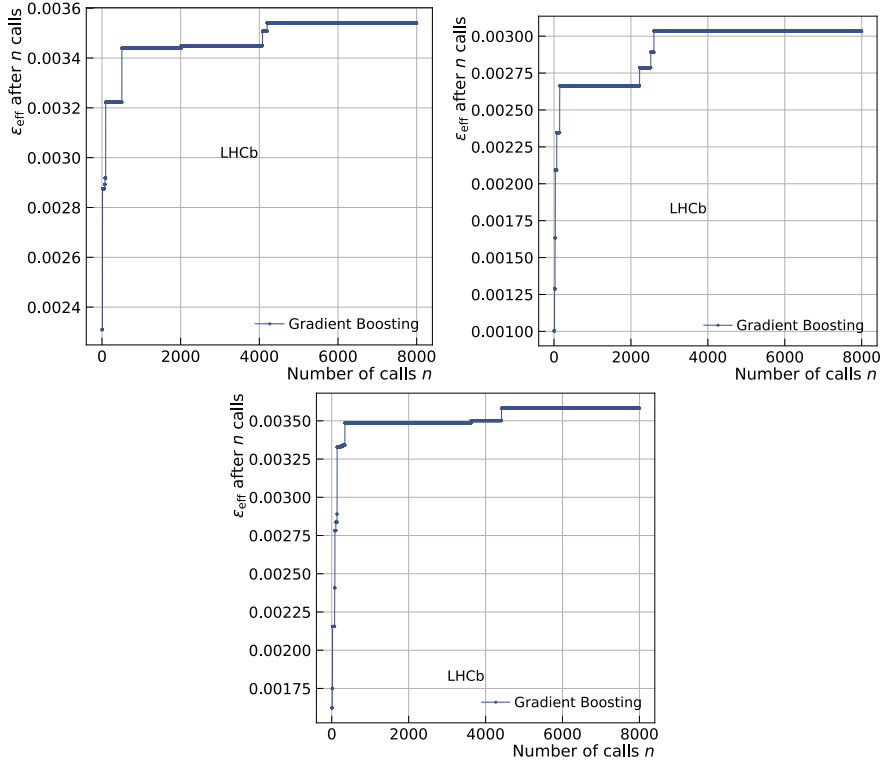


Figure 3.8 – Maximised value of the average tagging power as a function of the gradient boosted regression tree algorithm iteration for the Run 1 new (top left), Run 2 B2CC (top right) and Run 2 B2OC (bottom) implementations of the OSe tagger.

tagged candidates over all candidates classified as correctly tagged. The feature selection, BDT training and feature importance evaluation chain is repeated iteratively in order to exclude highly-correlated and poorly-important features, until the BDT performance starts to degrade significantly.

For each candidate, the BDT predicts the probability P that such candidates is correctly tagged. In order to obtain a mistag probability η , the following transformation is applied on both P and tagging decision d :

$$(\eta, d) \rightarrow \begin{cases} (P, -d) & \text{if } P \leq 0.5 \\ (1 - P, d) & \text{otherwise} \end{cases} \quad (3.18)$$

The distributions of η for training and test samples, splitted per target value, are shown in Fig. 3.15.

3.3. Optimisation of the opposite-side electron tagger

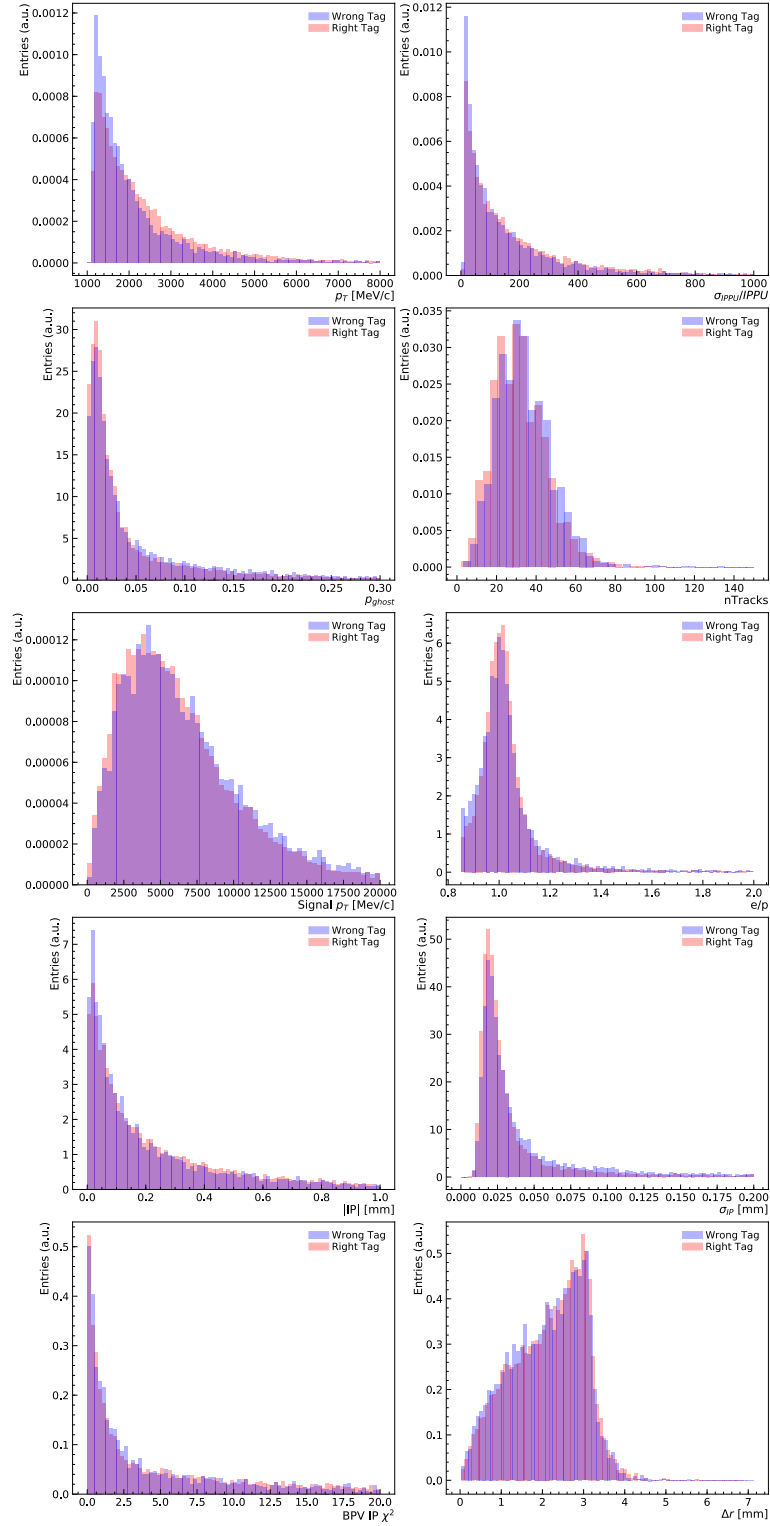


Figure 3.9 – Distributions (for the *sWeighted*, Run 1 $B^+ \rightarrow J/\psi K^+$ sample) of the input features of the BDT classifier, for candidates with a right (red) and wrong (blue) decision from the OSe tagger.

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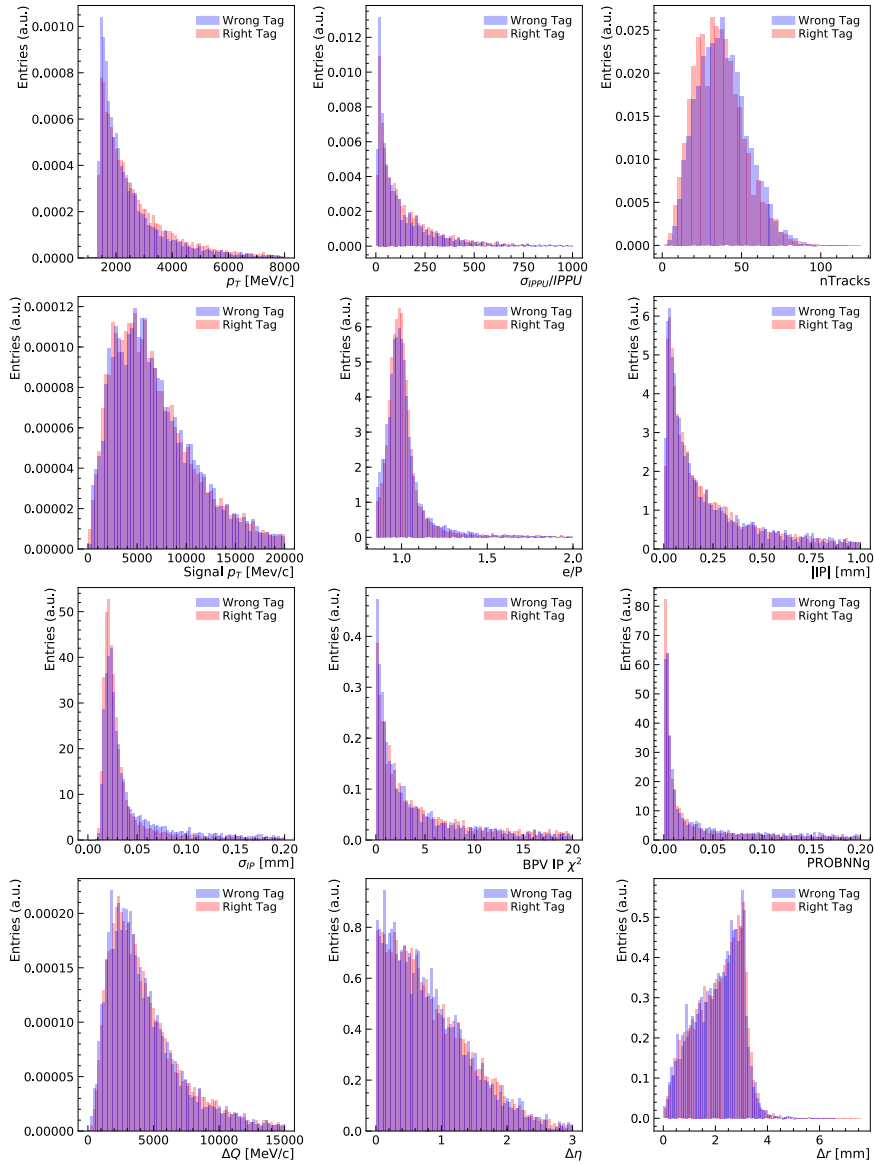


Figure 3.10 – Distributions (for the *sWeighted*, Run 2 $B^+ \rightarrow J/\psi K^+$ sample) of the input features of the BDT classifier, for candidates with a right (red) and wrong (blue) decision from the OSe tagger.

3.3. Optimisation of the opposite-side electron tagger

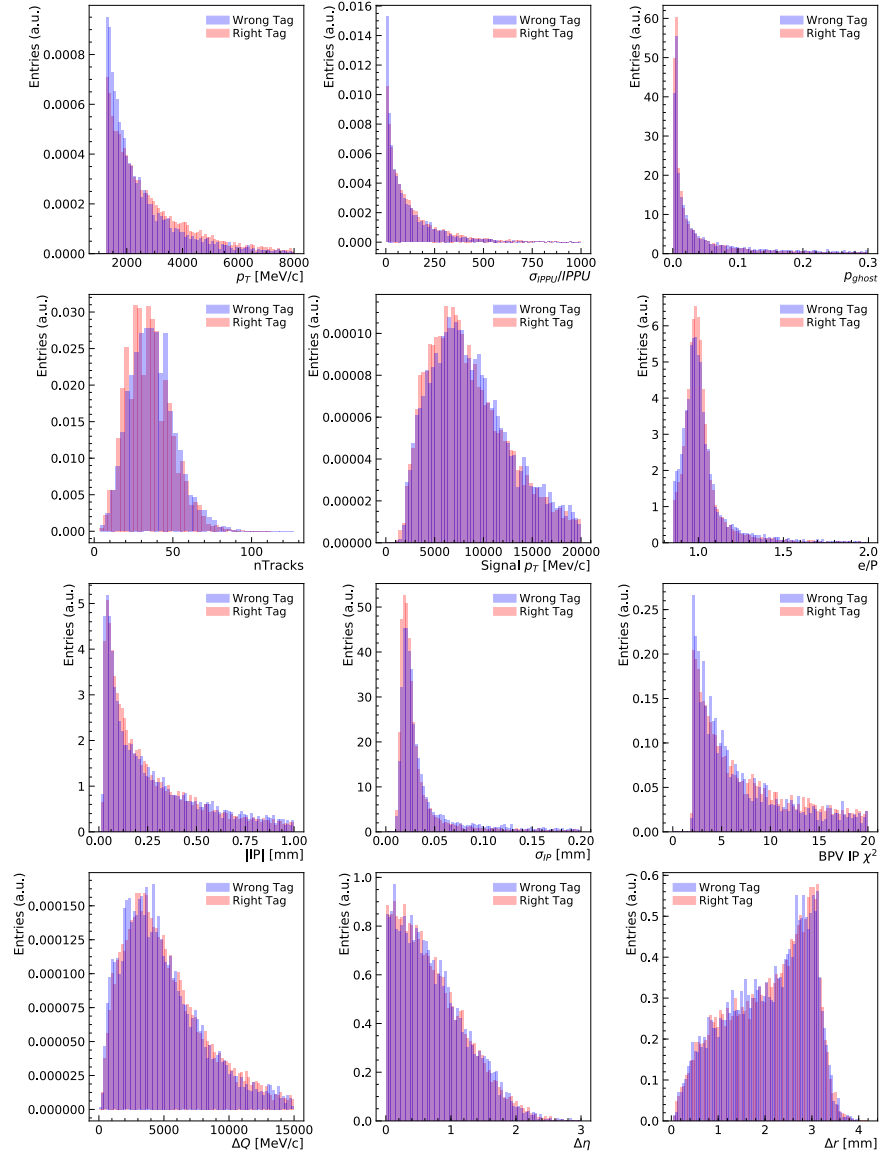


Figure 3.11 – Distributions (for the *sWeighted*, Run 2 $B^+ \rightarrow \bar{D}^0 \pi^+$ sample) of the input features of the BDT classifier, for candidates with a right (red) and wrong (blue) decision from the OSe tagger.

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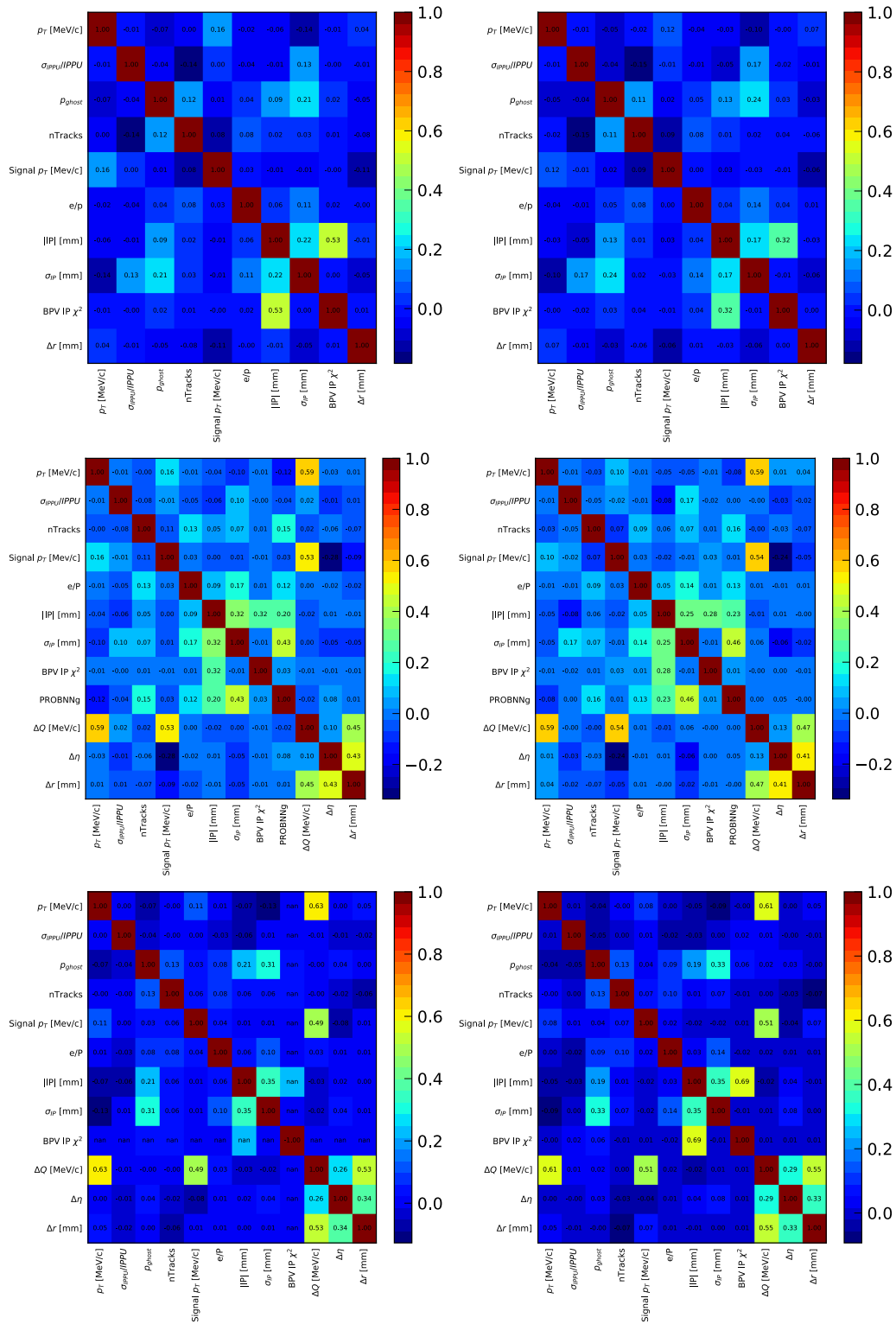


Figure 3.12 – Pearson correlation coefficients between the input features of the Run 1 new (top), Run 2 B2CC (middle) and Run 2 B2OC (bottom) BDT classifiers, for candidates with a correct (left) and wrong (right) decision from the OSe tagger.

3.3. Optimisation of the opposite-side electron tagger

Table 3.8 – Features considered for the BDT used to evaluate the predicted mistag of the OSe tagger. For each tuning (Run 1 new, Run 2 B2CC and Run 2 B2OC), the symbol $\checkmark(\times)$ indicates if a given feature is included (discarded).

Feature	Description	Run 1 new	Run 2 B2CC	Run 2 B2OC
nTracks	Number of reconstructed tracks	\checkmark	\checkmark	\checkmark
p_T	Transverse momentum of tagging track	\checkmark	\checkmark	\checkmark
σ_{IP}	IP uncertainty of tagging track	\checkmark	\checkmark	\checkmark
Signal p_T	Transverse momentum of B candidate	\checkmark	\checkmark	\checkmark
BPV IP χ^2	IP χ^2 of tagging track w.r.t B vertex	\checkmark	\checkmark	\checkmark
p_{ghost}	Ghost probability	\checkmark	\times	\checkmark
e/p	Inverse rigidity	\checkmark	\checkmark	\checkmark
Δr	difference in r coordinate between B and tagging track	\checkmark	\checkmark	\checkmark
$ IP $	Absolute value of tagging track IP Significance of the IP	\checkmark	\checkmark	\checkmark
$\sigma_{IPPU}/IPPU$	w.r.t pile-up vertex for tagging track	\checkmark	\checkmark	\checkmark
PROBNNg	Ghost probability from neural networks	\times	\checkmark	\times
$\Delta\eta$	Difference in pseudorapidity between B and tagging track	\times	\checkmark	\checkmark
ΔQ	Magnitude of difference in momenta between B and tagging track	\times	\checkmark	\checkmark

3.3.4 Performance evaluation

Performance on $B^+ \rightarrow J/\psi K^+$ and $B^+ \rightarrow \bar{D}^0 \pi^+$ data

Once the BDT is trained on the training sample, the mistag η is predicted for each candidate using the evaluation samples. Then, a two-fold evaluation is applied:

- the mistag calibration is determined on the first evaluation sample. The obtained calibration is then applied to the second evaluation sample, and a calibrated per-event tagging power is computed on the latter;
- same as above, but with the two evaluation samples swapped.

The calibrated per-event tagging power is computed by considering, for each tagged B candidate, only the tagging particle with the highest transverse momentum. The calibration model consists of a first order natural spline with a logistic link function. The result of these calibrations are shown in Fig. 3.16. The calibrated per-event tagging power is reported in Table 3.9.

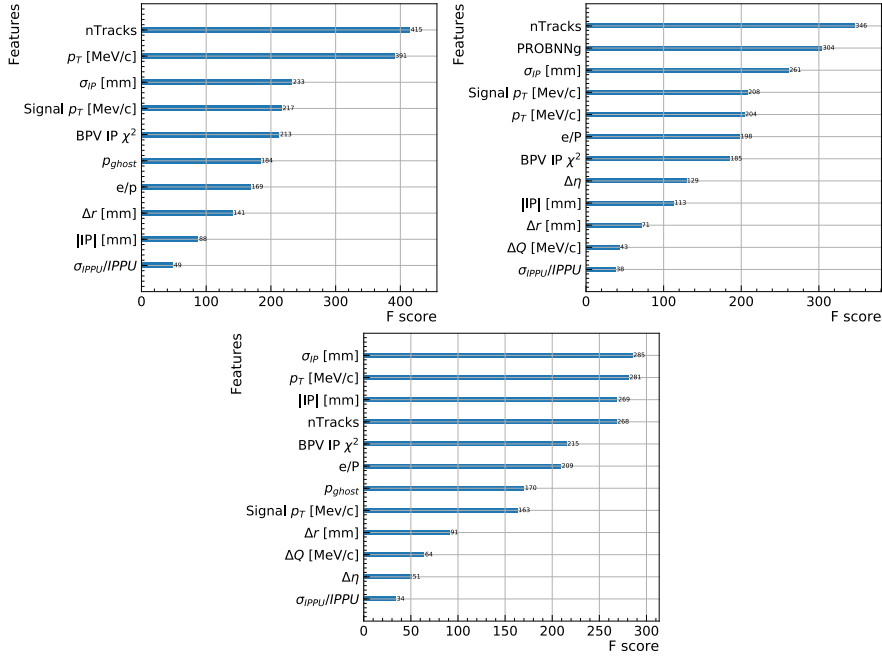


Figure 3.13 – Feature importance for the BDT classifiers of the Run 1 new (top left), Run 2 B2CC (top right) and Run 2 B2OC (bottom) implementations of the OSe tagger.

Performance on $B^0 \rightarrow D^- \pi^+$ data

The performance (tagging efficiency, mistag probability, tagging power) of the calibrated OSe tagger is evaluated on Run 1 (2012) and Run 2 (2016) *sWeighted* data samples of $B^0 \rightarrow D^- \pi^+$ decays. These decays ensure a robust estimation of the performance thanks to the high statistics collected at LHCb. Moreover, this channel was not exploited in the development of the OSe tagger, so that it constitutes an independent validation of these algorithms. The performance of the other OS taggers (OS_μ , OS_K , OS_{Charm} , OS_{Vtx} , and their combination) is presented as well in this section in order to provide a complete overview.

The calibration and the performance evaluation are done as follows:

- each sample (Run 1 and Run 2) is split randomly in two subsamples;
- the calibrations are found on one subsample for all OS taggers;
- the obtained calibrations are applied to the other subsample, and the calibrated performance is evaluated.
- the calibrated OS taggers are combined, the combination is calibrated in order to correct for effects due to correlations among taggers, and the performance of the calibrated combination is evaluated.

3.3. Optimisation of the opposite-side electron tagger

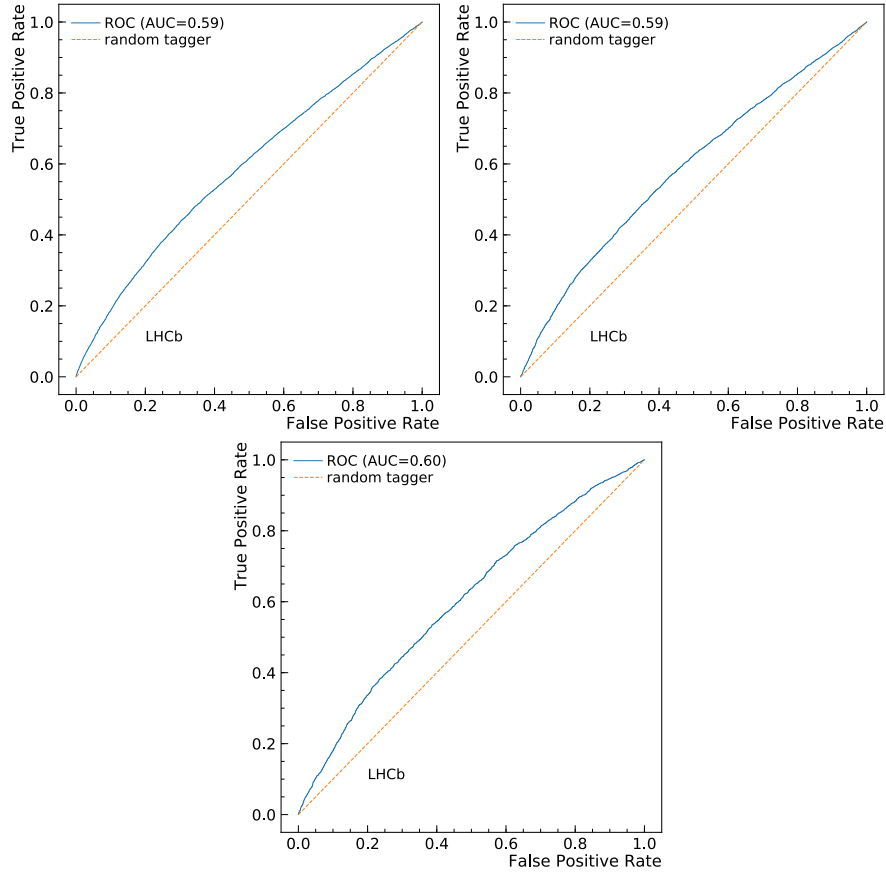


Figure 3.14 – True positive rate as a function of the false positive rate (ROC curves) for the BDT classifiers of the Run 1 new (top left), Run 2 B2CC (top right) and Run 2 B2OC (bottom) implementations of OSe. The obtained ROC curves are represented in blue, while the expected ROC curve in case of random tag decision is shown as a dashed orange line. For each BDT, the *Area Under the ROC Curve* (AUC score) is reported as well.

The calibrations are obtained via a time-dependent analysis of the $B^0 \rightarrow D^- \pi^+$ decays, where acceptance and resolution effects are neglected as described in Ref. [61]; moreover, the Cabibbo-suppressed decay mode $B^0 \rightarrow D^+ \pi^-$ is ignored as well. The chosen model $\omega(\eta)$ for each tagger is a GLM model with a logistic link function, and a first order spline as basis function. The results of the calibration and the mistag distribution of each OSe implementation are shown in Figs. 3.17 and 3.18; the calibration and the mistag of the corresponding OS combinations are also reported in Figs. 3.19 and 3.20.

The performance is reported in Tables 3.10 and 3.11. The Run 1 new tuning allows to gain a relative 9% in tagging power for the OSe tagger on Run 1 data; the corresponding, relative gain of the OS combination is 3%. The tagging power of OSVtx and OScharm,

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Table 3.9 – Calibrated, per-event tagging power ε_{eff} (in %) of the OSe algorithms obtained on the evaluation sets of each OSe implementation. The errors include both statistical uncertainty and uncertainties from the calibration procedure. The average is computed by assuming uncorrelated measurements.

Algorithm	set 1	set 2	average
Run 1 new	0.513 ± 0.040	0.496 ± 0.038	0.504 ± 0.028
Run 2 B2CC	0.324 ± 0.031	0.364 ± 0.033	0.343 ± 0.023
Run 2 B2OC	0.455 ± 0.043	0.434 ± 0.041	0.444 ± 0.030

Table 3.10 – Performance (tagging efficiency, average mistag and tagging power in %) of the OS taggers on *sWeighted* Run 1 $B^0 \rightarrow D^-\pi^+$ data. The numbers for OSe and the OS combination are shown separately for the Run 1 old and Run 1 new tunings. The first uncertainty is statistical and the second comes from the calibration.

Tagger	ε_{tag}	$\langle\omega\rangle$	ε_{eff}
OSVtx	22.026 ± 0.100	$37.295 \pm 0.030 \pm 0.376$	$1.422 \pm 0.009 \pm 0.084$
OSCharm	4.632 ± 0.050	$34.026 \pm 0.049 \pm 0.824$	$0.473 \pm 0.006 \pm 0.049$
OSe Run 1 old	3.028 ± 0.041	$30.570 \pm 0.113 \pm 0.963$	$0.457 \pm 0.008 \pm 0.045$
OSe Run 1 new	4.337 ± 0.049	$33.089 \pm 0.085 \pm 0.777$	$0.496 \pm 0.007 \pm 0.046$
OS μ Run 1	8.539 ± 0.067	$28.756 \pm 0.071 \pm 0.582$	$1.541 \pm 0.016 \pm 0.085$
OSK Run 1	18.800 ± 0.094	$36.724 \pm 0.031 \pm 0.417$	$1.325 \pm 0.009 \pm 0.083$
OS combination Run 1 old	39.004 ± 0.117	$34.679 \pm 0.035 \pm 0.273$	$3.662 \pm 0.020 \pm 0.131$
OS combination Run 1 new	39.733 ± 0.118	$34.576 \pm 0.035 \pm 0.270$	$3.781 \pm 0.021 \pm 0.133$

which were trained on Run 1 data, increases on Run 2 data compared to Run 1; for this reason, no specific optimisation for the Run 2 conditions is performed. The tagging power of OSe, OS μ , and OSK with the Run 1 tunings is lower on Run 2 data compared to Run 1. However, compared to the Run 1 tunings, the Run 2 tunings show a relative improvement in tagging power of about $\sim 160\%$ for OSe, and $\sim 6\%$ for OS μ and OSK on Run 2 data. This allows to recover similar performances as the ones obtained on Run 1 data with the Run 1 tunings, both for the individual taggers and their combination. Moreover, the Run 2 B2CC and B2OC tunings show consistent tagging powers on Run 2 data, meaning that the optimisation is robust against the different kinematics of the adopted decays.

3.3. Optimisation of the opposite-side electron tagger

Table 3.11 – Performance (tagging efficiency, average mistag and tagging power in %) of the OS taggers on *sWeighted* Run 2 $B^0 \rightarrow D^-\pi^+$ data. The numbers for OSe, OS μ , OSK, and the OS combination are shown separately for the Run 1, Run 2 B2CC, and Run 2 B2OC tunings. The first uncertainty is statistical and the second comes from the calibration.

Tagger	ϵ_{tag}	$\langle\omega\rangle$	ϵ_{eff}
OSVtx	20.834 ± 0.075	$36.139 \pm 0.029 \pm 0.301$	$1.601 \pm 0.009 \pm 0.070$
OSCharm	5.025 ± 0.040	$33.875 \pm 0.041 \pm 0.624$	$0.523 \pm 0.005 \pm 0.040$
OSe Run 1 old	1.868 ± 0.025	$34.300 \pm 0.096 \pm 0.941$	$0.184 \pm 0.003 \pm 0.022$
OSe Run 2 B2CC	4.451 ± 0.038	$33.352 \pm 0.081 \pm 0.608$	$0.493 \pm 0.006 \pm 0.036$
OSe Run 2 B2OC	3.333 ± 0.033	$30.917 \pm 0.075 \pm 0.702$	$0.486 \pm 0.006 \pm 0.036$
OS μ Run 1	8.343 ± 0.051	$30.357 \pm 0.042 \pm 0.466$	$1.288 \pm 0.010 \pm 0.061$
OS μ Run 2 B2CC	9.151 ± 0.053	$30.837 \pm 0.041 \pm 0.432$	$1.344 \pm 0.010 \pm 0.061$
OS μ Run 2 B2OC	8.040 ± 0.050	$29.174 \pm 0.043 \pm 0.463$	$1.395 \pm 0.010 \pm 0.062$
OSK Run 1	15.737 ± 0.067	$35.902 \pm 0.030 \pm 0.357$	$1.251 \pm 0.008 \pm 0.063$
OSK Run 2 B2CC	19.516 ± 0.073	$36.889 \pm 0.026 \pm 0.310$	$1.342 \pm 0.007 \pm 0.064$
OSK Run 2 B2OC	15.793 ± 0.067	$35.565 \pm 0.030 \pm 0.348$	$1.316 \pm 0.008 \pm 0.063$
OS combination Run 1 old	36.239 ± 0.088	$35.285 \pm 0.024 \pm 0.227$	$3.139 \pm 0.013 \pm 0.097$
OS combination Run 2 B2CC	40.154 ± 0.090	$35.123 \pm 0.025 \pm 0.210$	$3.555 \pm 0.014 \pm 0.100$
OS combination Run 2 B2OC	36.555 ± 0.089	$34.225 \pm 0.026 \pm 0.220$	$3.638 \pm 0.015 \pm 0.102$

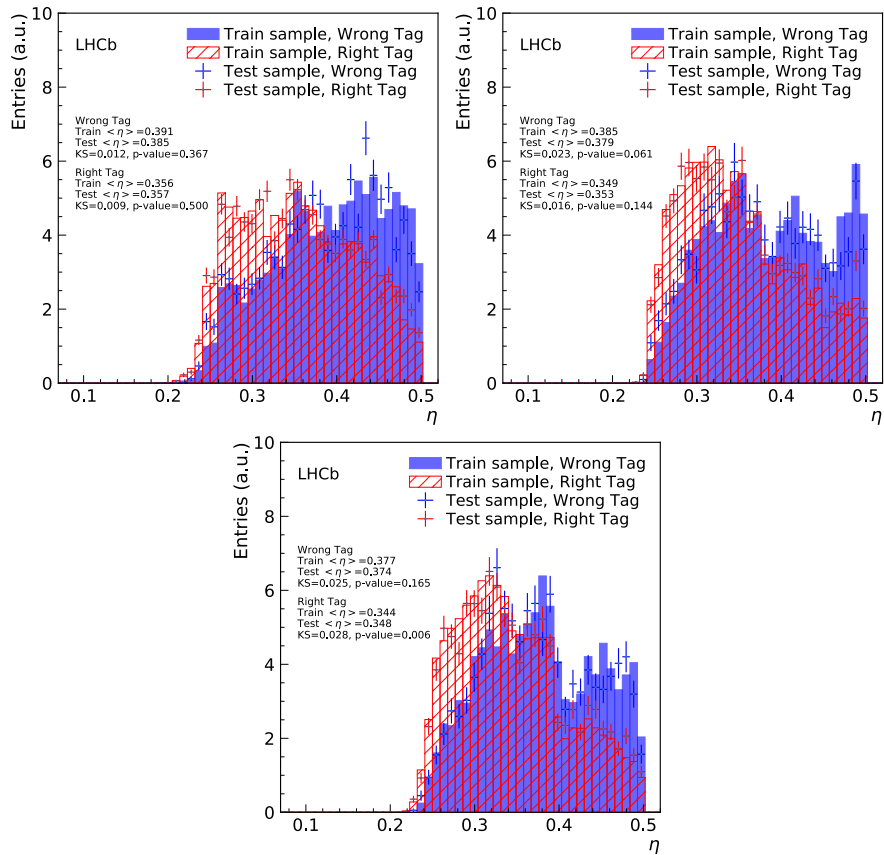


Figure 3.15 – s Weighted distributions of the mistag probability η predicted by the BDT classifiers of the Run 1 new (top left), Run 2 B2CC (top right) and Run 2 B2OC (bottom) versions of the OSe tagger. The blue-solid (red-hatched) histogram represents the training data for candidates having the wrong (right) tag decision. The blue (red) points indicate the test data for candidates with wrong (right) tag decision. The overtraining is checked, separately for candidates with wrong and right tag decision, by means of a Kolmogorov-Smirnov (KS) test to measure the compatibility between training data and test data. The conventional value of 0.05 is chosen as significance level to reject the hypothesis of compatibility.

3.3. Optimisation of the opposite-side electron tagger

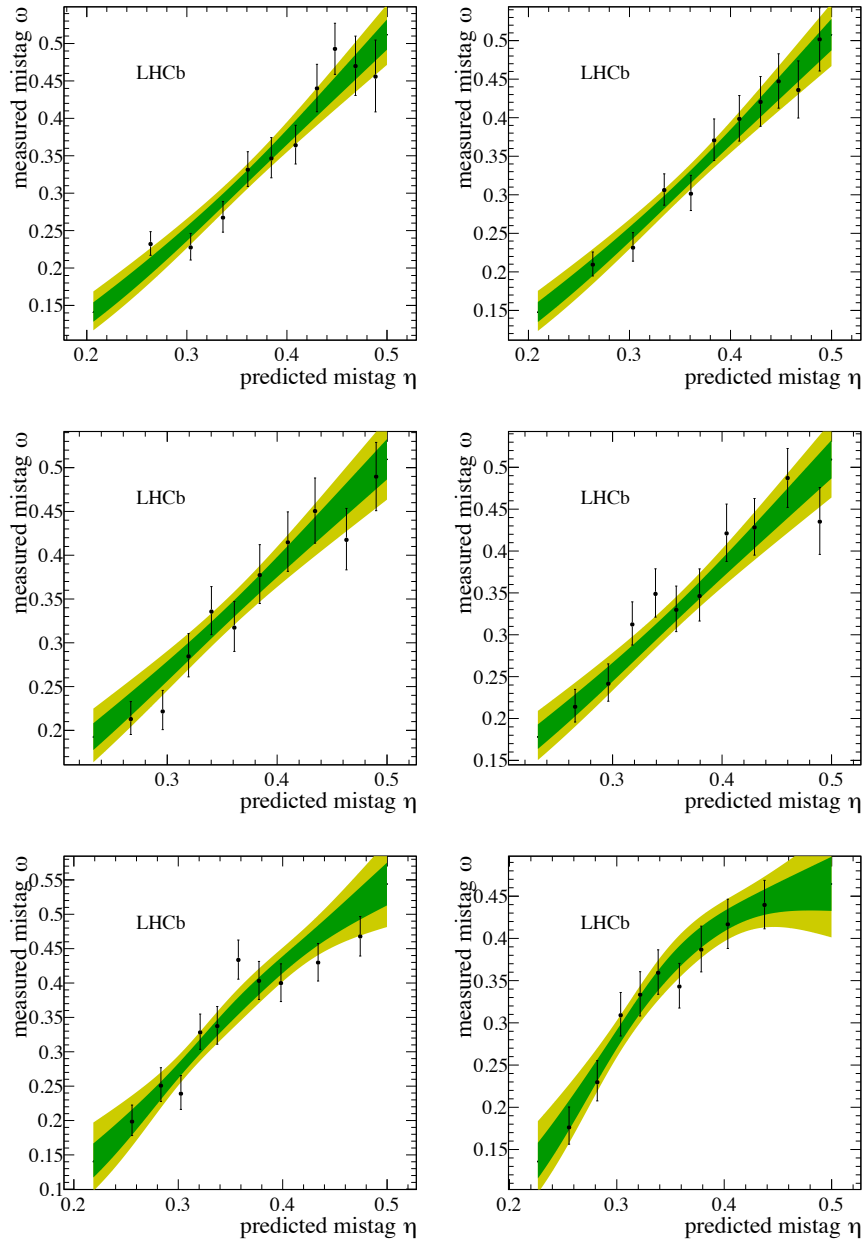


Figure 3.16 – OSe mistag calibration results for the (top) Run 1 new, (middle) Run 2 B2CC and (bottom) Run 2 B2OC optimisations. Left: calibration obtained on the second evaluation sample plotted together with the first evaluation sample. Right: calibration obtained on the first evaluation sample plotted together with the second evaluation sample. The $sWeighted$ data sample is shown as black points. The green (yellow) band indicates the 68% (95%) C.L. interval for the fitted calibration functions.

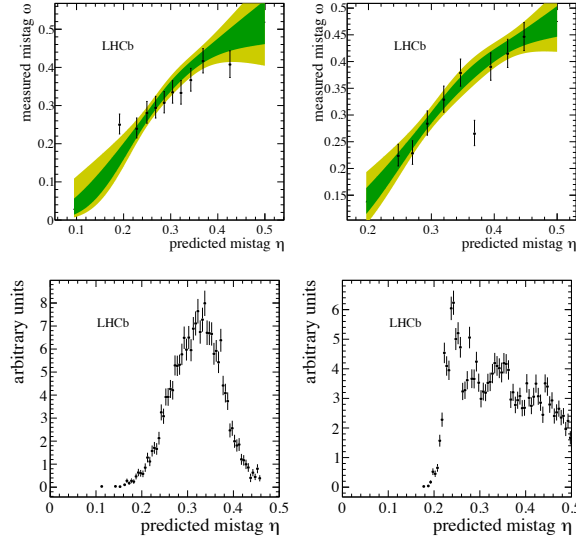


Figure 3.17 – Top: mistag calibration results on $sWeighted$ Run 1 $B^0 \rightarrow D^- \pi^+$ data for the Run 1 old (left) and Run 1 new (right) versions of the OSe tagger. The $sWeighted$ data sample is shown as black points. The green (yellow) band indicates the 68% (95%) C.L. interval for the fitted calibration functions. Bottom: distributions of the uncalibrated mistag η .

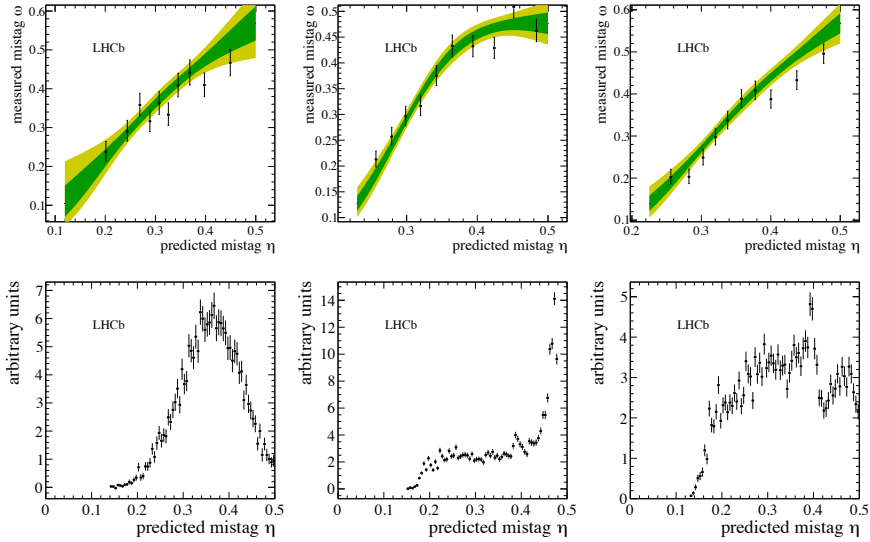


Figure 3.18 – Mistag calibration results on $sWeighted$ Run 2 $B^0 \rightarrow D^- \pi^+$ data for the OSe taggers. The results obtained with the Run 1 old (left), Run 2 B2CC (center), and Run 2 B2OC (right) tunings are shown. The $sWeighted$ data sample is shown as black points. The green (yellow) band indicates the 68% (95%) C.L. interval for the fitted calibration functions. Bottom: distributions of the uncalibrated mistag η .

3.3. Optimisation of the opposite-side electron tagger

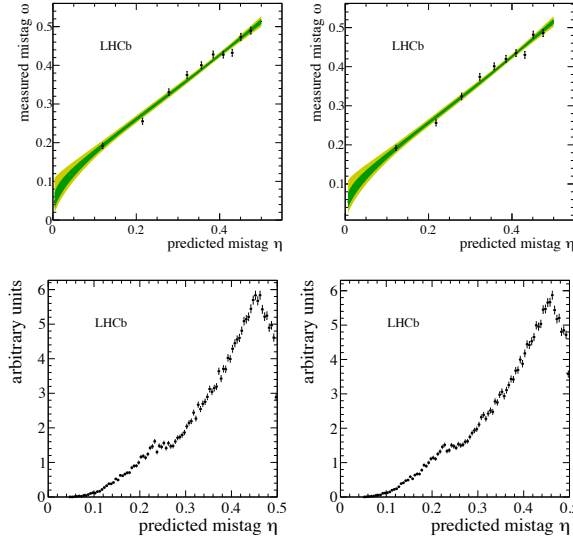


Figure 3.19 – Top: mistag calibration results on *sWeighted* Run 1 $B^0 \rightarrow D^- \pi^+$ data for the combination of the OS taggers. The results obtained with the Run 1 old (left) and Run 1 new (right) tunings of OSe are shown. The *sWeighted* data sample is shown as black points. The green (yellow) band indicates the 68% (95%) C.L. interval for the fitted calibration functions. Bottom: distributions of the uncalibrated mistag η .

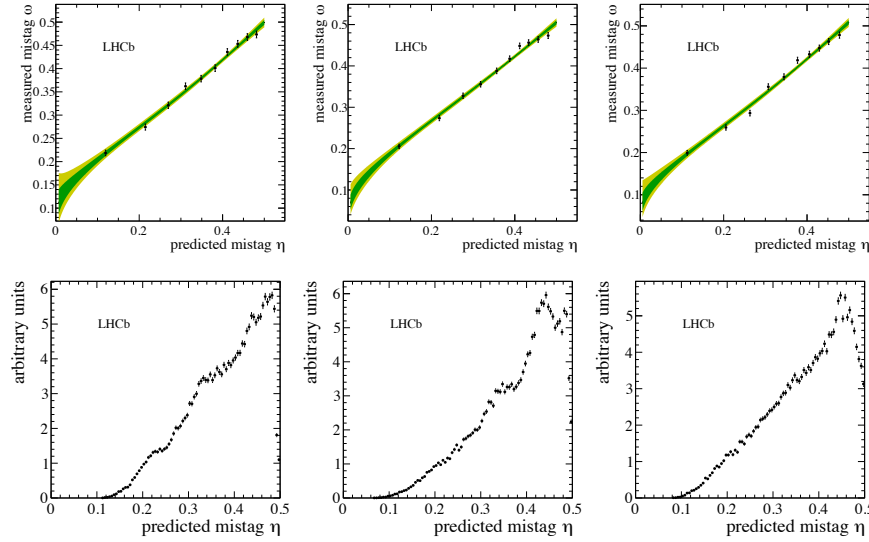


Figure 3.20 – Mistag calibration results on *sWeighted* Run 2 $B^0 \rightarrow D^- \pi^+$ data for the combination of the OS taggers. The results obtained with the Run 1 old (left), Run 2 B2CC (center), and Run 2 B2OC (right) tunings of OSe, OS μ , and OSK are shown. The *sWeighted* data sample is shown as black points. The green (yellow) band indicates the 68% (95%) C.L. interval for the fitted calibration functions. Bottom: distributions of the uncalibrated mistag η .

4 Selection of $B^0 \rightarrow D^{\mp} \pi^{\pm}$ decays

4.1 Data sample and preselection

The sample of data is passed through the following selection steps:

1. stripping and trigger requirements;
2. a *cut-based* preselection;
3. vetoes for misidentified backgrounds and wrongly associated *primary vertices* (PVs);
4. a multivariate classification (MVA);
5. a final randomised multiple candidate selection.

In what follows, the details of each step are provided.

4.1.1 Stripping and trigger requirements

Signal $B^0 \rightarrow D^\mp \pi^\pm$ candidates are reconstructed using a dedicated stripping line (called `B02DPiD2HHHBeauty2CharmLine`). Each event is required to have less than 500 long tracks. The criteria that the charged tracks have to fulfill are listed in Table 4.1. Three of these hadrons have to form a common vertex to build a D^\mp meson. Further requirements on the D^\mp combination are given in Table 4.2. The B^0 candidates are built by combining a D^\mp candidate and a bachelor particle if the requirements listed in Table 4.3 are fulfilled. Finally, a bagged boosted decision tree (BDT) classifier [65], which is trained on simulated data, is applied. A minimum value of 0.05 is required for the output value of the BDT. Stripped candidates are then filtered according to how they were selected at the trigger level: no specific requirements are made at L0; at HLT1, B candidates are required to be TOS from the `Hlt1TrackAllL0Decision` trigger line; at HLT2, B candidates are required to be TOS from one of the following lines: `Hlt2Topo2BodyBBDTDecision`, `Hlt2Topo3BodyBBDTDecision` or `Hlt2Topo4BodyBBDTDecision`. These trigger lines are described in detail in Refs. [66, 67].

Figure 4.1 shows the $D^\mp \pi^\pm$ and $K^\pm \pi^\mp \pi^\mp$ mass distributions of the reconstructed candidates after the stripping and trigger selections. In the $D^\mp \pi^\pm$ mass distribution the $B^0 \rightarrow D^\mp \pi^\pm$ signal peak is already visible at around $5280 \text{ MeV}/c^2$. The

Table 4.1 – Stripping requirements applied in the selection of charged tracks. The more stringent requirements given in brackets are for the bachelor track. The $\text{IP}\chi^2$ is the vertex-fit χ^2 difference for the PV reconstructed with and without the B^0 candidate.

track χ^2/ndof	< 3.0 (< 2.5)
momentum p	$> 1 \text{ GeV}/c$ ($> 5 \text{ GeV}/c$)
transverse momentum p_T	$> 100 \text{ MeV}/c$ ($> 500 \text{ MeV}/c$)
$\text{IP}\chi^2$ w.r.t. any PV	> 4.0
ghost probability	< 0.4

4.1. Data sample and preselection

Table 4.2 – Stripping requirements on the three-track combinations forming D^\mp candidates. DOCA is the Distance Of Closest Approach of the daughter particles w.r.t. each other, and DIRA indicates the cosine of the angle between the momentum of the D^\mp meson and the direction from the best PV to the decay vertex. The best PV is defined as the vertex with the lowest $\text{IP}\chi^2$.

$\sum p_T(hhh)$	$> 1800 \text{ MeV}/c$
DOCA	$< 0.5 \text{ mm}$
$m(hhh)$	$\in [1769.62, 2068.49] \text{ MeV}/c^2$
D^\mp vertex χ^2/ndof	< 10.0
D^\mp vertex separation χ^2 to any PV	> 36
D^\mp DIRA	> 0.0

Table 4.3 – Stripping requirements on the $D^\mp\pi^\pm$ combination (before the BDT requirement mentioned in the text).

B^0 vertex χ^2/ndof	< 10.0
reconstructed B^0 proper decay time t	$> 0.2 \text{ ps}$
$\text{IP}\chi^2$ w.r.t. the best PV	< 25.0
B^0 DIRA	> 0.999

structure at masses lower than the B peak originates from partially reconstructed $B^0 \rightarrow D\rho$ and $B^0 \rightarrow D^*\pi$ decays. The $K^\pm\pi^\mp\pi^\mp$ mass distribution features a clearly visible $D^\mp \rightarrow K^\pm\pi^\mp\pi^\mp$ peak at $1870 \text{ MeV}/c^2$ and a $D^{*\mp} \rightarrow D(K^\pm\pi^\mp)\pi^\mp$ peak around $2010 \text{ MeV}/c^2$.

4.1.2 Preselection and sample definitions

Additional preselection criteria (shown in Table 4.4) are applied offline. In order to obtain the correct correlations between the uncertainties on vertex positions, particle momenta, flight distances, decay times, and invariant masses, a Kalman filter, known as `DecayTreeFitter` (DTF) [68], is used. The decay-time related observables are derived from a DTF fit where the position of the primary vertex has been used to constrain the production vertex of the B^0 meson. To determine the momentum and the invariant mass of the B^0 meson, the invariant mass of the D^\mp meson is constrained to the central value of the PDG ($m_{D^\mp}^{\text{PDG}} = 1869.61 \text{ MeV}/c^2$ [22]) in a separate DTF fit. The $\text{PID}K$ variable of Eq. 2.1 is used to identify the kaon and the pions from the D^\mp decays, and to identify the bachelor pion from the B^0 decay. The requirement on the $\text{PID}K$ of the bachelor pion defines two samples of candidates: the so-called *pion sample* ($\text{PID}K \leq 5$) and the so-called *kaon sample* ($\text{PID}K > 5$). This distinction will be useful in the fit to the B^0 mass distribution described in Sec. 4.3 for determining the sample composition. All the following selection steps are applied to both the pion and kaon samples.

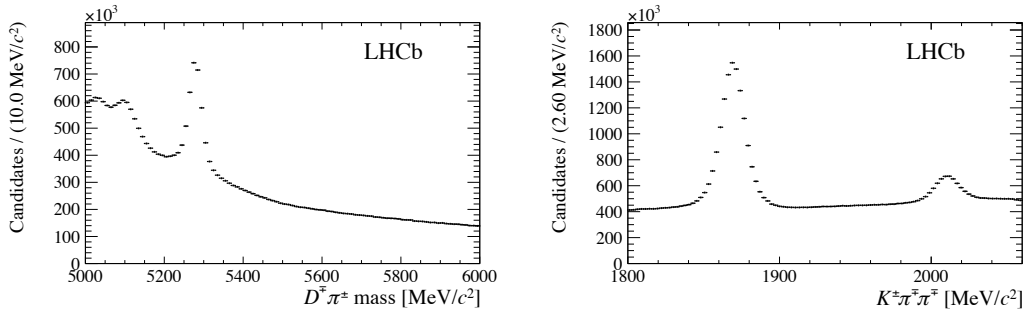


Figure 4.1 – $D^\mp \pi^\pm$ and $K^\pm \pi^\mp \pi^\mp$ mass distributions of the reconstructed $B^0 \rightarrow D^\mp \pi^\pm$, $D^\mp \rightarrow K^\pm \pi^\mp \pi^\mp$ candidates after the stripping and trigger selection.

Table 4.4 – Offline preselection requirements.

B^0 candidate decay time	> 0.2 ps
$ m(K^\pm \pi^\mp \pi^\mp) - m_{D^\mp}^{\text{PDG}} $	< 35 MeV/ c^2
PIDK for pions	$< +8$ from D^\mp
PIDK for kaon	> -2 from D^\mp

4.1.3 Vetoes against physics backgrounds

Misidentification of muons, kaons and protons as pions leads to exclusive backgrounds. These are suppressed by means of explicit *veto*s. In order to reduce contributions from semileptonic decays such as $B^0 \rightarrow D^- (\rightarrow K^+ \pi^- \pi^-) \mu^+ \nu_\mu$, the bachelor pion is required to have no hits in the muon chambers. A $p \rightarrow \pi$ mis-identification can lead to background contributions from $\Lambda_b^0 \rightarrow \Lambda_c^+ (\rightarrow K^- \pi^+ p) \pi^-$. To reduce these contributions, the proton mass hypothesis is applied separately to both pion candidates from the D^\mp final state. The invariant mass of the three hadrons is recalculated and if the candidate is inside a ± 30 MeV/ c^2 (± 50 MeV/ c^2) window around the Λ_c^+ mass, $m_{\Lambda_c^+} = 2286.46$ MeV/ c^2 [22], it is required to have $\text{PID}p < -8.0$ ($\text{PID}p < -5.0$). A plot showing the distributions before and after applying the veto is given in Fig. 4.2. This requirement shows a signal efficiency of $(93.48 \pm 0.06)\%$. The rejection of $\Lambda_b^0 \rightarrow \Lambda_c^+ \pi^-$ is checked with simulation. After stripping and preselection alone, $(99.720 \pm 0.004)\%$ of the $\Lambda_b^0 \rightarrow \Lambda_c^+ \pi^-$ decays are rejected, and this veto rejects $(76.6 \pm 0.6)\%$ of the remaining $\Lambda_b^0 \rightarrow \Lambda_c^+ \pi^-$ decays.

In the same way as protons may be misidentified as pions, it is possible for kaons to be misidentified as pions. Such a mis-identification would lead to background contributions from $B_s^0 \rightarrow D_s^- (\rightarrow K^+ K^- \pi^-) \pi^+$. To check for these contributions, the kaon mass hypothesis is applied in turn to each of the two pions from the D^\mp final state. The invariant mass of the resulting $K^\pm K^\mp \pi^\mp$ system is recalculated and plotted for two different ranges of the B_s^0 mass: the first range, from 5330 to 5400 MeV/ c^2 , covers the signal region of the B_s^0 meson as possible background contribution. The second range,

4.1. Data sample and preselection

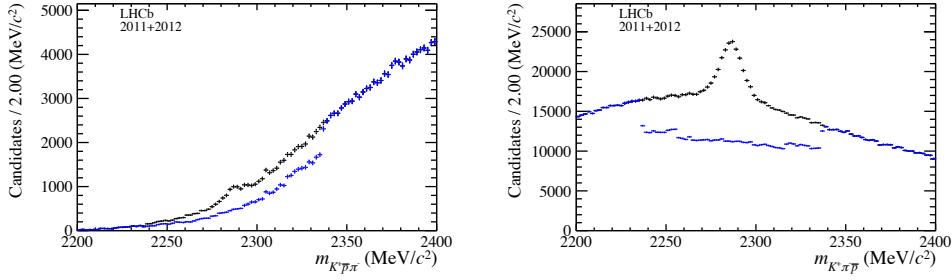


Figure 4.2 – Distributions of the invariant mass of the $K^\pm\pi^\mp p^\mp$ combinations where each of the two daughter pions of the D^\mp meson candidate is assigned in turn the proton mass. The distribution is given without (black) and with (blue) the Λ_c^\mp veto described in the text. On the left (right) the proton mass hypothesis is applied to the pion with the lower (higher) p_T .

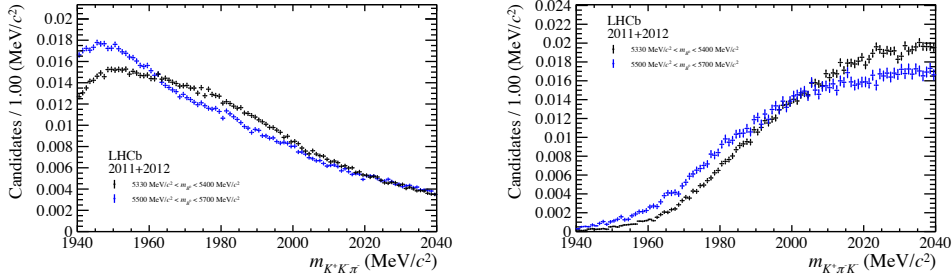


Figure 4.3 – Distributions of the invariant mass of the $K^\pm K^\mp \pi^\mp$ combinations where each of the two daughter pion of the D^\mp meson candidate is assigned in turn the kaon mass. The distributions are given in the B_s^0 meson mass range from 5330 to 5400 MeV/c^2 (black) and in the B_s^0 meson mass range from 5500 to 5700 MeV/c^2 (blue) after applying the Λ_c^\mp veto. On the left (right) the kaon mass hypothesis is applied to the pion with the lower (higher) p_T .

from 5500 to 5700 MeV/c^2 is the upper mass *sideband* for this possible background contamination. As can be seen in Fig. 4.3 the distribution in the B_s^0 signal region does not show any significant peaking structure compared to that in the upper B_s^0 mass sideband region. The visible differences are expected as the distributions arise from different kinematic ranges. To double check for possible resonant contributions from a kaon mis-identification, the decay of the D^\mp meson after applying the kaon mass hypothesis is investigated. Possible resonant decays of the D meson can take place via a K^* or ϕ resonance. These resonances would be visible in the $K\pi$ and KK invariant mass distributions, which are plotted for the same two ranges in Fig. 4.4. As the distributions in the signal and background ranges look compatible, the D_s^\mp contamination is negligible and no veto is applied. In the same way as for the D^\mp meson daughters, it is also possible that the bachelor pion candidate is actually a misidentified kaon. This mis-identification could

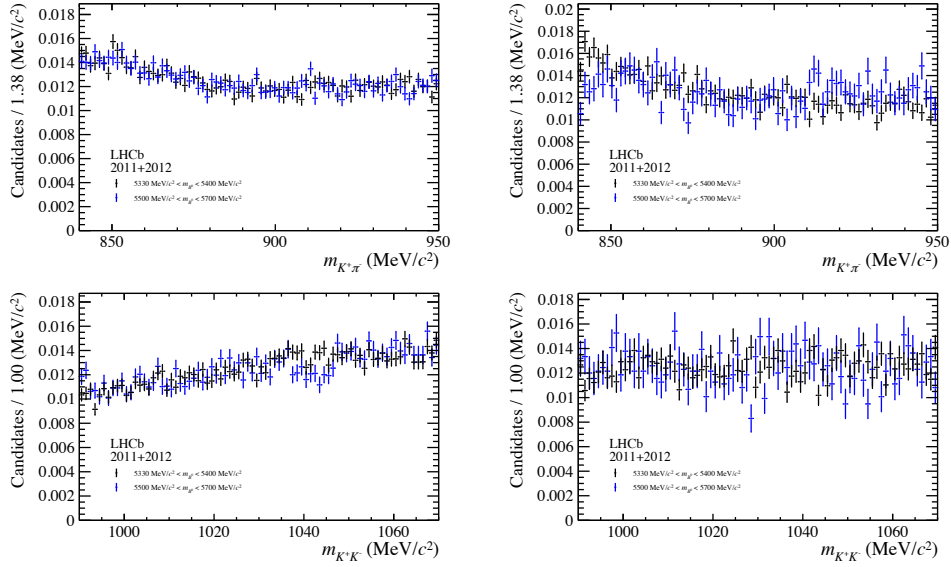


Figure 4.4 – Distributions of the invariant mass of the $K^\mp \pi^\pm$ combination (top) and the $K^\mp K^\pm$ combination (bottom), where each of the two daughter pion of the D^\mp meson candidate is assigned in turn the kaon mass. The distributions are shown in the B_s^0 meson mass range of $[5330, 5400]$ MeV/c^2 (black) and $[5500, 5700]$ MeV/c^2 (blue) as described in the text. Additionally, the D_s^\mp mass is required to be in $[1940, 2040]$ MeV/c^2 . On the left (right) the kaon mass hypothesis is applied to the pion with the lower (higher) p_T .

lead to background contributions of $B^0 \rightarrow D(\rightarrow K^\pm \pi^\mp) \pi^\mp K^\pm$. A similar background, $B^0 \rightarrow D(\rightarrow K^\pm \pi^\mp) K^\mp \pi^\pm$, could arise if one D^\mp meson daughter pion is misidentified as a kaon and combined with the bachelor pion. To check for this contribution, the kaon mass hypothesis is applied to the bachelor pion and the D meson daughter pions, and the invariant mass distributions for the four possible $K\pi$ systems are plotted after applying the MVA classifier described in Sec. 4.1.5 (Fig. 4.5). As the distributions show no significant peaking structures, this contribution is neglected and no specific cuts are applied.

4.1.4 Wrongly associated primary vertices

Given an average number of total pp interactions per bunch crossing of $\nu = 2.5$, a large fraction of events have more than one reconstructed PV. The PV to which the B^0 candidate has the smallest $\text{IP}\chi^2$ (*best* PV) is chosen as the B^0 production vertex.

In events where the association of the B^0 candidate to its best PV is wrong, the reconstructed decay time of this candidate will be incorrect. These wrongly associated candidates cause a large tail in the decay time distribution, which can be clearly observed in signal MC where the true decay time is known: giving each candidate a weight equal

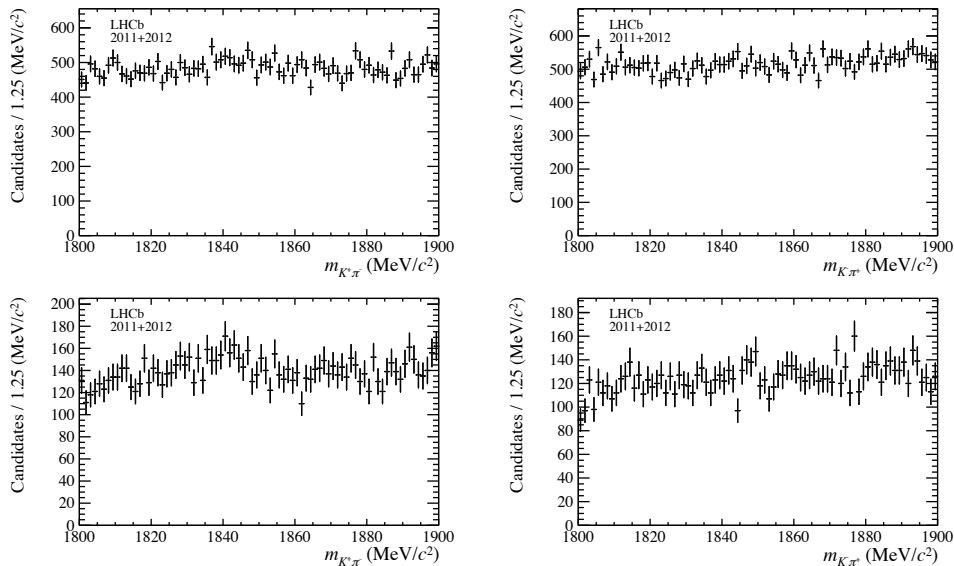


Figure 4.5 – Distributions of the invariant mass of the four possible $K^\mp\pi^\pm$ combinations. In the top (bottom) plots the bachelor pion is combined with the pion from the D^\mp meson with lower (higher) p_T . In the left (right) plots the kaon mass hypothesis is applied to the bachelor pion (pions from the D^\mp meson).

to $e^{t/\tau}$, where τ is the true lifetime, leads to an excess of candidates at high decay times. To remove these incorrect associations in MC, one can compare the z position of the associated PV with the z position of the true PV and reject the candidate if the distance between those positions exceeds 5 times its uncertainty (Fig. 4.6). In real data, the true PV is unknown, so a selection involving the B^0 impact parameter $\chi_{\text{DTF,PV}}^2$ is adopted instead. If there are multiple PVs in an event, the B^0 candidate is constrained by DTF to originate, in turn, from each of them; then, the $\chi_{\text{DTF,PV}}^2$ with respect to all the other PVs is computed, and the smallest one (called $\text{MinIP}\chi^2$) is considered. (Fig. 4.7). For events with only a single PV, $\text{MinIP}\chi^2$ is not defined. The main advantage of this $\text{MinIP}\chi^2$ variable is that all PVs are treated equally, without any biasing choice. The cut on $\text{MinIP}\chi^2$ is optimised to retain 98% of the truth-matched signal candidates in MC. The optimal requirement is then found at $\text{MinIP}\chi^2 > 16.5$. A plot showing the signal MC weighted decay time distribution after applying this cut is given in Fig. 4.7.

4.1.5 Development of an MVA classifier

The combinatorial background, consisting of candidates created from random combinations of tracks, is rejected by using a Boosted Decision Tree (BDT) classifier [69, 70]. The signal input to the training stage consists of signal MC candidates simulated under 2012 data-taking conditions, while the upper mass sideband above $5500 \text{ MeV}/c^2$ from the 2012

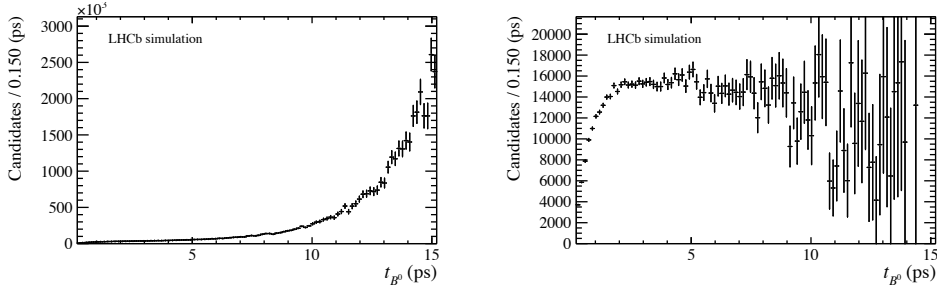


Figure 4.6 – Left: decay time distribution of signal MC events weighted with $e^{t/\tau}$, where τ is the true lifetime. At high decay times an excess of candidates can be observed. Right: same distribution after requiring that the absolute difference between the best PV z position and the true PV z position is within 5 times the best PV z position uncertainty. As the excess of candidates at high decay times vanishes (from left to right), it is concluded that this excess is due to candidates wrongly associated to their PV.

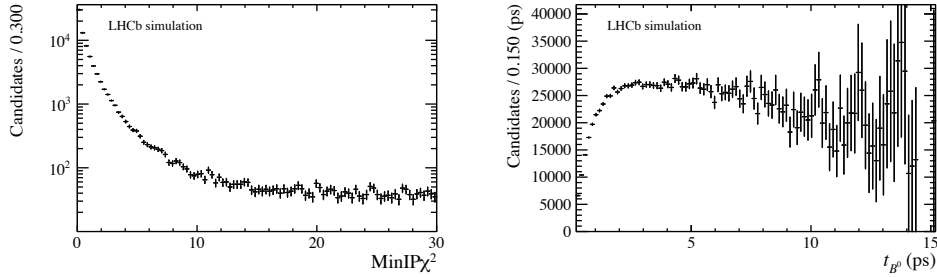


Figure 4.7 – Left: distribution of $\text{MinIP}\chi^2$ for signal MC events. Right: decay time distribution of signal MC events weighted with $e^{t/\tau}$, where τ is the true lifetime, after requiring that $\text{MinIP}\chi^2 > 16.5$.

data sample is used as template for the combinatorial background. The BDT is trained on one half of these samples, the other half being used to test its performance. Before the BDT training, all previous selection steps (the cut-based preselection, the mass vetoes and the wrongly associated PV veto) are applied. To reduce the number of input features, the ones with a correlation larger than 97% with any other feature are removed. The 16 final input features are listed in Table 4.5. The correlation matrices between the input features in the signal and the background samples are shown in Fig. 4.8, while the distributions of the input features can be found in Appendix C.1.

The BDT implementation of TMVA [71] is used. The BDT is built out of 1700 trees, with a depth limited to four. For each node, at least 2.5% of the training events have to be present. The chosen boosting method is the AdaBoost [72] algorithm with a boost factor $\beta = 0.5$. The number of trees and the maximal depth of trees have been increased iteratively until no significant increase of the performance without overtraining was

4.1. Data sample and preselection

Table 4.5 – List of input features used in the training of the BDT.

B^0 candidate	$\cos \text{ of } \sphericalangle \left[\frac{\overrightarrow{\text{Vertex}}(B^0) - \overrightarrow{\text{PV}}, \vec{p}(B^0)}{\text{vertex } \chi^2} \right]$ DTF χ^2 with best PV constraint
D^\mp candidate	$\cos \text{ of } \sphericalangle \left[\frac{\overrightarrow{\text{Vertex}}(D^\mp) - \overrightarrow{\text{Vertex}}(B^0), \vec{p}(D^\mp)}{\text{flight distance } \chi^2 \text{ w.r.t. } B^0 \text{ vertex}} \right]$ IP χ^2 w.r.t. B^0 vertex IP χ^2 w.r.t. best PV radial flight distance vertex χ^2/ndof transverse momentum
bachelor π^\pm	IP χ^2 w.r.t. the best PV transverse momentum track χ^2/ndof
D^\mp daughters	IP χ^2 w.r.t. best PV

observed. The BDT is tested on the events that are not used in the training. The plot of this overtraining check is given in Fig. 4.9.

4.1.6 BDT selection optimisation

To estimate the best requirement on the output of the BDT classifier, the statistical uncertainty of the CP coefficients derived from the analysis of simulated samples is used as the figure of merit (FoM). To determine the sensitivity, the preselection, the mass vetoes and the wrongly associated PV veto are applied and the BDT classifier is calculated for every candidate. The BDT cut point is then scanned with a step size of 0.01 from -0.15 to 0.10 and a step size of 0.05 in the outer regions. For each cut point, a simulated (*toy*) sample is generated. This sample contains the same signal and combinatorial background yields as determined from the real dataset via a maximum likelihood fit of the B^0 mass distribution. Finally, a time-dependent analysis of each toy dataset is performed in order to estimate the statistical uncertainty on S_f and $S_{\bar{f}}$. These statistical uncertainties as a function of the BDT cut are shown in Fig. 4.10. Based on these distributions, the BDT cut point is chosen to be at 0.05 .

4.1.7 Multiple candidates

After the stripping selection and trigger requirements, approximately 9% of the events contain at least two B^0 candidates, and 18 – 20% of all B^0 candidates share an event. If the offline selection is also applied, around 0.4% of the events contain multiple B^0



Figure 4.8 – Correlation matrices of the input features used in the training of the BDT for signal (top) and background (bottom).

candidates, and 0.8% of all B^0 candidates share an event. More details are given in Appendix C.2. In order to be consistent with the prescription used in the stripping and trigger requirements, only the best PV is chosen; all events in which the best PV is no longer present after the offline selection are removed. Finally, since the remaining B^0 candidates are considered to be equally likely signal candidates, a single B^0 candidate per event is chosen randomly following the prescription of Ref. [73], which prevents any unexpected bias due to a more specific choice.

4.1.8 Selection performance

The offline selection performances are listed in Table 4.6. They are determined by using data candidates of the 2012 sample with an invariant B^0 mass above $5500 \text{ MeV}/c^2$ to represent combinatorial background, and signal MC candidates (see Sec. 4.2) to represent the signal. Additionally, the BDT performances are quoted in Table 4.7 split by magnet polarity and year of data taking. Finally, in order to check the contribution of non-resonant $B^0 \rightarrow K^+ \pi^- \pi^- \pi^+$ decays, the B^0 and D^\mp invariant mass distributions are analysed after applying the full offline selection (with the exception of the B^0 and D^\mp mass cuts) in two ways. First, the D^\mp mass distribution is plotted for candidates falling in a B^0 mass signal window (Fig. 4.11). From this plot, the maximal contamination from non-resonant

4.2. Simulation and expected sample composition

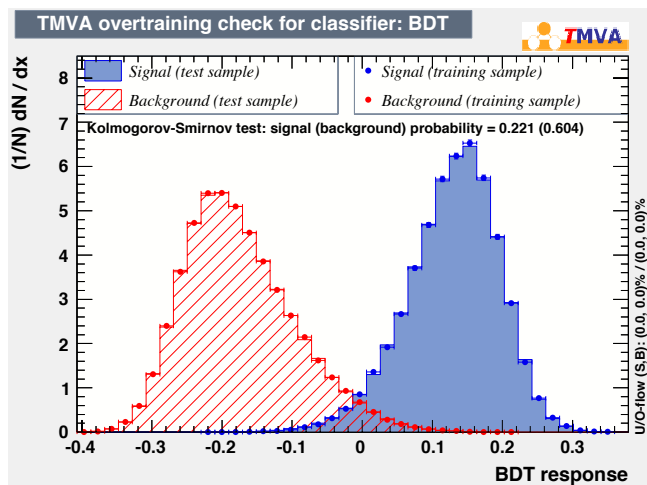


Figure 4.9 – Distributions of the BDT response on training and test samples.

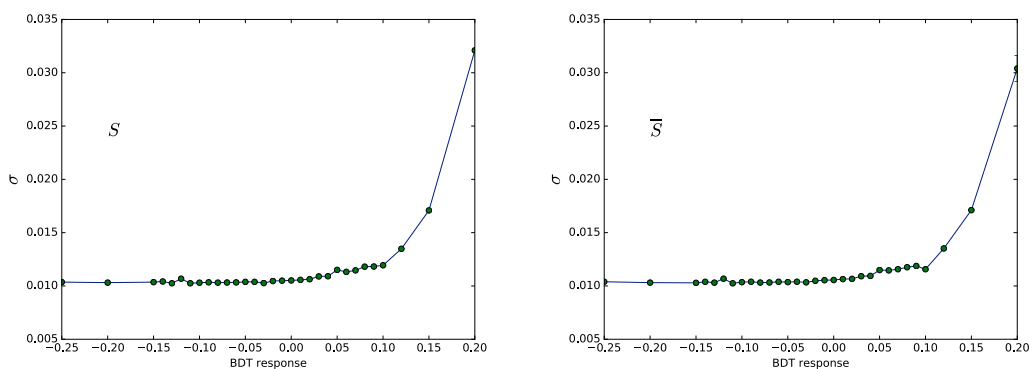


Figure 4.10 – Expected statistical uncertainty on of S_f (left) and $S_{\bar{f}}$ (right) as a function of the cut on the output of the BDT classifier, as obtained from simulated samples.

decays can be estimated to be roughly 1%. Then, the B^0 distribution after excluding the D^{\mp} signal window is plotted. To quantify the non-resonant $B^0 \rightarrow K^+\pi^-\pi^-\pi^+$ decays, the sum of an exponential and a Gaussian with a fixed shape is used to fit the resulting B^0 mass distribution, as shown in Fig. 4.12. As the fitted B^0 yield is 645 ± 242 , the non-resonant contribution is assumed to be negligible.

4.2 Simulation and expected sample composition

Simulated samples are used to (i) gain a detailed overview of all sources of b -hadron decays that contribute to the sample and (ii) model the relevant distributions studied in the analysis. Simulated data undergoes the same reconstruction and selection as real data.

Chapter 4. Selection of $B^0 \rightarrow D^\mp \pi^\pm$ decays

Table 4.6 – Signal efficiencies and background rejections of the different selection steps given with respect to the previous selection step. The preselection efficiency is computed w.r.t. the number of candidates passing trigger and stripping requirements. The last row shows the overall selection performance without the multiple candidate removal.

Selection step	ε_{sig}	$1 - \varepsilon_{\text{bkg}}$
preselection	$(93.61 \pm 0.06) \%$	$(85.20 \pm 0.02) \%$
Λ_c^\mp veto	$(93.48 \pm 0.06) \%$	$(9.85 \pm 0.03) \%$
semileptonic veto	$(98.96 \pm 0.03) \%$	$(7.66 \pm 0.03) \%$
wrongly associated PV veto	$(97.75 \pm 0.04) \%$	$(15.81 \pm 0.04) \%$
BDT selection	$(83.63 \pm 0.10) \%$	$(97.18 \pm 0.01) \%$
total	$(70.7 \pm 0.1) \%$	$(99.911 \pm 0.002) \%$

Table 4.7 – BDT performance for each magnet polarity (up and down) and year of data taking. The quoted efficiencies contain signal and background.

	2011, up	2011, down	2012, up	2012, down
# cand. before BDT	398 357	569 853	1 301 800	1 316 597
# cand. after BDT	210 844	285 137	601 345	609 880
$\varepsilon_{\text{sig+bkg}}$	$(50.67 \pm 0.08) \%$	$(50.04 \pm 0.07) \%$	$(46.19 \pm 0.04) \%$	$(46.32 \pm 0.04) \%$

Each sample is split into four subsamples according to magnet polarity (up or down) and year of data taking (2011 or 2012), in proportions similar to those present in real data.

The simulated samples used are listed in Table 4.8, together with the number of true signal events passing the final selection and the corresponding total efficiencies. The PID requirements on the bachelor pion are not applied in order to compute these efficiencies.

4.2.1 PIDK correction

The PIDK distributions in data and MC differ. To correct for that, the PIDK distributions in MC are resampled using the binned PIDK probability density functions of dedicated calibration samples. These calibration samples consist of kinematically clean $D^{*+} \rightarrow D^0(\rightarrow K^- \pi^+) \pi^+$ decays, for which no requirement on RICH information is applied in the reconstruction.

The need for this resampling is due to the fact that, if the same cut is applied on data and MC, the resulting distributions in other observables may differ if the PIDK distributions in data and MC are different. Moreover, a correct PIDK distribution in MC allows the proper estimation (on MC) of the efficiency or misidentification rate for a given PIDK cut, which is an essential ingredient in the fit to the B^0 invariant mass distribution (as

4.2. Simulation and expected sample composition

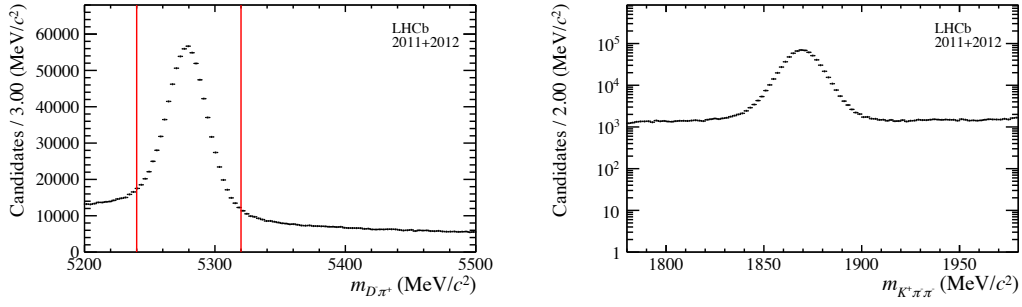


Figure 4.11 – Left: B^0 mass distribution with red vertical lines indicating the selected signal window. Right: resulting D^{\mp} mass distribution in the B^0 signal window.

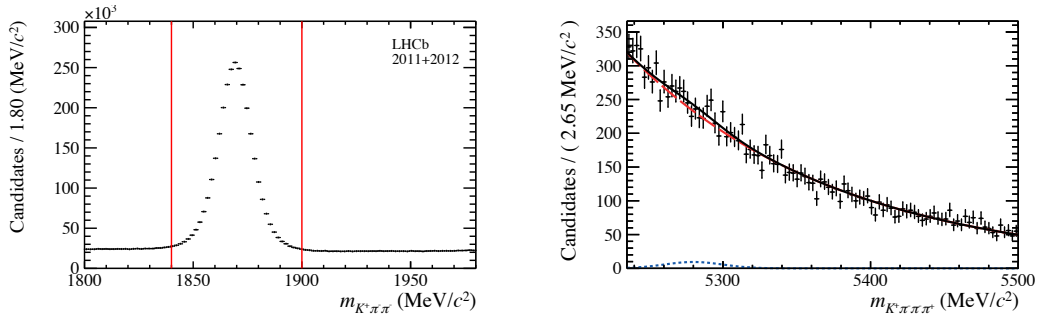


Figure 4.12 – Left: D^{\mp} mass distribution with red vertical lines indicating the excluded range. Right: B^0 mass distribution outside the D^{\mp} mass window with the fitting function overlaid.

described in Sec. 4.3).

The following strategy is adopted. A two-dimensional binning in momentum, p , and pseudorapidity, η , is defined. For each bin, the corresponding PIDK distribution in the calibration sample is built and for each event in the MC sample, a random PIDK value is sampled from the PIDK distribution associated with the corresponding bin in the calibration sample. More details are given in Appendix D.

Because of the Λ_c^{\mp} veto described in Sec. 4.1.3, the PID p variable for the D^{\mp} daughter particles is resampled as well in a similar manner using $\Lambda^0 \rightarrow p\pi^-$ decays as calibration channel.

The nominal binning used for the PID resampling is the following:

- momentum: 100 uniform bins between 2 GeV/ c and 200 GeV/ c , two equal bins between 200 GeV/ c and 300 GeV/ c ;
- pseudorapidity: one bin between 1.5 and 1.55, 69 uniform bins between 1.55 and 5.0.

Sample	Event type	N_{sel}	Efficiency [%]
$B^0 \rightarrow D^\mp \pi^\pm$	11164003	101096	1.966 ± 0.006
$B^0 \rightarrow D^- K^+$	11264011	19300	1.833 ± 0.013
$B^0 \rightarrow D^\mp \rho^\pm$	11164401	2408	0.1178 ± 0.0024
$B^0 \rightarrow D^{*\mp} \pi^\pm$	11164404	14901	0.721 ± 0.006
$B_s^0 \rightarrow D_s^- \pi^+$	13264021	7942	0.1531 ± 0.0017
$\Lambda_b^0 \rightarrow \Lambda_c^+ \pi^-$	15164001	325	0.0155 ± 0.0009
$B^0 \rightarrow D^- K^{*+}$	11164470	361	0.0358 ± 0.0019

Table 4.8 – Samples of simulated data used in the analysis, with the numbers of candidates N_{sel} after the final selection, and the selection efficiencies. Efficiencies include generator level, trigger, stripping, offline selection and tagging efficiencies. The $B^0 \rightarrow D^\mp \pi^\pm$ signal sample is generated with the parameters given in Appendix M.

In order to check the robustness of the method and evaluate the systematic uncertainties, two alternative binning schemes are defined:

- *narrow* binning: the number of bins in the uniform binning parts for p and η are increased to 140 and 80, respectively;
- *wide* binning: the number of bins in the uniform binning parts for p and η are decreased to 60 for both.

The $PIDK$ variable is resampled for the $B^0 \rightarrow D^\mp \pi^\pm$ and $B^0 \rightarrow D^\mp K^\pm$ Monte Carlo samples using these two alternative schemes as well. The result of this resampling is shown in Fig. 4.13 for $B^0 \rightarrow D^\mp \pi^\pm$ decays. The effect on all the other physics background is expected to be very small and thus neglected, since they are all located in different regions of the $D^\mp \pi^\pm$ invariant mass. More details on this will be given in Secs. 4.3 and 5.5.

4.2.2 Surviving physics backgrounds

Some physics background candidates that survive the selection chain described in the previous section are expected. In the pion sample, these are:

- $B^0 \rightarrow D^- K^+$: Peaking background due to the bachelor kaon being wrongly identified as a pion.
- $B^0 \rightarrow D^\mp \rho^\pm (\rightarrow \pi^\pm \pi^0)$: Low mass background due to a missing neutral pion in the reconstruction.
- $B^0 \rightarrow D^{*\mp} (\rightarrow D^\mp \gamma / \pi^0) \pi^\pm$: Low mass background due to a missing neutral particle in the reconstruction.

In the kaon sample, the following backgrounds are expected:

4.2. Simulation and expected sample composition

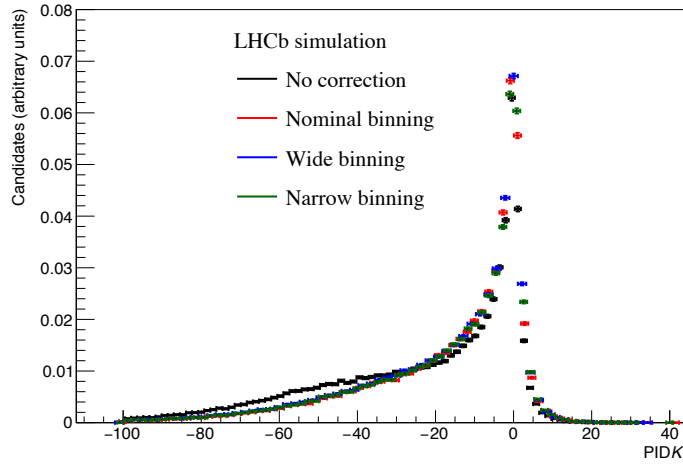


Figure 4.13 – PIDK distribution for simulated $B^0 \rightarrow D^\mp \pi^\pm$ decays without resampling (black), after the nominal resampling (red), after the resampling with the wide binning scheme (blue), and after the resampling with the narrow binning scheme (green).

- $B^0 \rightarrow D^\mp \pi^\pm$: Signal candidates having the bachelor pion wrongly identified as a kaon.
- $B^0 \rightarrow D^\mp \rho^\pm (\rightarrow \pi^\pm \pi^0)$: Low mass background where, in addition to the missing pion in the final state, a reconstructed pion is wrongly identified as a kaon.
- $B^0 \rightarrow D^- K^{*+} (\rightarrow \pi^0 K^+)$: Low mass background where the neutral pion is missing in the reconstruction.

The background fractions expected in the pion sample with respect to the $B^0 \rightarrow D^\mp \pi^\pm$ signal are reported in Table 4.9. These fractions are computed using the branching fractions of the expected decay as inputs and from the ratio of efficiencies estimated from MC and corrected as described in Sec. 4.2.1. Where relevant we consider also the ratio of the fragmentation probabilities of b quarks to different b hadrons, which are $34.0 \pm 2.1\%$, $21.8 \pm 4.7\%$ and $10.1 \pm 1.5\%$ for B^0 , B_s^0 and Λ_b^0 , respectively [74]. These expectations will be compared with the results from the fit to data described in the next section.

The $B_s^0 \rightarrow D_s^- \pi^+$ and $\Lambda_b^0 \rightarrow \Lambda_c^+ (\rightarrow K^- \pi^+ p) \pi^-$ backgrounds are suppressed to a negligible fraction by the offline selection described in Sec. 4.1, and are thus ignored in the description of the sample composition. Moreover, in the kaon sample, the $B^0 \rightarrow D^{*\mp} \pi^\pm$ and $B^0 \rightarrow D^{*-} K^+$ components, which are expected to be negligible, are ignored as well. More precisely, these components are taken into account by the PDF describing $B^0 \rightarrow D^- K^{*+} (\rightarrow \pi^0 K^+)$, since they are expected to sit in the same mass region.

Decay	\mathcal{B} from Ref. [32] [%]	ϵ_{bkg} [%]	f_{bkg} [%]
$B^0 \rightarrow D^- K^+$	0.00186 ± 0.00020	0.684 ± 0.008	2.61 ± 0.31
$B^0 \rightarrow D^\mp \rho^\pm (\rightarrow \pi^\pm \pi^0)$	0.071 ± 0.011	0.1149 ± 0.0024	16.7 ± 2.8
$B^0 \rightarrow D^\mp \pi^\pm (\rightarrow D^\mp \gamma / \pi^0) \pi^\pm$	0.0080 ± 0.0004	0.705 ± 0.006	11.6 ± 0.8
$\Lambda_b^0 \rightarrow \Lambda_c^+ (K^- \pi^+ p) \pi^-$	0.032 ± 0.004	0.0150 ± 0.0008	0.62 ± 0.24
$B_s^0 \rightarrow D_s^- \pi^+$	0.0164 ± 0.0014	0.1493 ± 0.0017	1.64 ± 0.32

Table 4.9 – Background contributions expected in the pion sample. Each fraction f_{bkg} is relative to the $B^0 \rightarrow D^\mp \pi^\pm$ yield. The $B^0 \rightarrow D^\mp \pi^\pm$ branching ratio and total selection efficiency in the pion sample are $(0.254 \pm 0.014)\%$ [32] and $(1.924 \pm 0.006)\%$, respectively.

4.3 Fits to the B^0 invariant mass

The *sPlot* technique [75] is applied in order to statistically isolate the signal contribution for the subsequent decay time fit. The $D^\mp \pi^\pm$ invariant mass, where the D^\mp mass is constrained to its known value in order to improve the mass resolution, is adopted as discriminating observable thanks to its small correlation with the B^0 decay time (see Appendix E).

In a first step, a binned extended maximum likelihood fit (“Fit A”) is performed in order to define the PDFs describing the signal and background components. The choice of a binned fit is justified by the very high statistics of the data sample. The invariant mass range of the fit is [5090, 6000] MeV/ c^2 . Only tagged candidates are considered, *i.e.* candidates with at least one nonzero tagging decision from the OS or SS taggers. The reason for this is that untagged candidates do not contribute to the sensitivity on the CP coefficients. The fit is performed simultaneously on the pion and kaon samples (see Sec. 4.1.2). This approach is adopted in order to control the contamination from the $B^0 \rightarrow D^- K^+$ background in the pion sample. The number of $B^0 \rightarrow DX$ candidates in the Y sample (with $X, Y = \pi, K$), $N_{B^0 \rightarrow DX}^Y$, can be defined via the following relations:

$$N_{B^0 \rightarrow D\pi}^K = \frac{\epsilon_{\text{PID}}(B^0 \rightarrow D\pi)_K}{\epsilon_{\text{PID}}(B^0 \rightarrow D\pi)_\pi} \times N_{B^0 \rightarrow D\pi}^\pi = \frac{1 - \epsilon_{\text{PID}}(B^0 \rightarrow D\pi)_\pi}{\epsilon_{\text{PID}}(B^0 \rightarrow D\pi)_\pi} \times N_{B^0 \rightarrow D\pi}^\pi, \quad (4.1)$$

$$N_{B^0 \rightarrow DK}^\pi = \frac{\epsilon_{\text{PID}}(B^0 \rightarrow DK)_\pi}{\epsilon_{\text{PID}}(B^0 \rightarrow DK)_K} \times N_{B^0 \rightarrow DK}^K = \frac{1 - \epsilon_{\text{PID}}(B^0 \rightarrow DK)_K}{\epsilon_{\text{PID}}(B^0 \rightarrow DK)_K} \times N_{B^0 \rightarrow DK}^K. \quad (4.2)$$

The quantities $\epsilon_{\text{PID}}(B^0 \rightarrow DX)_Y$ are the fractions of true $B^0 \rightarrow DX$ decays that are selected in the Y sample by applying the corresponding PIDK cut. These fractions (or efficiencies) are estimated on $B^0 \rightarrow D^\mp \pi^\pm$ and $B^0 \rightarrow D^- K^+$ MC samples where the

Table 4.10 – Fractions of true $B^0 \rightarrow D^\mp \pi^\pm$ and $B^0 \rightarrow D^- K^+$ decays that are selected in the π or K sample.

Decay	PIDK requirement	fraction
$B^0 \rightarrow D^\mp \pi^\pm$	< 5 (π sample)	$0.9790 \pm 0.0040(\text{stat}) \pm 0.0004(\text{syst})$
$B^0 \rightarrow D^\mp \pi^\pm$	> 5 (K sample)	$0.0211 \pm 0.0005(\text{stat}) \pm 0.0004(\text{syst})$
$B^0 \rightarrow D^- K^+$	< 5 (π sample)	$0.373 \pm 0.005(\text{stat}) \pm 0.008(\text{syst})$
$B^0 \rightarrow D^- K^+$	> 5 (K sample)	$0.627 \pm 0.007(\text{stat}) \pm 0.010(\text{syst})$

PIDK distributions are resampled from calibration data, as described in Sec. 4.2.1. A systematic uncertainty for these efficiencies is estimated by taking the largest discrepancy between the nominal value and the result obtained with the narrow and wide binning schemes introduced in Sec. 4.2.1. The results of these estimations are reported in Table 4.10.

Finally, an unbinned extended maximum likelihood fit (“Fit B”) is performed on data using the reduced mass interval $[5220, 5600]$ MeV/ c^2 in order to extract $sWeights$. In this second fit, all the parameters are fixed to the values found in Fit A, except for the signal and total background yields. The reduced mass window avoids diluting the $sWeights$ with background candidates having an invariant mass falling outside this window. This has the added advantage of reducing the dataset size used in the decay time fit.

4.3.1 Probability density functions

The PDFs used to describe both the pion and kaon sample components in Fit A are first estimated on MC samples. The parameters of the combinatorial background PDFs are instead determined directly from data. The PDFs used for the pion sample are:

- $B^0 \rightarrow D^\mp \pi^\pm$: sum of a double-sided Hypatia [76] and a Johnson SU [77] function (PDF $_{B^0 \rightarrow D\pi}^\pi$).
- $B^0 \rightarrow D^- K^+$: double-sided Hypatia function (PDF $_{B^0 \rightarrow DK}^\pi$).
- $B^0 \rightarrow D^\mp \rho^\pm$: Johnson SU function (PDF $_{B^0 \rightarrow D\rho}^\pi$).
- $B^0 \rightarrow D^{*\mp} \pi^\pm$: sum of a single-sided Crystal Ball function [78] and a Gaussian function (PDF $_{B^0 \rightarrow D^*\pi}^\pi$).
- Background: sum of an exponential function and a constant offset.

For the kaon sample they are:

- $B^0 \rightarrow D^\mp \pi^\pm$: double-sided Hypatia function (PDF $_{B^0 \rightarrow D\pi}^K$).
- $B^0 \rightarrow D^- K^+$: single-sided Hypatia function (PDF $_{B^0 \rightarrow DK}^K$).
- $B^0 \rightarrow D^\mp \rho^\pm$: double Gaussian function (PDF $_{B^0 \rightarrow D\rho}^K$).

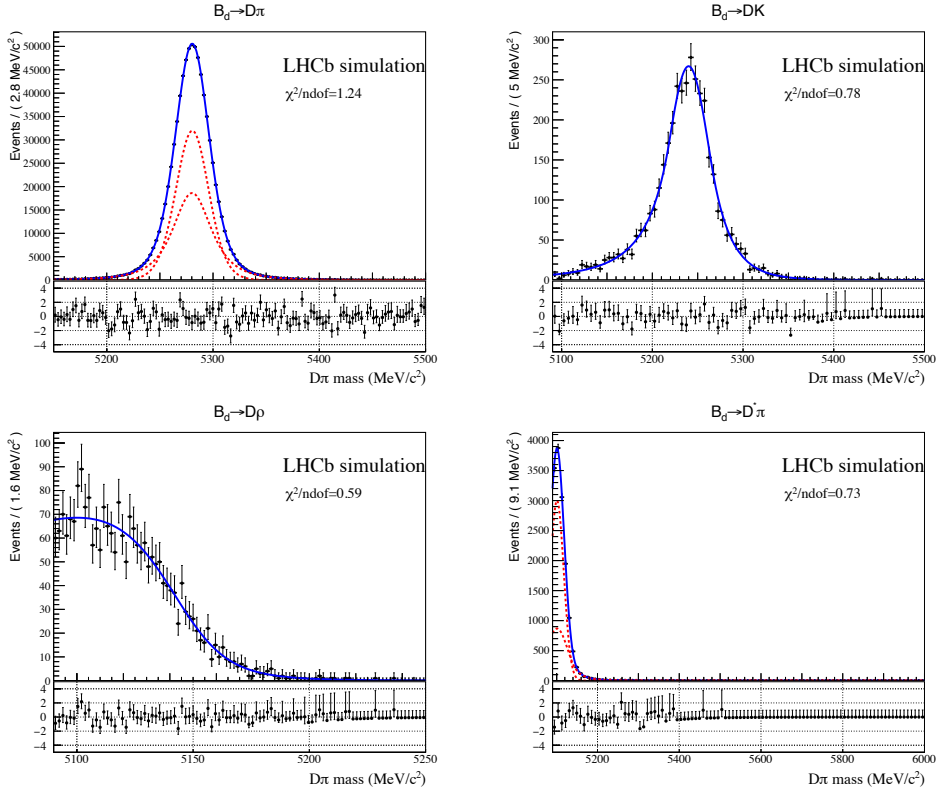


Figure 4.14 – $D^\mp \pi^\pm$ mass distributions of MC samples of $B^0 \rightarrow D^\mp \pi^\pm$ (top left), $B^0 \rightarrow D^- K^+$ (top right), $B^0 \rightarrow D^\mp \rho^\pm$ (bottom left) and $B^0 \rightarrow D^{*\mp} \pi^\pm$ (bottom right) events passing the $B^0 \rightarrow D^\mp \pi^\pm$ selection, with fits superimposed.

- $B^0 \rightarrow D^\mp K^{*\pm}$: Gaussian function ($\text{PDF}_{B^0 \rightarrow DK^*}^K$).
- Background: sum of an exponential function and a constant offset.

The definitions of all the PDFs listed above are reported in Appendix F. The fits to the MC samples are shown in Figs. 4.14 and 4.15. The parameters obtained from these fits that are then fixed in the data fits are listed in Tables 4.11 and 4.12.

4.3.2 Fit to data

In order to perform Fit A, two fitting functions are defined:

$$\begin{aligned}
 f_\pi(m) = & N_{B^0 \rightarrow D\pi}^\pi \text{PDF}_{B^0 \rightarrow D\pi}^\pi + N_{B^0 \rightarrow DK}^\pi \text{PDF}_{B^0 \rightarrow DK}^\pi \\
 & + N_{B^0 \rightarrow D^*\pi}^\pi \text{PDF}_{B^0 \rightarrow D^*\pi}^\pi + N_{B^0 \rightarrow D\rho}^\pi \text{PDF}_{B^0 \rightarrow D\rho}^\pi \\
 & + N_{\text{comb}}^\pi \text{PDF}_{\text{comb}}^\pi,
 \end{aligned} \tag{4.3}$$

4.3. Fits to the B^0 invariant mass

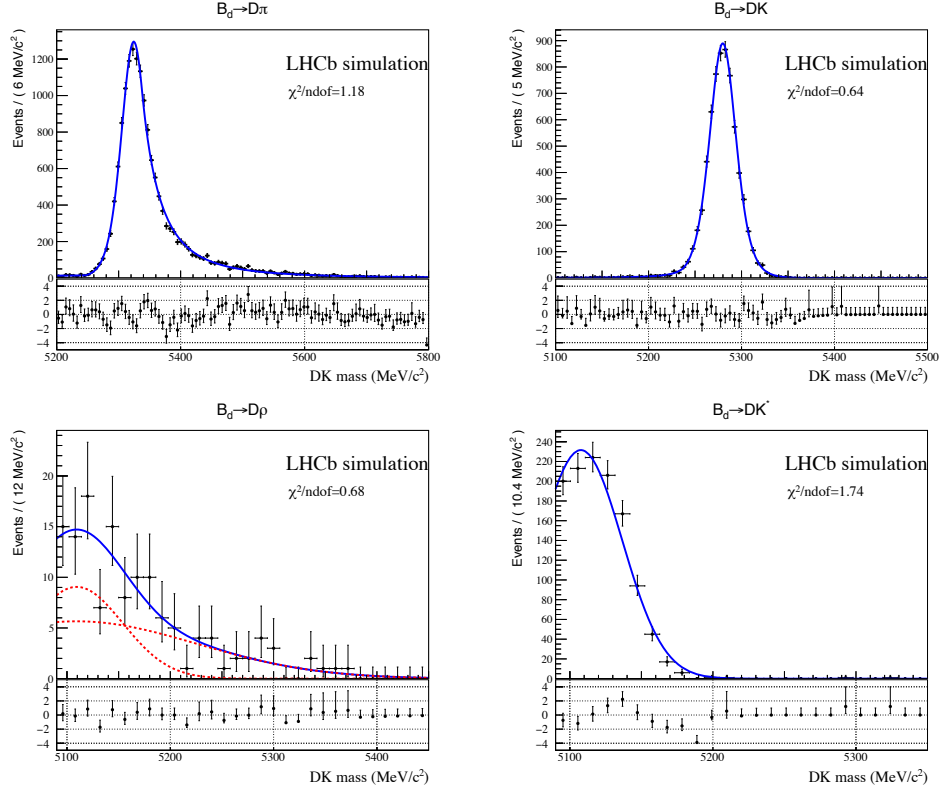


Figure 4.15 – $D^\mp K^\pm$ mass distributions of MC samples of $B^0 \rightarrow D^\mp \pi^\pm$ (top left), $B^0 \rightarrow D^- K^+$ (top right), $B^0 \rightarrow D^\mp \rho^\pm$ (bottom left) and $B^0 \rightarrow D^- K^{*+}$ (bottom right) events passing the $B^0 \rightarrow D^\mp \pi^\pm$ selection, with fits superimposed.

$$\begin{aligned}
 f_K(m) = & N_{B^0 \rightarrow DK}^K \text{PDF}_{B^0 \rightarrow DK}^K + N_{B^0 \rightarrow D\pi}^K \text{PDF}_{B^0 \rightarrow D\pi}^K \\
 & + N_{B^0 \rightarrow DK^*}^K \text{PDF}_{B^0 \rightarrow DK^*}^K + N_{B^0 \rightarrow D\rho}^K \text{PDF}_{B^0 \rightarrow D\rho}^K \\
 & + N_{\text{comb}}^K \text{PDF}_{\text{comb}}^K.
 \end{aligned} \tag{4.4}$$

Two extended likelihood functions are defined using data and PDFs related to both samples:

$$\mathcal{L}_x = \frac{e^{-N_{x,\text{exp}}} (N_{x,\text{exp}})^{N_{x,\text{obs}}}}{N_{x,\text{obs}}!} \prod_{i=1}^{N_{x,\text{obs}}} \frac{f_x(m_i)}{N_{x,\text{exp}}}, \quad x = \pi, K, \tag{4.5}$$

where $N_{\pi,\text{exp}} = N_{B^0 \rightarrow D\pi}^\pi + N_{B^0 \rightarrow DK}^\pi + N_{B^0 \rightarrow D^* \pi}^\pi + N_{B^0 \rightarrow D\rho}^\pi + N_{\text{comb}}^\pi$, $N_{K,\text{exp}} = N_{B^0 \rightarrow DK}^K + N_{B^0 \rightarrow D\pi}^K + N_{B^0 \rightarrow DK^*}^K + N_{B^0 \rightarrow D\rho}^K + N_{\text{comb}}^K$, and $N_{x,\text{obs}}$ is the number of observed candidates in the x sample. The product $\mathcal{L}_\pi \mathcal{L}_K$ is maximised during the fit.

The following strategy is adopted to perform Fit A:

- The mean and width parameters ($\mu_{B^0 \rightarrow D\pi}^\pi$, $\sigma_{H_{B^0 \rightarrow D\pi}}^\pi$, $\sigma_{J_{B^0 \rightarrow D\pi}}^\pi$, $\mu_{B^0 \rightarrow DK}^K$, $\sigma_{B^0 \rightarrow DK}^K$) of $\text{PDF}_{B^0 \rightarrow D\pi}^\pi$ and $\text{PDF}_{B^0 \rightarrow DK}^K$ are floated in the fit.
- The tail parameters ($a1_{B^0 \rightarrow D\pi}^\pi$, $a2_{B^0 \rightarrow D\pi}^\pi$, $n1_{B^0 \rightarrow D\pi}^\pi$, $n2_{B^0 \rightarrow D\pi}^\pi$) of $\text{PDF}_{B^0 \rightarrow D\pi}^\pi$ are constrained in the following way: $a1_{B^0 \rightarrow D\pi}^\pi$, $a2_{B^0 \rightarrow D\pi}^\pi$ are set to the values found on MC and both multiplied by a floating scale factor $sa_{B^0 \rightarrow D\pi}^\pi$; the same constraint is applied to $n1_{B^0 \rightarrow D\pi}^\pi$ and $n2_{B^0 \rightarrow D\pi}^\pi$, where the scale factor is labelled as $sn_{B^0 \rightarrow D\pi}^\pi$.
- The yield parameters $N_{B^0 \rightarrow D\pi}^K$ and $N_{B^0 \rightarrow DK}^\pi$ are constrained according to Eqs. 4.1 and 4.2. The efficiencies $\epsilon_{\text{PID}}(B^0 \rightarrow D\pi)_{D\pi}$ and $\epsilon_{\text{PID}}(B^0 \rightarrow DK)_{DK}$ are Gaussian-constrained independently in the fit, using the values reported in Table 4.10. The yield $N_{B^0 \rightarrow D\rho}^K$ is fixed to be 0.92 times the yield $N_{B^0 \rightarrow DK^*}^K$, the latter being floated in the fit. This is done according to the expected $B^0 \rightarrow D^\mp \pi^\pm$ to $B^0 \rightarrow D^\mp K^{\pm*}$ ratio in the kaon sample, which is 0.92 ± 0.21 . All the other yields appearing in Eqs. 4.3 and 4.4 are floated in the fit.
- The mean parameters ($\mu_{B^0 \rightarrow D^*\pi}^\pi$, $\mu_{B^0 \rightarrow D\rho}^{\pi/K}$) of $\text{PDF}_{B^0 \rightarrow D^*\pi}^\pi$, $\text{PDF}_{B^0 \rightarrow D\rho}^{\pi/K}$, are constrained to be shifted from $\mu_{B^0 \rightarrow D\pi}^\pi$ (in the π sample) and $\mu_{B^0 \rightarrow DK}^K$ (in the K sample) by the same amount found in MC. The shift of the component with respect to the $B^0 \rightarrow D^\mp \pi^\pm$ ($B^0 \rightarrow D^\mp K^\pm$) peak in the π (K) sample is denoted as $\Delta\mu_{\text{comb}}^{K/\pi}$. The mean parameters ($\mu_{B^0 \rightarrow DK^*}^K$, $\mu_{B^0 \rightarrow DK}^\pi$, $\mu_{B^0 \rightarrow D\pi}^K$) of $\text{PDF}_{B^0 \rightarrow DK^*}^K$, $\text{PDF}_{B^0 \rightarrow DK}^\pi$, $\text{PDF}_{B^0 \rightarrow D\pi}^K$ are floated in the fit.
- The exponent parameters ($c_{\text{comb}}^{\pi/K}$) and fractions ($f_{\text{comb}}^{\pi/K}$) of $\text{PDF}_{\text{comb}}^{\pi/K}$ are floated in the fit.

The projections of the fitted f_π and f_K in the $D^\mp \pi^\pm$ and $D^\mp K^\pm$ invariant mass observables (Fit A) are shown in Fig. 4.16, for the π and K data samples, respectively. A list of all the parameters fixed in Fit A is given in Tables 4.11 and 4.12. The fitted parameters (including yields and PID efficiencies) are listed in Table 4.13.

As cross-check, the fitted yields of $B^0 \rightarrow D^\mp K^\pm$, $B^0 \rightarrow D^\mp \rho^\pm$ and $B^0 \rightarrow D^{*\mp} \pi^\pm$ in the pion sample are compared with the expected yields, which are obtained, for each background, by multiplying the fitted $B^0 \rightarrow D^\mp \pi^\pm$ yield in the pion sample by the f_{bkg} fractions given in Table 4.9. These yields are reported in Table 4.14. There is a full agreement between expected and observed yields for the $B^0 \rightarrow D^\mp \rho^\pm$ and $B^0 \rightarrow D^{*\mp} \pi^\pm$ components, while the agreement for the $B^0 \rightarrow D^\mp K^\pm$ is at the level of 2.5σ .

4.3.3 *sWeight* calculation

After Fit A, all the floating shape parameters in $f_\pi(m)$ are fixed, all the background components in the π sample are combined into a single PDF, and the B^0 mass range is restricted to $[5220, 5600]$ MeV/ c^2 . Concretely, $f_\pi(m)$ is redefined as

$$f_\pi(m) = N_{B^0 \rightarrow D\pi}^\pi \text{PDF}_{B^0 \rightarrow D\pi}^\pi + N_{\text{bkg}}^\pi \text{PDF}_{\text{bkg}}^\pi. \quad (4.6)$$

4.3. Fits to the B^0 invariant mass

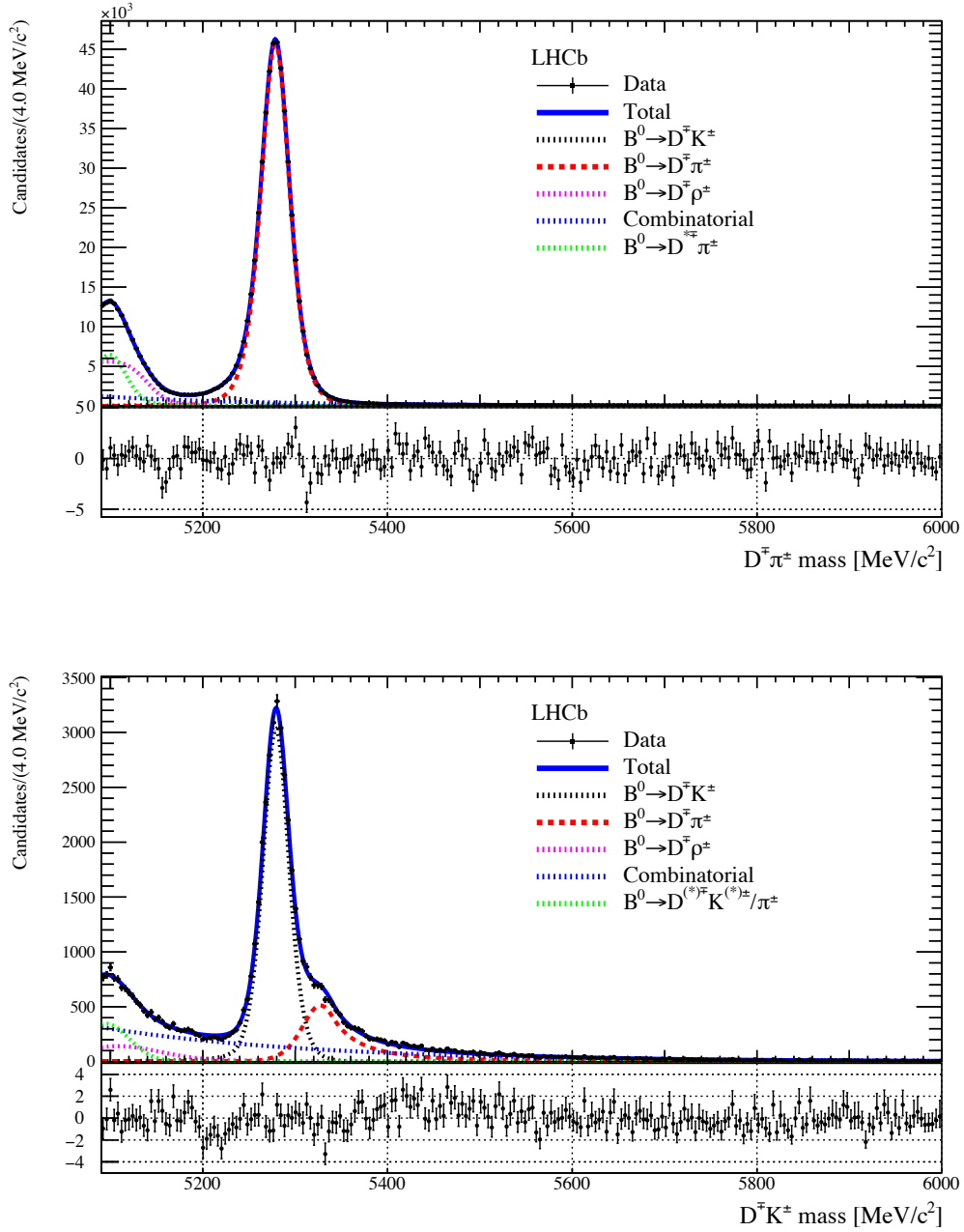


Figure 4.16 – Top: $D^{\mp}\pi^{\pm}$ mass distribution of the π sample. Bottom: $D^{\mp}K^{\pm}$ mass distribution of the K sample. The result of the simultaneous fit (Fit A) to both samples is superimposed. The plot below each histogram shows the normalised fit residuals (data minus fit divided by fit error).

Chapter 4. Selection of $B^0 \rightarrow D^\mp \pi^\pm$ decays

Table 4.11 – Parameters of $f_\pi(m)$ fixed or constrained in Fit A. All non-zero values are obtained from the fits to MC samples described in Sec. 4.3.1.

Parameter	Value	Status in Fit A
$a1_{B^0 \rightarrow D\pi}^\pi$	0.722 ± 0.091	constrained
$a2_{B^0 \rightarrow D\pi}^\pi$	0.96 ± 0.12	constrained
$n1_{B^0 \rightarrow D\pi}^\pi$	5.92 ± 0.92	constrained
$n2_{B^0 \rightarrow D\pi}^\pi$	5.83 ± 0.38	constrained
$\beta_{B^0 \rightarrow D\pi}^\pi$	0.0	fixed
$\lambda_{B^0 \rightarrow D\pi}^\pi$	-1.240 ± 0.060	fixed
$\zeta_{B^0 \rightarrow D\pi}^\pi$	0.0	fixed
$f_{B^0 \rightarrow D\pi}^\pi$	0.436 ± 0.060	fixed
$\sigma_{B^0 \rightarrow DK}^\pi$	$23.43 \pm 0.42 \text{ MeV}/c^2$	fixed
$a1_{B^0 \rightarrow DK}^\pi$	0.898 ± 0.025	fixed
$a2_{B^0 \rightarrow DK}^\pi$	1.092 ± 0.033	fixed
$n1_{B^0 \rightarrow DK}^\pi$	3.83 ± 0.40	fixed
$n2_{B^0 \rightarrow DK}^\pi$	22.0 ± 7.6	fixed
$\beta_{B^0 \rightarrow DK}^\pi$	0.0	fixed
$\lambda_{B^0 \rightarrow DK}^\pi$	-24 ± 10	fixed
$\zeta_{B^0 \rightarrow DK}^\pi$	0.0	fixed
$\nu_{B^0 \rightarrow D\rho}^\pi$	-2.01 ± 0.15	fixed
$\mu_{B^0 \rightarrow D\rho}^\pi$	$4828 \pm 80 \text{ MeV}/c^2$	constrained into $\Delta\mu_{B^0 \rightarrow D\rho}^\pi$
$\sigma_{B^0 \rightarrow D\rho}^\pi$	$550 \pm 190 \text{ MeV}/c^2$	fixed
$\tau_{B^0 \rightarrow D\rho}^\pi$	1.163 ± 0.090	fixed
$\alpha_{B^0 \rightarrow D^*\pi}^\pi$	-1.443 ± 0.031	fixed
$n_{B^0 \rightarrow D^*\pi}^\pi$	4.65 ± 0.30	fixed
$\mu_{B^0 \rightarrow D^*\pi}^\pi$	$5100.93 \pm 0.23 \text{ MeV}/c^2$	constrained into $\Delta\mu_{B^0 \rightarrow D^*\pi}^\pi$
$\sigma G_{B^0 \rightarrow D^*\pi}^\pi$	$16.52 \pm 0.20 \text{ MeV}/c^2$	fixed
$\sigma C_{B^0 \rightarrow D^*\pi}^\pi$	$25.84 \pm 0.48 \text{ MeV}/c^2$	fixed
$f_{B^0 \rightarrow D^*\pi}^\pi$	0.302 ± 0.011	fixed

The N_{bkg}^π parameter describes the total number of background events in the new range. The $\text{PDF}_{\text{bkg}}^\pi$ function is defined as:

$$\begin{aligned} \text{PDF}_{\text{bkg}}^\pi = & f_{\text{comb}}^\pi \text{PDF}_{\text{comb}}^\pi + f_{B^0 \rightarrow DK}^\pi \text{PDF}_{B^0 \rightarrow DK}^\pi + f_{B^0 \rightarrow D\rho}^\pi \text{PDF}_{B^0 \rightarrow D\rho}^\pi \\ & + (1 - f_{\text{comb}}^\pi - f_{B^0 \rightarrow DK}^\pi - f_{B^0 \rightarrow D\rho}^\pi) \text{PDF}_{B^0 \rightarrow D^*\pi}^\pi. \end{aligned} \quad (4.7)$$

For each background component in the π sample, the fraction f_j^π is determined by the following expression:

$$f_j^\pi = \frac{N_j^\pi \int_{5220 \text{ MeV}/c^2}^{5600 \text{ MeV}/c^2} \text{PDF}_j^\pi dm}{\sum_i N_i^\pi \int_{5220 \text{ MeV}/c^2}^{5600 \text{ MeV}/c^2} \text{PDF}_i^\pi dm}, \quad (4.8)$$

4.3. Fits to the B^0 invariant mass

Table 4.12 – Parameters of $f_K(m)$ fixed or constrained in Fit A. All non-zero values are obtained from the fits to MC samples described in Sec. 4.3.1.

Parameter	Value	Status in Fit A
$\sigma_{B^0 \rightarrow D\pi}^K$	$23.97 \pm 0.46 \text{ MeV}/c^2$	fixed
$a1_{B^0 \rightarrow D\pi}^K$	3.14 ± 0.14	fixed
$a2_{B^0 \rightarrow D\pi}^K$	0.569 ± 0.039	fixed
$n1_{B^0 \rightarrow D\pi}^K$	0.05 ± 0.11	fixed
$n2_{B^0 \rightarrow D\pi}^K$	2.81 ± 0.12	fixed
$\beta_{B^0 \rightarrow D\pi}^K$	0.0	fixed
$\lambda_{B^0 \rightarrow D\pi}^K$	-3.77 ± 0.57	fixed
$\zeta_{B^0 \rightarrow D\pi}^K$	0.0	fixed
$\sigma_{B^0 \rightarrow DK}^K$	$17.32 \pm 0.26 \text{ MeV}/c^2$	fixed
$a_{B^0 \rightarrow DK}^K$	2.34 ± 0.19	fixed
$n_{B^0 \rightarrow DK}^K$	1.56 ± 0.33	fixed
$\beta_{B^0 \rightarrow DK}^K$	0.0	fixed
$\lambda_{B^0 \rightarrow DK}^K$	-3.45 ± 0.34	fixed
$\zeta_{B^0 \rightarrow DK}^K$	0.0	fixed
$f_{B^0 \rightarrow D\rho}^K$	0.58 ± 0.17	fixed
$\mu_{B^0 \rightarrow D\rho}^K$	$5\,109 \pm 24 \text{ MeV}/c^2$	constrained into $\Delta\mu_{B^0 \rightarrow D\rho}^K$
$\sigma1_{B^0 \rightarrow D\rho}^K$	$117 \pm 18 \text{ MeV}/c^2$	fixed
$\sigma2_{B^0 \rightarrow D\rho}^K$	$45 \pm 16 \text{ MeV}/c^2$	fixed

where the indices i and j run over the combinatorial, $B^0 \rightarrow D^- K^+$ and $B^0 \rightarrow D^\mp \rho^\pm$ background components in the π sample, and N_j^π is the number of events of component j in the old mass range.

An unbinned extended maximum likelihood fit (Fit B) is then performed to the π sample only. The only floating parameters are the yields $N_{B^0 \rightarrow D\pi}^\pi$ and N_{bkg}^π . The result of the fit is reported in Table 4.15. The yields of the signal component and the total background as obtained in Fit B are compatible with the yields of the signal and the sum of the yields of each background as obtained in Fit A after integrating the PDFs in the Fit B mass range.

Fit B is used as starting point to apply the *sPlot* technique and extract *sWeights* used to subtract the total background component from the π sample. The projection of the fitted $f_\pi(m)$ in Fit B and a comparison between the weighted and unweighted datasets projected over the B^0 decay time and D^\mp invariant mass observables are reported in Fig. 4.17.

Table 4.13 – Results of Fit A.

Parameter	Fitted value
$\mu_{B^0 \rightarrow DK}^\pi$	$5\,228.62 \pm 0.92$
$\sigma_{B^0 \rightarrow DK}^K$	17.17 ± 0.15
$\mu_{B^0 \rightarrow DK^*}^K$	$5\,094.8 \pm 3.9$
$\sigma_{B^0 \rightarrow DK^*}^K$	25.5 ± 2.6
$c_{\text{comb}}^{1\pi}$	-0.00576 ± 0.00017
$c_{\text{comb}}^{2\pi}$	-0.0010 ± 0.0010
f_{comb}^π	0.899 ± 0.025
c_{comb}^K	-0.004397 ± 0.000066
$\mu_{B^0 \rightarrow DK}^K$	$5\,279.19 \pm 0.14$
$\mu_{B^0 \rightarrow D\pi}^\pi$	$5\,278.360 \pm 0.032$
$sa_{B^0 \rightarrow D\pi}^\pi$	0.684 ± 0.022
$sn_{B^0 \rightarrow D\pi}^\pi$	2.71 ± 0.80
$\sigma_{H_{B^0 \rightarrow D\pi}^\pi}$	37.69 ± 0.69
$\sigma_{J_{B^0 \rightarrow D\pi}^\pi}$	17.01 ± 0.17
$\mu_{B^0 \rightarrow D\pi}^K$	$5\,327.32 \pm 0.78$
$\epsilon_{\text{PID}}(B^0 \rightarrow DK)_K$	0.6197 ± 0.0079
$\epsilon_{\text{PID}}(B^0 \rightarrow D\pi)_\pi$	0.98048 ± 0.00041
$N_{B^0 \rightarrow DK}^K$	$28\,820 \pm 242$
$N_{B^0 \rightarrow DK^*}^K$	$3\,164 \pm 110$
$N_{B^0 \rightarrow D\rho}^\pi$	$73\,766 \pm 1239$
$N_{B^0 \rightarrow D^*\pi}^\pi$	$52\,494 \pm 819$
N_{comb}^K	$17\,469 \pm 341$
N_{comb}^π	$56\,230 \pm 1336$
$N_{B^0 \rightarrow D\pi}^\pi$	$483\,398 \pm 1040$

Table 4.14 – Expected and fitted yields for the physical background components in the pion sample.

Decay	Expected yield [10^4]	Fitted yield [10^4]
$B^0 \rightarrow D^- K^+$	1.26 ± 0.15	1.65 ± 0.05
$B^0 \rightarrow D^\mp \rho^\pm$	8.1 ± 1.4	7.38 ± 0.12
$B^0 \rightarrow D^{*\mp} \pi^\pm$	5.6 ± 0.4	5.25 ± 0.08

Table 4.15 – Results of Fit B (second column) and yields calculated by integrating the PDFs fitted in Fit A in the mass range used for Fit B (third column).

Parameter	Fitted value (Fit B)	Fitted value (Fit A)
$N_{B^0 \rightarrow D\pi}^\pi$	$479\,045 \pm 732$	$483\,398 \pm 1040$
N_{bkg}^π	$34\,381 \pm 300$	$34\,615 \pm 664$

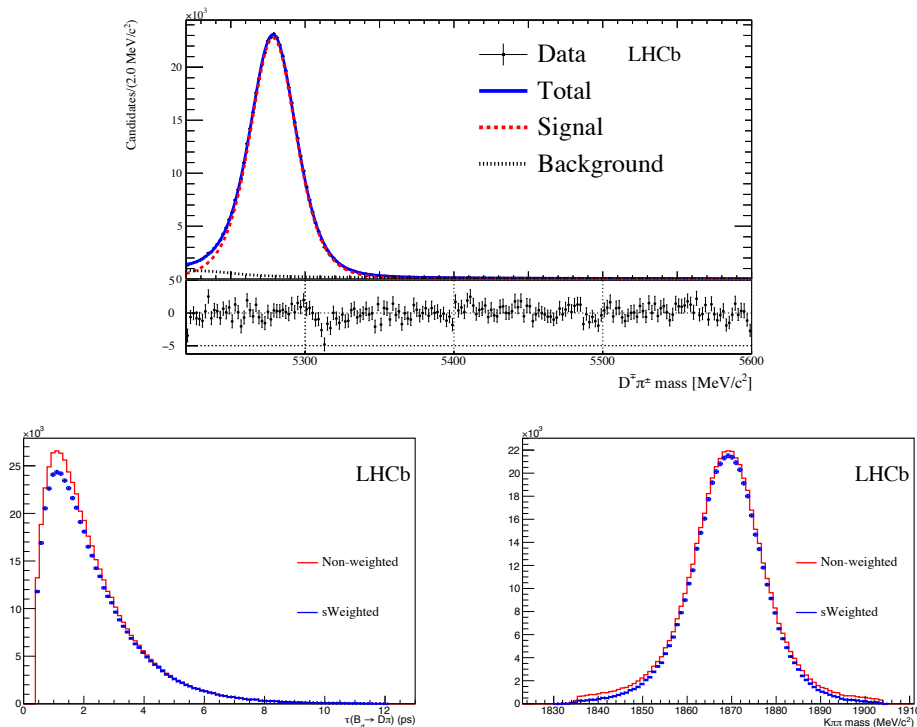


Figure 4.17 – Top: $D^{\mp}\pi^{\pm}$ mass distribution of the π sample with the results of Fit B superimposed. The plot below the histogram shows the normalised fit residuals (data minus fit divided by fit error). Bottom: B^0 decay time (left) and $K^{\pm}\pi^{\mp}\pi^{\pm}$ mass (right) distributions of the π sample in the B^0 mass region [5220, 5600] MeV/c^2 , with (blue) and without (red) $sWeights$ from Fit B.

4.3.4 Fits of subsamples

In order to validate the data sample, selection and fit procedure, Fit A is repeated in smaller subsamples. These subsamples are divided per year of data taking (2011, 2012), magnet polarity (up, down) and final state ($D^-\pi^+$, $D^+\pi^-$). In order to cope with the reduced statistics in the 2011 subsample, the combinatorial background PDF of the K sample ($\text{PDF}_{\text{comb}}^K$ is taken as a simple exponential (instead of an exponential plus a constant function)). The projections of the fitted PDFs describing the pion and kaon samples for each data subsample are shown in Figs. 4.18 and 4.19, respectively.

Fit B strategy is also repeated exactly as before for each subsample. The corresponding signal and background fitted yields are listed in Table 4.16. The sum of the signal yields for each subsample is compatible with the signal yield from the fit of total sample (reported in Table 4.15), which is $(4.790 \pm 0.007) \times 10^5$. The asymmetry between the yields of the $D^-\pi^+$ and $D^+\pi^-$ samples is 0.0100 ± 0.0015 , which is in agreement with the detection asymmetry between π^+ and π^- obtained in this analysis and by previous

Chapter 4. Selection of $B^0 \rightarrow D^\mp \pi^\pm$ decays

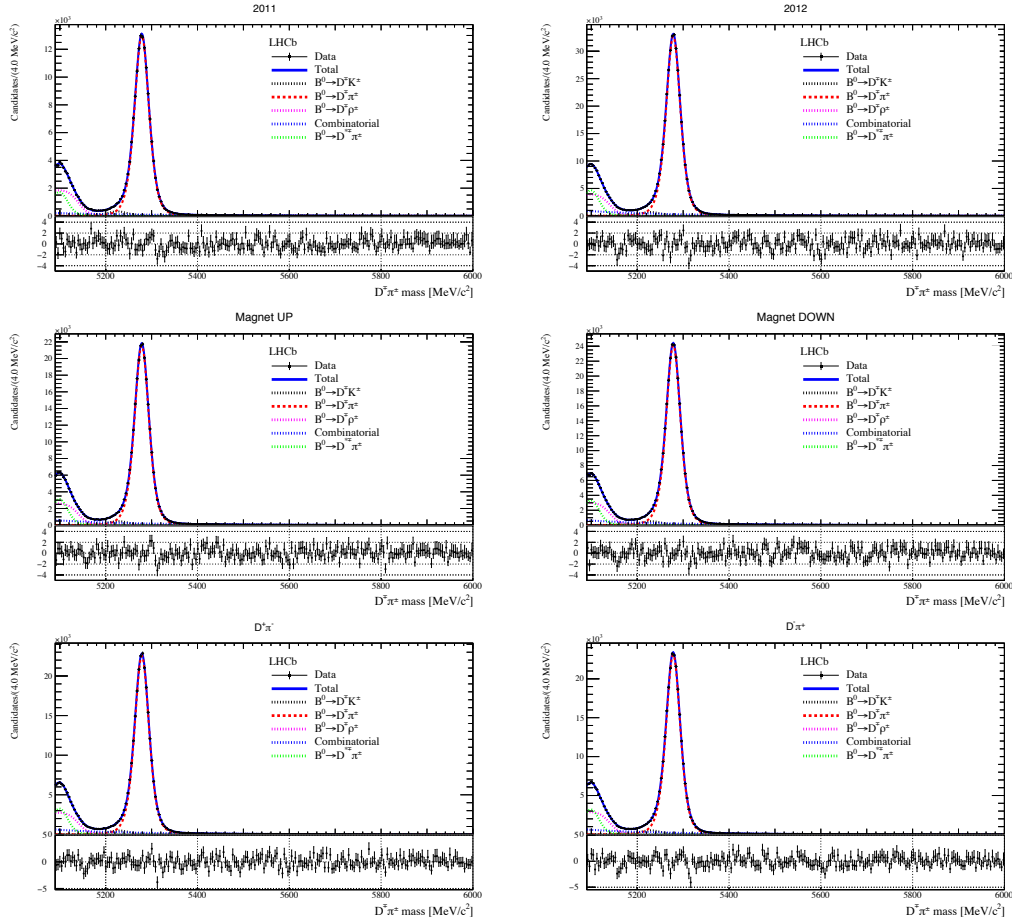


Figure 4.18 – $D^\mp \pi^\pm$ mass distributions of the pion sample for each data subsample, with the result of Fit A superimposed. The plots below the histograms show the normalised fit residuals.

measurements (more details given in Sec. 5.3). The ratio between the fitted yields on the 2011 and 2012 samples is compatible with the different collected luminosities and data taking conditions between the two years (twice as much luminosity is collected in 2012 compared to 2011, and the b -production cross-section is increased by a factor $\sim 8/7$ in 2012 because of the increase of the centre-of-mass energy). Moreover, the ratio of the yields obtained with the magnet up and down samples is in agreement with the ratio of luminosities collected with the two magnet polarities (same luminosity in 2012, +30 % more magnet down data in 2011).

4.3. Fits to the B^0 invariant mass

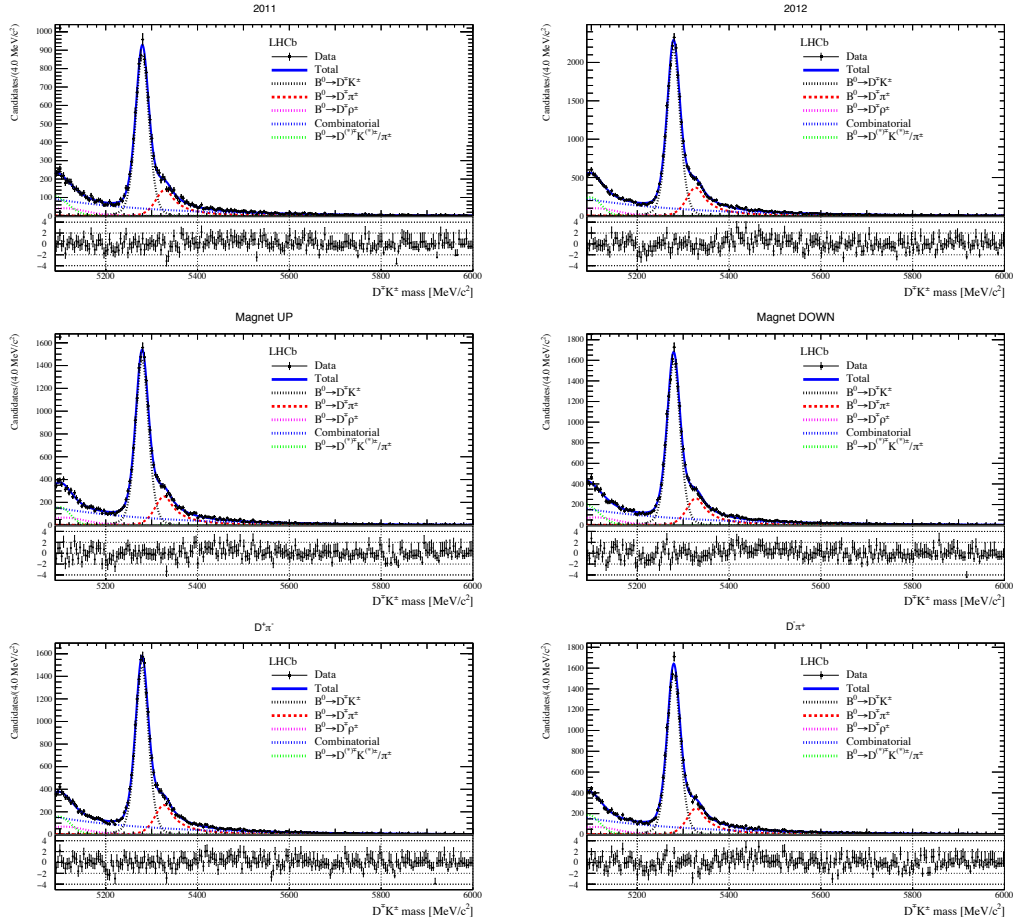


Figure 4.19 – $D^{\mp}K^{\pm}$ mass distributions of the pion sample for each data subsample, with the result of Fit A superimposed. The plots below the histograms show the normalised fit residuals.

Table 4.16 – Signal yields (in units of 10^5) in the pion sample for each subsample, obtained from Fit B.

2011	2012	Sum
1.383 ± 0.004	3.424 ± 0.006	4.807 ± 0.007
Magnet Up	Magnet Down	Sum
2.263 ± 0.005	2.523 ± 0.005	4.786 ± 0.007
$D^-\pi^+$	$D^+\pi^-$	Sum
2.421 ± 0.005	2.373 ± 0.005	4.794 ± 0.007

5 Measurement of CP violation in
 $B^0 \rightarrow D^{\mp} \pi^{\pm}$ decays

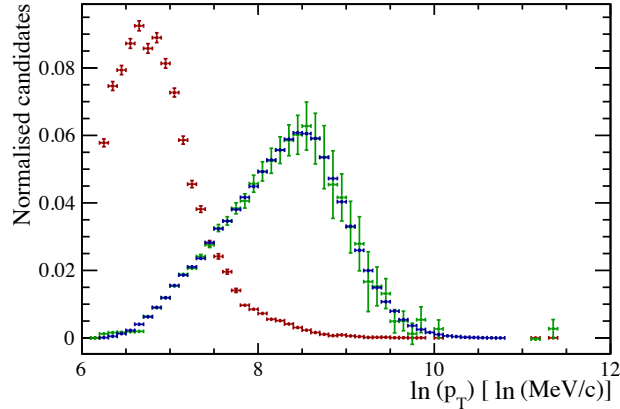


Figure 5.1 – Normalised $\ln(p_T)$ distributions of $sWeighted B^0 \rightarrow D^\mp \pi^\pm$ signal (blue), prompt $D^\mp + track$ before reweighting (red) and after reweighting (green).

5.1 Decay-time resolution

The decay-time resolution is determined from a sample of fake B^0 candidates formed from a prompt D^\mp candidate combined with a track originating from the PV. This sample is referred to as “ $D^\mp + track$ ”. The candidates are taken from the B02DKLTUBD2HHH stripping line. They are subjected to the same offline selection as that of the signal sample without a BDT cut and with two additional requirements: a single reconstructed PV per event is required in order to reduce wrong PV associations, and the D^\mp IP χ^2 with respect to the PV is required to be less than 9 to reduce the D^\mp contribution from B^0 decays. The combined stripping and offline selection yields 51 053 candidates. True $D^\mp + track$ candidates are unfolded from combinatorial background and nonresonant decays by means of $sWeights$ computed via a fit to the $K^\pm \pi^\mp \pi^\mp$ invariant mass distribution.

5.1.1 Companion track momentum reweighting

The decay time resolution is found to be dependent upon the companion track p_T which is considerably lower on average for the $D^\mp + track$ candidates than it is for genuine $B^0 \rightarrow D^\mp \pi^\pm$ signal. This is corrected for by reweighting the prompt sample to have the same $\ln(p_T)$ distribution as that of the $B^0 \rightarrow D^\mp \pi^\pm$ signal. The logarithm is taken to compress the high- p_T tails, and a binning scheme is chosen to have an equal number of signal events per bin. The $\ln(p_T)$ spectra for signal, prompt, and reweighted prompt candidates is shown in Fig. 5.1. By following the steps described in Sec. 5.1.2, prior to reweighting the average proper time resolution is determined to be ~ 71 fs, and after reweighting the resolution is found to be consistent with the value of 54 fs that was obtained in other B meson time-dependent analyses [79].

5.1.2 Resolution determination from decay-time error parameterisation

In order to study potential second order corrections to the decay-time error distribution, fits to the decay-time distribution of the D^\mp +track sample in bins of the per-event decay time error are performed. The decay-time error is obtained from DTF by propagating the uncertainty on the B^0 four-momentum. The binning scheme is chosen such that the sum of the *sWeights* associated to D^\mp +track candidates in each bin is equal. The fit is similar to that used to determine the resolution in Ref. [80], consisting of three components: a delta function convolved with a Gaussian resolution function accounts for the genuine prompt D^\pm +track component; a pair of exponential functions convolved with the same Gaussian function accounts for signal candidates coming from b -hadron decays, and a Gaussian function with a large width accounts for wrong-PV associated backgrounds. The time constant of the exponentials and the mean of the wrong-PV component are fixed from a global fit to the sample, while the mean and width of the resolution, the width of the wrong-PV component and the relative fractions of the prompt, wrong-PV and from- b components are all free parameters in the fits to each decay-time error bin. A likelihood fit is performed in 20 bins of the decay-time error from which the measured resolution $\langle\sigma\rangle_i$ is determined. The results of these fits are presented in Table 5.1, and a representative fit is shown in Fig. 5.2. A χ^2 fit is performed to the obtained values of the per-bin average error and resolution of the form:

$$\langle\sigma\rangle_i = \langle\sigma\rangle + p_1 (\langle\delta\rangle_i - \langle\delta\rangle) + p_2 (\langle\delta\rangle_i - \langle\delta\rangle)^2 \quad (5.1)$$

where $\langle\delta\rangle$ is the average per-event decay time error of the whole (unbinned) sample, while $\langle\delta\rangle_i$ is the average per-event decay time error in each bin. In the prompt D^\mp +track sample $\langle\delta\rangle$ is determined to be 0.0307 ± 0.0097 ps, in good agreement with the *sWeighted* $B^0 \rightarrow D^\mp \pi^\pm$ signal sample value of 0.034 ± 0.011 ps. The fit determines the average resolution, $\langle\sigma\rangle$, in addition to the trend. This fit is shown in Fig. 5.2, the result of which is presented in Table 5.2. The global average resolution is determined from this fit to be $\langle\sigma\rangle = 0.05491 \pm 0.00038$ ps. The procedure is found to be stable and yields compatible results with fits to 10 bins (0.05523 ± 0.00041 ps) and 30 bins (0.05464 ± 0.00037 ps).

This method, which accounts for second-order corrections to the decay time error, is used to define the width of a single Gaussian in the decay time fit to data, $\mathcal{R}(t-t') = G(t-t', \langle\sigma\rangle)$, with $\langle\sigma\rangle = 0.05491 \pm 0.00038$ ps. The uncertainty quoted here is statistical. Systematic uncertainties will be considered in Sec. 5.5.2.

5.2 Time-dependent selection efficiency

Because of some of the selection criteria described in Sec. 4.1, the B^0 decay time distribution is biased, *i.e.* different from the shape it would have with a constant selection

Chapter 5. Measurement of CP violation in $B^0 \rightarrow D^\mp \pi^\pm$ decays

Table 5.1 – Measured resolution $\langle\sigma\rangle_i$ obtained from a fit to the p_T corrected $sPlot$ of the decay-time distribution in bins of per-event decay-time error, δ , for prompt D^\mp +track signal. The average per-event decay time error $\langle\delta\rangle_i$ in each bin is also reported.

Bin i	lower edge	$\langle\delta\rangle_i$	$\langle\sigma\rangle_i$
0	0.01	0.0142 ± 0.0016	0.01731 ± 0.00053
1	0.0165376	0.01801 ± 0.00075	0.02439 ± 0.00089
2	0.0192247	0.02038 ± 0.00063	0.0286 ± 0.0011
3	0.0214493	0.02248 ± 0.00052	0.0347 ± 0.0011
4	0.0232264	0.02388 ± 0.00036	0.0384 ± 0.0013
5	0.0245968	0.02528 ± 0.00033	0.0422 ± 0.0014
6	0.0257605	0.02641 ± 0.00034	0.0449 ± 0.0014
7	0.0269093	0.02753 ± 0.00033	0.0489 ± 0.0015
8	0.0280345	0.02857 ± 0.00028	0.0489 ± 0.0015
9	0.0290414	0.02955 ± 0.00030	0.0525 ± 0.0018
10	0.0301189	0.03054 ± 0.00024	0.0552 ± 0.0019
11	0.0309259	0.03138 ± 0.00027	0.0582 ± 0.0017
12	0.0318409	0.03229 ± 0.00032	0.0594 ± 0.0016
13	0.0328907	0.03347 ± 0.00036	0.0641 ± 0.0015
14	0.0341106	0.03482 ± 0.00039	0.0643 ± 0.0014
15	0.0354999	0.03638 ± 0.00052	0.0658 ± 0.0014
16	0.0372226	0.03830 ± 0.00063	0.0719 ± 0.0012
17	0.0395386	0.04096 ± 0.00086	0.0736 ± 0.0012
18	0.0424521	0.0447 ± 0.0014	0.0786 ± 0.0011
19	0.0473915	0.0561 ± 0.0095	0.0933 ± 0.0010

Table 5.2 – Average per-event decay-time error $\langle\delta\rangle$, and resolution parameters p_1 , p_2 and $\langle\sigma\rangle$ obtained from a fit to the per-bin decay time error.

Parameter	Result
$\langle\delta\rangle$	0.0307 ± 0.0097 ps
p_1	2.031 ± 0.022
p_2	-19.30 ± 1.6 ps ⁻¹
$\langle\sigma\rangle$	0.05491 ± 0.00038 ps

5.2. Time-dependent selection efficiency

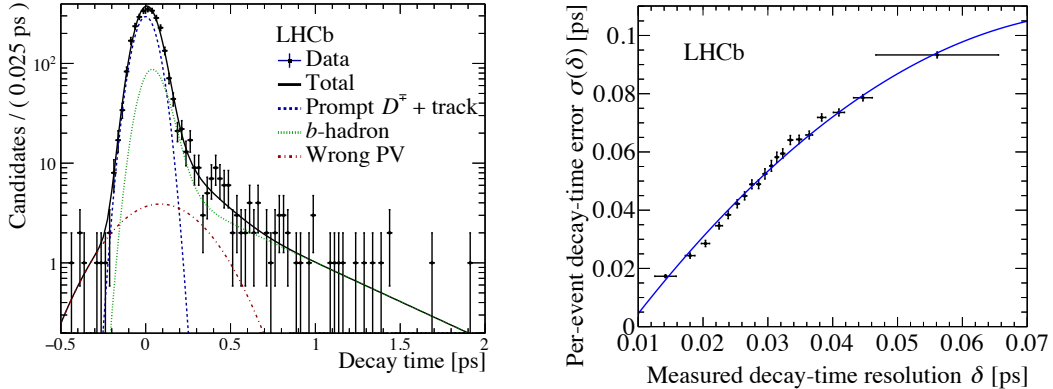


Figure 5.2 – Left: p_T -corrected and background-subtracted decay-time distribution of the $D^+ + \text{track}$ sample for the 15th bin ($[0.0341, 0.0355]$ ps) in per-event decay time error. The fit result is overlaid as the black solid curve: the wrong-PV, from- b , and prompt components are shown as the red dot-dashed, green dotted, and blue dashed curves, respectively. The numerical results are presented in Table 5.1. Right: measured resolution as a function of the average per-event decay time error determined from fits to the decay time in bins of decay time error. The horizontal bars are the standard deviation of the average per-event decay time error in each bin. The overlaid fit is described in the text.

efficiency. This efficiency, called here and after “acceptance”, is a function of the reconstructed proper time. In particular, it goes very rapidly to zero at low decay times due to the impact parameter requirements which exclude short-lived B^0 candidates; then, it reaches a “plateau” at intermediate decay times; finally, it drops at high decay times due to the acceptance of the VELO reconstruction.

The acceptance function $a(t)$ is parameterised using splines defined analytically as described in Ref. [81]. These splines are cubic polynomials defined in sub-ranges of the decay time. The boundaries of each sub-range, called “knots”, are located at 0.4, 0.5, 1.0, 1.5, 2.0, 2.3, 2.6, 3.0, 4.0, 10.0, and 12.0 ps. The location of the 11 knots and the higher density of knots at low decay times, where the acceptance is a strongly-varying function of t , ensure that the resulting acceptance is sufficiently smooth. For each knot t_i , a coefficient v_i is defined, which is the actual value of the acceptance function $a(t_i)$. In order to fix the overall scale of the acceptance function, the v_{10} coefficient is set to 1.0. Moreover, since statistical fluctuations at high decay times may strongly affect v_{11} , the latter is constrained to be the linear extrapolation from the previous two coefficients:

$$v_{11} = v_{10} + \frac{v_{10} - v_9}{t_{10} - t_9} \times (t_{11} - t_{10}). \quad (5.2)$$

The knot positions and the number of knots are optimized in order to fit the $B^0 \rightarrow D^\mp \pi^\pm$

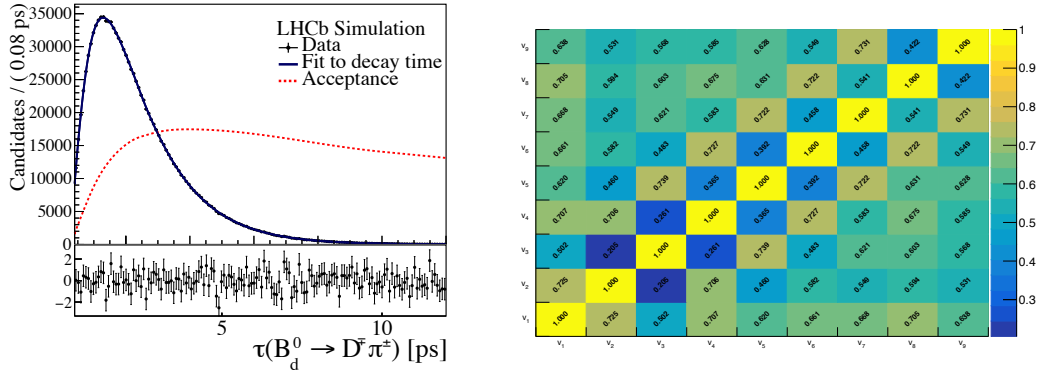


Figure 5.3 – Left: distribution of the reconstructed decay time of simulated and selected $B^0 \rightarrow D^\mp \pi^\pm$ decays (data points), with fit model superimposed (blue curve), and fitted acceptance function (red dotted curve). Right: correlation matrix of the nine fitted acceptance parameters.

Table 5.3 – Acceptance parameters fitted on the signal Monte Carlo sample.

Parameter name	Fitted value
v_1	0.1961 ± 0.0016
v_2	0.3348 ± 0.0032
v_3	0.6159 ± 0.0057
v_4	0.8667 ± 0.0073
v_5	0.9982 ± 0.0086
v_6	1.0747 ± 0.0091
v_7	1.1051 ± 0.0094
v_8	1.1590 ± 0.0086
v_9	1.188 ± 0.014

Monte Carlo decay time distribution with sufficient fit quality. The PDF adopted in this fit is proportional to:

$$a(t) \int dt' \mathcal{R}(t - t') e^{-t'/\tau_d}, \quad (5.3)$$

where $\mathcal{R}(t - t')$ is the average resolution model discussed in Sec. 5.1 and τ_d is the B^0 lifetime value used in the Monte Carlo generation. All acceptance coefficients are floating in the fit, while resolution and lifetime are fixed.

The fit projection is shown in Fig. 5.3 together with the correlation matrix obtained from the fit, whereas the fitted coefficients are listed in Table 5.3.

5.3 Decay-time fit to data

The CP coefficients S_f and $S_{\bar{f}}$ are determined from an unbinned maximum likelihood fit where each candidate is weighted with the $sWeights$ extracted from the mass fit described in Sec. 4.3.3. Hence, the total PDF is given solely by the PDF $f(t, d_{OS}, d_{SS}, \eta_{OS}, \eta_{SS}, q)$ describing the signal distribution. This is proportional to

$$a(t) \int dt' \mathcal{R}(t-t') P(t'|d_{OS}, d_{SS}, \eta_{OS}, \eta_{SS}, q) E^{OS}(\eta_{OS}) E^{SS}(\eta_{SS}) D^{OS}(d_{OS}) D^{SS}(d_{SS}) Q(q), \quad (5.4)$$

where $R(t-t')$ is the Gaussian resolution function, $a(t)$ is the acceptance function, $E^{OS}(\eta_{OS})$ and $E^{SS}(\eta_{SS})$ are the PDFs of the predicted mistag probability of the taggers, $D^{OS}(d_{OS})$ and $D^{SS}(d_{SS})$ are the PDFs of the decision of the taggers, and $Q(q)$ is the PDF of the final states. The function $P(t|d_{OS}, d_{SS}, \eta_{OS}, \eta_{SS}, q)$ represents the expected decay-time distribution for a B^0 or a \bar{B}^0 decaying into a $D^-\pi^+$ or $D^+\pi^-$ final state. This is conditional on the tagging decisions d_i , the mistag probabilities η_i and the final state q , and it contains the decay rates of Eqs. 1.48–1.51. A detailed description of the time PDF including the tagging parameters, as well as the detection and production asymmetries, is given in Appendix G. The function maximised during the fit is the logarithm of the likelihood obtained from the PDF given in Eq. 5.4,

$$\ln \mathcal{L} = s \sum_{i=1}^N s_W^i \ln f(t_i, d_{OS}^i, d_{SS}^i, \eta_{OS}^i, \eta_{SS}^i, q^i). \quad (5.5)$$

where s_W^i are the $sWeights$, N is the number of candidates in the fitted sample, and s is a correction factor given by

$$s = \frac{\sum_{i=1}^N s_W^i}{\sum_{i=1}^N (s_W^i)^2}. \quad (5.6)$$

This factor s takes into account the dilution due to the background subtraction with the $sWeights$, so that correctly-estimated uncertainties from the fit are obtained [82].

In the PDF, $\Delta\Gamma$ is fixed to zero. Moreover, the C_f ($C_{\bar{f}}$) coefficient is fixed to 1 (-1) because the value of $r_{D\pi}^2$ is such that the sensitivity to C_f ($C_{\bar{f}}$) is negligible. The tagging efficiency differences $\Delta\varepsilon^i$ ($i = OS, SS$) are found to be consistent with zero in the $B^0 \rightarrow D^-\pi^+$ Monte Carlo sample: for this reason, these coefficients are fixed to zero in the fit. Systematic uncertainties will be considered in Sec. 5.5 for all these assumptions.

The following physics parameters are Gaussian-constrained to their measured values,

$$\tau = 1/\Gamma = 1.518 \pm 0.004 \text{ ps}, \quad (5.7)$$

$$\Delta m = 0.5050 \pm 0.0023 \text{ ps}^{-1}, \quad (5.8)$$

where τ is taken as the world average value [74], and Δm is the LHCb measurement from semileptonic B^0 decays [83] (the world average of Δm is not included because it uses an analysis performed by the LHCb collaboration on Run 1 $B^0 \rightarrow D^\mp \pi^\pm$ decays as input).

The free parameters of the fit are:

- the S_f and $S_{\bar{f}}$ coefficients;
- the detection asymmetry¹

$$A_D = \frac{\varepsilon(f) - \varepsilon(\bar{f})}{\varepsilon(f) + \varepsilon(\bar{f})}, \quad (5.9)$$

where $\varepsilon(f)$ ($\varepsilon(\bar{f})$) is the detection efficiency of the final state f (\bar{f});

- the production asymmetry

$$A_P = \frac{\sigma(\bar{B}^0) - \sigma(B^0)}{\sigma(\bar{B}^0) + \sigma(B^0)}, \quad (5.10)$$

where $\sigma(\bar{B}^0)$ ($\sigma(B^0)$) is the inclusive \bar{B}^0 (B^0) production cross-section in the LHCb acceptance;

- the calibration parameters for both OS and SS taggers;
- the tagging efficiencies $\varepsilon_{\text{tag}}^{\text{OS}}$ and $\varepsilon_{\text{tag}}^{\text{SS}}$;
- the time acceptance coefficients.

The value of the parameters obtained from the fit to data are listed in Table 5.4. The correlation matrix of the parameters is reported in Appendix H. The projection of the PDF on the decay-time distribution is shown in Fig 5.4, while Fig 5.5 shows the B^0 - \bar{B}^0 asymmetries of Eq. 1.46 (distorted by experimental effects) for the two final states.

It is possible to define the CP asymmetries between Cabibbo-favoured (CF) and Cabibbo-suppressed (CS) rates as follows:

$$A_{\text{CF}} = \frac{\Gamma_{B^0 \rightarrow f}(t) - \Gamma_{\bar{B}^0 \rightarrow \bar{f}}(t)}{\Gamma_{B^0 \rightarrow f}(t) + \Gamma_{\bar{B}^0 \rightarrow \bar{f}}(t)}, \quad A_{\text{CS}} = \frac{\Gamma_{\bar{B}^0 \rightarrow f}(t) - \Gamma_{B^0 \rightarrow \bar{f}}(t)}{\Gamma_{\bar{B}^0 \rightarrow f}(t) + \Gamma_{B^0 \rightarrow \bar{f}}(t)}, \quad (5.11)$$

where $f = D^- \pi^+$ and $\bar{f} = D^+ \pi^-$. These asymmetries above (distorted by experimental effects) are plotted together with data in Fig. 5.6. The projections of the fitting function considering the decay-time distribution of the four independent decays rates, $B^0 \rightarrow D^- \pi^+$, $\bar{B}^0 \rightarrow D^- \pi^+$, $B^0 \rightarrow D^+ \pi^-$ and $\bar{B}^0 \rightarrow D^+ \pi^-$, for OS and SS tagged candidates, are

¹This definition has the opposite sign compared to the one in Ref. [84].

5.3. Decay-time fit to data

Table 5.4 – Results of the decay time fit.

Parameter	Fitted value	Comment
S_f	0.058 ± 0.021	Statistical uncertainty when fitting w/o Gauss-const. and PIDK syst. is 0.0198
$S_{\bar{f}}$	0.038 ± 0.021	Statistical uncertainty when fitting w/o Gauss-const. and PIDK syst. is 0.0199
A_P	-0.0064 ± 0.0028	Compare with -0.0100 ± 0.0047 (Eq. 5.13)
A_D	0.0086 ± 0.0019	Compare with Ref. [84] (Sec. 5.4)
Γ	$0.6587 \pm 0.0017 \text{ ps}^{-1}$	Gaussian-constrained to $0.6588 \pm 0.0017 \text{ ps}^{-1}$
Δm	$0.5054 \pm 0.0022 \text{ ps}^{-1}$	Gaussian-constrained to $0.5050 \pm 0.0023 \text{ ps}^{-1}$
p_0^{OS}	-0.152 ± 0.021	OS tagger calibration parameters
p_1^{OS}	-0.035 ± 0.024	
p_2^{OS}	-0.0070 ± 0.0089	
p_3^{OS}	-0.32 ± 0.11	
p_4^{OS}	-0.47 ± 0.49	
Δp_0^{OS}	-0.079 ± 0.049	
Δp_1^{OS}	0.141 ± 0.036	
Δp_2^{OS}	-0.024 ± 0.013	
Δp_3^{OS}	-0.26 ± 0.16	
Δp_4^{OS}	-0.52 ± 0.71	
p_0^{SS}	-0.041 ± 0.021	SS tagger calibration parameters
p_1^{SS}	-0.012 ± 0.022	
Δp_0^{SS}	-0.085 ± 0.044	
Δp_1^{SS}	0.043 ± 0.033	
$\varepsilon_{\text{tag}}^{\text{OS}}$	0.43237 ± 0.00077	Fraction of OS tagged candidates (relative to tagged candidates only)
$\varepsilon_{\text{tag}}^{\text{SS}}$	0.93046 ± 0.00040	Fraction of SS tagged candidates (relative to tagged candidates only)
v_1	0.3192 ± 0.0062	Time acceptance coefficients
v_2	0.494 ± 0.010	
v_3	0.793 ± 0.016	
v_4	0.994 ± 0.019	
v_5	1.093 ± 0.021	
v_6	1.117 ± 0.021	
v_7	1.140 ± 0.021	
v_8	1.175 ± 0.019	
v_9	1.154 ± 0.026	

shown in Figs. 5.7 and 5.8, respectively. The 2D contour plots for the CP coefficients S_f and $S_{\bar{f}}$ and for the detection and production asymmetry are shown in Fig. 5.9.

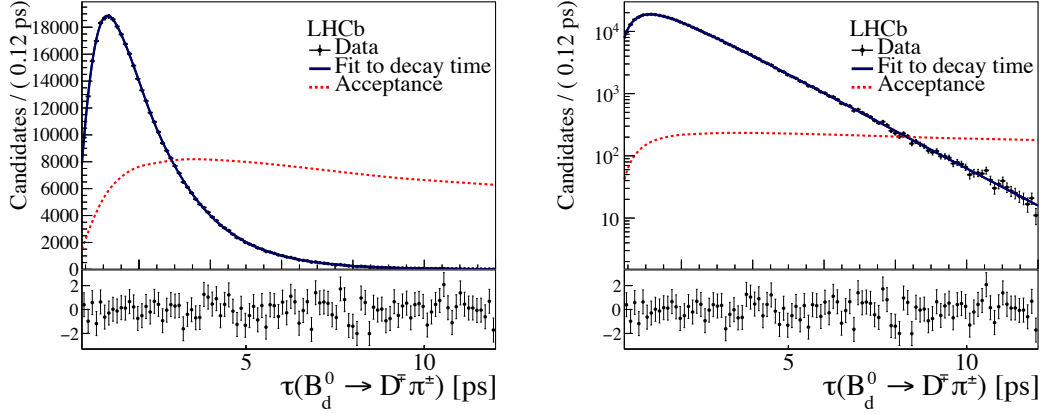


Figure 5.4 – Distribution of the reconstructed decay time of $sWeighted B^0 \rightarrow D^\mp \pi^\pm$ decays (data points), with fit model superimposed (blue curve), and fitted acceptance function (red dotted curve), in linear (left) and logarithmic (right) scale.

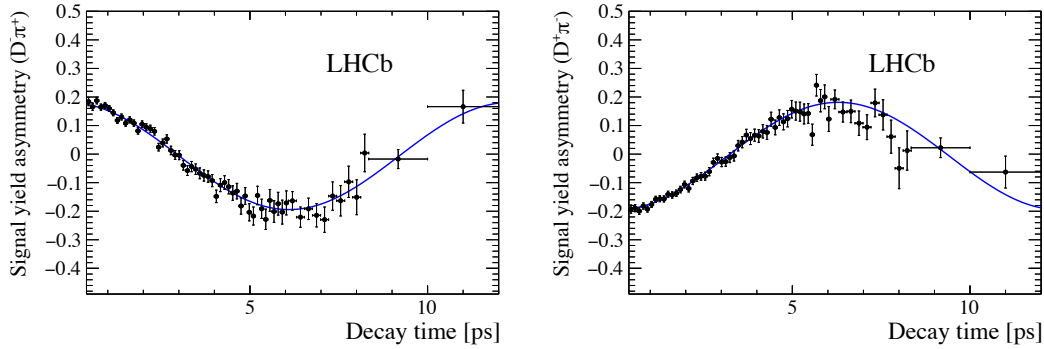


Figure 5.5 – Time-dependent asymmetry between B^0 and \bar{B}^0 decays (data points) for the $D^-\pi^+$ (left) and $D^+\pi^-$ (right) final states, with fit model superimposed (blue curve).

Using the flavour tagging calibrations obtained from the fit, the tagging performance of the signal sample is computed. The average of the total squared dilution is $(6.554 \pm 0.017)\%$. Taking into account also untagged candidates, *i.e.* considering the tagging efficiency of $(85.23 \pm 0.05)\%$, the tagging power is $(5.59 \pm 0.01)\%$.

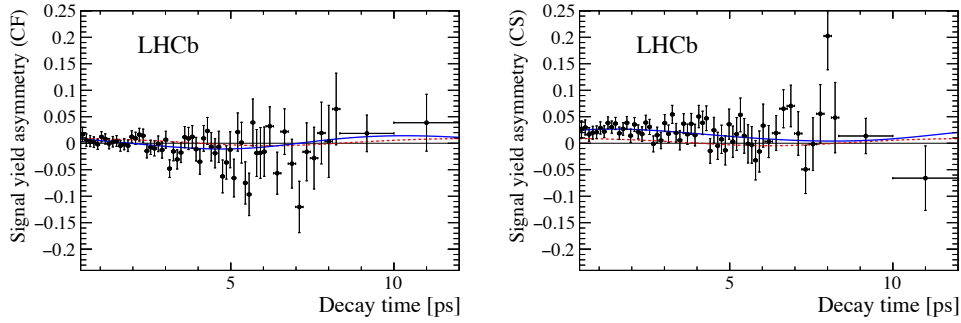


Figure 5.6 – Decay-time-dependent signal-yield asymmetries for (left) Cabibbo-favoured and (right) Cabibbo-suppressed decay topologies, defined in Eq. 5.11. The points with error bars are the data, the blue solid curve is the fit model, and the red dotted curve indicates the fit model when $S_{\bar{f}} \equiv -S_f$ (*i.e.* CP conservation in the interference between mixing and decay) is required.

5.4 Fit validation

5.4.1 Check of nuisance parameters

The values of the nuisance parameters obtained in the fit (production/detection asymmetries, flavour tagging calibrations) are compared with available external measurements.

The production asymmetry A_P is compared with the LHCb measurement of Ref. [84]. The production asymmetry is computed by weighting the production asymmetry measured from this paper in bins of p_T and η , $A_{P,i}$, with the signal fractions $\varepsilon_i = \frac{s_i}{\sum_j s_j}$ in each bin i of $B^0 \rightarrow D^{\mp}\pi^{\pm}$ data,

$$A_P = \sum_i \varepsilon_i A_{P,i}, \quad (5.12)$$

where s_i is the sum of the *sWeights* in bin i . This yields

$$A_P = -0.0100 \pm 0.0047 \text{ (stat)} \pm 0.0004 \text{ (syst)}, \quad (5.13)$$

compatible within 0.65σ with the value obtained from the $B^0 \rightarrow D^{\mp}\pi^{\pm}$ decay-time fit.

The detection asymmetry A_D is also obtained from Ref. [84], and is measured using $B^0 \rightarrow J/\psi K^{*0}$ to be 0.0098 ± 0.0046 and 0.0056 ± 0.0030 for 2011 and 2012, respectively, and using $B_s^0 \rightarrow D_s^- \pi^+$ to be 0.0143 ± 0.0086 and 0.0103 ± 0.0058 for 2011 and 2012, respectively. The central value obtained from the fit to $B^0 \rightarrow D^{\mp}\pi^{\pm}$ is in agreement with this set of results.

The values of the parameters of the tagging calibrations are compared with those found in the control samples as described in Secs. 3.2.1 and 3.2.2. The strategy presented in

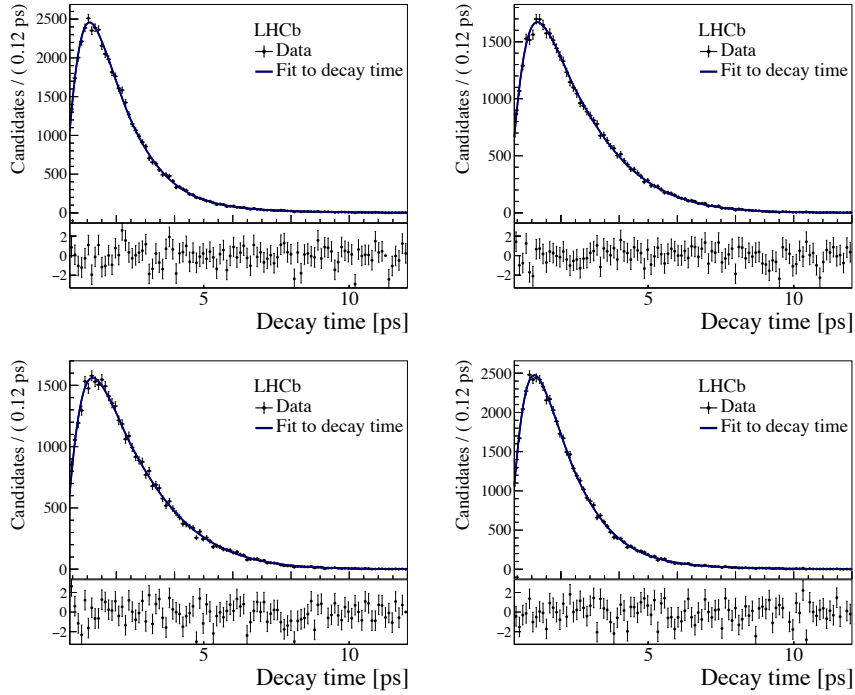


Figure 5.7 – Decay-time distributions of the $sWeighted$ data samples for (top left) $B^0 \rightarrow D^- \pi^+$, (top right) $\bar{B}^0 \rightarrow D^- \pi^+$, (bottom left) $B^0 \rightarrow D^+ \pi^-$ and (bottom right) $\bar{B}^0 \rightarrow D^+ \pi^-$ for OS inclusively tagged candidates. The fit result is superimposed as the blue curves.

Sec. 3.2 is followed, and no perfect portability of the calibrations is assumed a priori, as Sec. 3.2.1 has shown this not to be the case. However, the parameter values found in the signal fit are expected to be similar to those from the control channels given in Tables 3.2 and 3.4. A full comparison with a χ^2 test that takes into account the correlation between the parameters is performed. The discrepancy (corresponding to the χ^2 minimum) is around 0.91σ for the OS tagger, and 0.29σ for the SS tagger. The parameters which present the largest disagreement are Δp_3^{OS} for the OS tagger and the Δp_0^{SS} for the SS tagger.

5.4.2 Fits in data subsamples

A check of the stability of the results against the different data taking conditions is performed by repeating the fit in four subsamples of the data, namely data taken with magnet “up” and “down” polarities, and data taken in 2011 and 2012. The $sWeights$ for each subsample are obtained via the mass fits described in Sec. 4.3.4. The detailed results of these time fits are reported in Appendix I. A comparison between the fitted values for S_f and $S_{\bar{f}}$ obtained in each subsample is shown in Fig. 5.10. In all cases, the parameters

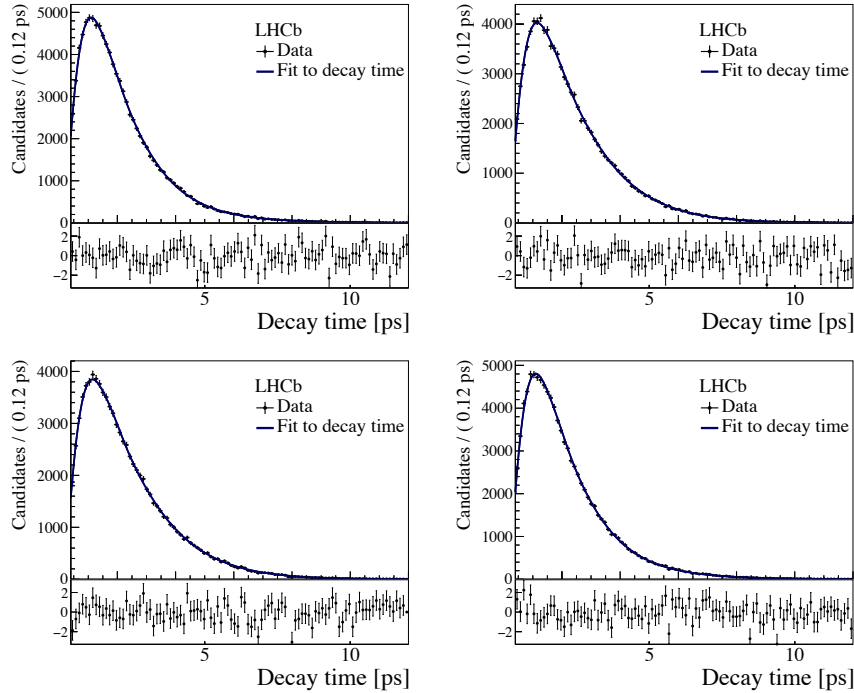


Figure 5.8 – Decay-time distributions of the *sWeighted* data sample for (top left) $B^0 \rightarrow D^-\pi^+$, (top right) $\bar{B}^0 \rightarrow D^-\pi^+$, (bottom left) $B^0 \rightarrow D^+\pi^-$ and (bottom right) $\bar{B}^0 \rightarrow D^+\pi^-$ for SS inclusively tagged candidates. The fit result is superimposed as the blue curves.

are in agreement and have good p -values, the smallest being the one between the values of $S_{\bar{f}}$ from the magnet polarity splits (2.7%). In addition, the average between the fitted values in each split (black line in Fig. 5.10) is always very close to the central value from the nominal fit (red hatched band).

The stability of the results against the tagging algorithm adopted in the fit are also checked. In this case, the data sample with *sWeights* obtained from the nominal mass fit (Sec. 4.3) is split in three independent subsamples according to the tagging decision:

- candidates tagged exclusively by the OS tagger;
- candidates tagged exclusively by the SS tagger;
- candidates tagged by both the OS and SS taggers.

The values of S_f and $S_{\bar{f}}$ obtained in these subsamples are compared in Fig. 5.11. All values are compatible. More details are given in Appendix I. Given the difference of the tagging algorithms and their calibrations, the stability of the results in this test provides additional confidence on the strategy adopted of floating the calibration parameters in the fit.

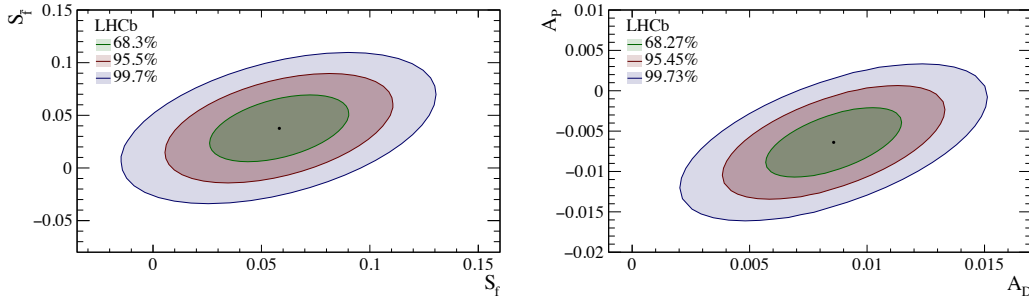


Figure 5.9 – Plots for $(S_f, S_{\bar{f}})$ (left) and (A_P, A_D) (right) showing the one, two and three sigma contours. The shown uncertainties include the full statistical uncertainty and the systematic uncertainty due to the Gaussian constraints on the mixing frequency Δm and the B^0 decay width Γ .

The stability of the results against the B^0 kinematics and global properties of the event is tested. More specifically, the decay time fit is repeated in bins of the following variables:

- transverse momentum of the B^0 (4 bins);
- number of reconstructed primary vertices (3 bins);
- total number of reconstructed tracks (3 bins);
- difference in pseudorapidity ($\Delta\eta$) between D meson and bachelor pion (4 bins).

The motivation for these tests is that flavour tagging calibration parameters depend on the above observables; as a consequence, the fitted values for the S_f and $S_{\bar{f}}$ coefficients might also show a significant trend in these variables because of the correlation with the flavour tagging calibrations. Moreover, the difference in pseudorapidity is sensitive to potential misalignments in the detectors which might affect the measured value of CP asymmetries. The values of S_f and $S_{\bar{f}}$ obtained in these subsamples are compared in Fig. 5.12, whereas more details are given in Appendix I. All values are compatible, and no significant dependence of S_f and $S_{\bar{f}}$ on the studied variables is observed.

Finally, the time fit is repeated separately for TOS candidates on L0Hadron and all the other candidates. The values of S_f and $S_{\bar{f}}$ obtained in these subsamples are compared in Fig. 5.13. All values are compatible, and no significant dependence of S_f and $S_{\bar{f}}$ is observed. More details can be found in Appendix I.

5.4.3 Time fits to bootstrapped Monte Carlo samples

The fit is also validated using Monte Carlo simulation. The $B^0 \rightarrow D^\mp \pi^\pm$ simulated sample is *bootstrapped* [85] (*i.e.*, resampled allowing repetition of the same candidates) 1000 times. Each bootstrapped sample contains the same signal yield as obtained from the nominal mass fit (Table 4.15), corrected for the *sWeights* dilution factor of Eq. 5.6 to

5.4. Fit validation

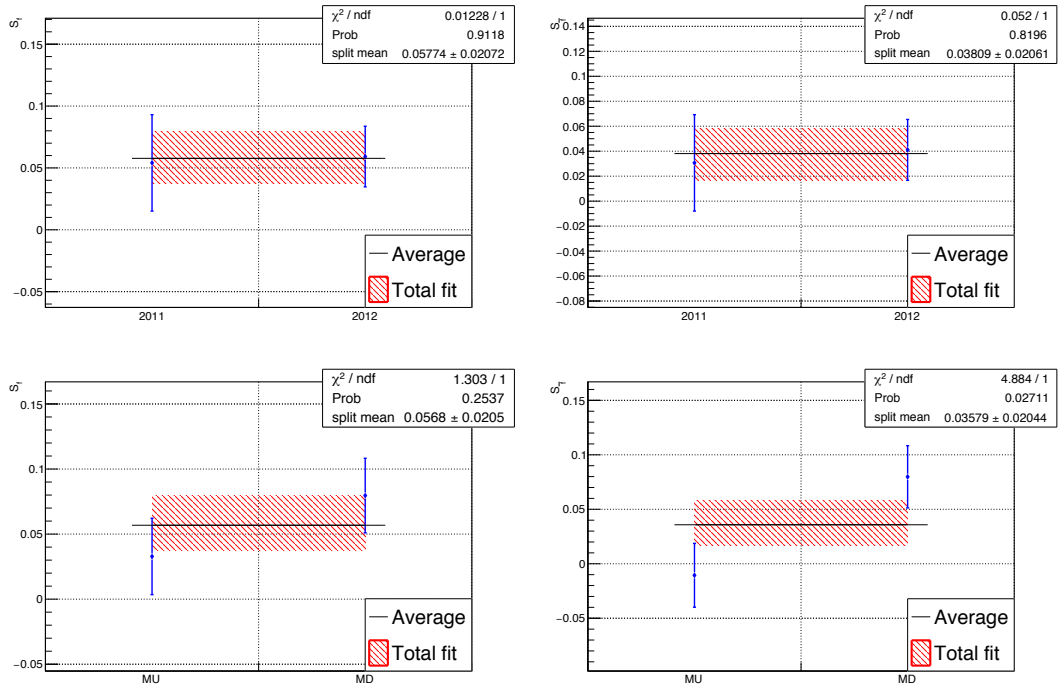


Figure 5.10 – Fitted values of S_f (left) and $S_{\bar{f}}$ (right) as a function of the data-taking year (top) and magnet polarity (bottom). The red hatched band shows the values obtained from the nominal fit of the full sample. The horizontal black line is the result of a χ^2 fit to obtain the weighted average of the results of each subsample.

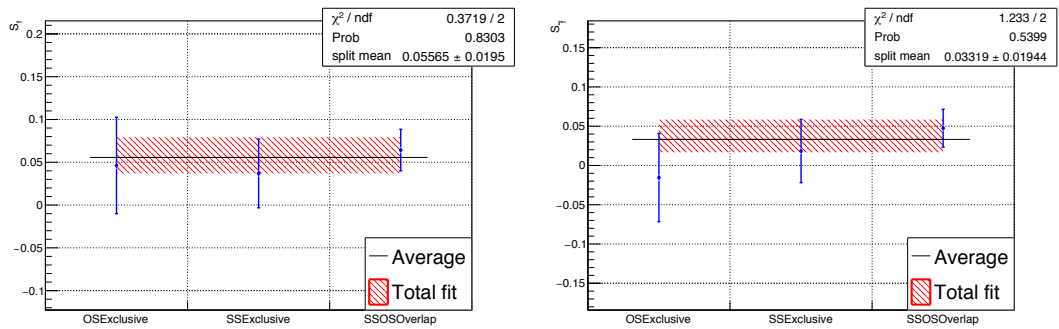


Figure 5.11 – Fitted values of S_f (left) and $S_{\bar{f}}$ (right) when candidates tagged exclusively by OS or SS, or both simultaneously are considered. The red hatched band shows the values obtained from the nominal fit of the full sample. The horizontal black line is the result of a χ^2 fit to obtain the weighted average of the results of each subsample.

Chapter 5. Measurement of CP violation in $B^0 \rightarrow D^\mp \pi^\pm$ decays

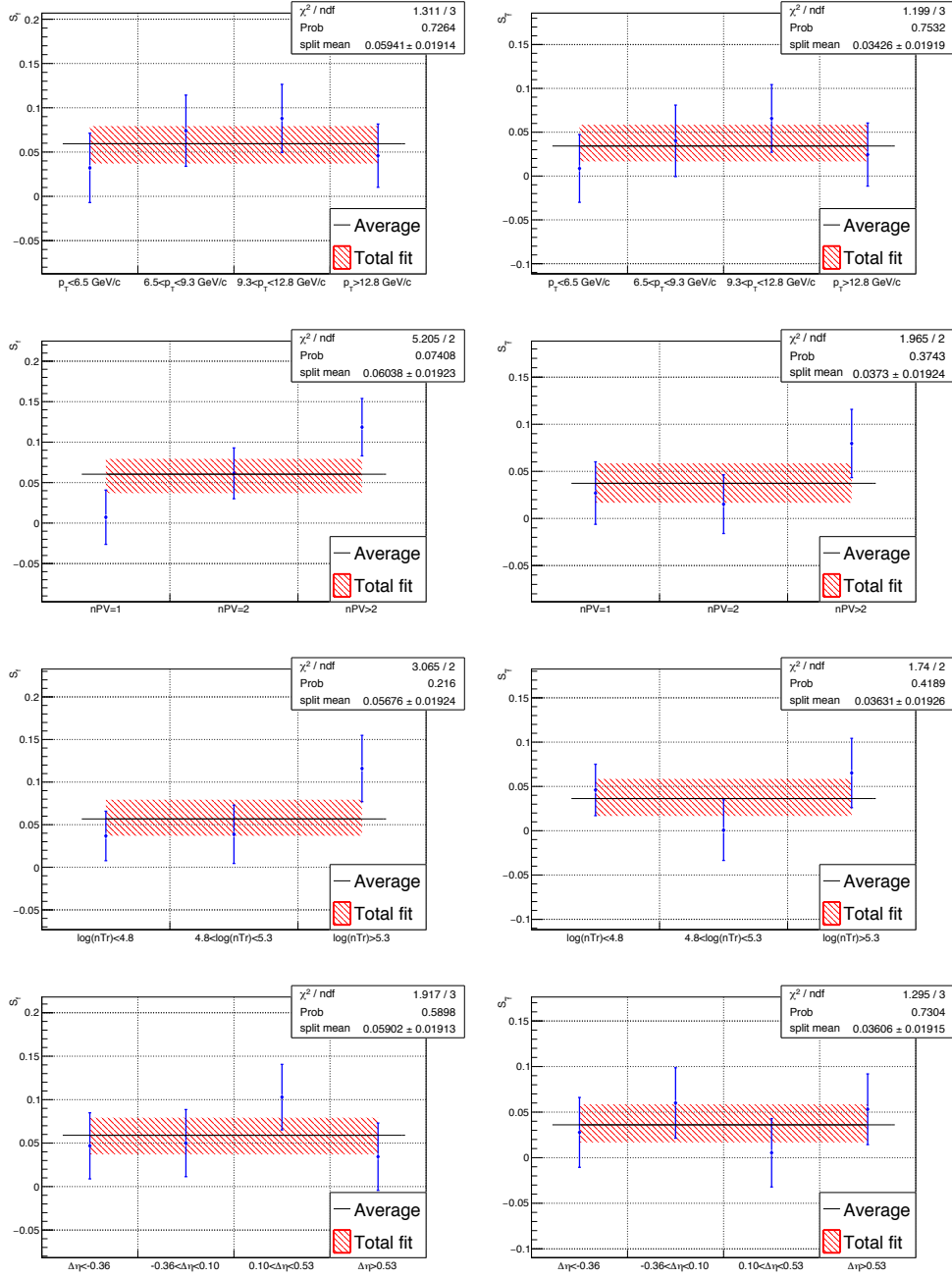


Figure 5.12 – Fitted values of S_f (left) and $S_{\bar{f}}$ (right) when the decay time fit is performed in bins of (from top to bottom) the transverse momentum of the B^0 , number of primary vertices, number of tracks and difference in pseudorapidity between the D meson and the bachelor pion. The red hatched band shows the values obtained from the nominal fit of the full sample. The horizontal black line is the result of a χ^2 fit to obtain the weighted average of the results of each subsample.

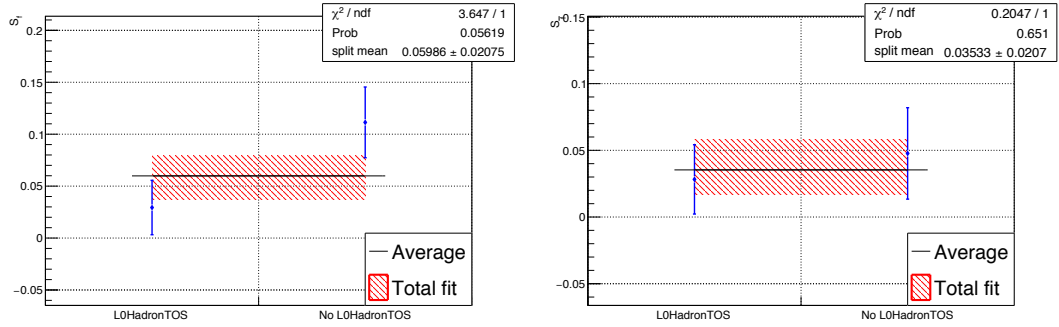


Figure 5.13 – Fitted values of S_f (left) and $S_{\bar{f}}$ (right) when the decay time fit is performed separately for TOS candidates on L0Hadron and all the other candidates. The red hatched band shows the values obtained from the nominal fit of the full sample. The horizontal black line is the result of a χ^2 fit to obtain the weighted average of the results of each subsample.

have the same effective yield as in the data time fit.

Each sample is then fitted using exactly the same strategy as described in Sec. 5.3, except that the central value of the Gaussian constraints on the Γ and Δm parameters is randomly drawn from a Gaussian distribution centred on the Monte Carlo generation value given in Appendix M ($1/\Gamma = 1.519$ ps and $\Delta m = 0.510$ ps $^{-1}$) with a standard deviation equal to the width of the constraint (± 0.004 ps for $1/\Gamma$ and ± 0.0023 ps $^{-1}$ for Δm). This allows fluctuations of the Γ and Δm measurements, and avoids underestimation of the fitted uncertainties.

The distributions of the fitted value, uncertainty, pull and residual² of S_f and $S_{\bar{f}}$ are shown in Fig. 5.14. Other fitted parameters are reported in Appendix J. Each of these distribution is fitted with a Gaussian function. The width of the fitted pull distributions are close to unity, meaning that the uncertainty coming from the likelihood fit is correctly estimated. The mean value of the distribution of the uncertainties of each parameter is close to the value of the uncertainty found in the fit to data. The on-average better precision found in the fit to MC is due to the higher tagging performance of the simulation.

The distribution of the residuals of the S_f parameter shows a mean of 0.0071 ± 0.0006 , corresponding to one third of the statistical uncertainty of the fit to data; for $S_{\bar{f}}$, the mean is -0.0013 ± 0.0006 , which corresponds to about 6% of the statistical uncertainty of the fit to data.

Several configurations are implemented to test the bootstrap study and its results, and

²The residual is defined as fitted value minus generated value, whereas the pull is the residual divided by the fitted uncertainty.

to try to address the origin of these biases. The fits to the bootstrapped samples are repeated in the following different configurations:

- using the true flavour of the B^0 candidate instead of the tagging decision and mistag probability: no biases are found on S_f and $S_{\bar{f}}$;
- using a *toy* (or *cheated*) tagger, as explained in Appendix K: no biases are found on S_f and $S_{\bar{f}}$;
- using the calibration parameters obtained in the signal MC sample using the true flavour information (see Sec. 3.2.1 and Sec. 3.2.2): no biases are found on S_f and $S_{\bar{f}}$;
- fixing the calibration parameters to the values obtained from the MC samples of the control channels: biases of the order 1σ on S_f and $S_{\bar{f}}$ are found;
- applying Gaussian-constraints on the calibration parameters using the values obtained from the MC samples of the control channels: biases of the order of half the statistical uncertainty of S_f and $S_{\bar{f}}$ are found;

This study confirms that the strategy of floating the calibration parameters in the fit is the optimal choice. Other than the biases related to the flavour-tagging calibrations, the origin of the small bias observed on the S_f parameter in the nominal configuration could not be clarified. To confirm this bias, the study is repeated by fitting additional 1000 bootstrapped samples using an independent fitter. The mean of the distribution of the residuals in this second study is confirmed to be of the same size, namely 0.0064 ± 0.0007 for S_f and -0.0024 ± 0.0007 for $S_{\bar{f}}$. Hence, the weighted average of the small residuals on S_f (0.0068 ± 0.0005) and $S_{\bar{f}}$ (-0.0018 ± 0.0005) of both studies are considered as systematic uncertainties. As described in Appendix L, the correlation between the systematic uncertainties on S_f and $S_{\bar{f}}$ associated to the fit biases reported here is 0.4.

5.5 Systematics

The identified systematic uncertainties are listed in Table 5.5 in decreasing order of their size. Their quadratic sum is 0.011 and 0.007 for S_f and $S_{\bar{f}}$, respectively. A description of each systematic effect is given in the following subsections. The “fit biases” are the residuals observed in the Monte Carlo bootstrap study discussed in Sec. 5.4.3.

5.5.1 Systematic uncertainties from Gaussian constraints

Systematic uncertainties due to external measurements used in the PDF are accounted for through Gaussian constraints in the likelihood. These parameters are the mixing frequency, Δm , and the B^0 lifetime, τ . The fit has been repeated by fixing the Gaussian-constrained parameters to their central values, in order to not propagate the uncertainty of these parameters to the statistical uncertainties of the fit. The statistical uncertainties

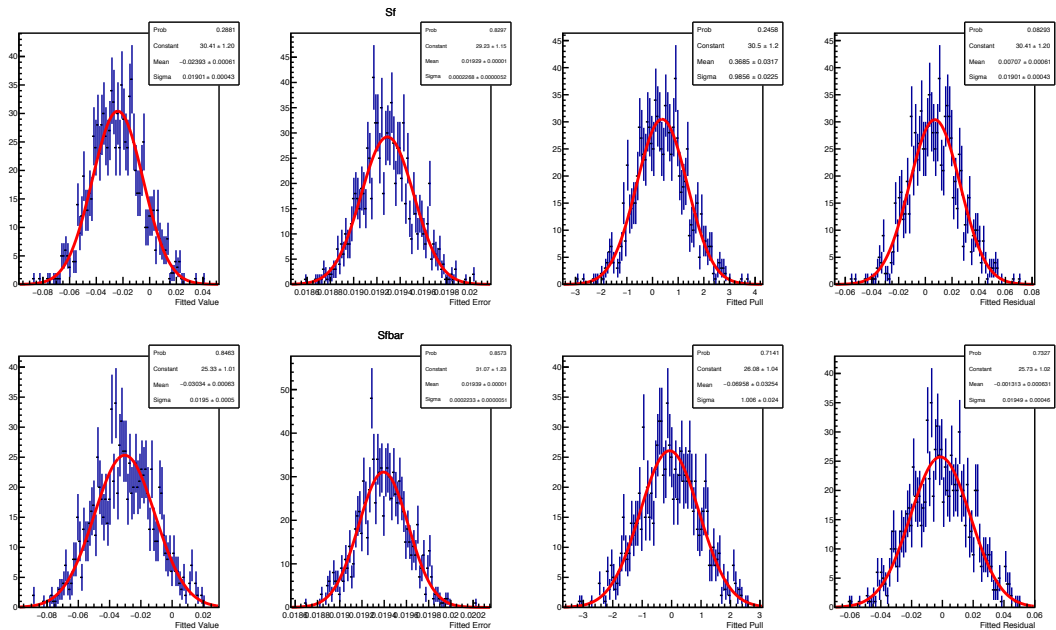


Figure 5.14 – Distributions of the fitted value, error, pull and residual for S_f (top) and $S_{\bar{f}}$ (bottom) obtained from fits to bootstrapped Monte Carlo samples. Each distribution is fitted with a Gaussian function. Pulls and residuals are computed by taking the Monte Carlo generation values as reference (Appendix M).

of S_f and $S_{\bar{f}}$ with Δm fixed are 0.0198 and 0.0199, respectively, whereas the statistical correlation is 0.6. Considering the difference in quadrature between the uncertainty from the nominal fit and that from this fit, the systematic uncertainty due to Δm are 0.0073 and 0.0061 for S_f and $S_{\bar{f}}$, respectively. The correlation between the systematic uncertainties due to Δm on S_f and $S_{\bar{f}}$ is -1 , as described in Appendix L. The fit with τ fixed shows that the systematics uncertainty due to this parameter is negligible.

Systematic uncertainty associated with the PIDK efficiencies (Table 4.10) are taken into account in the mass fit by means of Gaussian constraints on these parameters (Sec. 4.3.2). The mass fit is repeated by neglecting these uncertainties in the Gaussian constraints. Then, the time fit is performed with this new set of $sWeights$. The difference in quadrature between the uncertainty from this fit and that from the nominal fit gives the systematic due to the binning scheme in the PIDK resampling, which is 0.0008 for both S_f and $S_{\bar{f}}$.

5.5.2 Systematic uncertainties estimated with pseudoexperiments

When computing the systematic uncertainties with pseudoexperiments (or *toys*), a sample with the same size as the data is generated by sampling the PDF with parameters fixed to the value found in the data fit. The values of S_f and $S_{\bar{f}}$ are fixed to those used in the

Table 5.5 – Systematic uncertainties on the CP parameters S_f and $S_{\bar{f}}$.

Source	S_f	$S_{\bar{f}}$
uncertainty of Δm	0.0073	0.0061
fit biases	0.0068	0.0018
background subtraction	0.0042	0.0023
flavour tagging calibration models	0.0011	0.0015
flavour tagging efficiency asymmetries	0.0012	0.0015
PIDK efficiencies	0.0008	0.0008
acceptance model	0.0007	0.0007
assumption $\Delta\Gamma = 0$	0.0007	0.0007
assumption $C_f = -C_{\bar{f}} = 1$	0.0006	0.0006
decay time resolution	0.0012	0.0008
total systematic uncertainty	0.0111	0.0073
statistical uncertainties	0.0198	0.0199

generation of the Monte Carlo sample (Appendix M). In the generation of the samples the PDF is modified to consider alternative models according to the source of systematic uncertainty under investigation. The generated sample is then fitted with the nominal model. For each parameter, the mean of the distribution of the residuals from 1000 toys is taken as a symmetric systematic uncertainty. If the mean is consistent with zero within $\pm 1 \sigma$, the error on the mean is taken instead. The systematic uncertainties estimated with this toy-based method are the following:

- the flavour tagging calibration model;
- the flavour tagging efficiency asymmetries;
- the acceptance model;
- the decay time resolution;
- the assumption $C_f = -C_{\bar{f}} = 1$;
- the assumption $\Delta\Gamma = 0$.

Flavour tagging calibration model

Toys are generated using for the SS calibration the nominal model with a first order polynomial, and for the OS the model is reduced by one degree as compared to the nominal one. In the fit, the calibration models of both taggers are increased by one degree compared to what was used in the generation step. The distribution of the residuals of S_f and $S_{\bar{f}}$ are shown in Fig. 5.15. The residuals are not compatible with zero and therefore they are assigned as systematic uncertainties.

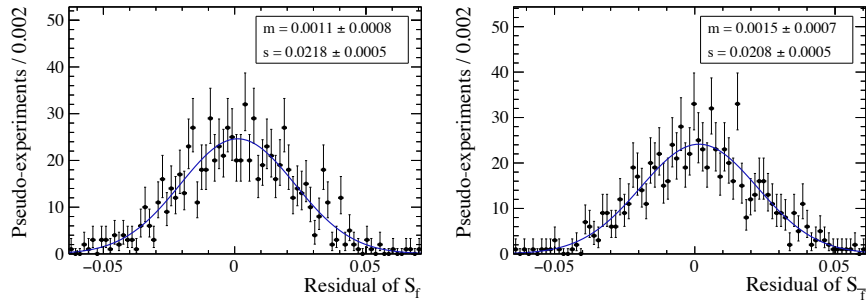


Figure 5.15 – Distribution of S_f (left) and $S_{\bar{f}}$ (right) residuals for the determination of the systematic uncertainty due to the tagging calibration models.

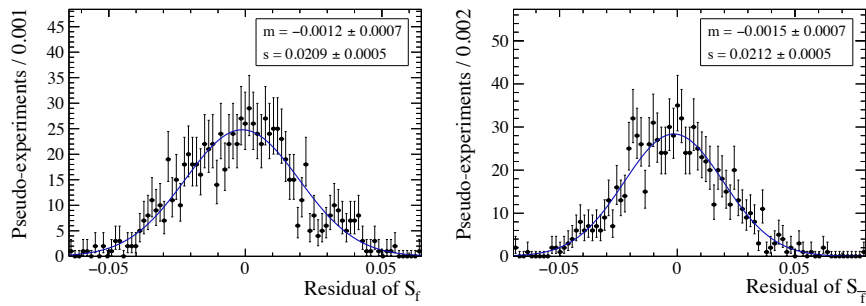


Figure 5.16 – Distribution of S_f (left) and $S_{\bar{f}}$ (right) residuals for the determination of the systematic uncertainty due to the assumption on the flavour tagging efficiency asymmetry.

Flavour tagging efficiency asymmetries

Toys are generated with the flavour tagging asymmetries set to their (negative) estimate from simulation minus their uncertainty, namely -0.14% and -0.13% for the OS and SS tagger, respectively. The distributions of the residuals of S_f and $S_{\bar{f}}$ are shown in Fig. 5.16. The residuals are not compatible with zero and therefore they are assigned as systematic uncertainties.

Acceptance model

The acceptance model is modified in the generation by replacing the nominal knots for the spline function with new knots, namely at 0.4, 0.45, 0.8, 1.3, 2.5, 6.0, and 12.0 ps. The distribution of the residuals of S_f and $S_{\bar{f}}$ are shown in Fig. 5.17. Residuals consistent with zero are found and therefore the uncertainty on the residuals is assigned as systematic uncertainty.

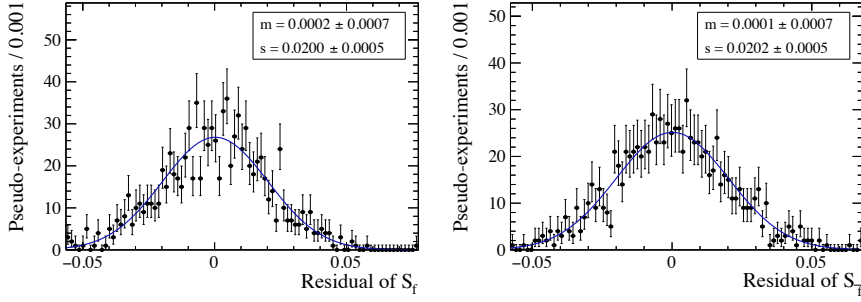


Figure 5.17 – Distribution of S_f (left) and $S_{\bar{f}}$ (right) residuals for the determination of the systematic uncertainty due to the acceptance model.

Decay time resolution

Toys are generated with time resolutions 20 fs larger and 20 fs smaller than the nominal value of 55 fs. The distributions of the fitted value of S_f and $S_{\bar{f}}$ are shown in Fig. 5.18. The largest residual is considered as overall systematic uncertainty.

Fixed C_f

Toys are generated with $C_f = -C_{\bar{f}}$ set to the average of the measurements by Belle and BaBar minus the largest uncertainty among the two measurements, namely 0.993 [28, 29]. The distributions of the residuals of S_f and $S_{\bar{f}}$ are shown in Fig. 5.19. Residuals consistent with zero are found, therefore the uncertainty on the residuals is assigned as systematic uncertainty.

Fixed $\Delta\Gamma$

Toys are generated with $\Delta\Gamma$ set to the world average value plus its uncertainty, namely 0.0079 ps^{-1} [86]. Moreover, the D_f and $D_{\bar{f}}$ coefficients (defined in Eqs. 1.35-1.38) have been fixed to their expected values of -0.0103 and -0.0155 , the same used in the Monte Carlo production of the $B^0 \rightarrow D^- \pi^+$ sample (Appendix M). The distribution of the residuals of S_f and $S_{\bar{f}}$ are shown in Fig. 5.20. Residuals consistent with zero are found, therefore the uncertainty on the residuals is assigned as systematic uncertainty.

5.5.3 Systematics related to the background subtraction

Systematic uncertainties can arise from the choice of the mass fit strategy adopted to calculate $sWeights$ (Sec. 4.3). Fit B, used to compute the $sWeights$, is repeated in the full mass window ($[5090, 6000] \text{ MeV}/c^2$) instead of the narrow signal region ($[5220, 5600] \text{ MeV}/c^2$).

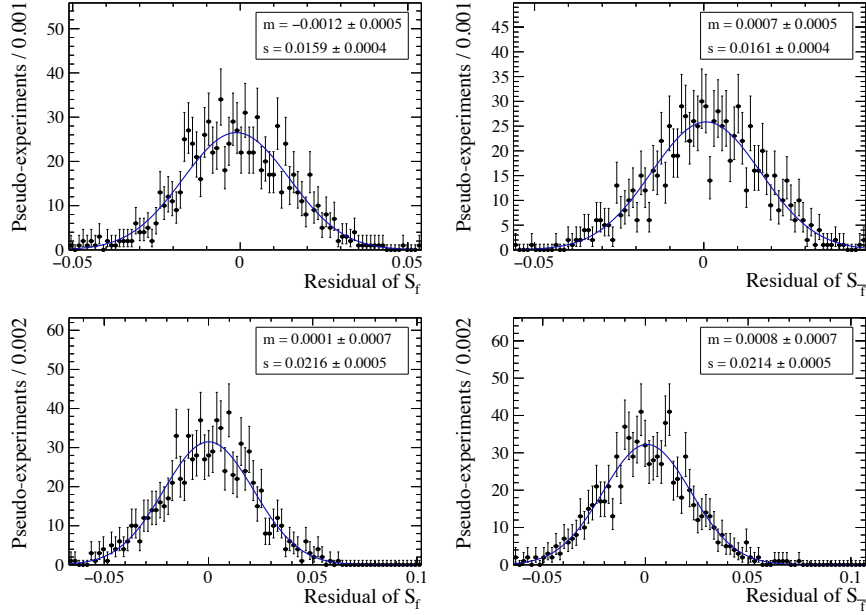


Figure 5.18 – Distribution of S_f (left) and $S_{\bar{f}}$ (right) residuals for the determination of the systematic uncertainty due to the resolution model. Top: 75 fs resolution model. Bottom: 35 fs resolution model.

In this way, the resulting sample is enriched in background events. The aim of this test is to estimate how much background events (with negative $sWeights$) affect the result for S_f and $S_{\bar{f}}$ in the final decay time fit. The fitted total background yield in this new mass fit configuration is $199\,767 \pm 481$, compared to $34\,102 \pm 299$ in the nominal fit configuration (Table 4.15). The projection of the PDF used for Fit B in the wide mass range is shown in Fig. 5.21. The new $sWeights$ are then used in a decay time fit performed on the full sample following the same strategy as reported in Sec. 5.3. The correlated disagreement, defined as the difference between the fit results divided by the difference in quadrature between the fitted uncertainties, between the result of this fit and that of the nominal fit is 2.3σ for S_f and 1.8σ for $S_{\bar{f}}$. Because of this discrepancy, the difference between the newly obtained S_f and $S_{\bar{f}}$ coefficients and the nominal values is taken as systematic uncertainty, yielding 0.0042 and 0.0023 for S_f and $S_{\bar{f}}$ respectively. The correlation between the systematic uncertainties on S_f and $S_{\bar{f}}$ is estimated to be 0.7 , as shown in Appendix L.

Another test is made by repeating the mass fit with a different strategy:

- a $PIDK < 0$ cut (instead of $PIDK < 5$) is applied on the pion PID in order to define the pion sample;
- both Fit A and Fit B are performed in the narrow signal region ($[5220, 5600] \text{ MeV}/c^2$);
- during Fit A, only the pion sample is considered (no simultaneous fit in kaon and

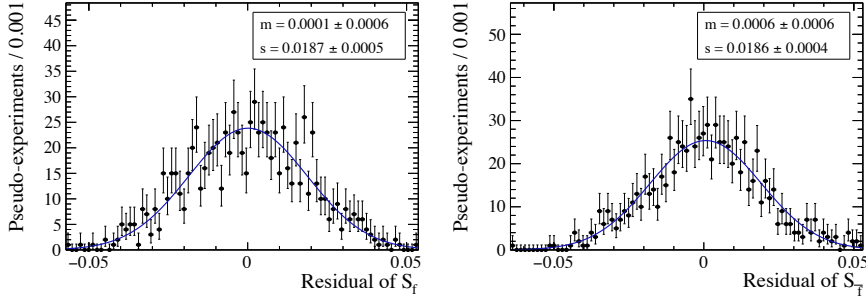


Figure 5.19 – Distribution of S_f (left) and $S_{\bar{f}}$ (right) residuals for the determination of the systematic uncertainty due to the assumption $C_f = -C_{\bar{f}} = 1$.

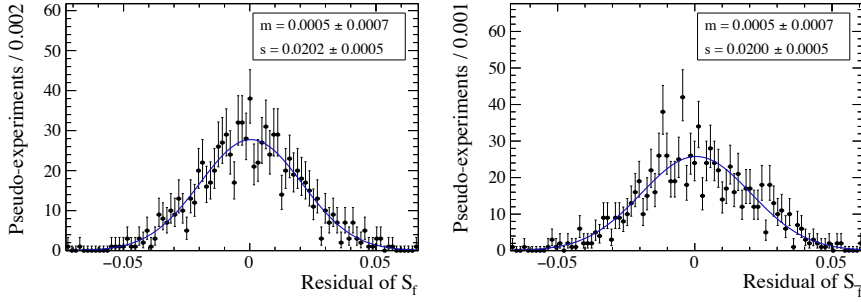


Figure 5.20 – Distribution of S_f (left) and $S_{\bar{f}}$ (right) residuals for the determination of the systematic uncertainty due the assumption $\Delta\Gamma = 0$.

pion samples is performed);

- only $B^0 \rightarrow D^\mp K^\pm$ and combinatorial background are considered, whereas all the other physical background are neglected;
- the $B^0 \rightarrow D^\mp K^\pm$ yield is Gaussian constrained to be 0.0101 ± 0.0012 of the signal yield, based on the selection efficiencies (including the $PIDK < 0$ cut) found on Monte Carlo.

The signal and total background yield obtained in this fit are $406\,818 \pm 674$ and $23\,938 \pm 266$ respectively. The projection of the PDF used for Fit A and Fit B in this configuration is shown in Fig. 5.22. A decay time fit is performed on the resulting sample with $sWeights$ by following the same strategy as reported in Sec. 5.3. The correlated discrepancy between the result of this fit and that of the nominal fit is 0.4σ and 1.6σ for S_f and $S_{\bar{f}}$ respectively. Given the good level of agreement, and the fact that systematic uncertainties on the PID efficiencies are already considered, no further systematics are assigned.

As additional cross-check, the decay time fit is repeated for $B^0 \rightarrow D^\mp \pi^\pm$ candidates restricted in the $[5250, 5330] \text{ MeV}/c^2$ invariant mass region, very close to the $B^0 \rightarrow D^\mp \pi^\pm$ signal peak position. No $sWeights$ are applied on this subsample. The correlated

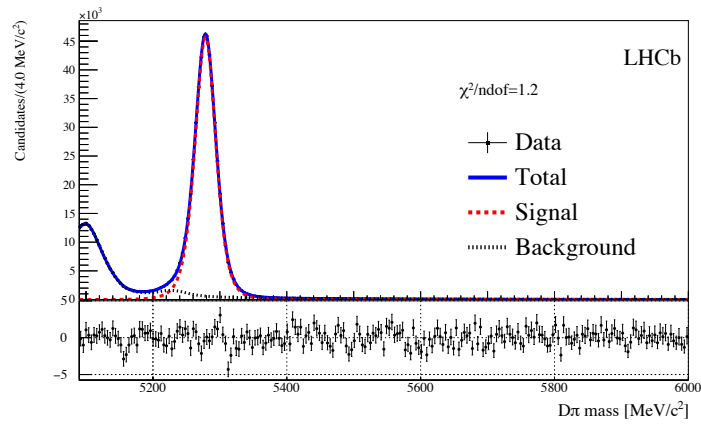


Figure 5.21 – $D^{\mp}\pi^{\pm}$ mass distribution of the π sample with the results of Fit B in the large mass window superimposed.

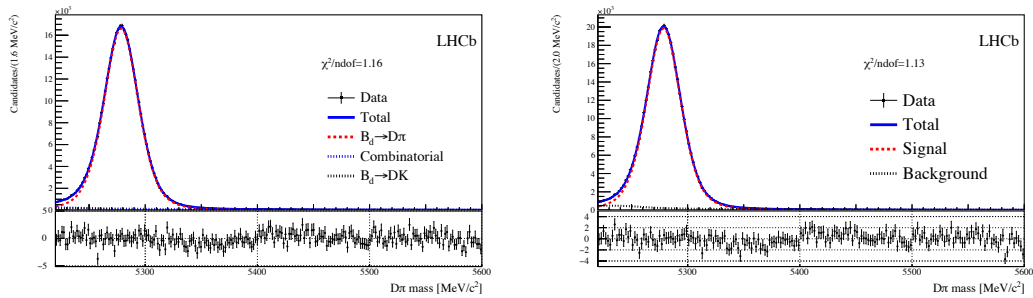


Figure 5.22 – $D^{\mp}\pi^{\pm}$ mass distribution of the alternative π sample defined by the cut $\text{PIDK} < 0$ on the bachelor pion with the result of Fit A (left) and Fit B (right) superimposed.

disagreement between the result of this fit and that of the nominal fit is 0.2σ and 1.3σ for S_f and $S_{\bar{f}}$ respectively. Given the good level of agreement, no further systematics are assigned, and the following conclusions are drawn:

- the amount of combinatorial and $B^0 \rightarrow D^{\mp}K^{\pm}$ backgrounds in the signal region is very small, and their presence doesn't affect significantly the fitted S_f and $S_{\bar{f}}$ coefficients as these are compatible with the nominal fit result;
- any systematics due to a wrong modelling of signal and/or background PDF in the $B^0 \rightarrow D^{\mp}\pi^{\pm}$ signal peak region is negligible, since the fitted value obtained from the nominal fit (with *sWeights*) and this alternative fit (with no mass fit at all) are compatible.

6 Summary and interpretation of the results

Table 6.1 – Comparison of the measurements of S_f and $S_{\bar{f}}$. The first uncertainty is statistical, the second is systematic.

	S_f [%]	$S_{\bar{f}}$ [%]
Belle [36]	$+6.8 \pm 2.9 \pm 1.2$	$+3.1 \pm 3.0 \pm 1.2$
Babar [33]	$-2.3 \pm 4.8 \pm 1.4$	$+4.3 \pm 4.8 \pm 1.4$
LHCb (this analysis) [39]	$+5.8 \pm 2.0 \pm 1.1$	$+3.8 \pm 2.0 \pm 0.7$

6.1 Interpretation of the results

In this thesis, a time-dependent analysis of the decay $B^0 \rightarrow D^\mp \pi^\pm$ in order to extract the CP observables S_f and $S_{\bar{f}}$ was presented. The values obtained are

$$S_f = 0.058 \pm 0.020(\text{stat}) \pm 0.011(\text{syst}), \quad (6.1)$$

$$S_{\bar{f}} = 0.038 \pm 0.020(\text{stat}) \pm 0.007(\text{syst}). \quad (6.2)$$

The statistical and systematic correlations between S_f and $S_{\bar{f}}$ are +60% and -41%, respectively, and the total correlation is +44%. These values are in agreement with and more precise than the measurements from Belle and Babar [36, 33]. A direct comparison is shown in Table 6.1.

The measurements of S_f and $S_{\bar{f}}$ are interpreted in terms of the angle γ and the strong phase δ , by following the definition of S_f and $S_{\bar{f}}$ given in Sec. 1.4:

$$S_f = -\frac{2r_{D\pi} \sin[\delta - (\gamma + 2\beta)]}{1 + r_{D\pi}^2}, \quad S_{\bar{f}} = \frac{2r_{D\pi} \sin[\delta + (\gamma + 2\beta)]}{1 + r_{D\pi}^2}. \quad (6.3)$$

A frequentist method, called PLUGIN and described in Ref. [87], is adopted to derive confidence intervals for γ and δ by using external inputs for β and $r_{D\pi}$. Given the observed values $\vec{A}^{\text{obs}} = (S_f, S_{\bar{f}})$ and the parameters $\vec{\alpha} = (\gamma, \delta)$, a $\chi^2(\vec{\alpha})$ function is built as

$$\chi^2(\vec{\alpha}) = -2 \ln \mathcal{L}(\vec{\alpha} | \vec{A}^{\text{obs}}) \propto \left(\vec{A}(\vec{\alpha}) - \vec{A}^{\text{obs}} \right)^T V^{-1} \left(\vec{A}(\vec{\alpha}) - \vec{A}^{\text{obs}} \right), \quad (6.4)$$

where V is the covariance matrix of S_f and $S_{\bar{f}}$. The best fit point, $\vec{\alpha}_{\text{min}}$, is the one that minimises the expression of Eq. 6.4. The p -value, or 1-CL, is computed for each possible value of each component of $\vec{\alpha}$ (γ and δ) as follows:

1. a value for a given component of $\vec{\alpha}$ (*e.g.* $\gamma = \gamma_0$) is chosen, and the associated new minimum $\vec{\alpha}'_{\text{min}}(\gamma_0)$ is found;
2. the corresponding test statistics $\Delta\chi^2 = \chi^2(\vec{\alpha}'_{\text{min}}(\gamma_0)) - \chi^2(\vec{\alpha}_{\text{min}})$ is built;
3. pseudoexperiments are generated to sample values for \vec{A} , called \vec{A}_j , from $\mathcal{L}(\vec{\alpha}'_{\text{min}}(\gamma_0) | \vec{A})$;
4. for each \vec{A}_j , a new value for the test statistics, $\Delta\chi_j^2$, is computed by replacing

- $\vec{A}^{\text{obs}} \rightarrow \vec{A}_j$ in Eq. 6.4, and by taking the difference between the minimised values of χ^2 with respect to $\vec{\alpha}$, once with γ as free parameter, and once with $\gamma = \gamma_0$;
5. the p -value is computed as the fraction of pseudoexperiments for which $\Delta\chi^2 < \Delta\chi_j^2$.

In all the minimisation steps described above, the values $\beta = (22.2 \pm 0.7)^\circ$ [86] and $r_{D\pi} = (1.82 \pm 0.12 \pm 0.36(\text{SU}(3)))\%$ are used as external Gaussian constraints. The latter is calculated from the branching fraction of $B^0 \rightarrow D_s^+ \pi^-$ decays, assuming SU(3) symmetry, following the relation of Refs. [29, 28]:

$$r_{D\pi} = \tan\theta_c \frac{f_D}{f_{D_s}} \sqrt{\frac{\mathcal{B}(B^0 \rightarrow D_s^+ \pi^-)}{\mathcal{B}(B^0 \rightarrow D^- \pi^+)}} \quad (6.5)$$

where $\tan\theta_c = 0.23101 \pm 0.00032$ [21] is the tangent of the Cabibbo angle, $\frac{f_{D_s}}{f_D} = 1.173 \pm 0.003$ [88, 89, 90] is the ratio of decay constants, $\mathcal{B}(B^0 \rightarrow D_s^+ \pi^-) = (2.16 \pm 0.26) \times 10^{-5}$ [22] and $\mathcal{B}(B^0 \rightarrow D^- \pi^+) = (2.52 \pm 0.13) \times 10^{-3}$ [22]. An additional 20% relative error is added on $r_{D\pi}$ to account for uncertainties due to possible non-factorizable SU(3)-breaking effects, as reported in Ref. [91].

The angle γ is determined to be in the interval $[5, 86]^\circ \cup [185, 266]^\circ$ and δ to be in the interval $[-41, 41]^\circ \cup [140, 220]^\circ$, both at the 68% CL. These intervals are illustrated in Fig. 6.1. In Fig. 6.2, contours are shown in the two-dimensional plane (γ, δ) .

In addition to γ and δ , the interval of $|\sin(2\beta + \gamma)|$ is determined as well. This quantity does not rely on any external input for β , and it is thus experimentally cleaner. This interval is found to be $[0.77, 1.0]$ at the 68% CL as shown in Fig. 6.3. The absolute value of $\sin(2\beta + \gamma)$ is considered because the decay-time fit cannot resolve the ambiguity on the sign, *i.e.* the same p -value is found for $2\beta + \gamma$ and $-(2\beta + \gamma)$.

The intervals for γ , δ and $|\sin(2\beta + \gamma)|$ are also determined by assuming a SU(3)-breaking uncertainty of 0%, 20% and 100% on the value of $r_{D\pi}$. These are presented in Figs. 6.4 and 6.5.

6.2 Summary and perspectives

The $B^0 \rightarrow D^\mp \pi^\pm$ analysis presented in this thesis was performed on the full LHCb Run 1 (2011–2012) dataset, corresponding to an integrated luminosity of 3 fb^{-1} .

As can be seen from Eqs. 6.1 and 6.2, the dominant contribution in the error budget of S_f and $S_{\bar{f}}$ is the statistical uncertainty. For this reason, the precision of the measurement can be easily improved by using a larger data sample.

The statistics expected to be collected during Run 2 (2015–2018) is $\sim 6 \text{ fb}^{-1}$, with a luminosity of $4 \times 10^{32} \text{ cm}^{-2} \text{ s}^{-1}$. The centre-of-mass energy during Run 2 (13 TeV) is

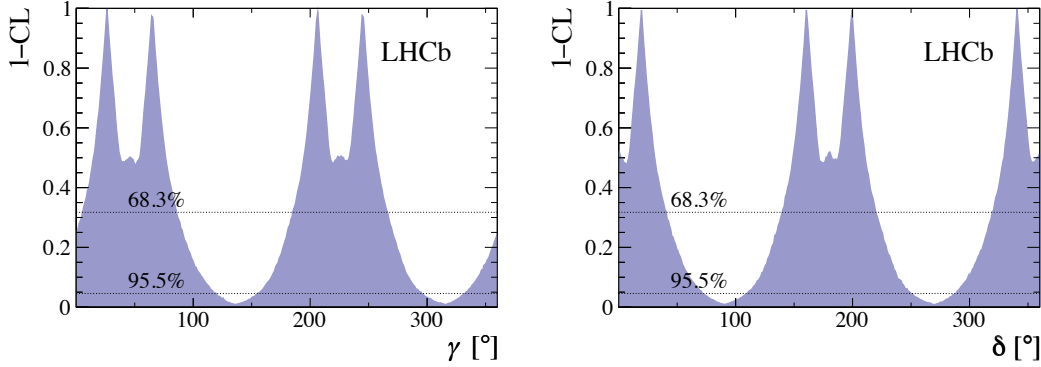


Figure 6.1 – p -value, or $1-\text{CL}$, as a function of γ (left) and δ (right) obtained using the measured values of S_f and $S_{\bar{f}}$.

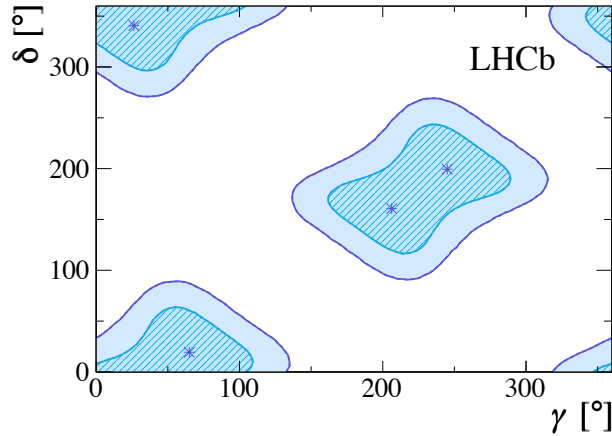


Figure 6.2 – Contours in the two-dimensional plane (γ, δ) obtained using the measured values of S_f and $S_{\bar{f}}$. The crosses indicate the preferred values ($1-\text{CL}=1$). The blue hatched (solid) areas correspond to 39% (87%) confidence level.

about twice the Run 1 value; for this reason, the B^0 production cross-section is also increased approximately by a factor two compared to Run 1. So, the increase in the number of reconstructed $B^0 \rightarrow D^\mp \pi^\pm$ decays in Run 2 compared to Run 1 is about $\sim 2 \times \frac{6}{3} = 4$, which corresponds to a factor $\sim \sqrt{4} = 2$ of decrease in statistical uncertainty.

After Run 2, a two-year shutdown (2019–2020) will allow a major upgrade of several LHCb detectors to take place. Particularly relevant for the $B^0 \rightarrow D^\mp \pi^\pm$ analysis are the VELO upgrade [92], which will imply an improvement of hit efficiency and IP resolution thanks to the pixel geometry, the Upstream Tracker [93] (replacing the TT), which will improve the particle acceptance, and the scintillating fibre tracker [93] (replacing IT and OT), which is designed to cope with the expected higher occupancy due to the higher luminosity. Moreover, the L0 trigger will be removed, and a 40 MHz readout electronics

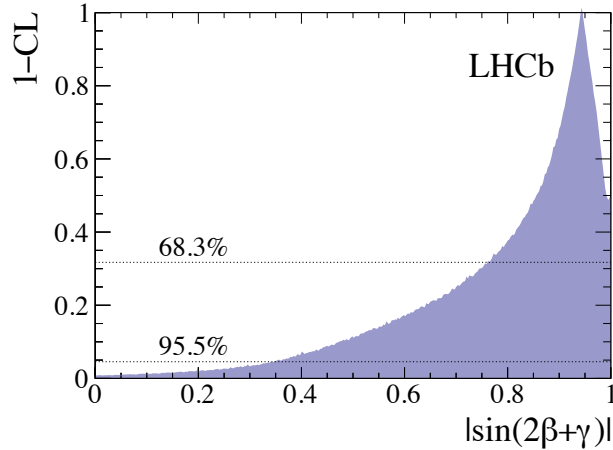


Figure 6.3 – p -value, or 1-CL, as a function of $|\sin(2\beta + \gamma)|$ using the measured values of S_f and $S_{\bar{f}}$.

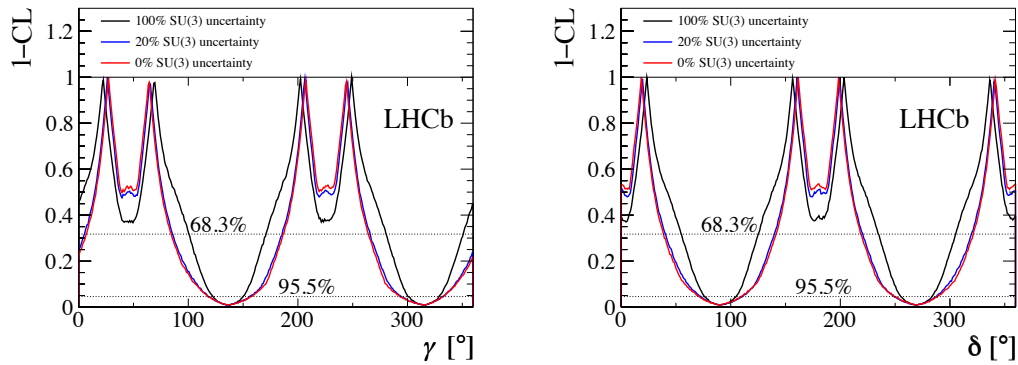


Figure 6.4 – p -value, or 1-CL, as a function of (left) γ and (right) δ for assumptions of 0%, 20% and 100% for the SU(3) uncertainty on the parameter $r_{D\pi}$.

will be installed; a fully software-implemented trigger will thus be adopted. According to Ref. [94], the resulting trigger efficiencies for fully hadronic decays as $B^0 \rightarrow D^\mp \pi^\pm$ is expected to increase by a factor ~ 2 compared to Run 1. The LHC will restart in 2021 with an increased luminosity of $2 \times 10^{33} \text{ cm}^{-2}\text{s}^{-1}$, and pp collisions at a centre-of-mass energy of 14 TeV will be delivered during Run 3 (2021–2023) and Run 4 (2026–2029). At the end of Run 4, the total amount of data collected with the upgraded detector is expected to reach $\sim 50 \text{ fb}^{-1}$. The expected increase in the number of reconstructed $B^0 \rightarrow D^\mp \pi^\pm$ decays in Runs 3 and 4 is about $\sim 2 \times 2 \times \frac{50}{3} \sim 70$, which corresponds to a factor $\sim \sqrt{70} \sim 8$ of decrease in statistical uncertainty compared to the Run 1 result.

Finally, a second upgrade of LHCb is under discussion, which would allow the experiment to operate at a luminosity of $1 - 2 \times 10^{34} \text{ cm}^{-2}\text{s}^{-1}$ starting in 2031 [95]. The expected pp collision statistics that will be collected in this high-luminosity scenario corresponds to

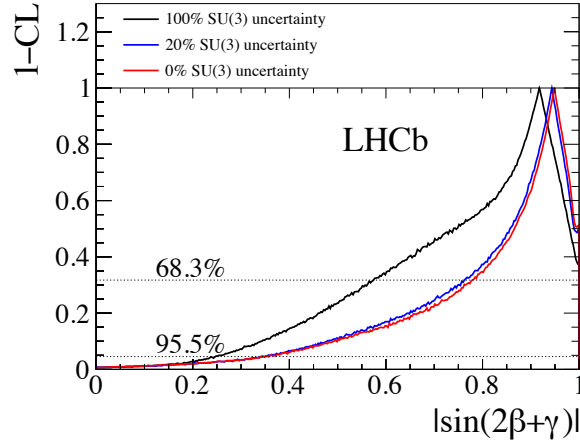


Figure 6.5 – $1-\text{CL}$ as a function of $|\sin(2\beta + \gamma)|$ for assumptions of 0%, 20% and 100% for the $\text{SU}(3)$ uncertainty on the parameter $r_{D\pi}$.

$\sim 300 \text{ fb}^{-1}$. Given the same centre-of-mass energy of 14 TeV, the increase in the number of $B^0 \rightarrow D^\mp \pi^\pm$ decays will be about $\sim 2 \times 2 \times \frac{300}{3} \sim 400$, corresponding to a reduction of the Run 1 statistical uncertainty by a factor $\sim \sqrt{400} \sim 20$.

These extrapolations are made by assuming the same tagging power as obtained on Run 1 data; future developments of flavour tagging algorithms are thus crucial to further improve these projections. As described in Sec. 3.3.4, the performance of OS taggers on Run 2 data is compatible with, and not worse than, the one obtained on Run 1 data thanks to the reoptimisation campaign on Run 2 data, while the tagging power of the $\text{SS}\pi$ and $\text{SS}p$ taggers is increased. Preliminary studies on simulated samples showed that both SS taggers and Run 2-optimised OS taggers have similar performances in Run 3 as the ones of Run 2; further improvements can be achieved by tuning these taggers specifically on Run 3 (and beyond) conditions. In parallel, a new approach, called *inclusive tagger*, is under development. This algorithm, consisting of a deep neural network, exploits all tracks and vertices reconstructed in the events in order to provide a tagging decision and a mistag estimate. Preliminary results on Run 1 $B^+ \rightarrow J/\psi K^+$ data indicate a tagging power of the order of $\sim 7 - 8\%$, which would represent a relative increase of $\sim 40 - 60\%$ compared to the combination of the standard OS and SS taggers.

Concerning the systematic uncertainties on S_f and $S_{\bar{f}}$, the external constraint on Δm , which is the dominant contribution to the systematic error budget, will also benefit from the high statistics collected by LHCb, since the world-leading measurement of this parameter is already obtained by the LHCb collaboration from semileptonic B^0 decays [83], and this result will be updated with new data. In addition to Δm , other sources of systematic uncertainty ($\Delta\Gamma$, decay-time resolution ...) are expected to be reduced thanks to the increased statistics foreseen for the next decades.

The precision on the values of γ and δ extracted from the measured values for S_f and $S_{\bar{f}}$ will benefit from the increased knowledge on β and $r_{D\pi}$. The β angle will be measured with unprecedented precision also thanks to the Belle II experiment [96, 97], which will start its operations between 2018 and 2019. The precision on the $r_{D\pi}$ parameter will increase thanks to the improvement in the measurements of the $B^0 \rightarrow D_s^+ \pi^-$ and $B^0 \rightarrow D^- \pi^+$ branching ratios, according to Eq. 6.5. Moreover, if the $B_s^0 \rightarrow D^+ K^-$ decay will be observed, an independent estimation of $r_{D\pi}$ will be available by following the relation

$$r_{D\pi} = \frac{f_\pi}{f_K} \sqrt{\frac{\mathcal{B}(B_s^0 \rightarrow D^+ K^-)}{\mathcal{B}(B^0 \rightarrow D^- \pi^+)}} \quad (6.6)$$

where SU(3) symmetry is assumed as for Eq. 6.5 and $\frac{f_\pi}{f_K} = 1.1956(10)_{-18}^{+26}$ [89]. A search for $B_s^0 \rightarrow D^+ K^-$ with Run 1 data gave a null result [98].

The value of γ that will be extracted with future time-dependent analyses of $B^0 \rightarrow D^\mp \pi^\pm$ decays, and similarly of $B_s^0 \rightarrow D_s^\mp K^\pm$ decays, will be one of the inputs for the global combination of direct measurements performed by LHCb. The current sensitivity, which is of the order of $\sim 5-6^\circ$ [87], is expected to go down at the level of $\sim 0.4^\circ$ with $\sim 300 \text{ fb}^{-1}$ of collected data [99]. This sensitivity will allow to test the expected relationship between γ and the CP coefficients given by Eq. 6.3, which is affected by small theoretical uncertainties, and to compare these direct measurements of γ with the indirect determination from other CKM parameters: a discrepancy in any of these tests will be a clear signature of new physics beyond the SM.



Appendices

A Opposite-side tagging studies

A.1 Mass fit of $B^+ \rightarrow \bar{D}^0\pi^+$

A fit to the mass distribution of B^+ candidates is done to calculate $sWeights$, used in the subsequent steps of the analysis to subtract the backgrounds surviving the selection. A two-step procedure similar to that adopted for the $B^0 \rightarrow D^\mp\pi^\pm$ analysis (“Fit A” in a wide mass window to account for all backgrounds, and “Fit B” in a subset to calculate the weights, as described in Sec. 4.3, where the same mass windows are adopted) is used for the fit of the $B^+ \rightarrow \bar{D}^0\pi^+$ candidates. The projection of the total PDF on the π sample and the K sample (“Fit A”) is shown in Fig. A.1, as well as the projection of the total PDF in the reduced mass interval (“Fit B”). The π sample and K sample are defined by the PID requirement on the companion track, $PIDK < 5$ and $PIDK > 5$, respectively.

The background components expected in the π sample for the $B^+ \rightarrow \bar{D}^0\pi^+$ mass fit are listed below, together with the PDF used for each component:

- $B^+ \rightarrow \bar{D}^0\pi^+$: double-sided Hypatia function.
- $B^+ \rightarrow \bar{D}^0K^+$: double-sided Hypatia function.
- $B^0 \rightarrow \bar{D}^0\pi^+\pi^-$: Crystal ball function plus Gaussian function.
- $B^+ \rightarrow \bar{D}^{0*}\pi^+$: Johnson SU function plus Gaussian function.
- Combinatorial background: single exponential function.

The list of components expected in the K sample is the following:

- $B^+ \rightarrow \bar{D}^0\pi^+$: double-sided Hypatia function.
- $B^+ \rightarrow \bar{D}^0K^+$: single-sided Hypatia function.
- $B^+ \rightarrow \bar{D}^{0*}\pi^+$: Crystal ball function plus exponential function.
- $B^+ \rightarrow \bar{D}^0K^{+*}$: Gaussian function.
- Combinatorial background: single exponential function.

All the PDFs listed above are defined in Appendix F.

The values for the fitted parameters floated in the fit are reported in Table A.1 for Fit A and Table A.2 for Fit B. The naming convention for each parameter is similar to the one used in Sec. 4.3.2.

A.2 Reweighting of $B^+ \rightarrow \bar{D}^0\pi^+$ to $B^0 \rightarrow D^\mp\pi^\pm$

In order to improve the OS calibration portability, a multi-dimensional reweighting of the $sWeighted$ $B^+ \rightarrow \bar{D}^0\pi^+$ distributions is made to match the $B^0 \rightarrow D^\mp\pi^\pm$ kinematics.

The reweighting is made in two steps. In the first step, the variables considered in the reweighting are the transverse momentum, the pseudo-rapidity η and the decay time τ_B of

A. Opposite-side tagging studies

Table A.1 – Results of the $B^+ \rightarrow \bar{D}^0 \pi^+$ mass fit (Fit A).

Parameter name	Fitted value
$\mu_{B^0 \rightarrow D\pi\pi}^\pi$	5132.61 ± 0.23
$s\sigma_{B^0 \rightarrow D\pi\pi}^\pi$	0.780 ± 0.015
$\sigma_{B^+ \rightarrow DK}^K$	19.47 ± 0.30
$\sigma_{B^+ \rightarrow DK}^\pi$	16.62 ± 0.69
$\mu_{B^+ \rightarrow DK^*}^K$	4960 ± 150
$\sigma_{B^+ \rightarrow DK^*}^K$	88 ± 37
c_{comb}^K	-0.001305 ± 0.000035
$c_{\text{comb}}^{1\pi}$	-0.001279 ± 0.000022
$\mu_{B^+ \rightarrow D\pi}^K$	5283.18 ± 0.22
$\mu_{B^+ \rightarrow D\pi}^\pi$	5283.880 ± 0.046
$sa_{B^+ \rightarrow D\pi}^\pi$	0.804 ± 0.016
$\mu_{B^+ \rightarrow D\pi}^K$	5325.4 ± 1.2
$sn_{B^+ \rightarrow D\pi}^\pi$	2.70 ± 0.94
$\sigma_{B^+ \rightarrow D\pi}^\pi$	22.850 ± 0.054
$N_{B^0 \rightarrow D\pi\pi}^\pi$	27245 ± 430
$N_{B^+ \rightarrow DK}^K$	18030 ± 296
N_{LM}^K	5154 ± 944
$N_{B^+ \rightarrow D^*\pi}^K$	5704 ± 1350
$N_{B^+ \rightarrow D^*\pi}^\pi$	41871 ± 578
N_{comb}^K	58761 ± 555
N_{comb}^π	146824 ± 793
$N_{B^+ \rightarrow D\pi}^\pi$	322597 ± 812

Table A.2 – Results of the $B^+ \rightarrow \bar{D}^0 \pi^+$ mass fit (Fit B).

Parameter	Fitted value
N_{bkg}^π	$85\,687 \pm 377$
$N_{B^+ \rightarrow D\pi}^\pi$	$319\,974 \pm 612$

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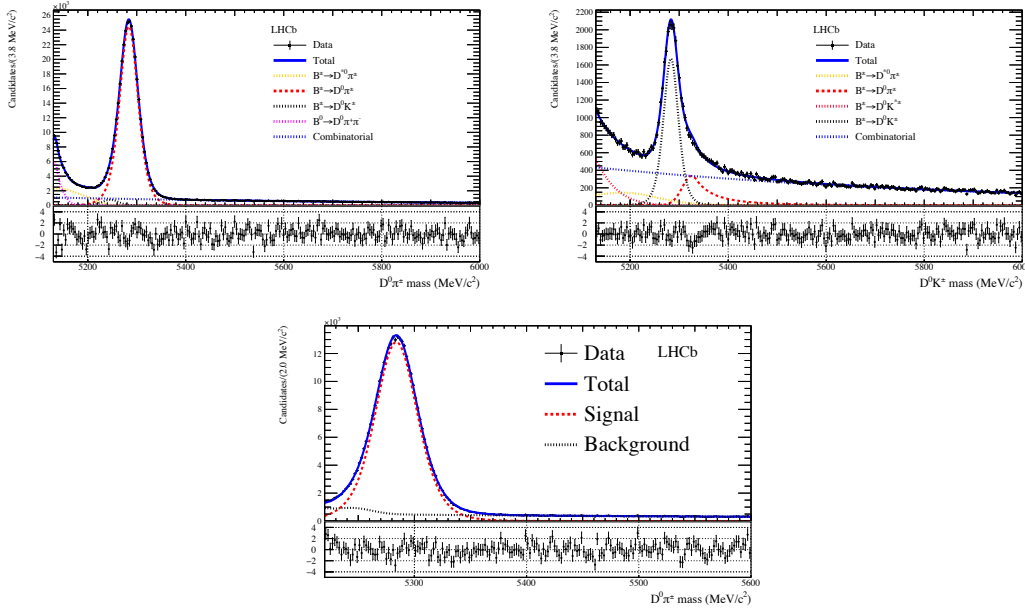


Figure A.1 – Top left: $\bar{D}^0\pi^+$ mass distribution of the π sample. Top right: $\bar{D}^0\pi^+$ mass distribution of the K sample. The result of the simultaneous fit (Fit A) to both samples is superimposed. Bottom: $\bar{D}^0\pi^+$ mass distribution of the π sample with the results of Fit B superimposed.

the B candidate, as well as the number of tracks and the number of primary vertices of the events. A BDT-based approach is followed in order to cope with the high dimensionality of the space as well as with the correlations among variables [100]. A comparison between weighted and unweighted distributions is provided in Figs. A.2 and A.3.

In the second step, a new weight is computed by comparing the two-dimensional distributions of the D meson decay time and HLT2 trigger composition between $B^+ \rightarrow \bar{D}^0\pi^+$ and $B^0 \rightarrow D^\mp\pi^\pm$ after *sWeights* and the weights from the first step are applied. The HLT2 composition observable is a categorical variable which describes which HLT2 trigger line has been fired by the B candidate:

- `Hlt2Topo2BodyBBDTDecision` only (value 0);
- `Hlt2Topo3BodyBBDTDecision` or `Hlt2Topo4BodyBBDTDecision` only (value 1);
- overlap of the first two categories (value 2).

This reweighting is done separately from the first one in order to avoid a too fine partition of the samples, which would result in very low statistics in less populated bins. The reason why HLT2 trigger and D decay time are reweighted simultaneously is that these two observables are correlated. The result of this second reweighting is shown in Fig. A.4.

A.3 GOF tests for OS calibration on $B^+ \rightarrow \bar{D}^0\pi^+$ data

The number of free parameters (10) used in the GLM model for the OS calibration (Sec. 3.2.1) is the minimum number to obtain satisfactory GOF metrics. The GOF tests are performed automatically by the EPM; the metrics include the Pearson χ^2 , the deviance G^2 , the Cressie-Read (CR) metric and the Le Cessie-van Houwelingen-Copas-Hosmer metric (S), all described in Ref. [61].

All these tests return a normally distributed score: this means that the score is equal to the distance (measured in standard deviations) from the perfect case, which is a null score. A comparison between the GOF scores obtained for the nominal calibration (10 free parameters) and a simplified model (8 free parameters) is shown in Table A.3. In a simplified model, all scores are more than ~ 3 standard deviations away from a perfect fit, whereas the scores for the nominal model are ~ 2 standard deviations at most. For this reason, a calibration with 10 free parameters are chosen, and the model cannot be simplified further.

Table A.3 – GOF scores of two OS calibration fits of the reweighted $B^+ \rightarrow \bar{D}^0\pi^+$ dataset.

GOF metric	Score (10 parameters)	Score (8 parameters)
χ^2	-2.2	4.1
G^2	0.7	-3.9
CR	-1.7	2.9
S	1.8	-4.3

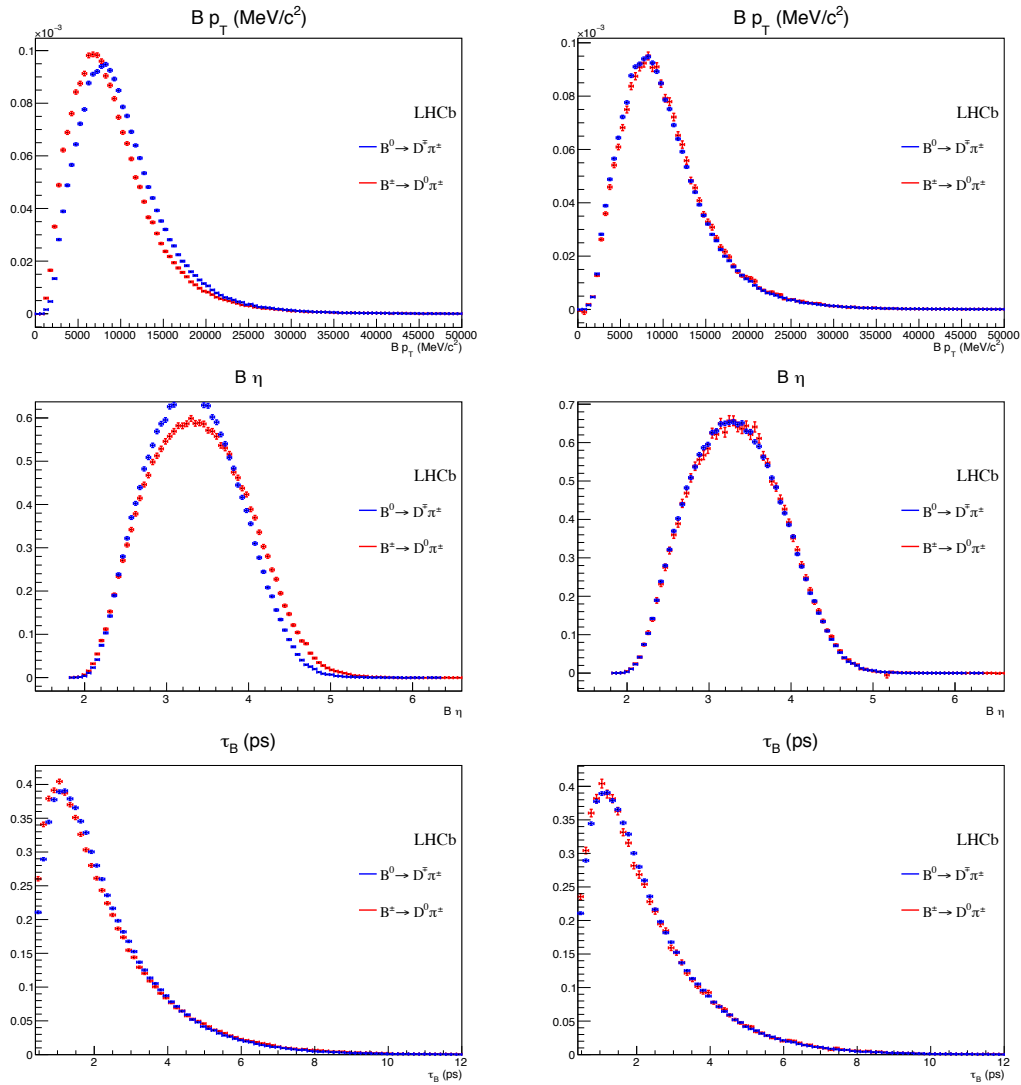


Figure A.2 – Normalised *sWeighted* distributions of the transverse momentum, the pseudo-rapidity η and the decay time τ_B of the B^0 and B^+ mesons. Left: unweighted distributions. Right: distributions after reweighting the $B^+ \rightarrow \bar{D}^0 \pi^+$ events.

A. Opposite-side tagging studies

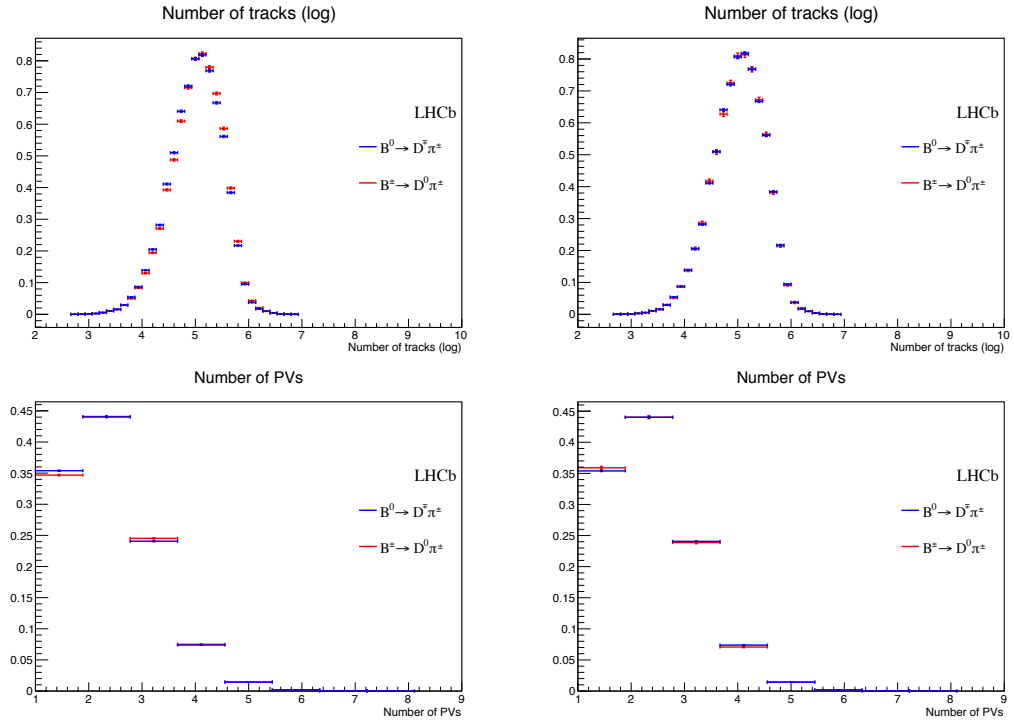


Figure A.3 – Normalised s -Weighted distributions of the number of tracks and PVs in a B^0 or B^+ event. Left: unweighted distribution. Right: distributions after reweighting the $B^+ \rightarrow \bar{D}^0 \pi^+$ events.

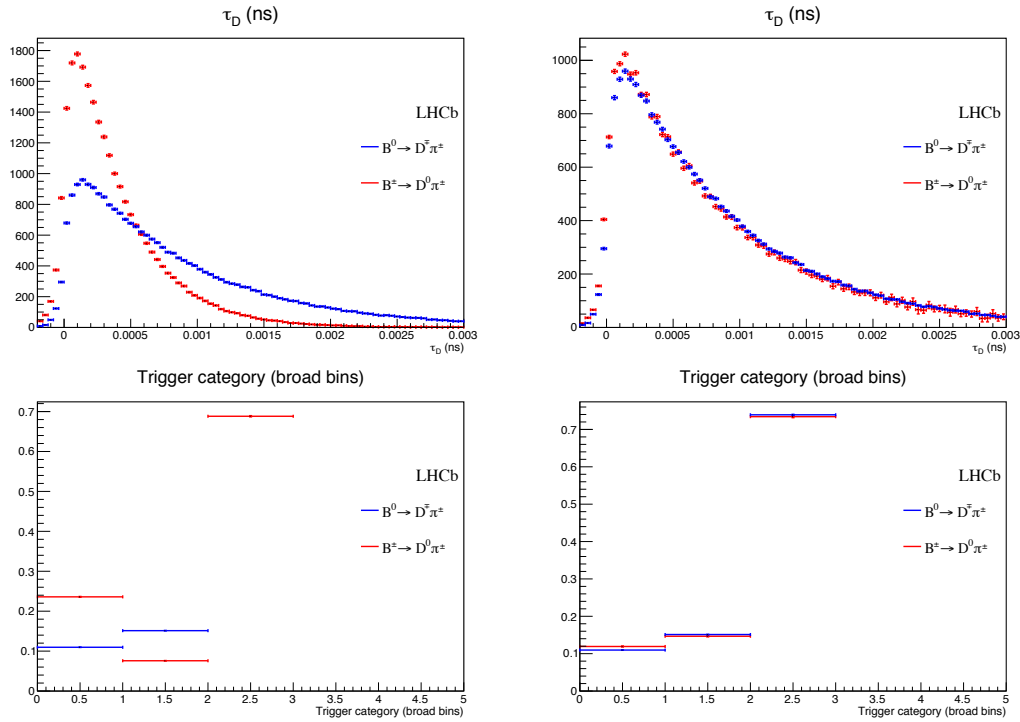


Figure A.4 – Normalised *sWeighted* distributions of the D^+ and \bar{D}^0 mesons decay time and HLT2 trigger composition, where the weight obtained from the first reweighting step is also applied. Left: unweighted distributions. Right: distributions after reweighting the $B^+ \rightarrow \bar{D}^0 \pi^+$ events.

B Opposite-side electron optimisation

The correlation between the predicted mistag η and the input features is shown in Figs. B.1, B.2 and B.3 for the Run 1 new, Run 2 B2CC and Run 2 B2OC implementations of the OSe tagger, respectively. This correlation, known as *partial dependence*, allows to check the impact of each feature on the classifier, in addition to the F-score (Sec. 3.3.3).

An important step in the BDT development is the hyperparameter tuning. In particular, the maximum depth (md) of each tree of the ensemble and the number of trees (nt) are optimised. In order to do so, a cross-validation+bootstrapping procedure is followed:

- For a given set of maximum depth and number of trees values, the training set is bootstrapped 10 times. Each bootstrapped sample is then divided in three exclusive subsamples.
- The first subsample is used to train a BDT. The BDT is then transformed into a mistag probability, and a calibration is performed on the second subsample (a simple second order logistic function is used). Finally, the calibration is applied on the third sample, where the per-event tagging power is computed. The ROC AUC is also obtained as additional performance metric.
- The above procedure is repeated by permutating the 3 samples. This means that, in total, there are $3 \times 10 = 30$ approximately independent estimations of the BDT performance for each set of hyperparameters. The average tagging power and ROC AUC values are finally computed over the 30 estimations, together with the standard error on the mean.

The result for the Run 1 new OSe algorithm is shown in Fig. B.4. The performance is weakly dependent on the hyperparameters. For this reason, the maximum depth and the number of trees are fixed to 3 and 300 respectively, in order to reduce complexity. A similar result is observed for the Run 2 B2OC and Run 2 B2CC OSe algorithms.

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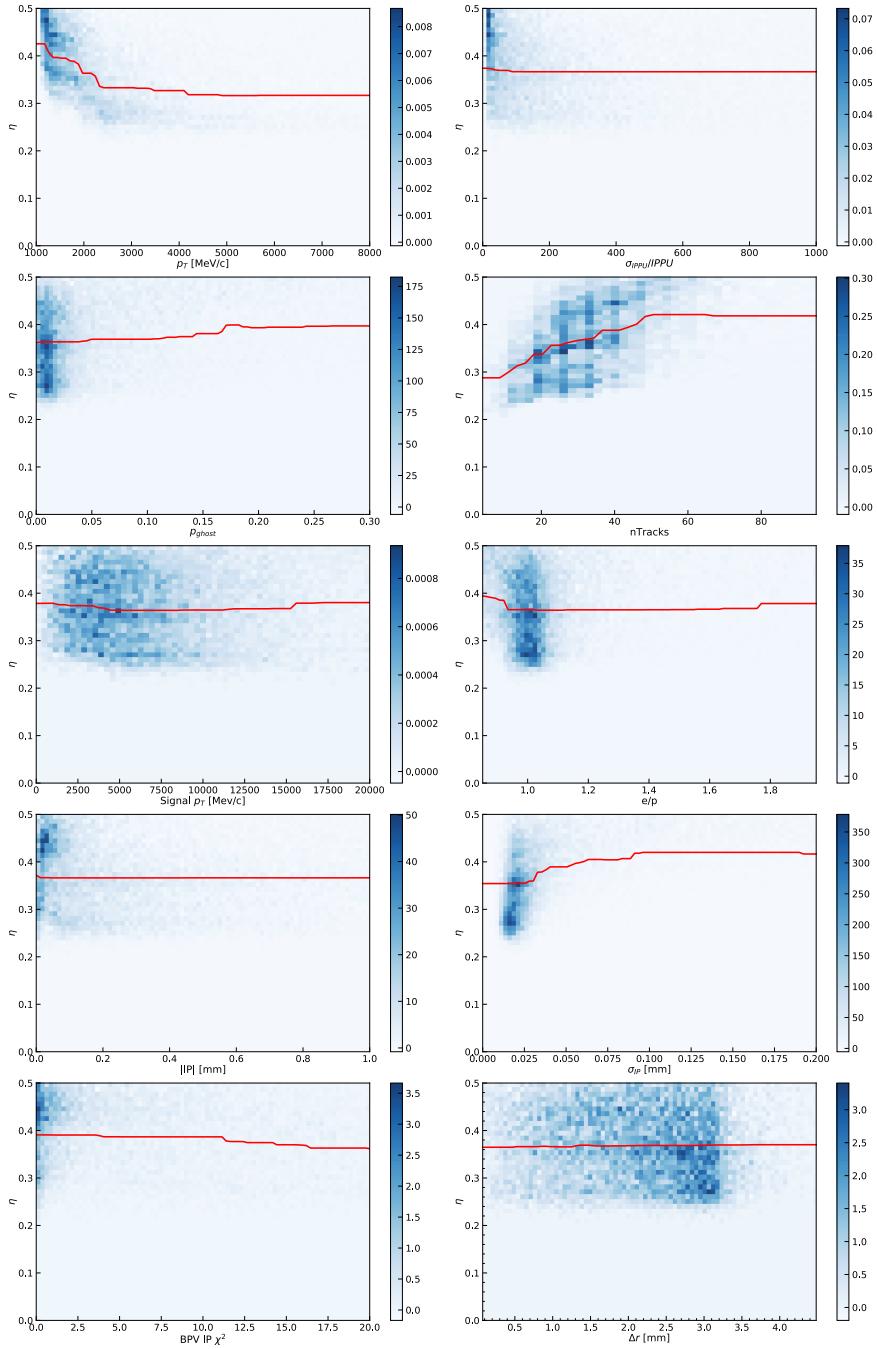


Figure B.1 – Partial dependence of the predicted mistag η (OSe Run 1 tagger) for each feature used as BDT input, marginalised over any other feature. The blue two-dimensional distributions represent the *sWeighted* data, whereas the red line shows the average η for each feature bin.

B. Opposite-side electron optimisation

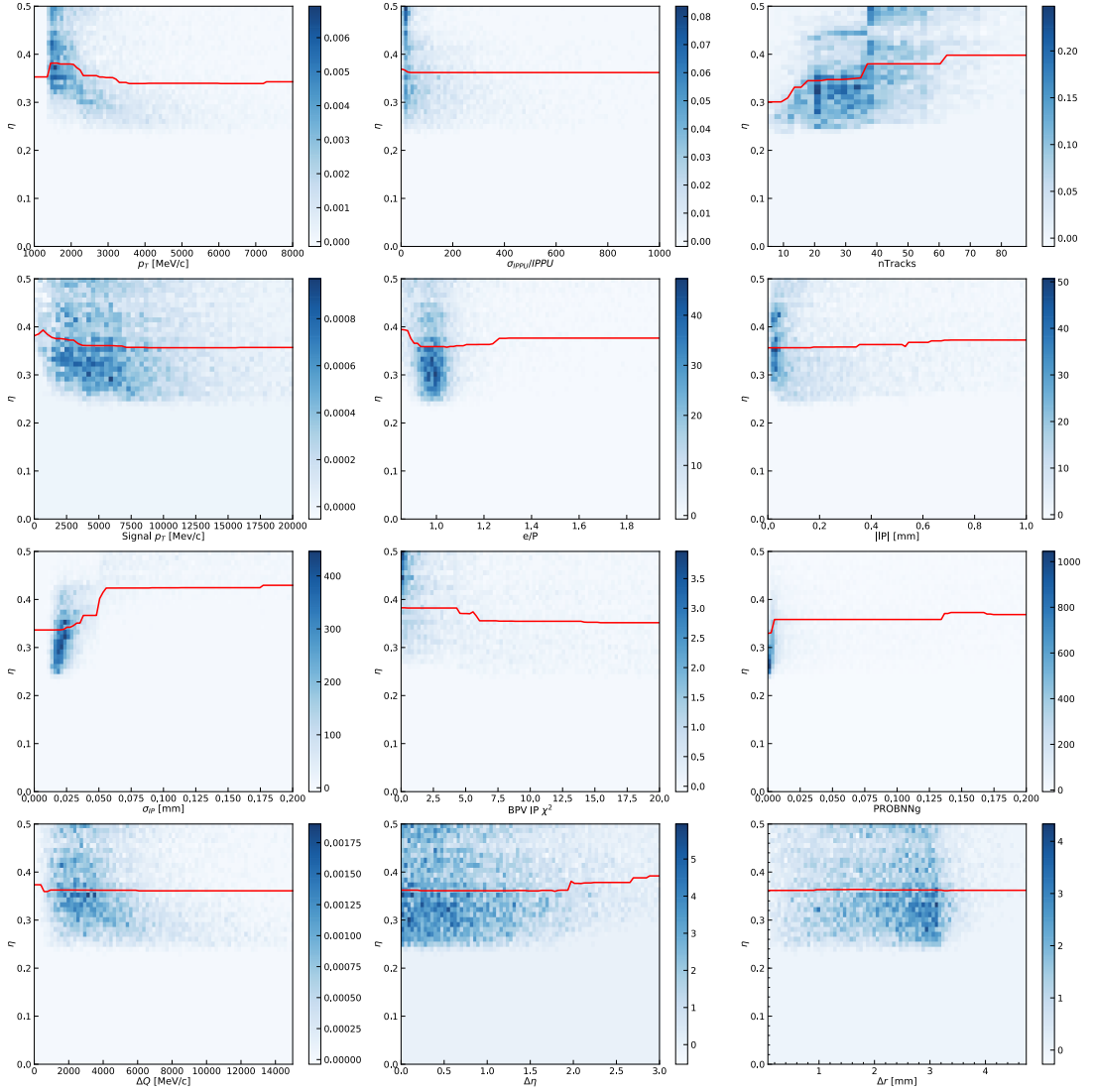


Figure B.2 – Partial dependence of the predicted mistag η (OSe Run 2 B2CC tagger) for each feature used as BDT input, marginalised over any other feature. The blue two-dimensional distributions represent the *sWeighted* data, whereas the red line shows the average η for each feature bin.

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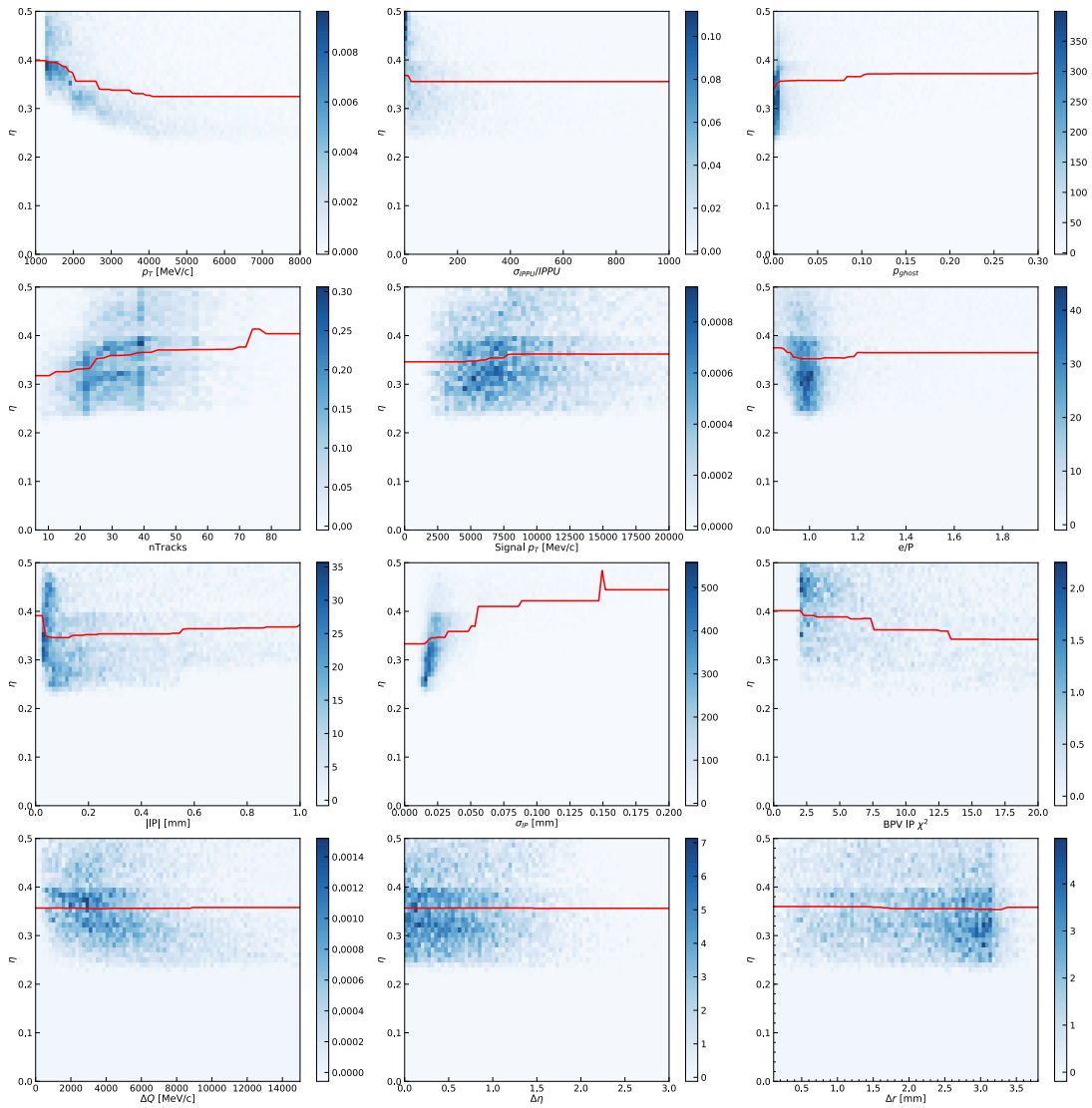


Figure B.3 – Partial dependence of the predicted mistag η (OSe Run 2 B2OC tagger) for each feature used as BDT input, marginalised over any other feature. The blue two-dimensional distributions represent the *sWeighted* data, whereas the red line shows the average η for each feature bin.

B. Opposite-side electron optimisation

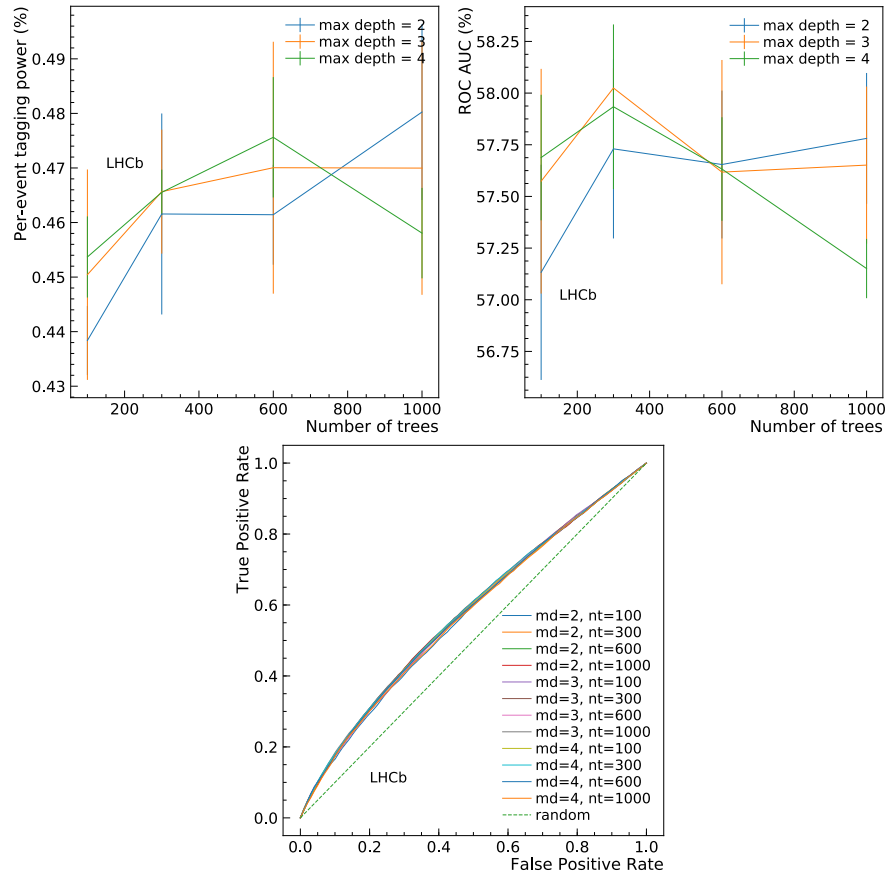


Figure B.4 – Tagging power (top left), ROC AUC (top right), and ROC curve (bottom) for each set of hyperparameters considered for the tuning of the BDT classifier used for the mistag estimation by the Run 1 new version of the OSe tagger (cross-validation).

C $B^0 \rightarrow D^\mp \pi^\pm$ selection studies

C.1 BDT input features

The distributions of the input features for the BDT (listed and defined in Table 4.5) are shown in Figs. C.1, C.2 and C.3.

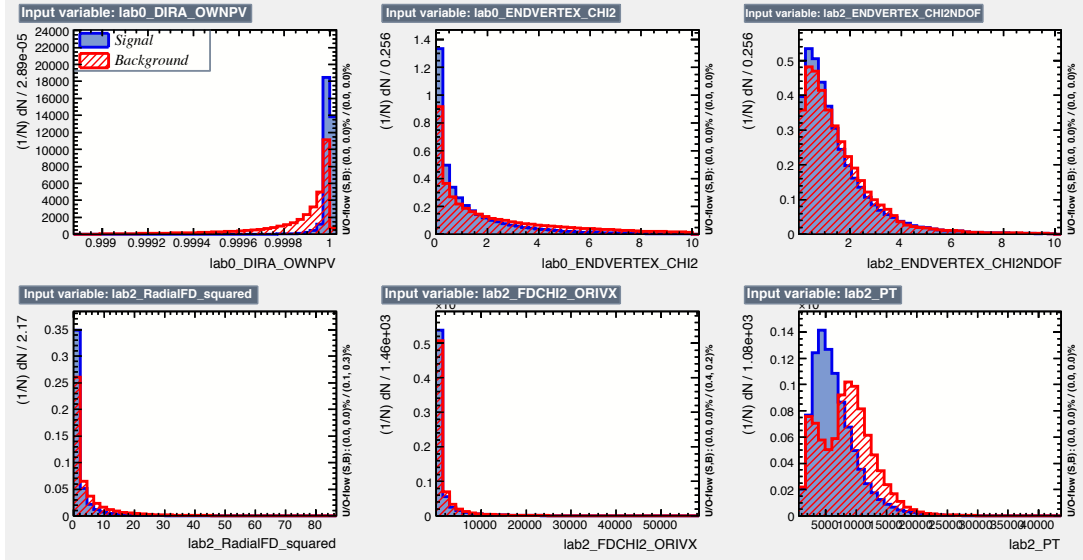


Figure C.1 – Input features used in the BDT training. From top left to bottom right: cosine of the direction angle of the B^0 , χ^2 of the B^0 vertex, χ^2/ndof of the D^\mp vertex, D^\mp radial flight distance, D^\mp flight distance χ^2 with respect to the B^0 vertex and transverse momentum of the D^\mp .

C.2 Multiple candidates

Table C.1 gives a summary of the multiple candidates left after stripping and trigger selection, while Table C.2 reports the number of multiple candidates after stripping, trigger and offline selection.

C. $B^0 \rightarrow D^\mp \pi^\pm$ selection studies

Table C.1 – Statistical information on multiple $B^0 \rightarrow D^\mp \pi^\pm$ candidates left after stripping and trigger selection.

	2011	2012		
fraction of candidates that are not unique in a given event	18.3 %	19.5 %		
fraction of candidates to be discarded to maintain one candidate per event	10.1 %	11.0 %		
fraction of events with multiple candidates	9.0 %	9.6 %		
	#cands	#events	#cands	#events
	1	5940804	1	16407228
	2	483991	2	1426286
	3	73902	3	226205
	4	20093	4	62640
	5	6132	5	19213
	6	2505	6	8044
	7	1087	7	3326
	8	528	8	1686
	9	251	9	839
	10	146	10	461
	11	78	11	279
	12	40	12	178
	13	28	13	109
	14	32	14	85
	15	10	15	53
	16	12	16	24
	17	7	17	16
	18	4	18	20
	19	5	19	9
	20	1	20	11
	21	3	21	5
	22	2	22	3
	23	0	23	2
	24	2	24	2
	25	1	25	1
	26	0	26	4
	30	0	30	1
	33	0	33	1
	40	1	40	0
	41	0	41	1

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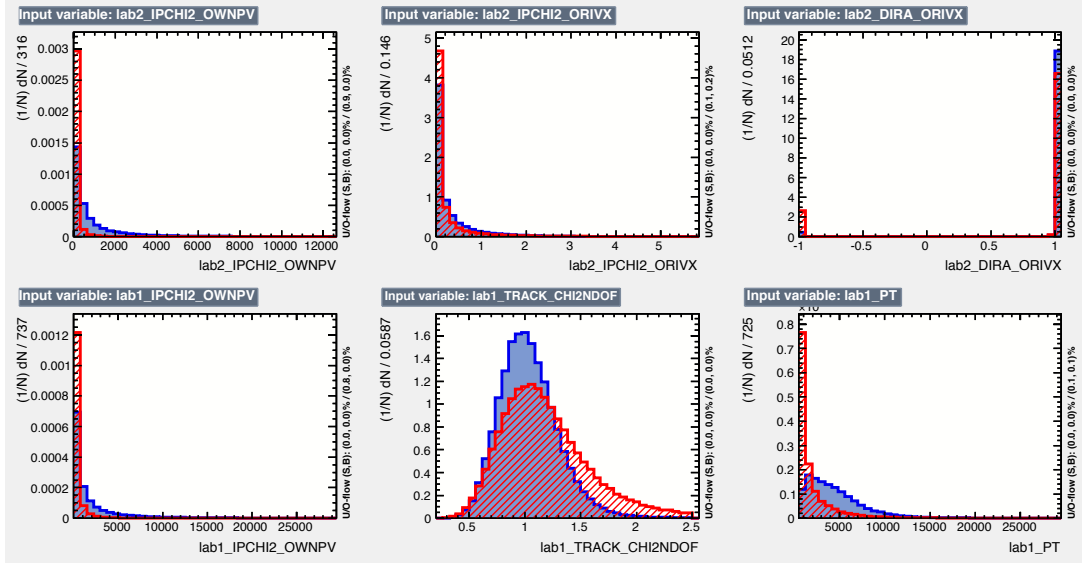


Figure C.2 – Input features used in the BDT training. From top left to bottom right: D^\mp IP χ^2 with respect to the associated PV and the B^0 vertex, cosine of the direction angle of the D^\mp , the IP χ^2 with respect to the associated PV of the bachelor pion, track χ^2/ndof of the bachelor pion and the transverse momentum of the bachelor pion.

Table C.2 – Statistical information on multiple $B^0 \rightarrow D^\mp \pi^\pm$ candidates left after stripping, trigger and offline selection.

	2011	2012		
fraction of candidates pairs that are not unique in an event	0.8 %	0.8 %		
fraction of candidates to be discarded to maintain one candidate per event	0.4 %	0.4 %		
fraction of events with multiple candidates	0.4 %	0.4 %		
	#cands	#events	#cands	#events
	1	483074	1	1200956
	2	1886	2	4962
	3	38	3	98
	4	4	4	9
	5	1	5	3

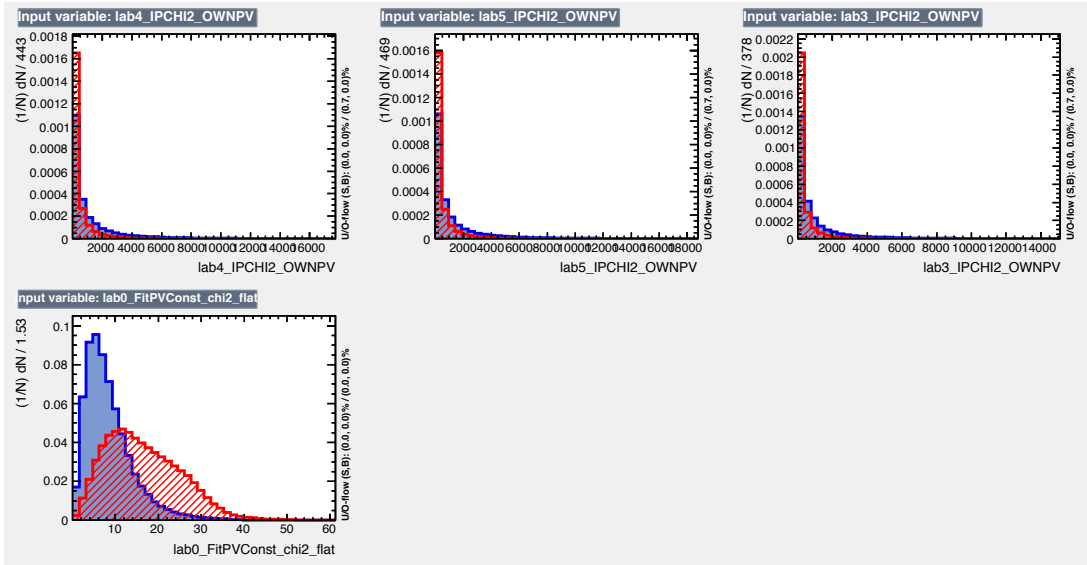


Figure C.3 – Input features used in the BDT training. From top left to bottom right: IP χ^2 of the associated primary vertex of the three D^\mp daughters and the χ^2 of the decay tree fit with PV constraint.

D Particle identification plots

The p and η distributions, as well as the PIDK efficiency and misidentification rates before the resampling in bins of p and η are shown in Figs. D.1 and D.2 for the bachelor particle of the signal, and in Fig. D.3 for the D^\mp daughters.

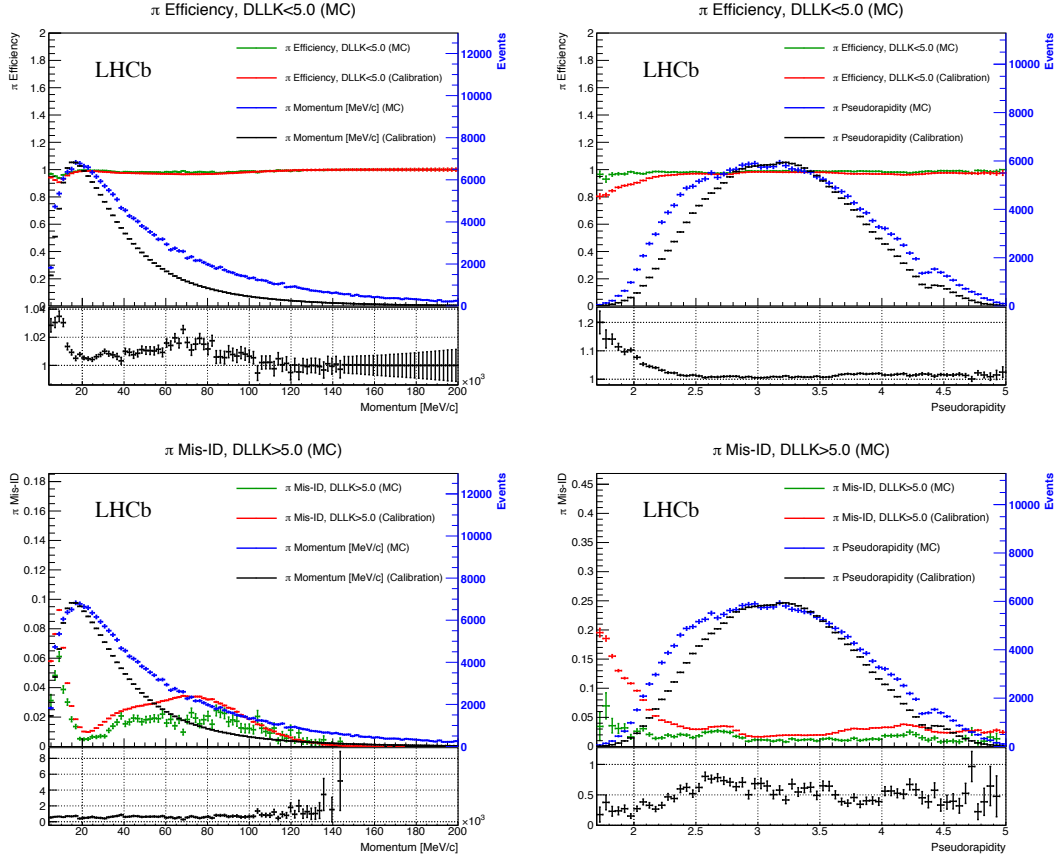


Figure D.1 – Efficiencies of the requirements $PIDK < 5$ (top) and $PIDK > 5$ (bottom) for bachelor pions as a function of momentum p (left) and pseudorapidity η (right), both for $B^0 \rightarrow D^\mp \pi^\pm$ signal MC (green) and calibration mode (red). The superimposed histograms show the p and η distributions of the MC signal (blue) and calibration (black) samples. The ratio of the efficiency or misidentification rate between the MC signal and data calibration samples is shown in the lower pad (black).

D. Particle identification plots

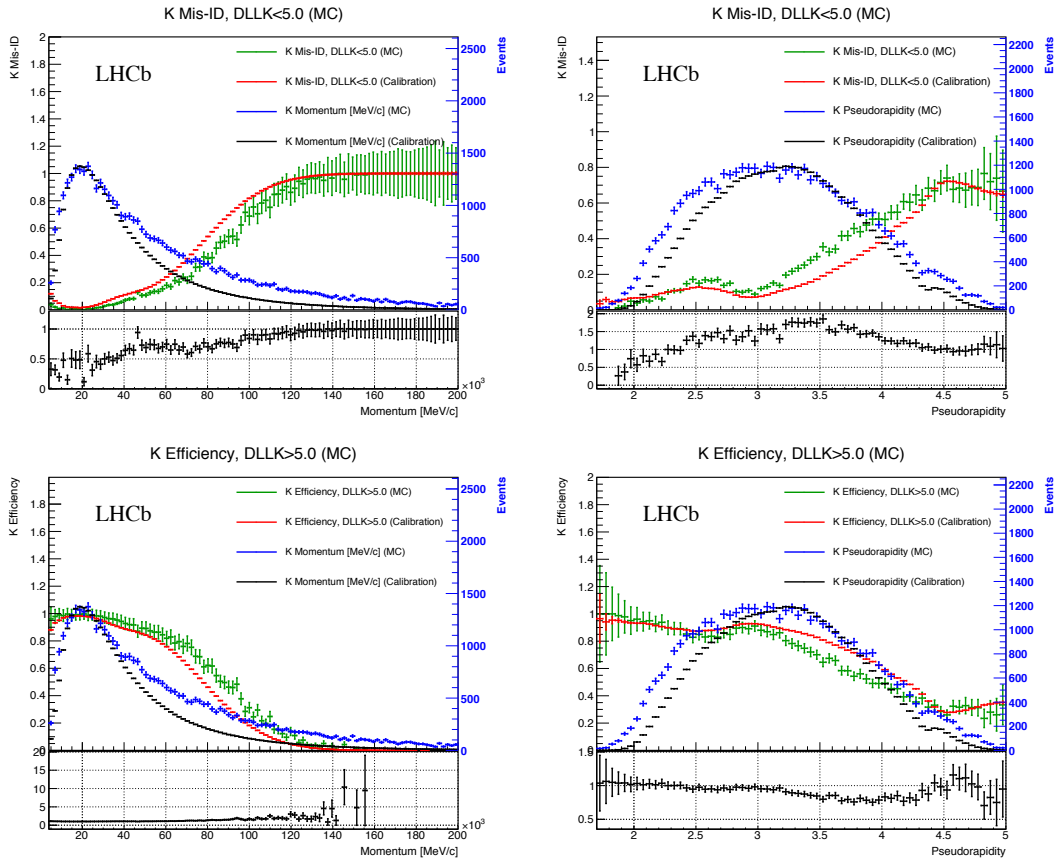


Figure D.2 – Efficiencies of the requirements $PIDK < 5$ (top) and $PIDK > 5$ (bottom) for bachelor kaons as a function of momentum p (left) and pseudorapidity η (right), both for $B^0 \rightarrow D^\mp K^\pm$ MC (green) and calibration mode (red). The superimposed histograms show the p and η distributions of the $B^0 \rightarrow D^\mp K^\pm$ MC (blue) and calibration (black) samples. The ratio of the efficiency or misidentification rate between the MC signal and data calibration samples is shown in the lower pad (black).

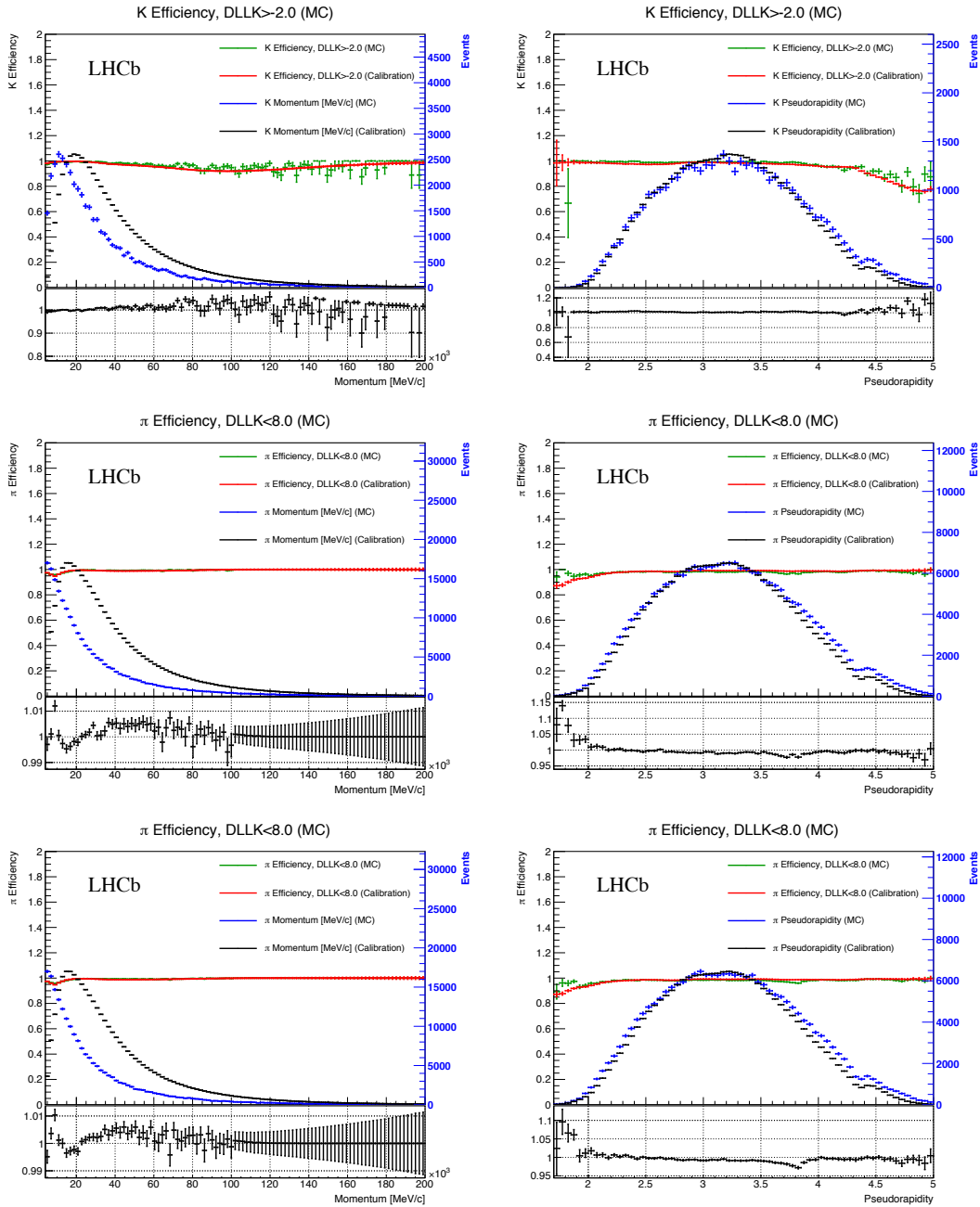


Figure D.3 – PIDK efficiencies for the kaon (top) and the two pions (middle, bottom) produced in the D^{\mp} decay as a function of momentum p (left) and pseudorapidity η (right), both for $B^0 \rightarrow D^{\mp} \pi^{\pm}$ signal MC (green) and calibration mode (red). The superimposed histograms show the p and η distributions of the MC signal (blue) and calibration (black) samples. The ratio of the efficiency between the MC signal and data calibration samples is shown in the lower pad (black).

E Correlation between B^0 mass and decay time

The small correlation between the B^0 invariant mass and decay time is shown by comparing the distribution of the decay time in bins of the invariant mass after applying the full selection. This is done separately for signal and background. For the signal distribution simulated data is used and the decay time is shown in six bins of the invariant mass. (Fig. E.1). In order to account for the combinatorial background, the upper mass sideband is chosen as a proxy. Figure E.2 shows the decay time in four bins of the invariant mass. The physics background contribution in the signal region is considered to be small enough, so that even a large correlation does not matter. Given the small differences for all distributions, the correlations between decay time and invariant mass is assumed to be small enough to justify the use of the invariant mass in the *sPlot* [75] technique for disentangling signal from background.

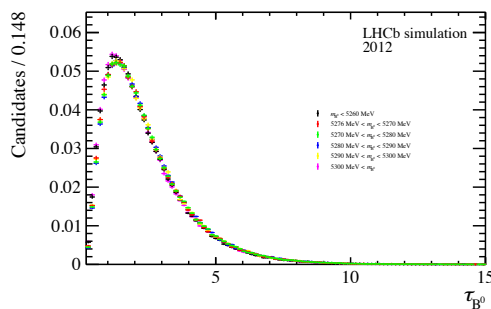


Figure E.1 – Normalised MC signal decay time distributions in six bins of the reconstructed invariant mass.

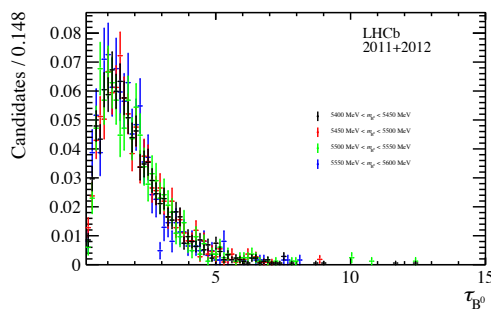


Figure E.2 – Upper mass sideband decay time distribution in four bins of the invariant mass. The shapes are shown normalised.

F PDF definitions

Throughout this section, each parameter p defined inside a PDF used for the mass fit is labelled as p_c^s , where $s = \pi, K$ indicates the sample and $c = B^0 \rightarrow D^{\mp} \pi^{\pm}, B^0 \rightarrow D^{\mp} \rho^{\pm} \dots$ indicates the component. For sake of clarity, the s and c labels are dropped in the equations that follow. The mass *observable* is always indicated as m . The \propto symbol indicates that all PDFs are defined up to a normalisation constant, which depends on the interval chosen for m .

- **Exponential function**

$$E(m, c) \propto e^{-cm}.$$

- **Gaussian function**

$$G(m, \mu, \sigma) \propto e^{-\frac{(m-\mu)^2}{2\sigma^2}}.$$

- **Double Gaussian function**

$$DG(m, \mu, \sigma_1, \sigma_2, f) \propto \frac{f}{\sigma_1} e^{-\frac{(m-\mu)^2}{2\sigma_1^2}} + \frac{(1-f)}{\sigma_2} e^{-\frac{(m-\mu)^2}{2\sigma_2^2}}.$$

- **Single-sided Crystal ball function**

Having defined

$$A = \left(\frac{n}{|\alpha|} \right)^n e^{-\frac{|\alpha|^2}{2}}, \quad B = \frac{n}{|\alpha|} - |\alpha|,$$

and

$$t = \begin{cases} \frac{m-\mu}{\sigma}, & \text{if } \alpha \geq 0, \\ -\frac{m-\mu}{\sigma}, & \text{if } \alpha < 0, \end{cases}$$

the single-sided Crystal Ball function [78] is expressed as follows:

$$CB(m, \mu, \sigma, \alpha, n) \propto \begin{cases} e^{-\frac{1}{2}t^2(m, \mu, \sigma)}, & \text{if } t \geq -|\alpha|, \\ A(\alpha, n) [B(\alpha, n) - t(m, \mu, \sigma)]^{-n}, & \text{if } t < -|\alpha|. \end{cases}$$

- **Double-sided Hypatia function**

Having defined

$$h(m, \mu, \sigma, \lambda, \zeta, \beta) \propto \left((m - \mu)^2 + A_\lambda^2(\zeta) \sigma^2 \right)^{\frac{1}{2}\lambda - \frac{1}{4}} e^{\beta(m-\mu)} K_{\lambda - \frac{1}{2}} \left(\zeta \sqrt{1 + \left(\frac{m - \mu}{A_\lambda(\zeta) \sigma} \right)^2} \right),$$

and its first derivative with respect to m , h' , then the double-sided Hypatia function H [76] is expressed as

$$H(m, \mu, \sigma, \lambda, \zeta, \beta, a_1, n_1, a_2, n_2) \propto \begin{cases} h(m, \mu, \sigma, \lambda, \zeta, \beta) & \text{if } \frac{m-\mu}{\sigma} > -a_1 \text{ and } \frac{m-\mu}{\sigma} < a_2, \\ \frac{h(\mu-a_1\sigma, \mu, \sigma, \lambda, \zeta, \beta)}{\left(1-m/\left(n\frac{h(\mu-a_1\sigma, \mu, \sigma, \lambda, \zeta, \beta)}{h'(\mu-a_1\sigma, \mu, \sigma, \lambda, \zeta, \beta)}-a_1\sigma\right)\right)^{n_1}} & \text{if } \frac{m-\mu}{\sigma} \leq -a_1, \\ \frac{h(\mu-a_2\sigma, \mu, \sigma, \lambda, \zeta, \beta)}{\left(1-m/\left(n\frac{h(\mu-a_2\sigma, \mu, \sigma, \lambda, \zeta, \beta)}{h'(\mu-a_2\sigma, \mu, \sigma, \lambda, \zeta, \beta)}-a_2\sigma\right)\right)^{n_2}} & \text{if } \frac{m-\mu}{\sigma} \geq a_2. \end{cases}$$

The K_λ functions are special Bessel functions of the third kind, whereas A_λ is defined as

$$A_\lambda^2 = \frac{\zeta K_\lambda(\zeta)}{K_{\lambda+1}(\zeta)}.$$

- **Single-sided Hypatia function**

A single-sided Hypatia function is obtained from a double-sided Hypatia function in the limit $a_2 \rightarrow +\infty$, $n_2 = 0$ (and by labelling a_1 and n_1 as a and n , respectively).

- **Johnson SU function**

Having defined the parameters

$$\begin{aligned} w &= e^{\tau^2}, \\ \omega &= -\nu\tau, \\ c &= \frac{1}{\sqrt{\frac{1}{2}(w-1)(w \cosh 2\omega + 1)}}, \\ z &= \frac{m - (\mu + c + \sigma\sqrt{w} \sinh \omega)}{c\sigma}, \\ r &= -\nu + \frac{\sinh^{-1} z}{\tau}, \end{aligned}$$

the Johnson SU function [77] is expressed as

$$J(m, \mu, \sigma, \nu, \tau) \propto \frac{1}{2\pi c(\nu, \tau)\sigma} e^{-\frac{1}{2}r(m, \mu, \sigma, \nu, \tau)^2} \frac{1}{\tau \sqrt{z(m, \mu, \sigma, \nu, \tau)^2 + 1}}.$$

G Signal PDF for the decay-time fit

The PDF $P(t|d_{\text{OS}}, d_{\text{SS}}, \eta_{\text{OS}}, \eta_{\text{SS}}, q)$ describing the B^0 decay time distribution is proportional, in the most general case, to

$$e^{-\frac{t}{\tau}} \left[C_{\text{cosh}}^{\text{eff}} \cosh \left(\frac{\Delta\Gamma t}{2} \right) + C_{\text{sinh}}^{\text{eff}} \sinh \left(\frac{\Delta\Gamma t}{2} \right) + C_{\text{cos}}^{\text{eff}} \cos(\Delta m t) + C_{\text{sin}}^{\text{eff}} \sin(\Delta m t) \right]. \quad (\text{G.1})$$

The four ‘‘effective’’ CP coefficients inside Eq. G.1 depend on the final state ($q = f = D^- \pi^+$ or $q = \bar{f} = D^+ \pi^-$), the tagging decision ($d_i = -1, 0, +1$, for $i = \text{OS}, \text{SS}$), the mistag and the tagging efficiency of the OS and SS taggers, and the asymmetries (production, detection and tagging efficiency). In the ideal case (no asymmetries, zero mistag, 100% tagging efficiency), these effective coefficients become the CP coefficients already introduced in Sec. 1.3.2.

The OS and SS taggers are combined on the fly during the time fit. All the steps to build the final PDF are described in details below.

The tagging efficiency $\varepsilon_{\text{tag}}^i$ of the i th tagger ($i = \text{OS}, \text{SS}$) is corrected for the tagging efficiency difference $\Delta\varepsilon^i$ as follows:

$$\varepsilon_B^i = \varepsilon_{\text{tag}}^i - \frac{1}{2} \Delta\varepsilon^i, \quad (\text{G.2})$$

$$\varepsilon_{\bar{B}}^i = \varepsilon_{\text{tag}}^i + \frac{1}{2} \Delta\varepsilon^i, \quad (\text{G.3})$$

where ε_B^i and $\varepsilon_{\bar{B}}^i$ are the two different tagging efficiencies for B^0 and \bar{B}^0 , respectively. Moreover, the OS and SS mistag are calibrated taking into account asymmetries between B^0 and \bar{B}^0 according to Eq. 3.10. We will refer to them here as ω_B^i and $\omega_{\bar{B}}^i$ (for $i = \text{OS}, \text{SS}$). It’s now convenient to define the quantity $\Delta^\pm(d_{\text{OS}}, d_{\text{SS}})$ in the following way according to the OS and SS tagging decisions.

- If $d_{\text{OS}} = d_{\text{SS}} = 0$ (untagged B^0 candidate):

$$\Delta^\pm = (1 - \varepsilon_B^{\text{OS}} - \varepsilon_{\bar{B}}^{\text{SS}} + \varepsilon_B^{\text{OS}} \varepsilon_{\bar{B}}^{\text{SS}}) \pm (1 - \varepsilon_{\bar{B}}^{\text{OS}} - \varepsilon_B^{\text{SS}} + \varepsilon_{\bar{B}}^{\text{OS}} \varepsilon_B^{\text{SS}}). \quad (\text{G.4})$$

- If $d_i = 0, d_j \neq 0$ and $i \neq j$ (B^0 candidate tagged by one tagger):

$$\Delta^\pm = \frac{1}{2} \varepsilon_B^j \left[1 - \varepsilon_B^i + d_j \left(1 - \varepsilon_B^i - 2\omega_B^j + 2\omega_B^j \varepsilon_B^i \right) \right] \pm \frac{1}{2} \varepsilon_{\bar{B}}^j \left[1 - \varepsilon_{\bar{B}}^i + d_j \left(1 - \varepsilon_{\bar{B}}^i - 2\omega_{\bar{B}}^j + 2\omega_{\bar{B}}^j \varepsilon_{\bar{B}}^i \right) \right]. \quad (\text{G.5})$$

- If $d_{\text{OS}} = d_{\text{SS}} = 1$ (B^0 candidate tagged by both taggers):

$$\begin{aligned}
 \Delta^\pm &= \frac{1}{4} \varepsilon_B^{\text{SS}} \varepsilon_B^{\text{OS}} \left[1 + d_{\text{SS}} \left(1 - 2\omega_B^{\text{SS}} \right) + d_{\text{OS}} \left(1 - 2\omega_B^{\text{OS}} \right) \right. \\
 &\quad \left. + d_{\text{OS}} d_{\text{SS}} \left(1 - 2\omega_B^{\text{SS}} - 2\omega_B^{\text{OS}} + 4\omega_B^{\text{SS}} \omega_B^{\text{OS}} \right) \right] \\
 &\quad \pm \frac{1}{4} \varepsilon_B^{\text{SS}} \varepsilon_B^{\text{OS}} \left[1 + d_{\text{SS}} \left(1 - 2\omega_B^{\text{SS}} \right) + d_{\text{OS}} \left(1 - 2\omega_B^{\text{OS}} \right) \right. \\
 &\quad \left. + d_{\text{OS}} d_{\text{SS}} \left(1 - 2\omega_B^{\text{SS}} - 2\omega_B^{\text{OS}} + 4\omega_B^{\text{SS}} \omega_B^{\text{OS}} \right) \right].
 \end{aligned} \tag{G.6}$$

Finally, the effective CP coefficients can be written as

$$C_{\text{sin,cos}}^{\text{eff}} = \begin{cases} (1 + A_{\text{D}}) C_{\text{sin,cos}}^{\text{phys}} (\Delta^- - A_{\text{P}} \Delta^+) & \text{for } q = f = D^- \pi^+, \\ -(1 - A_{\text{D}}) C_{\text{sin,cos}}^{\text{phys}} (\Delta^- - A_{\text{P}} \Delta^+) & \text{for } q = \bar{f} = D^+ \pi^-, \end{cases} \tag{G.7}$$

$$C_{\text{sinh,cosh}}^{\text{eff}} = \begin{cases} (1 + A_{\text{D}}) C_{\text{sinh}}^{\text{phys}} (\Delta^+ - A_{\text{P}} \Delta^-) & \text{for } q = f = D^- \pi^+, \\ (1 - A_{\text{D}}) C_{\text{sinh}}^{\text{phys}} (\Delta^+ - A_{\text{P}} \Delta^-) & \text{for } q = \bar{f} = D^+ \pi^-, \end{cases} \tag{G.8}$$

where A_{D} and A_{P} are the production and detection asymmetries defined in Eqs. 5.9 and 5.10, respectively.

H Correlation matrix of the $B^0 \rightarrow D^\mp \pi^\pm$ decay-time fit

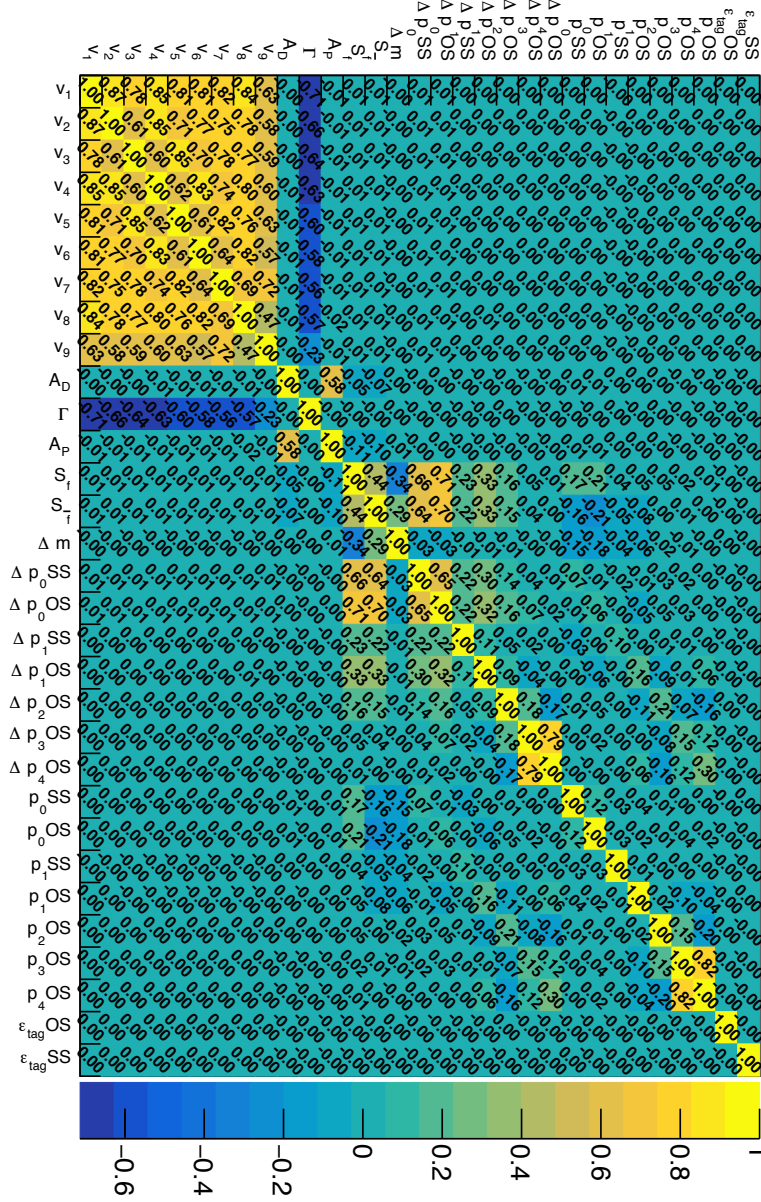


Figure H.1 – Correlation matrix of the $B^0 \rightarrow D^\mp \pi^\pm$ decay-time fit.

I. Decay-time fits of $B^0 \rightarrow D^{\mp}\pi^{\pm}$ data subsamples

I Decay-time fits of $B^0 \rightarrow D^{\mp}\pi^{\pm}$ data subsamples

Table I.1 – Results of the decay-time fits to the 2011, 2012, magnet-up and down subsamples.

Parameter	2011	2012	Magnet Up	Magnet Down
S_f	0.054 ± 0.039	0.059 ± 0.024	0.032 ± 0.029	0.079 ± 0.029
$S_{\bar{f}}$	0.031 ± 0.039	0.041 ± 0.024	-0.010 ± 0.029	0.080 ± 0.029
A_P	-0.0022 ± 0.0054	-0.0079 ± 0.0033	-0.0056 ± 0.0041	-0.0072 ± 0.0039
A_D	0.0137 ± 0.0036	0.0065 ± 0.0022	0.0075 ± 0.0027	0.0096 ± 0.0026
Γ	0.6588 ± 0.0017	0.6587 ± 0.0017	0.6587 ± 0.0017	0.6588 ± 0.0017
Δm	0.5047 ± 0.0023	0.5058 ± 0.0023	0.5054 ± 0.0023	0.5050 ± 0.0023
p_0^{OS}	-0.126 ± 0.038	-0.166 ± 0.025	-0.143 ± 0.030	-0.160 ± 0.029
p_1^{OS}	-0.095 ± 0.044	-0.011 ± 0.028	-0.054 ± 0.034	-0.014 ± 0.033
p_2^{OS}	-0.025 ± 0.018	-0.000 ± 0.010	-0.013 ± 0.013	-0.004 ± 0.012
p_3^{OS}	-0.03 ± 0.21	-0.43 ± 0.14	-0.20 ± 0.18	-0.40 ± 0.15
p_4^{OS}	1.08 ± 0.94	-1.41 ± 0.70	-0.11 ± 0.92	-0.63 ± 0.64
Δp_0^{OS}	-0.117 ± 0.092	-0.065 ± 0.057	-0.103 ± 0.070	-0.063 ± 0.068
Δp_1^{OS}	0.064 ± 0.067	0.172 ± 0.042	0.187 ± 0.051	0.101 ± 0.050
Δp_2^{OS}	-0.036 ± 0.025	-0.018 ± 0.015	-0.020 ± 0.019	-0.031 ± 0.017
Δp_3^{OS}	-0.08 ± 0.30	-0.29 ± 0.20	-0.17 ± 0.26	-0.34 ± 0.21
Δp_4^{OS}	0.8 ± 1.4	-1.5 ± 1.2	-0.2 ± 1.5	-0.72 ± 0.89
p_0^{SS}	-0.046 ± 0.039	-0.041 ± 0.025	0.020 ± 0.030	-0.097 ± 0.030
p_1^{SS}	0.005 ± 0.042	-0.019 ± 0.026	-0.022 ± 0.031	-0.004 ± 0.032
Δp_0^{SS}	-0.142 ± 0.083	-0.061 ± 0.052	-0.102 ± 0.063	-0.076 ± 0.061
Δp_1^{SS}	0.068 ± 0.062	0.030 ± 0.039	0.048 ± 0.046	0.036 ± 0.046
$\varepsilon_{\text{tag}}^{\text{OS}}$	0.4270 ± 0.0014	0.43458 ± 0.00091	0.4321 ± 0.0011	0.4326 ± 0.0011
$\varepsilon_{\text{tag}}^{\text{SS}}$	0.92753 ± 0.00076	0.93163 ± 0.00046	0.92978 ± 0.00058	0.93107 ± 0.00054
v_1	0.350 ± 0.011	0.3082 ± 0.0066	0.3199 ± 0.0078	0.3186 ± 0.0075
v_2	0.535 ± 0.018	0.479 ± 0.011	0.506 ± 0.014	0.483 ± 0.013
v_3	0.886 ± 0.029	0.760 ± 0.018	0.786 ± 0.021	0.800 ± 0.020
v_4	1.055 ± 0.034	0.973 ± 0.021	1.013 ± 0.026	0.978 ± 0.024
v_5	1.186 ± 0.038	1.060 ± 0.023	1.096 ± 0.028	1.089 ± 0.026
v_6	1.176 ± 0.037	1.095 ± 0.023	1.140 ± 0.028	1.096 ± 0.026
v_7	1.231 ± 0.039	1.108 ± 0.023	1.156 ± 0.028	1.126 ± 0.027
v_8	1.267 ± 0.035	1.143 ± 0.021	1.189 ± 0.026	1.163 ± 0.024
v_9	1.203 ± 0.050	1.135 ± 0.029	1.174 ± 0.037	1.136 ± 0.035

Appendices

Table I.2 – Results of the decay-time fits when only OS tagged, SS tagged or OS and SS candidates are considered exclusively.

Parameter	OS only	SS only	both OS and SS
S_f	0.046 ± 0.056	0.037 ± 0.040	0.064 ± 0.024
$S_{\bar{f}}$	-0.015 ± 0.056	0.019 ± 0.040	0.048 ± 0.024
A_P	0.009 ± 0.010	-0.0110 ± 0.0036	-0.0022 ± 0.0044
A_D	0.0100 ± 0.0066	0.0086 ± 0.0024	0.0083 ± 0.0030
Γ	0.6588 ± 0.0017	0.6587 ± 0.0017	0.6587 ± 0.0017
Δm	0.5051 ± 0.0023	0.5055 ± 0.0023	0.5049 ± 0.0023
p_0^{OS}	-0.164 ± 0.051		-0.148 ± 0.022
p_1^{OS}	-0.017 ± 0.058		-0.036 ± 0.024
p_2^{OS}	-0.007 ± 0.023		-0.0077 ± 0.0091
p_3^{OS}	-0.16 ± 0.26		-0.34 ± 0.11
p_4^{OS}	0.2 ± 1.4		-0.55 ± 0.50
Δp_0^{OS}	-0.10 ± 0.13		-0.059 ± 0.055
Δp_1^{OS}	0.120 ± 0.089		0.154 ± 0.038
Δp_2^{OS}	-0.071 ± 0.033		-0.015 ± 0.013
Δp_3^{OS}	-0.34 ± 0.37		-0.24 ± 0.16
Δp_4^{OS}	0.9 ± 2.1		-0.59 ± 0.71
p_0^{SS}		-0.023 ± 0.028	-0.072 ± 0.031
p_1^{SS}		-0.027 ± 0.027	0.008 ± 0.034
Δp_0^{SS}		-0.102 ± 0.075	-0.112 ± 0.058
Δp_1^{SS}		0.014 ± 0.041	0.075 ± 0.050
v_1	0.276 ± 0.014	0.3155 ± 0.0070	0.3344 ± 0.0087
v_2	0.403 ± 0.024	0.500 ± 0.012	0.504 ± 0.015
v_3	0.668 ± 0.040	0.793 ± 0.019	0.821 ± 0.023
v_4	0.903 ± 0.050	0.992 ± 0.023	1.017 ± 0.028
v_5	1.070 ± 0.059	1.083 ± 0.024	1.112 ± 0.030
v_6	1.045 ± 0.058	1.100 ± 0.024	1.159 ± 0.031
v_7	1.154 ± 0.062	1.137 ± 0.025	1.142 ± 0.030
v_8	1.138 ± 0.054	1.165 ± 0.022	1.199 ± 0.028
v_9	1.117 ± 0.086	1.148 ± 0.032	1.171 ± 0.040

I. Decay-time fits of $B^0 \rightarrow D^\mp \pi^\pm$ data subsamples

Table I.3 – Results of the decay-time fits in bins of the reconstructed B^0 transverse momentum (GeV/ c).

Parameter	$p_T \in [0, 6.5)$	$p_T \in [6.5, 9.3)$	$p_T \in [9.3, 12.8)$	$p_T \in [12.8, +\infty)$
S_f	0.032 ± 0.039	0.074 ± 0.040	0.088 ± 0.038	0.046 ± 0.036
$S_{\bar{f}}$	0.008 ± 0.039	0.041 ± 0.041	0.066 ± 0.039	0.025 ± 0.036
A_P	-0.0052 ± 0.0053	-0.0041 ± 0.0053	-0.0158 ± 0.0054	-0.0013 ± 0.0056
A_D	0.0092 ± 0.0033	0.0140 ± 0.0035	0.0008 ± 0.0037	0.0098 ± 0.0040
Γ	0.6587 ± 0.0017	0.6587 ± 0.0017	0.6588 ± 0.0017	0.6588 ± 0.0017
Δm	0.5048 ± 0.0023	0.5054 ± 0.0023	0.5056 ± 0.0023	0.5045 ± 0.0023
p_0^{OS}	-0.088 ± 0.041	-0.123 ± 0.040	-0.149 ± 0.039	-0.237 ± 0.037
p_1^{OS}	-0.008 ± 0.045	0.044 ± 0.048	-0.052 ± 0.046	-0.117 ± 0.042
p_2^{OS}	0.006 ± 0.016	-0.027 ± 0.019	-0.005 ± 0.021	-0.004 ± 0.016
p_3^{OS}	-0.60 ± 0.20	-0.44 ± 0.23	0.06 ± 0.31	-0.19 ± 0.20
p_4^{OS}	-1.61 ± 0.83	-0.9 ± 1.3	-0.8 ± 1.5	0.07 ± 0.81
Δp_0^{OS}	-0.094 ± 0.099	-0.104 ± 0.099	-0.064 ± 0.090	-0.062 ± 0.082
Δp_1^{OS}	0.074 ± 0.066	0.183 ± 0.073	0.218 ± 0.071	0.078 ± 0.064
Δp_2^{OS}	-0.040 ± 0.022	-0.012 ± 0.027	-0.042 ± 0.032	-0.001 ± 0.023
Δp_3^{OS}	-0.20 ± 0.27	-0.58 ± 0.37	-0.08 ± 0.58	-0.00 ± 0.29
Δp_4^{OS}	-0.0 ± 1.1	-3.2 ± 1.9	-2.8 ± 2.6	0.4 ± 1.2
p_0^{SS}	-0.198 ± 0.084	-0.223 ± 0.044	-0.100 ± 0.038	0.146 ± 0.032
p_1^{SS}	-0.142 ± 0.081	-0.045 ± 0.051	-0.034 ± 0.040	-0.014 ± 0.033
Δp_0^{SS}	-0.24 ± 0.13	-0.075 ± 0.087	-0.047 ± 0.080	-0.091 ± 0.074
Δp_1^{SS}	-0.03 ± 0.11	-0.041 ± 0.071	0.047 ± 0.059	0.073 ± 0.053
$\varepsilon_{\text{tag}}^{\text{OS}}$	0.4461 ± 0.0015	0.4210 ± 0.0014	0.4176 ± 0.0014	0.4452 ± 0.0015
$\varepsilon_{\text{tag}}^{\text{SS}}$	0.92426 ± 0.00078	0.94508 ± 0.00066	0.93922 ± 0.00070	0.91274 ± 0.00084
v_1	0.0492 ± 0.0018	0.2480 ± 0.0073	0.463 ± 0.014	0.779 ± 0.025
v_2	0.1206 ± 0.0048	0.433 ± 0.014	0.702 ± 0.023	1.030 ± 0.036
v_3	0.493 ± 0.015	0.757 ± 0.024	0.937 ± 0.032	1.307 ± 0.047
v_4	0.749 ± 0.021	0.955 ± 0.029	1.127 ± 0.037	1.344 ± 0.048
v_5	0.902 ± 0.025	1.079 ± 0.033	1.162 ± 0.038	1.432 ± 0.051
v_6	0.981 ± 0.027	1.033 ± 0.032	1.211 ± 0.039	1.389 ± 0.049
v_7	0.993 ± 0.027	1.114 ± 0.033	1.203 ± 0.039	1.406 ± 0.050
v_8	1.062 ± 0.025	1.142 ± 0.030	1.215 ± 0.035	1.403 ± 0.045
v_9	1.033 ± 0.038	1.112 ± 0.045	1.245 ± 0.053	1.365 ± 0.063

Appendices

Table I.4 – Results of the decay-time fits in bins of the number of reconstructed primary vertices (PV) in the event.

Parameter	1 PV	2 PV	> 2 PV
S_f	0.007 ± 0.034	0.061 ± 0.031	0.118 ± 0.035
$S_{\bar{f}}$	0.027 ± 0.033	0.016 ± 0.031	0.080 ± 0.036
A_P	-0.0055 ± 0.0047	-0.0112 ± 0.0043	-0.0009 ± 0.0049
A_D	0.0064 ± 0.0032	0.0053 ± 0.0029	0.0152 ± 0.0033
Γ	0.6587 ± 0.0017	0.6587 ± 0.0017	0.6588 ± 0.0017
Δm	0.5052 ± 0.0023	0.5051 ± 0.0023	0.5052 ± 0.0023
p_0^{OS}	-0.193 ± 0.034	-0.177 ± 0.032	-0.067 ± 0.037
p_1^{OS}	-0.011 ± 0.043	-0.033 ± 0.035	-0.042 ± 0.042
p_2^{OS}	-0.034 ± 0.018	0.011 ± 0.013	-0.002 ± 0.016
p_3^{OS}	-0.37 ± 0.20	-0.35 ± 0.17	-0.11 ± 0.20
p_4^{OS}	-0.05 ± 0.93	-1.25 ± 0.97	-0.4 ± 1.1
Δp_0^{OS}	-0.189 ± 0.079	-0.122 ± 0.074	0.120 ± 0.087
Δp_1^{OS}	0.119 ± 0.065	0.088 ± 0.053	0.273 ± 0.065
Δp_2^{OS}	-0.055 ± 0.026	-0.020 ± 0.019	0.000 ± 0.023
Δp_3^{OS}	0.16 ± 0.28	-0.53 ± 0.26	-0.17 ± 0.30
Δp_4^{OS}	1.9 ± 1.3	-2.6 ± 1.4	-1.2 ± 1.6
p_0^{SS}	-0.031 ± 0.035	-0.057 ± 0.032	-0.031 ± 0.037
p_1^{SS}	0.020 ± 0.038	-0.034 ± 0.034	-0.015 ± 0.039
Δp_0^{SS}	-0.102 ± 0.072	-0.175 ± 0.066	0.065 ± 0.078
Δp_1^{SS}	-0.009 ± 0.056	0.046 ± 0.049	0.094 ± 0.057
$\varepsilon_{\text{tag}}^{\text{OS}}$	0.4394 ± 0.0013	0.4330 ± 0.0012	0.4240 ± 0.0013
$\varepsilon_{\text{tag}}^{\text{SS}}$	0.93394 ± 0.00065	0.93134 ± 0.00059	0.92557 ± 0.00071
v_1	0.3144 ± 0.0084	0.3265 ± 0.0082	0.3150 ± 0.0090
v_2	0.484 ± 0.014	0.491 ± 0.014	0.510 ± 0.016
v_3	0.754 ± 0.022	0.802 ± 0.022	0.829 ± 0.026
v_4	0.935 ± 0.026	1.012 ± 0.026	1.040 ± 0.031
v_5	1.027 ± 0.029	1.093 ± 0.029	1.170 ± 0.035
v_6	1.038 ± 0.029	1.121 ± 0.029	1.204 ± 0.035
v_7	1.087 ± 0.030	1.144 ± 0.029	1.197 ± 0.035
v_8	1.103 ± 0.027	1.173 ± 0.026	1.262 ± 0.032
v_9	1.118 ± 0.041	1.150 ± 0.038	1.203 ± 0.046

I. Decay-time fits of $B^0 \rightarrow D^\mp \pi^\pm$ data subsamples

Table I.5 – Results of the decay-time fits in bins of the logarithm of the total number of reconstructed tracks ($\ln N_{\text{tr}}$) in the event.

Parameter	$\ln N_{\text{tr}} \in [0, 4.8)$	$\ln N_{\text{tr}} \in [4.8, 5.2)$	$\ln N_{\text{tr}} \in [5.2, +\infty)$
S_f	0.036 ± 0.029	0.038 ± 0.034	0.116 ± 0.039
$S_{\bar{f}}$	0.046 ± 0.029	0.000 ± 0.034	0.066 ± 0.039
A_P	-0.0031 ± 0.0046	-0.0071 ± 0.0046	-0.0093 ± 0.0047
A_D	0.0092 ± 0.0031	0.0073 ± 0.0031	0.0091 ± 0.0031
Γ	0.6587 ± 0.0017	0.6588 ± 0.0017	0.6588 ± 0.0017
Δm	0.5056 ± 0.0023	0.5049 ± 0.0023	0.5050 ± 0.0023
p_0^{OS}	-0.077 ± 0.031	-0.174 ± 0.035	-0.221 ± 0.038
p_1^{OS}	-0.050 ± 0.041	-0.038 ± 0.038	-0.003 ± 0.042
p_2^{OS}	-0.047 ± 0.020	0.013 ± 0.015	0.018 ± 0.016
p_3^{OS}	-0.15 ± 0.28	-0.27 ± 0.19	-0.22 ± 0.20
p_4^{OS}	-0.3 ± 1.4	-1.33 ± 0.94	-1.13 ± 0.97
Δp_0^{OS}	-0.077 ± 0.071	-0.190 ± 0.081	0.056 ± 0.091
Δp_1^{OS}	0.105 ± 0.062	0.099 ± 0.059	0.239 ± 0.065
Δp_2^{OS}	-0.030 ± 0.031	-0.042 ± 0.021	-0.010 ± 0.023
Δp_3^{OS}	0.26 ± 0.59	-0.11 ± 0.26	-0.71 ± 0.29
Δp_4^{OS}	-1.6 ± 2.3	1.8 ± 1.4	-3.8 ± 1.5
p_0^{SS}	0.015 ± 0.031	-0.028 ± 0.035	-0.126 ± 0.039
p_1^{SS}	-0.038 ± 0.032	-0.037 ± 0.037	0.028 ± 0.044
Δp_0^{SS}	-0.100 ± 0.063	-0.057 ± 0.073	-0.089 ± 0.083
Δp_1^{SS}	0.003 ± 0.047	0.074 ± 0.054	0.069 ± 0.063
$\varepsilon_{\text{tag}}^{\text{OS}}$	0.4082 ± 0.0012	0.4370 ± 0.0013	0.4523 ± 0.0013
$\varepsilon_{\text{tag}}^{\text{SS}}$	0.90933 ± 0.00073	0.93877 ± 0.00061	0.94361 ± 0.00059
v_1	0.2904 ± 0.0075	0.3406 ± 0.0091	0.3303 ± 0.0089
v_2	0.450 ± 0.013	0.521 ± 0.015	0.518 ± 0.016
v_3	0.728 ± 0.021	0.814 ± 0.024	0.847 ± 0.025
v_4	0.932 ± 0.025	1.010 ± 0.028	1.051 ± 0.030
v_5	1.008 ± 0.027	1.116 ± 0.031	1.167 ± 0.033
v_6	1.033 ± 0.028	1.144 ± 0.032	1.186 ± 0.033
v_7	1.066 ± 0.028	1.140 ± 0.032	1.226 ± 0.034
v_8	1.083 ± 0.025	1.219 ± 0.029	1.235 ± 0.030
v_9	1.135 ± 0.039	1.123 ± 0.042	1.211 ± 0.044

Appendices

Table I.6 – Results of the decay-time fits in bins of the difference in pseudorapidity $\Delta\eta$ between the D^\mp meson and the bachelor pion.

Parameter	$\Delta\eta \in (-\infty, -0.36)$	$\Delta\eta \in [-0.36, 0.10)$	$\Delta\eta \in [0.10, 0.53)$	$\Delta\eta \in [0.53, +\infty)$
S_f	0.047 ± 0.038	0.050 ± 0.039	0.103 ± 0.038	0.034 ± 0.039
$S_{\bar{f}}$	0.028 ± 0.038	0.060 ± 0.039	0.006 ± 0.038	0.054 ± 0.039
A_P	-0.0092 ± 0.0053	-0.0121 ± 0.0054	0.0004 ± 0.0053	-0.0051 ± 0.0053
A_D	0.0072 ± 0.0036	0.0043 ± 0.0037	0.0118 ± 0.0036	0.0106 ± 0.0035
Γ	0.6587 ± 0.0017	0.6587 ± 0.0017	0.6588 ± 0.0017	0.6588 ± 0.0017
Δm	0.5052 ± 0.0023	0.5047 ± 0.0023	0.5046 ± 0.0023	0.5059 ± 0.0023
p_0^{OS}	-0.127 ± 0.039	-0.145 ± 0.039	-0.157 ± 0.039	-0.188 ± 0.041
p_1^{OS}	-0.032 ± 0.049	-0.033 ± 0.044	-0.044 ± 0.045	-0.010 ± 0.045
p_2^{OS}	-0.037 ± 0.023	0.011 ± 0.017	0.003 ± 0.018	-0.016 ± 0.015
p_3^{OS}	-0.01 ± 0.32	-0.12 ± 0.22	0.01 ± 0.23	-0.95 ± 0.20
p_4^{OS}	-0.7 ± 2.0	0.16 ± 0.92	0.66 ± 0.96	-2.93 ± 0.91
Δp_0^{OS}	-0.105 ± 0.092	-0.082 ± 0.093	0.001 ± 0.089	-0.140 ± 0.093
Δp_1^{OS}	0.185 ± 0.074	0.276 ± 0.069	0.029 ± 0.067	0.082 ± 0.067
Δp_2^{OS}	-0.021 ± 0.037	-0.054 ± 0.026	0.007 ± 0.025	-0.032 ± 0.022
Δp_3^{OS}	0.16 ± 0.79	-0.31 ± 0.30	-0.10 ± 0.31	-0.56 ± 0.28
Δp_4^{OS}	-2.5 ± 2.9	-0.1 ± 1.3	-0.4 ± 1.3	-1.2 ± 1.3
p_0^{SS}	-0.061 ± 0.040	-0.009 ± 0.038	-0.029 ± 0.040	-0.071 ± 0.041
p_1^{SS}	0.018 ± 0.044	-0.060 ± 0.039	0.055 ± 0.043	-0.054 ± 0.044
Δp_0^{SS}	-0.143 ± 0.083	-0.121 ± 0.082	0.024 ± 0.081	-0.098 ± 0.084
Δp_1^{SS}	0.157 ± 0.066	0.033 ± 0.058	-0.038 ± 0.063	0.017 ± 0.063
$\varepsilon_{\text{tag}}^{\text{OS}}$	0.4370 ± 0.0015	0.4344 ± 0.0015	0.4346 ± 0.0015	0.4236 ± 0.0014
$\varepsilon_{\text{tag}}^{\text{SS}}$	0.93563 ± 0.00072	0.92611 ± 0.00077	0.92349 ± 0.00078	0.93669 ± 0.00071
v_1	0.2796 ± 0.0082	0.377 ± 0.011	0.396 ± 0.012	0.2429 ± 0.0072
v_2	0.439 ± 0.015	0.565 ± 0.019	0.591 ± 0.020	0.404 ± 0.013
v_3	0.768 ± 0.025	0.872 ± 0.029	0.879 ± 0.030	0.679 ± 0.022
v_4	0.882 ± 0.028	1.061 ± 0.034	1.105 ± 0.035	0.952 ± 0.029
v_5	1.078 ± 0.033	1.152 ± 0.037	1.180 ± 0.038	0.982 ± 0.031
v_6	1.009 ± 0.031	1.181 ± 0.037	1.202 ± 0.038	1.095 ± 0.033
v_7	1.112 ± 0.034	1.171 ± 0.037	1.208 ± 0.038	1.083 ± 0.033
v_8	1.095 ± 0.029	1.222 ± 0.034	1.200 ± 0.034	1.191 ± 0.031
v_9	1.130 ± 0.046	1.128 ± 0.048	1.312 ± 0.052	1.065 ± 0.045

I. Decay-time fits of $B^0 \rightarrow D^\mp \pi^\pm$ data subsamples

Table I.7 – Results of the decay-time fits performed for TOS candidates on L0Hadron and all the other candidates.

Parameter	L0HadronTOS	not L0HadronTOS
S_f	0.029 ± 0.026	0.111 ± 0.034
$S_{\bar{f}}$	0.029 ± 0.026	0.048 ± 0.034
A_P	-0.0092 ± 0.0035	-0.0002 ± 0.0049
A_D	0.0079 ± 0.0024	0.0093 ± 0.0032
Γ	0.6588 ± 0.0017	0.6587 ± 0.0017
Δm	0.5053 ± 0.0023	0.5049 ± 0.0023
p_0^{OS}	-0.196 ± 0.027	-0.083 ± 0.034
p_1^{OS}	-0.066 ± 0.030	-0.005 ± 0.039
p_2^{OS}	-0.006 ± 0.012	-0.010 ± 0.015
p_3^{OS}	-0.20 ± 0.16	-0.46 ± 0.18
p_4^{OS}	-0.66 ± 0.79	-0.70 ± 0.80
Δp_0^{OS}	-0.145 ± 0.060	0.044 ± 0.084
Δp_1^{OS}	0.122 ± 0.046	0.151 ± 0.059
Δp_2^{OS}	-0.032 ± 0.017	-0.012 ± 0.020
Δp_3^{OS}	-0.34 ± 0.25	0.04 ± 0.25
Δp_4^{OS}	-1.8 ± 1.3	0.7 ± 1.1
p_0^{SS}	0.013 ± 0.024	-0.229 ± 0.046
p_1^{SS}	-0.024 ± 0.025	-0.029 ± 0.050
Δp_0^{SS}	-0.108 ± 0.053	-0.061 ± 0.083
Δp_1^{SS}	0.012 ± 0.038	0.100 ± 0.070
$\varepsilon_{\text{tag}}^{\text{OS}}$	0.40498 ± 0.00094	0.4834 ± 0.0014
$\varepsilon_{\text{tag}}^{\text{SS}}$	0.93331 ± 0.00048	0.92502 ± 0.00071
v_1	0.4111 ± 0.0092	0.1889 ± 0.0052
v_2	0.608 ± 0.015	0.331 ± 0.010
v_3	0.909 ± 0.022	0.633 ± 0.019
v_4	1.087 ± 0.025	0.863 ± 0.024
v_5	1.174 ± 0.027	0.981 ± 0.028
v_6	1.188 ± 0.027	1.018 ± 0.028
v_7	1.200 ± 0.027	1.051 ± 0.029
v_8	1.234 ± 0.024	1.095 ± 0.026
v_9	1.218 ± 0.034	1.063 ± 0.039

J Decay-time fit validation with bootstrapping

In this appendix, further results of the Monte Carlo time fit validation are shown in addition to Fig. 5.14. The results for the production asymmetry, detection asymmetry, and acceptance parameters are not shown because no independent reference value for these parameters exists in Monte Carlo.

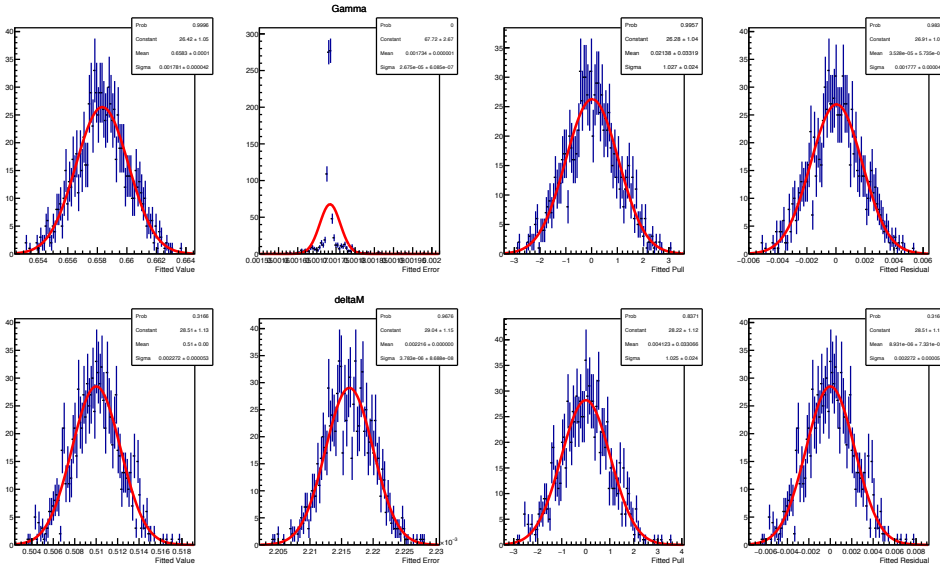


Figure J.1 – Distributions of the fitted value, error, pull and residual for Γ (top) and Δm (bottom). Each distribution is fitted with a Gaussian function. The result of the Gaussian fit is shown for the fitted error as well, even though uncertainties are not always Gaussian. Pulls and residuals are computed by taking the Monte Carlo generation value as reference (see Appendix M).

J. Decay-time fit validation with bootstrapping

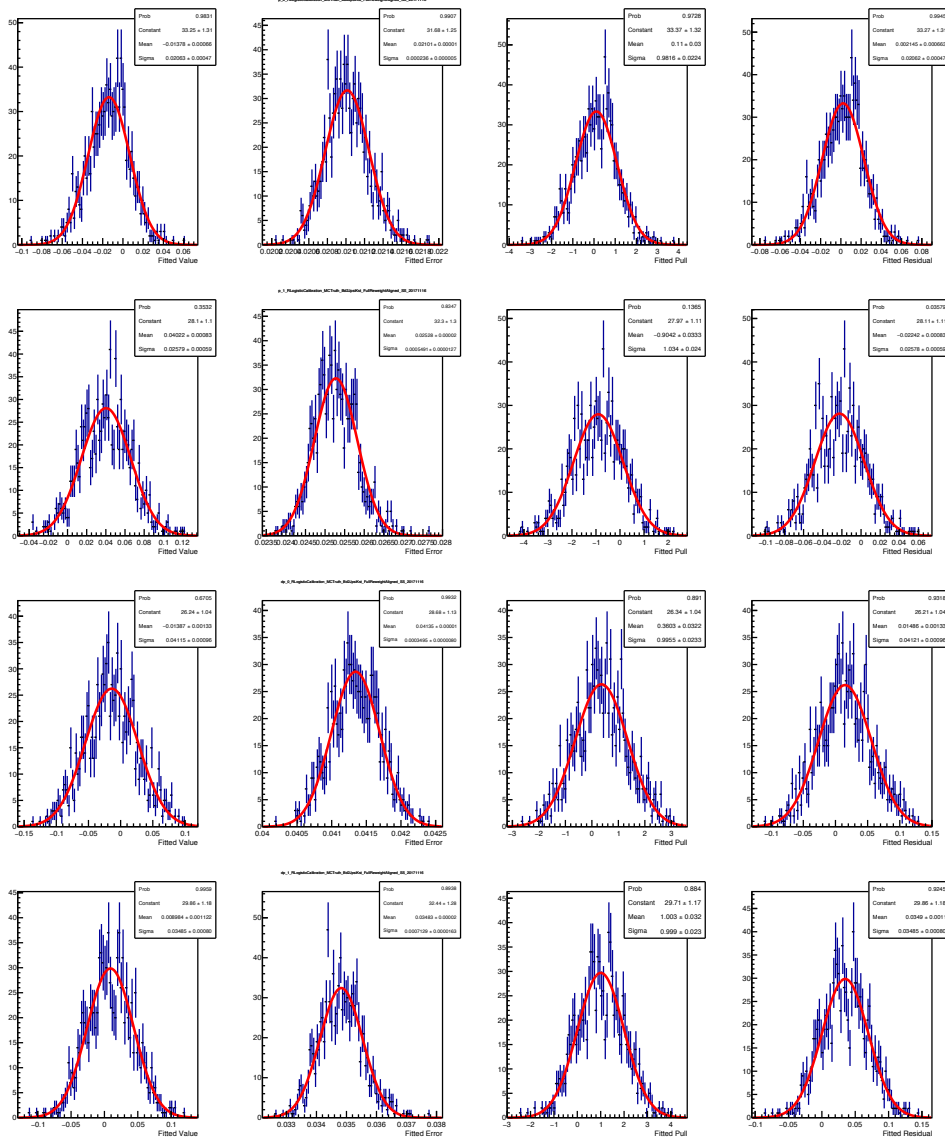


Figure J.2 – Distributions of the fitted value, error, pull and residual for the SS tagger calibration parameters p_0^{SS} , p_1^{SS} , Δp_0^{SS} , and Δp_1^{SS} (from top to bottom). Each distribution is fitted with a Gaussian function. The result of the Gaussian fit is shown for the fitted error as well, even though uncertainties are not always Gaussian. Pulls and residuals are computed by taking the values found on the $B^0 \rightarrow J/\psi K^{*0}$ Monte Carlo calibration as reference (see Table 3.5).

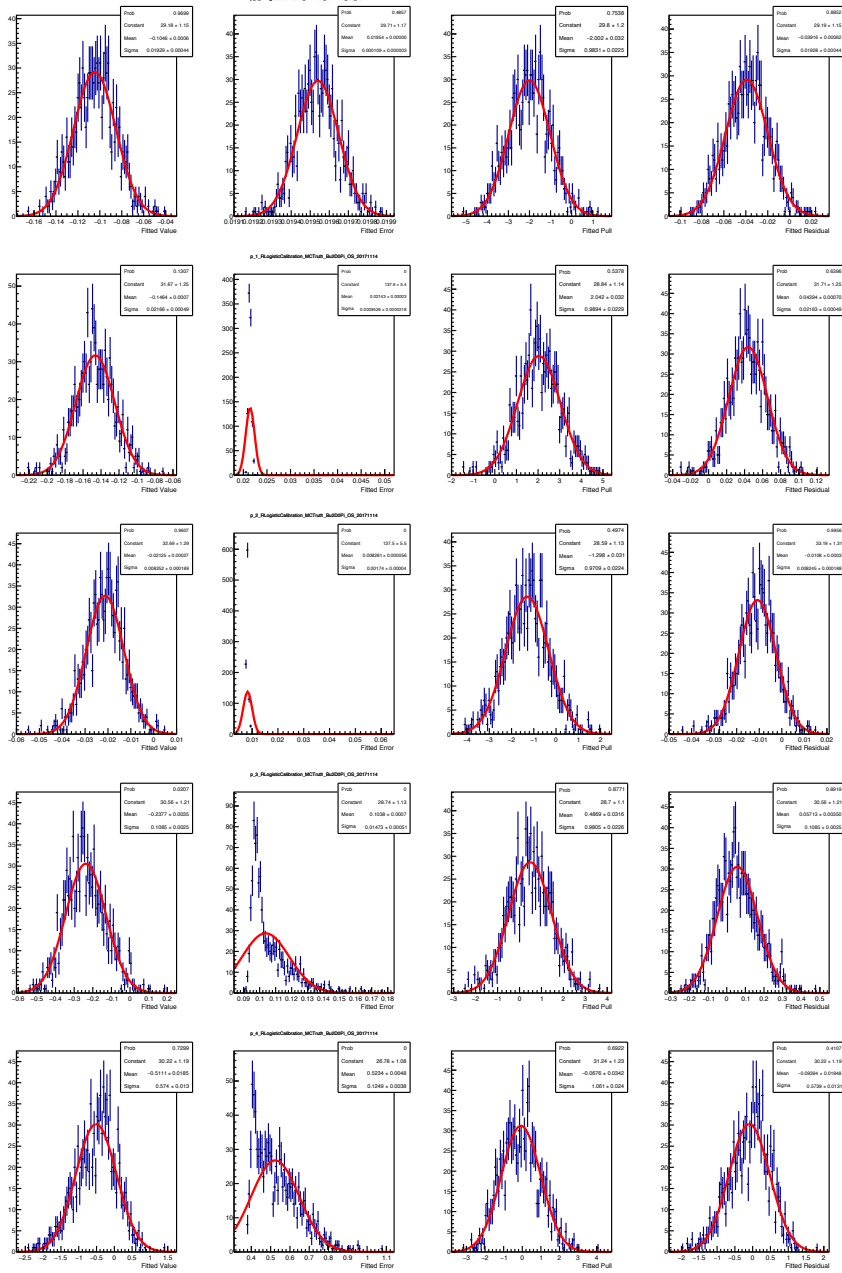


Figure J.3 – Distributions of the fitted value, error, pull and residual for the OS tagger calibration parameters p_0^{OS} , p_1^{OS} , p_2^{OS} , p_3^{OS} , and p_4^{OS} (from top to bottom). Each distribution is fitted with a Gaussian function. The result of the Gaussian fit is shown for the fitted error as well, even though uncertainties are not always Gaussian. Pulls and residuals are computed by taking the values found on the $B^+ \rightarrow \bar{D}^0\pi^+$ Monte Carlo calibration as reference (see Table 3.3).

J. Decay-time fit validation with bootstrapping

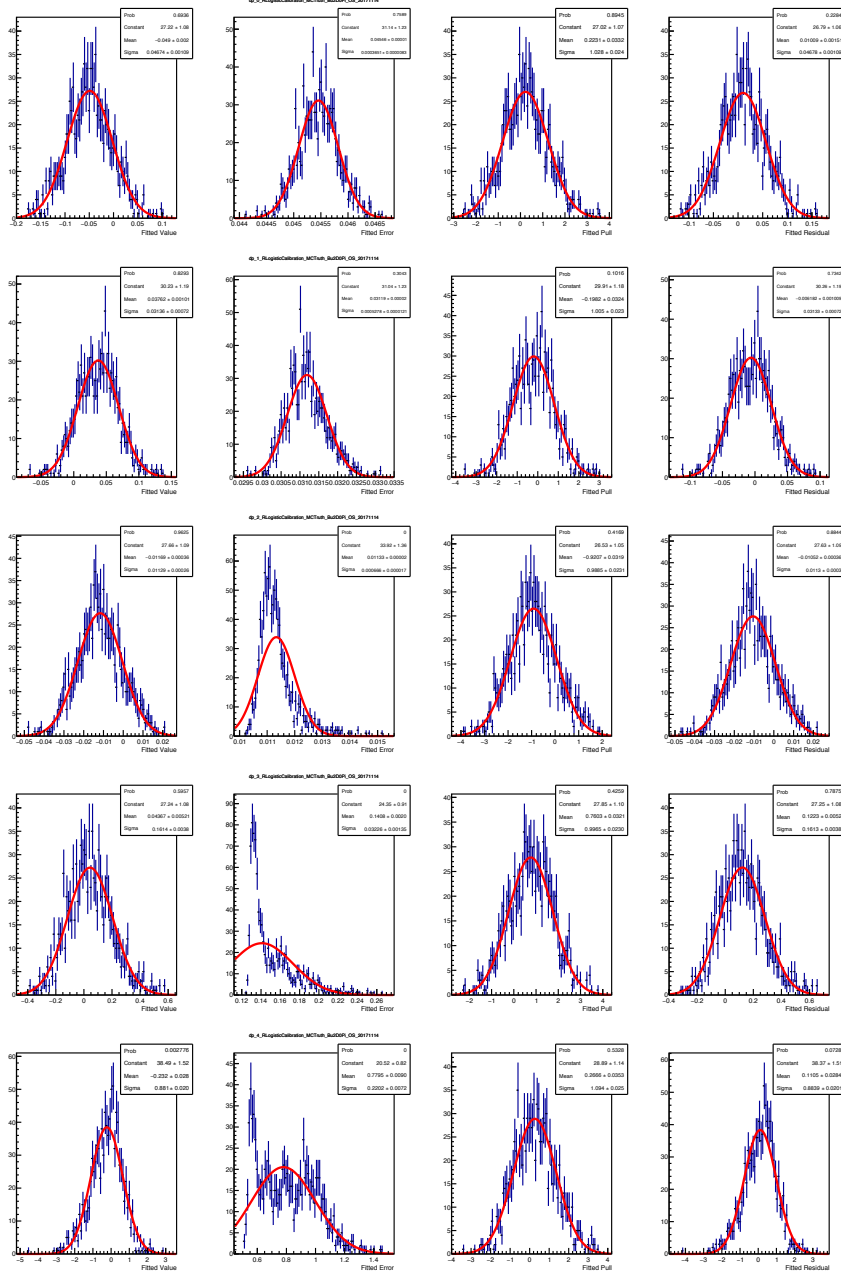


Figure J.4 – Distributions of the fitted value, error, pull and residual for the OS tagger calibration parameters Δp_0^{OS} , Δp_1^{OS} , Δp_2^{OS} , Δp_3^{OS} , and Δp_4^{OS} (from top to bottom). Each distribution is fitted with a Gaussian function. The result of the Gaussian fit is shown for the fitted error as well, even though uncertainties are not always Gaussian. Pulls and residuals are computed by taking the values found on the $B^+ \rightarrow \bar{D}^0 \pi^+$ Monte Carlo calibration as reference (see Table 3.3).

K Test of the decay-time fit via a toy tagger

The *toy tagger* used to perform the test mentioned in Sec. 5.4.3 is created as follows. First, a mistag (η) PDF is created from the *sWeighted* η distribution of the OS tagger on data. This template is created as a `RoohistPdf`. Then, for each candidate in each bootstrapped Monte Carlo sample, a value of η is drawn from this PDF. The decision of the toy tagger is initially taken from the true ID of the B^0 meson, which is always correct by definition. In order to emulate wrong tagging decisions, a random number r_i is generated for the i^{th} B^0 candidate between 0 and 1. If η_i is the mistag assigned to this candidate, the tagging decision d_i may need to be flipped (and thus made wrong) according to the following criterion:

$$d_i \rightarrow \begin{cases} -d_i & \text{if } r_i \leq \eta_i \\ d_i & \text{otherwise} \end{cases}. \quad (\text{K.1})$$

During the time fit, the mistag calibration is simply taken as a linear function (Eq. 3.9) with $p_0 = \langle \eta \rangle = 0.370029$ (taken from the adopted template) and $p_1 = 1$, which means $\omega = \eta$ for all candidates. In fact, the per-event mistag η is the true mistag ω probability by construction. In this way, it is possible to test the time fit with a per-event mistag without relying on any approximation or uncertainty coming from the calibration procedure. Moreover, the tagging efficiency is 100% by construction.

The distributions of the fitted value, error, pull and residual for the relevant parameters are shown in Fig. K.1. All pull distributions have means compatible with 0 and widths compatible with 1, meaning that the maximum likelihood estimation of the parameters is unbiased and returns correct uncertainties.

K. Test of the decay-time fit via a toy tagger

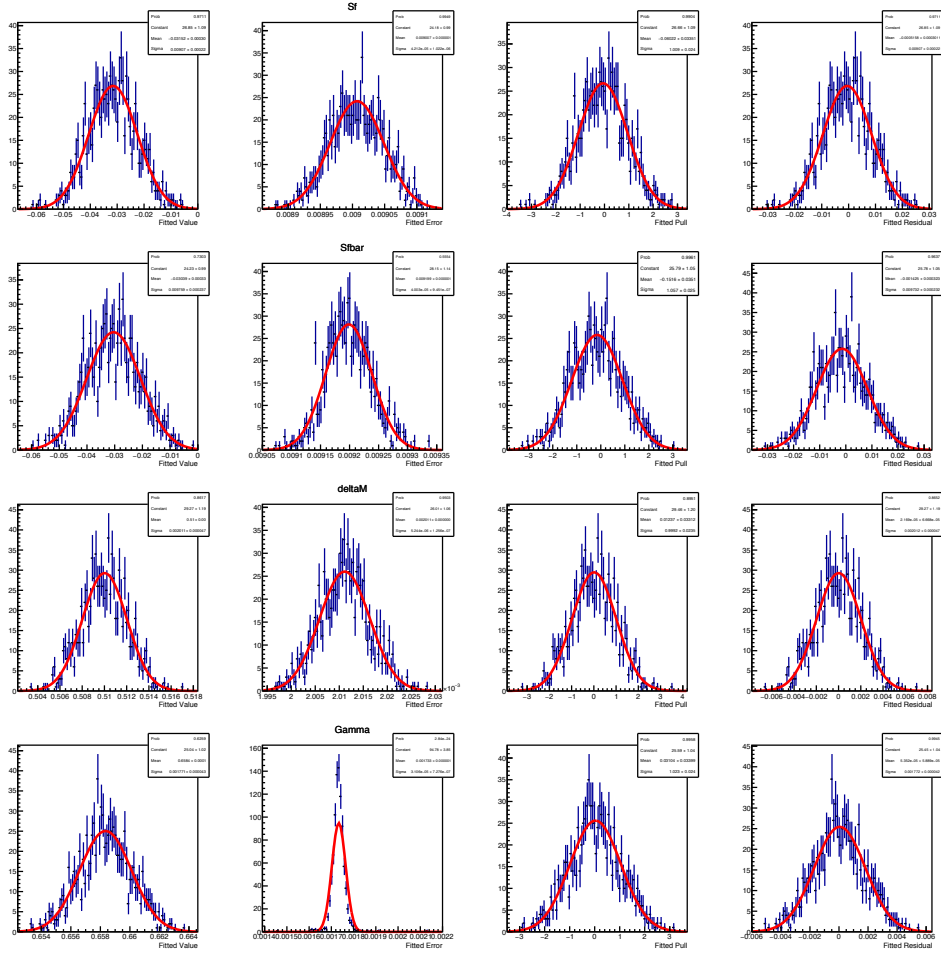


Figure K.1 – Distributions of the fitted value, error, pull and residual for the main parameters (S_f , $S_{\bar{f}}$, Δm , and Γ , from top to bottom) fitted on bootstrapped Monte Carlo samples with a toy tagger. Each distribution is fitted with a Gaussian function. Pulls and residuals are computed by taking the Monte Carlo generation value as reference.

L Correlation between systematic uncertainties

As reported in Table 5.5, the main systematic uncertainties on S_f and $S_{\bar{f}}$ are the ones due to the constraint on Δm , the fit biases, and the background subtraction. The correlations between the systematics uncertainties on S_f and $S_{\bar{f}}$ due to these sources are described in details in this appendix. The correlation of the total systematic resulting from these three contributions is -0.41 . The correlation between other sources of systematics is neglected.

L.1 Correlation of Δm systematics

The correlation of systematics uncertainties due to Δm between S_f and $S_{\bar{f}}$ is estimated by comparing the nominal fit result, and the result obtained with Δm fixed as described in Sec. 5.5.1. This correlation $\rho_{S_f, S_{\bar{f}}}^{\Delta m}$ is simply computed from the difference in statistical covariance between the two fit results:

$$\rho_{S_f, S_{\bar{f}}}^{\Delta m} = \frac{\sigma_{S_f}^{\text{nominal}} \sigma_{S_{\bar{f}}}^{\text{nominal}} \rho_{S_f, S_{\bar{f}}}^{\text{nominal}} - \sigma_{S_f}^{\Delta m \text{ fixed}} \sigma_{S_{\bar{f}}}^{\Delta m \text{ fixed}} \rho_{S_f, S_{\bar{f}}}^{\Delta m \text{ fixed}}}{\sqrt{(\sigma_{S_f}^{\text{nominal}})^2 - (\sigma_{S_f}^{\Delta m \text{ fixed}})^2} \sqrt{(\sigma_{S_{\bar{f}}}^{\text{nominal}})^2 - (\sigma_{S_{\bar{f}}}^{\Delta m \text{ fixed}})^2}}. \quad (\text{L.1})$$

The value of $\rho_{S_f, S_{\bar{f}}}^{\Delta m}$ so obtained is -1 , meaning that S_f and $S_{\bar{f}}$ are fully anticorrelated driven by the uncertainty on Δm . This can be understood with the following argument. The \bar{B}^0 versus B^0 time-dependent asymmetry for the $f = D^- \pi^+$ and $\bar{f} = D^+ \pi^-$ final states can be written as

$$\begin{aligned} -S_f \sin(\Delta mt) + \cos(\Delta mt) &= \cos(\Delta mt + \delta_f), \\ -S_{\bar{f}} \sin(\Delta mt) - \cos(\Delta mt) &= -\cos(\Delta mt - \delta_{\bar{f}}), \end{aligned}$$

respectively, where $C_f = -C_{\bar{f}} = 1$ is assumed, while $\delta_f = \sin^{-1}(S_f)$ and $\delta_{\bar{f}} = \sin^{-1}(S_{\bar{f}})$. Since S_f and $S_{\bar{f}}$ are small, the approximations $\delta_f \sim S_f$ and $\delta_{\bar{f}} \sim S_{\bar{f}}$ can be done. As shown in Fig. 3.3, the sensitivity to S_f and $S_{\bar{f}}$ is obtained around the zero of the asymmetries, where the cosine terms disappear. If t_0 is one of the zeroes, the following condition must hold:

$$\cos(\Delta mt_0 + S_f) = \cos(\pi - \Delta mt_0 + S_{\bar{f}}) = 0. \quad (\text{L.2})$$

So, if Δm is shifted by some systematic amount, S_f and $S_{\bar{f}}$ have to be shifted in opposite directions (and thus be anticorrelated) in order to satisfy L.2.

L.2 Correlation of systematics due to fit biases

As described in Sec. 5.4.3, a bias on S_f and $S_{\bar{f}}$ is observed from bootstrapped MC samples, and the size of this bias is assigned as systematic uncertainty. The associated correlation is estimated from the two-dimensional distribution of $(S_f^{\text{fit}} - S_f^{\text{gen}})$ versus

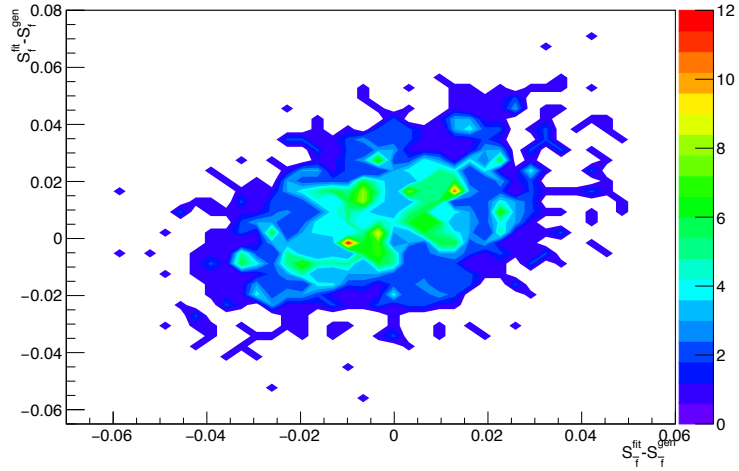


Figure L.1 – Two-dimensional distribution of $(S_f^{\text{fit}} - S_f^{\text{gen}})$ versus $(S_{\bar{f}}^{\text{fit}} - S_{\bar{f}}^{\text{gen}})$ obtained from the fits to bootstrapped MC samples (Sec. 5.4.3).

$(S_{\bar{f}}^{\text{fit}} - S_{\bar{f}}^{\text{gen}})$ obtained on the same set of bootstrapped MC samples, as shown in Fig. L.1. The resulting correlation is 0.4.

L.3 Correlation of systematics due to background subtraction

A systematic due to *sWeighting* and background subtraction is assigned by repeating Fit B in a wider mass range, as described in Sec. 5.5.3. In order to estimate the correlation of this systematic uncertainty between S_f and $S_{\bar{f}}$, the data sample is bootstrapped in a similar way as done for the Monte Carlo (Sec. 5.4.3). Then, *sWeights* are obtained twice on each sample, once with the nominal strategy, and once by selecting a wide mass range in Fit B. Finally, the time fit is performed on each sample using both the sets of *sWeights* computed in the previous step. The correlation is estimated from the two-dimensional distribution of $(S_f^{\text{nominal}} - S_f^{\text{wide mass}})$ versus $(S_{\bar{f}}^{\text{nominal}} - S_{\bar{f}}^{\text{wide mass}})$, which is shown in Fig. L.2. The resulting correlation is 0.7.

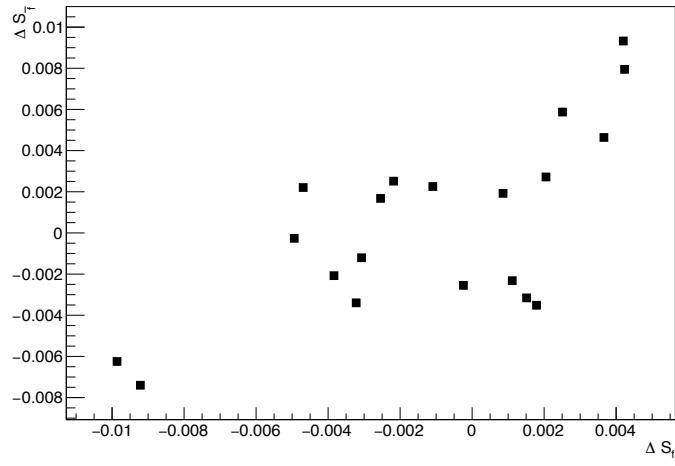


Figure L.2 – Two-dimensional distribution of $S_f^{\text{nominal}} - S_f^{\text{wide mass}}$ (ΔS_f) versus $S_f^{\text{nominal}} - S_f^{\text{wide mass}}$ (ΔS_f) obtained from fits to bootstrapped data samples, where Fit B is performed with the nominal strategy or with a wider mass range for each sample.

M Inputs for the $B^0 \rightarrow D^{\mp}\pi^{\pm}$ simulation

Table M.1 – Parameter values used in the generation of the $B^0 \rightarrow D^{\mp}\pi^{\pm}$ MC sample, and resulting values for the CP coefficients.

Parameter	Generation value
Δm	0.51 ps^{-1}
$\tau = 1/\Gamma$	1.519068 ps
$\Delta\Gamma/\Gamma$	0
$ q/p $	1
$\arg(q/p)$	-0.764
$ A_f $	0.0849
$\arg(A_f)$	0.002278
$ \bar{A}_f $	0.00137
$\arg(\bar{A}_f)$	-1.128958
$ A_{\bar{f}} $	0.00137
$\arg(A_{\bar{f}})$	1.3145
$ \bar{A}_{\bar{f}} $	0.0849
$\arg(\bar{A}_{\bar{f}})$	0.002278
S_f	-0.0305
$S_{\bar{f}}$	-0.0282
C_f	0.9995
$C_{\bar{f}}$	-0.9995
D_f	-0.0103
$D_{\bar{f}}$	-0.0155

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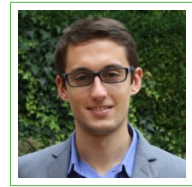
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Date of birth December 14, 1989
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Education

- 2014-2018 **Ph.D. in Physics**, EPFL, Lausanne.
Thesis: "Measurement of time-dependent CP violation in $B^0 \rightarrow D^\mp \pi^\pm$ decays and optimisation of flavour tagging algorithms at LHCb"
Supervisors: Prof. O Schneider (EPFL), Dr. C. Fitzpatrick (EPFL)
- 2011-2014 **MS.c. in Physics**, University of Ferrara, Italy, Grade: 110/110 cum laude.
Thesis: "A Study of Semileptonic B Decays in Orbitally Excited D Mesons at LHCb"
Supervisor: Dr. C. Bozzi (INFN Ferrara)
- 2008-2011 **BS.c. in Physics**, University of Ferrara, Italy, Grade: 110/110 cum laude.
Thesis: "Studio dei parametri caratteristici di dispositivi optoelettronici per il rivelatore di muoni dell'esperimento SuperB"
Supervisors: Prof. R. Calabrese (University of Ferrara), Dr. W. Baldini (INFN Ferrara), Dr. G. Cibinetto (INFN Ferrara)

Selected publications

- LHCb collaboration, R. Aaij *et al.*, *Measurement of CP violation in $B^0 \rightarrow D^\mp \pi^\pm$ decays*, JHEP **06** (2018) 084, arXiv:1805.03448 [hep-ex]
- LHCb collaboration, R. Aaij *et al.*, *Measurement of CP asymmetry in $B_s^0 \rightarrow D_s^\mp K^\pm$ decays*, JHEP **03** (2018) 059, arXiv:1712.07428 [hep-ex]
- V. Battista, *b-flavour tagging in pp collisions at LHCb*, Il Nuovo Cimento **C4** (2016), 10.1393/ncc/i2016-16335-5
- V. Battista *et al.*, *A study of spillover clusters and ghost tracks in the Silicon Tracker with 25 ns bunch spacing*, CERN-LHCb-INT-2016-010 (2016)

Teaching assistantship

- 2014-2017 Laboratory of nuclear physics, exercises on nuclear and particle physics (EPFL)
2012 Exercises on general physics, data analysis and programming (University of Ferrara)

Scholarships

- Summer 2012 CERN Summer Student Programme
Calibration of silicon-based sensors for the monitoring of the radiation field in the LHCb cavern
Supervisors: Dr. G. Corti (CERN), Dr. M. Karacson (CERN)
- Summer 2011 DOE/INFN Summer Exchange Program, SLAC
Optimisation of the selection of $D_s^{+*} \rightarrow D_s^+(\rightarrow \mu^+\nu_\mu)\gamma$ decays using simulated events for the BaBar experiment
Supervisor: Dr. M. Convery (SLAC)

Research activity

CP violation measurements with B^0 and B_s^0 decays

I performed measurements of the CKM angle γ with time-dependent analyses of $B^0 \rightarrow D^\mp \pi^\pm$ and $B_s^0 \rightarrow D_s^\mp K^\pm$ using data collected with the LHCb detector. The result obtained with $B^0 \rightarrow D^\mp \pi^\pm$ is more precise than previous measurements from BaBar and Belle collaborations. The analysis of $B_s^0 \rightarrow D_s^\mp K^\pm$ is the first measurement of γ using B_s^0 decays.

Flavour tagging algorithms

I implemented and optimised flavour tagging algorithms used by the LHCb collaboration to infer the flavour of neutral B mesons for time-dependent analyses.

Silicon Tracker

I was responsible for the calibration and the maintenance of the Silicon Tracker detector of LHCb.

Talks at conferences

- Apr 2018 *Incontri di Fisica delle Alte Energie*, Milano,
"Misure di violazione di CP in decadimenti $B^0 \rightarrow D^\mp \pi^\pm$ a LHCb"
- Aug 2017 *CHIPP & Swiss Physical Society Meeting*, Geneva
"Flavour tagging in pp collisions at LHCb"
- May 2017 *Technology and Instrumentation in Particle Physics*, Beijing
"Operation of the LHCb silicon tracking and vertexing systems in LHC Run-2"
- Feb 2017 *Lake Louise Winter Institute*, Lake Louise
" CP violation in B decays at LHCb"
- Aug 2016 *CHIPP Annual Plenary Meeting*, Lugano
"Time-dependent measurement of CP violation in $B^0 \rightarrow D^- \pi^+$ decays at LHCb"
- Mar 2016 *Les Rencontres de Physique de la Vallée d'Aoste*, La Thuile
" b -flavour tagging in pp collisions at LHCb"
- Apr 2015 *Incontri di Fisica delle Alte Energie*, Roma,
"Violazione di CP e misura di $|V_{ub}|$ tramite decadimenti semileptonici di adroni B a LHCb"

Software skills

C, C++, Python, Linux shell, ROOT, RooFit, TMVA, Matlab, Mathematica, Keras, scikit-learn, MySQL, Tableau, L^AT_EX

Languages

English (fluent, C2), French (advanced, B2) and Italian (mother tongue)

Science outreach

- 2016 Collaboration to the organisation of the "EPFL Open Days". An exhibit of simple experiments was organised for the general public in order to explain the principles of particle physics and detectors.
- 2015-2016 Collaboration with an Italian science popularisation journal (*Quaderni di scienza e scienziati molisani*).

