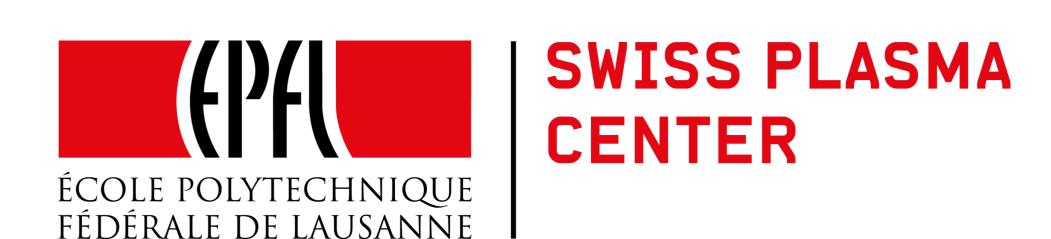
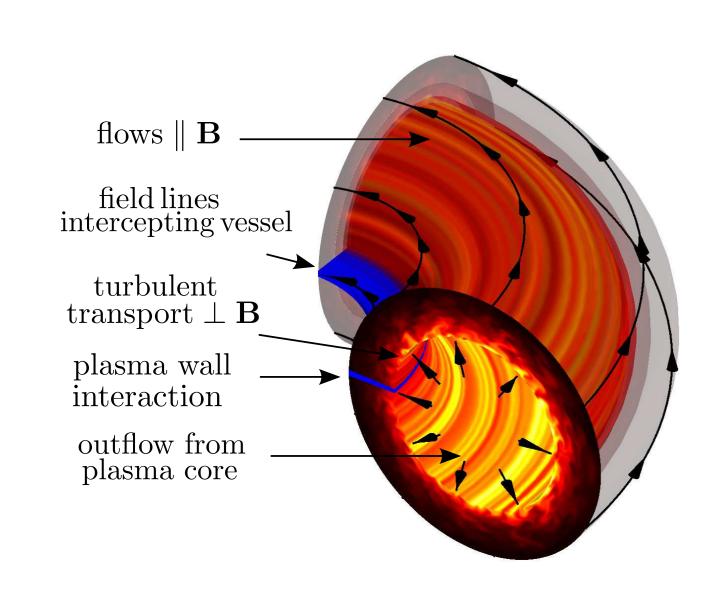
Progress in simulating SOL plasma turbulence with the GBS code

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Introduction



- ▶ In the tokamak SOL, magnetic field lines intersect the walls of the fusion device
- ► Heat and particles flow along magnetic field lines and are exhausted to the vessel
- ► Turbulence amplitude and size comparable to **steady-state** values
- ▶ Neutral particles interact with the plasma

The Global Braginskii Solver (GBS) code: a 3D, flux-driven, global turbulence code in limited geometry used to study plasma turbulence in the SOL

[Ricci et al., PPCF 2012; Halpern et al., JCP 2016]

▶ GBS solves 3D fluid equations for electrons and ions, Poisson's and Ampere's equations, and a kinetic equation for neutral atoms.

The Global Braginskii Solver (GBS) code

Two-fluid drift-reduced Braginskii equations, $k_{\perp}^2 \gg k_{\parallel}^2$, $d/dt \ll \omega_{ci}$

$$\frac{\partial n}{\partial t} = -\frac{\rho_{*}^{-1}}{B}[\phi, n] + \frac{2}{B}[C(p_{e}) - nC(\phi)] - \nabla \cdot (nv_{\parallel e}\mathbf{b}) + \mathcal{D}_{n}(n) + S_{n} + n_{n}\nu_{iz} - n\nu_{rec}$$

$$\frac{\partial \Omega}{\partial t} = -\frac{\rho_{*}^{-1}}{B}\nabla_{\perp} \cdot [\phi, \omega] - \nabla_{\perp} \cdot \left[\nabla_{\parallel}(v_{\parallel i}\omega)\right] + B^{2}\nabla \cdot (j_{\parallel}\mathbf{b}) + 2BC(p) + \frac{B}{3}C(G_{i}) + \mathcal{D}_{\Omega}(\Omega) - \frac{n_{n}}{n}\nu_{cx}\Omega$$

$$\frac{\partial U_{\parallel e}}{\partial t} = -\frac{\rho_{*}^{-1}}{B}[\phi, v_{\parallel e}] - v_{\parallel e}\nabla_{\parallel}v_{\parallel e} + \frac{m_{i}}{m_{e}}\left[\frac{\nu j_{\parallel}}{n} + \nabla_{\parallel}\phi - \frac{\nabla_{\parallel}p_{e}}{n} - 0.71\nabla_{\parallel}T_{e} - \frac{2}{3n}\nabla_{\parallel}G_{e}\right] + \mathcal{D}_{v_{\parallel e}}(v_{\parallel e})$$
(3)

$$\frac{\partial v_{\parallel i}}{\partial t} = -\frac{\rho_{*}^{-1}}{B} [\phi, v_{\parallel i}] - v_{\parallel i} \nabla_{\parallel} v_{\parallel i} - \frac{\nabla_{\parallel} \rho}{n} - \frac{2}{3n} \nabla_{\parallel} G_{i} + \mathcal{D}_{v_{\parallel i}} (v_{\parallel i}) + \frac{n_{n}}{n} (v_{iz} + v_{cx}) (v_{\parallel n} - v_{\parallel i})
\frac{\partial T_{e}}{\partial t} = -\frac{\rho_{*}^{-1}}{B} [\phi, T_{e}] - v_{\parallel e} \nabla_{\parallel} T_{e} + \frac{4T_{e}}{3B} \left[\frac{C(p_{e})}{n} + \frac{5}{2} C(T_{e}) - C(\phi) \right] + \frac{2T_{e}}{3n} \left[0.71 \nabla \cdot (j_{\parallel} \mathbf{b}) - n \nabla \cdot (v_{\parallel e} \mathbf{b}) \right]
+ \mathcal{D}_{T_{e}} (T_{e}) + \mathcal{D}_{T_{e}}^{\parallel} (T_{e}) + S_{T_{e}} + \frac{n_{n}}{n} v_{iz} \left[-\frac{2}{3} E_{iz} - T_{e} + \frac{m_{e}}{m_{i}} v_{\parallel e} \left(v_{\parallel e} - \frac{4}{3} v_{\parallel n} \right) \right] - \frac{n_{n}}{n} v_{en} \frac{m_{e}}{m_{i}} \frac{2}{3} v_{\parallel e} (v_{\parallel n} - v_{\parallel e})$$
(5)

$$+ \mathcal{D}_{T_{e}}(T_{e}) + \mathcal{D}_{T_{e}}^{\parallel}(T_{e}) + S_{T_{e}} + \frac{n_{n}}{n} \nu_{iz} \left[-\frac{2}{3} E_{iz} - T_{e} + \frac{m_{e}}{m_{i}} v_{\parallel e} \left(v_{\parallel e} - \frac{4}{3} v_{\parallel n} \right) \right] - \frac{n_{n}}{n} \nu_{en} \frac{m_{e}}{m_{i}} \frac{2}{3} v_{\parallel e} (v_{\parallel n} - v_{\parallel e})$$

$$\frac{\partial T_{i}}{\partial t} = -\frac{\rho_{*}^{-1}}{B} [\phi, T_{i}] - v_{\parallel i} \nabla_{\parallel} T_{i} + \frac{4 T_{i}}{3B} \left[\frac{C(p_{e})}{n} - \frac{5}{2} \tau C(T_{i}) - C(\phi) \right] + \frac{2 T_{i}}{3n} \left[\nabla \cdot (j_{\parallel} \mathbf{b}) - n \nabla \cdot (v_{\parallel i} \mathbf{b}) \right]$$

$$+ \mathcal{D}_{T_{i}}(T_{i}) + \mathcal{D}_{T_{i}}^{\parallel}(T_{i}) + S_{T_{i}} + \frac{n_{n}}{n} (\nu_{iz} + \nu_{ex}) \left[\tau^{-1} T_{n} - T_{i} + \frac{1}{3\tau} (v_{\parallel n} - v_{\parallel i})^{2} \right]$$

$$(6)$$

$$\rho_{\star} = \rho_{s}/R, \quad \nabla_{\parallel} f = \mathbf{b}_{0} \cdot \nabla f + \frac{\beta_{e0}}{2} \frac{\rho_{\star}^{-1}}{B} [\psi, f], \quad p = n(T_{e} + \tau T_{i}), \quad U_{\parallel e} = v_{\parallel e} + \frac{\beta_{e0}}{2} \frac{m_{i}}{m_{e}} \psi, \quad \Omega = \nabla \cdot \omega = \nabla \cdot (n \nabla_{\perp} \phi + \tau \nabla_{\perp} p_{i})$$

- ► Equations implemented in GBS, a **flux-driven** plasma turbulence code with limited geometry to study SOL heat and particle transport
- ▶ System completed with first-principles boundary conditions applicable at the magnetic pre-sheath entrance where the magnetic field lines intersect the limiter [Loizu et al., PoP 2012]
- ▶ Parallelized using domain decomposition (MPI and OpenMP), excellent parallel scalability up to \sim 10000 cores
- ▶ Gradients and curvature discretized using finite differences, Poisson Brackets using Arakawa scheme, integration in time using Runge Kutta method
- ► Code **fully verified** using method of manufactured solutions [Riva *et al.*, PoP 2014]
- ▶ Note: $L_{\perp} \rightarrow \rho_s$, $L_{\parallel} \rightarrow R_0$, $t \rightarrow R_0/c_s$, $\nu = ne^2 R_0/(m_i \sigma_{\parallel} c_s)$ normalization

The Poisson and Ampere equations

- ▶ Generalized Poisson equation, $\nabla \cdot (n\nabla_{\perp}\phi) = \Omega \tau \nabla^2_{\perp} p_i$
- ▶ Ampere's equation from Ohm's law, $\left(\nabla_{\perp}^2 \frac{\beta_{e0}}{2} \frac{m_i}{m_e} n\right) v_{\parallel e} = \nabla_{\perp}^2 U_{\parallel e} \frac{\beta_{e0}}{2} \frac{m_i}{m_e} n v_{\parallel i}$
- ► Stencil based **parallel multigrid** implemented in GBS
- ▶ The elliptic equations are separable in the parallel direction leading to **independent 2D solutions** for each perpendicular plane

The kinetic neutral atoms equation

$$\frac{\partial f_{\mathsf{n}}}{\partial t} + \vec{v} \cdot \frac{\partial f_{\mathsf{n}}}{\partial \vec{x}} = -\nu_{\mathsf{iz}} f_{\mathsf{n}} - \nu_{\mathsf{cx}} n_{\mathsf{n}} \left(\frac{f_{\mathsf{n}}}{n_{\mathsf{n}}} - \frac{f_{\mathsf{i}}}{n_{\mathsf{i}}} \right) + \nu_{\mathsf{rec}} f_{\mathsf{i}}$$
(7

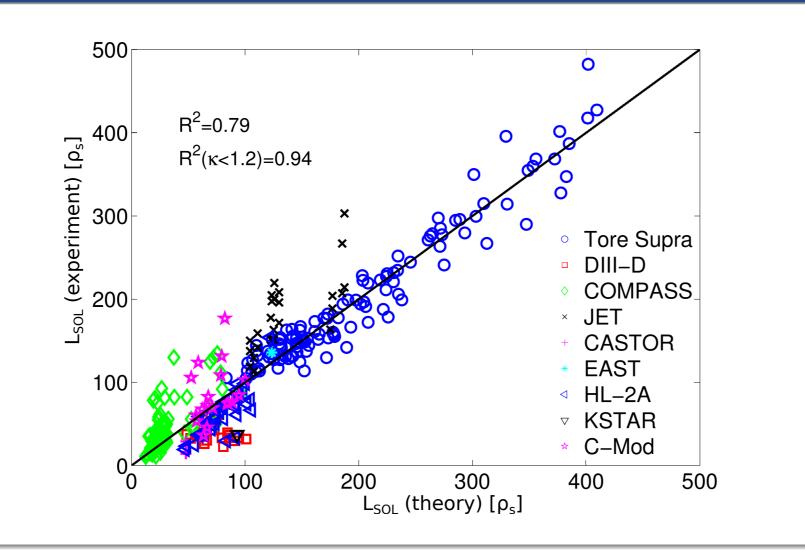
- ▶ **Method of characteristics** to obtain the formal solution of *f*_n [Wersal *et al.*, NF 2015]
- ▶ Two assumptions, $\tau_{\text{neutral losses}} < \tau_{\text{turbulence}}$ and $\lambda_{\text{mfp, neutrals}} \ll L_{\parallel, \text{plasma}}$, leading to a 2D steady state system for each perpendicular plane
- ▶ Linear integral equation for neutral density obtained by integrating f_n over \vec{v}
- ► Spatial discretization leading to a linear system of equations

$$\begin{bmatrix} n_{n} \\ \Gamma_{out} \end{bmatrix} = \begin{bmatrix} K_{p \to p} & K_{b \to p} \\ K_{p \to b} & K_{b \to b} \end{bmatrix} \cdot \begin{bmatrix} n_{n} \\ \Gamma_{out} \end{bmatrix} + \begin{bmatrix} n_{n,rec} \\ \Gamma_{out,rec} + \Gamma_{out,i} \end{bmatrix}$$
(8)

▶ This system is solved for neutral density, n_n , and neutral particle flux at the boundaries, Γ_{out} , with the threaded LAPACK solver.

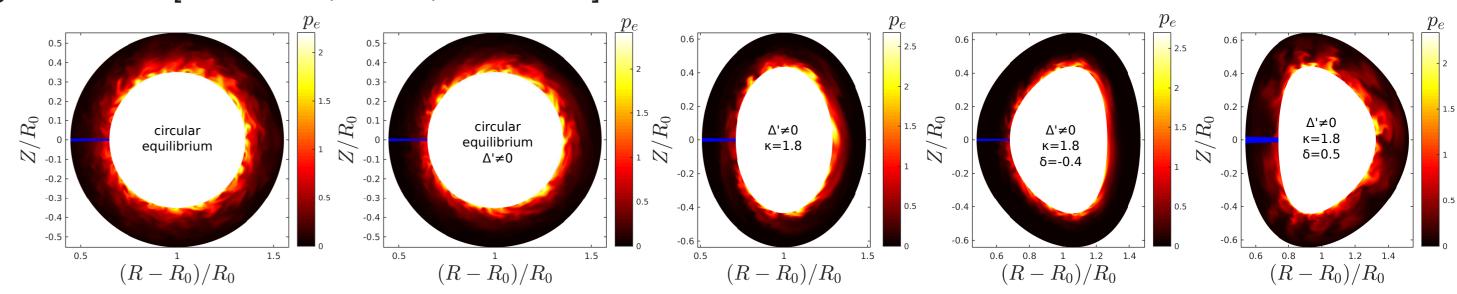
Achievements of GBS

- Characterization of non-linear turbulent regimes in the SOL [Mosetto et al., PoP 2015]
- ► SOL width scaling as a function of dimensionless / engineering plasma parameters [Halpern et al., PPCF 2016]
- ► Origin and nature of intrinsic toroidal plasma rotation in the SOL [Loizu et al., PoP 2014]
- ► Mechanisms regulating SOL equilibrium electrostatic potential [Loizu et al., PPCF 2013]



Plasma shaping effects on SOL turbulence

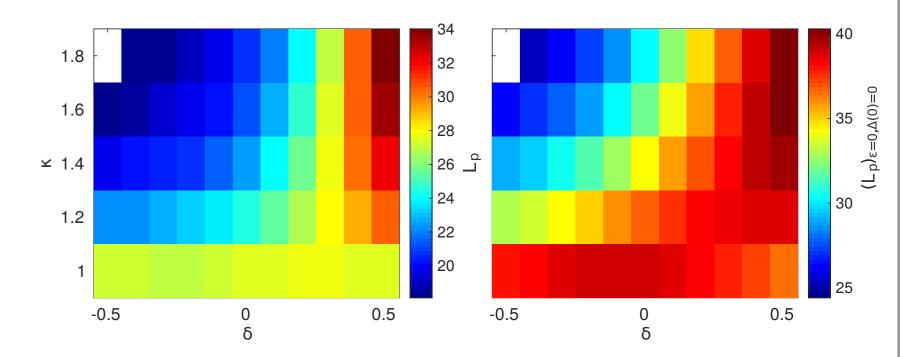
▶ Fully-turbulent non-linear simulations with same physical parameters, in different magnetic geometries [Riva et al., PPCF, submitted]



- ▶ Mitigation of turbulence by Δ' , κ , and negative δ ; enhancement of turbulence by positive δ
- ▶ Good agreement between non-linear simulations and Gradient Removal theory [Ricci et al., PoP 2013]

(κ,δ)	Non $-$ linear sim.	Gradient Removal Theory	Non - linear sim.	Gradient Removal Theory
	$\epsilon \simeq$ 0.25, Δ (0) \simeq 7	$\epsilon \simeq$ 0.25, Δ (0) \simeq 7	$\epsilon=0,\Delta(0)=0$	$\epsilon=0,\Delta(0)=0$
(1.0, 0.0)	25 ± 1	27.4	37 ± 2	38.9
(1.8, 0.0)	20 ± 1	20.7	26 ± 3	30.3
(1.8, -0.3)	15 ± 1	18.1	20 ± 1	26.2
(1.8, 0.3)	23 ± 1	26.8	43 ± 3	36.8

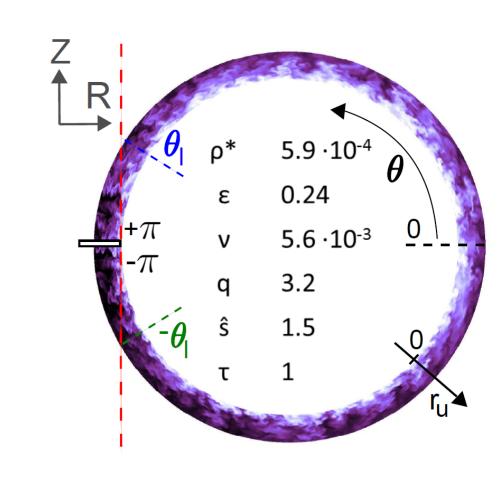
- **Linear scan** over κ and δ allows to predict the SOL width for non-circular magnetic geometries
- ▶ It is possible to **generalize** the analytical first-principle $L_{\mathcal{D}}$ scaling to include shaping effects

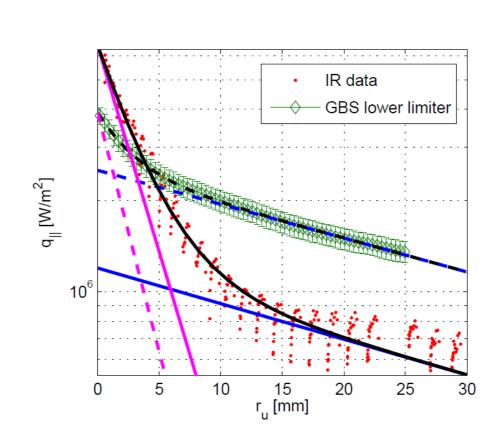


Simulation of TCV SOL

► GBS simulation of TCV discharge # 49170

- ► Full size simulation with realistic TCV input parameters
- Simulation parameters at the LCFS given by flush-mounted Langmuir probes
- ► Comparison with infrared imagining of heat flux [Nespoli et *al.*, JNM 2015]

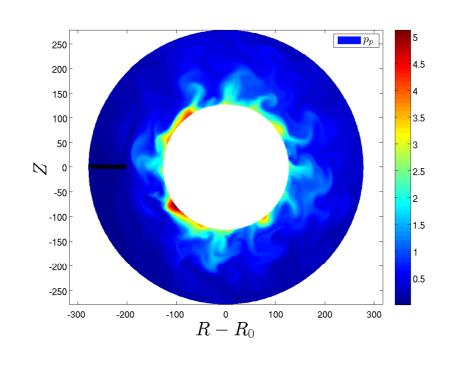


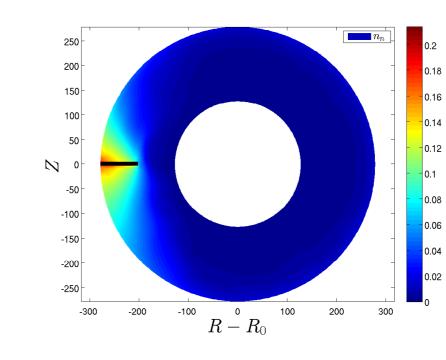


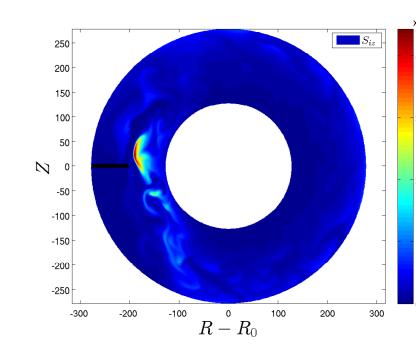
- ▶ Double scale length in heat flux profile as in TCV measurements
- Good agreement for what concerns the scale lengths
- ▶ Heat flux fall-off in the near SOL smaller with respect to experiments
- ▶ Non ambipolar current at the limiter observed in near SOL as in the experiment

Simulation with neutral atoms and closed flux surface region

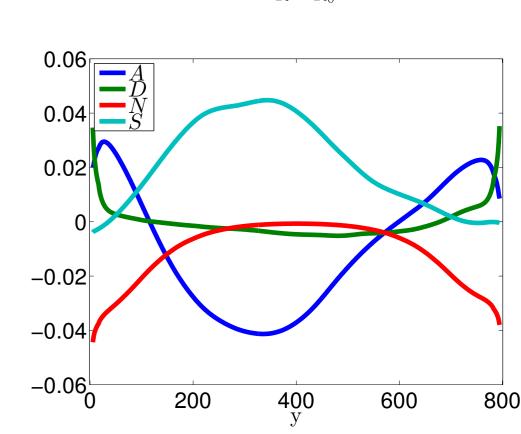
- ▶ Self-consistent GBS simulations with neutral dynamics that include closed flux surface region
- ▶ Neutral density peaks around the limiter due to recycling and ionization follows plasma fluctuations







- ► SOL quasi-steady state balance in the electron temperature equation
- ► The perpendicular drifts (S) and the neutral interaction terms (N) are balanced by the parallel advection (A) and the parallel diffusion (D) [Wersal et al., NF 2015]



Summary and Outlook

- ▶ GBS is a tool to carry out SOL turbulence simulations of medium size tokamaks
- ▶ Recent developments concern the implementation of shaping effects, neutral atom dynamics, the open-closed field lines interface, and validation agains TCV measurements
- ▶ A more flexible algorithm to simulate diverted SOL is being implemented



