Single-Layer MoS$_2$ Electronics

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CONSPECTUS: Atomic crystals of two-dimensional materials consisting of single sheets extracted from layered materials are gaining increasing attention. The most well-known material from this group is graphene, a single layer of graphite that can be extracted from the bulk material or grown on a suitable substrate. Its discovery has given rise to intense research effort culminating in the 2010 Nobel Prize in physics awarded to Andre Geim and Konstantin Novoselov. Graphene however represents only the proverbial tip of the iceberg, and increasing attention of researchers is now turning towards the veritable zoo of so-called “other 2D materials”. They have properties complementary to graphene, which in its pristine form lacks a bandgap: MoS$_2$, for example, is a semiconductor, while NbSe$_2$ is a superconductor. They could hold the key to important practical applications and new scientific discoveries in the two-dimensional limit. This family of materials has been studied since the 1960s, but most of the research focused on their tribological applications: MoS$_2$ is best known today as a high-performance dry lubricant for ultrahigh-vacuum applications and in car engines. The realization that single layers of MoS$_2$ and related materials could also be used in functional electronic devices where they could offer advantages compared with silicon or graphene created a renewed interest in these materials. MoS$_2$ is currently gaining the most attention because the material is easily available in the form of a mineral, molybdenite, but other 2D transition metal dichalcogenide (TMD) semiconductors are expected to have qualitatively similar properties.

In this Account, we describe recent progress in the area of single-layer MoS$_2$-based devices for electronic circuits. We will start with MoS$_2$ transistors, which showed for the first time that devices based on MoS$_2$ and related TMDs could have electrical properties on the same level as other, more established semiconducting materials. This allowed rapid progress in this area and was followed by demonstrations of basic digital circuits and transistors operating in the technologically relevant gigahertz range of frequencies, showing that the mobility of MoS$_2$ and TMD materials is sufficiently high to allow device operation at such high frequencies.

Monolayer MoS$_2$ and other TMDs are also direct band gap semiconductors making them interesting for realizing optoelectronic devices. These range from simple phototransistors showing high sensitivity and low noise, to light emitting diodes and solar cells. All the electronic and optoelectronic properties of MoS$_2$ and TMDs are accompanied by interesting mechanical properties with monolayer MoS$_2$ being as stiff as steel and 30X stronger. This makes it especially interesting in the context of flexible electronics where it could combine the high degree of mechanical flexibility commonly associated with organic semiconductors with high levels of electrical performance. All these results show that MoS$_2$ and TMDs are promising materials for electronic and optoelectronic applications.

INTRODUCTION

The most important basic building block of modern electronic circuits is the field-effect transistor (FET), mostly used as a switch in digital circuits. A typical FET is composed of source and drain regions that serve as contacts for the thin area connecting them, the channel. It is covered with a thin dielectric material, the gate dielectric, and capped with a metal electrode, the top gate. This electrode is used to electrostatically control the conductivity of the channel by changing the charge carrier concentration. If the transistor is to be used as a switch, the conductivity should be changed using the gate electrode from a value with high resistance, corresponding to the open position of the switch and called the off state, to a highly conductive on state, corresponding to the closed position of the switch. An ideal switch should also be able to instantly switch between the on and off states.

A good transistor channel material should satisfy several criteria. A material with high charge carrier mobility will allow fast operation and high on state current. Even though important, the mobility is not the only essential material property for transistors, as can be seen in the case of graphene with a high intrinsic mobility but no bandgap, resulting in transistors that cannot be turned off. In order to achieve high on/off ratios, a material with a bandgap is needed, with higher
bandgaps resulting in lower off state currents. This is however limited by the fact that materials with higher bandgaps show lower conductivity because of smaller charge carrier concentration. The mobility and bandgap in general also follow an empiric relationship with the mobility decreasing as the bandgap is increased.\(^1\) In order to achieve a sharp switching characteristic (subthreshold swing), a material that allows a high degree of electrostatic control is needed.

As the size of transistors is decreased, the subthreshold swing increases, causing an increase in the off state current. The undesired change in the distribution of electric fields causes short channel effects and can be modeled using the Poisson equation,\(^2\) resulting in a natural length \(\lambda\) defined as

\[
\lambda = \sqrt{\frac{\varepsilon_{\text{ox}}}{N_{\text{eff}}}}d_{\text{ch}}d_{\text{ox}}
\]

where \(\varepsilon_{\text{ch}}\) and \(\varepsilon_{\text{ox}}\) are dielectric constants of the channel and the oxide and \(d_{\text{ch}}\) and \(d_{\text{ox}}\) are the thicknesses of the semiconducting channel and the oxide. \(N\) is the effective gate number, equal to 1 for a single gate, 2 for a dual-gate geometry where the semiconducting channel is sandwiched between a pair of gate electrodes, and 3 for a trigate device where the gate surrounds the channel from three sides. A device can be considered free of short-channel effects if the gate is at least six times longer than \(\lambda\).

The natural length can be decreased by decreasing the gate oxide thickness, increasing the gate oxide dielectric constant, and also modifying the channel material and replacing it with a thinner layer that also has a lower dielectric constant. Two-dimensional materials that represent the ultimate limit of miniaturization in the vertical dimension are therefore very interesting in this context because they would allow a very high degree of electrostatic control. One such example is transition metal dichalcogenides (TMDs), with single-layer MoS\(_2\) being the typical representative of semiconducting TMDs.\(^3\)

Even though some of the semiconducting properties of TMDs have been known for decades, with reports of successful exfoliation of thin-film MoS\(_2\) published by Frindt et al. in 1963,\(^4\) they have not been seriously considered for applications in semiconductor electronics until recently. Some of the first electrical measurements on semiconducting TMDs date back to the 1960s, when Fivaz and Moser showed that the carrier mobility of MoS\(_2\), MoSe\(_2\), and WSe\(_2\) exceed 100 cm\(^2\)/Vs in the direction along the layers.\(^5\) In 2004, the first transistors were fabricated on the surface of bulk WSe\(_2\) crystals, showing...
high mobility and ambipolar behavior. The room-temperature on/off ratio was however rather low (∼10) because of bulk conduction. This was soon followed by first devices based on single-layer MoS2, which showed lower mobility (∼3 cm2/(V s)) than the bulk material and less than optimal switching characteristics for single layers.7

MoS2 TRANSISTORS

The first implementation of a switchable single-layer MoS2 transistor was demonstrated by Radisavljevic et al.8 (Figure 1). In this device, 6.5 Å thick (Figure 1b) MoS2 serves as the semiconducting channel. It is deposited on SiO2 and covered by a 30 nm thick layer of HfO2, which serves as the top-gate dielectric (Figure 1c). The transfer curves of the device are shown in Figure 1d. The transistor exhibits a current on/off ratio exceeding 108 at room temperature (Figure 1d). Besides this very high on/off ratio, the device exhibits off-state currents smaller than 100 fA (25 fA/μm). The high degree of electrostatic control is also reflected by the subthreshold slope $S = (d(\log I_d)/dV_{tg})^{-1}$, which is as low as 74 mV/dec. The pioneering work of Radisavljevic et al. showed that devices based on 2D semiconductors could reach competitive performance levels and was an important step toward the realization of electronics and low stand-by power integrated circuits based on two-dimensional materials, with the additional potential for applications in flexible, transparent electronics.

In a subsequent study, Radisavljevic and Kis9 studied the intrinsic mobility and conductivity of single-layer MoS2 sandwiched in the dual-gated geometry (Figure 2). These material properties could not be deduced with the two-terminal geometry of the initial study8 and require careful investigation with four-terminal and Hall-effect measurements (inset, Figure 2a).

Figure 2. Mobility in single-layer MoS2 encapsulated in HfO2. (a) Conductivity $\sigma$ and conductance $G$ of MoS2 as a function of top-gate voltage. Above $V_{tg} \approx 1–2$ V, single-layer MoS2 enters a metallic state, with increasing conductivity as the temperature decreases. Inset: Optical image of the device9 (scale bar 5 μm). (b) Field-effect mobility $\mu$ as a function of temperature. Above $\sim$100 K, $\mu$ decreases due to phonon scattering. Adapted from ref 9. Copyright 2013 Nature Publishing Group.
transition (MIT). Figure 2b depicts the temperature dependence of the extracted field-effect mobility for a single layer of MoS$_2$. In the range of 100 and 300 K, the mobility is phonon-limited and can be fitted to the power law $\mu \propto T^{-\gamma}$, with $\gamma$ in the range of 0.55–0.78. This result is much smaller than the theoretical prediction for a single-layer ($\gamma \approx 1.52$, ref 10) or bulk crystals ($\gamma \approx 2.6$, ref 5). The latter indicates that in addition to the quenching of the homopolar phonon mode, other mechanisms might influence the mobility of single-layer MoS$_2$ encapsulated in the dual-gated geometry.\textsuperscript{11}

In a later study, Lembke and Kis\textsuperscript{12} demonstrated high-performance single-layer MoS$_2$ transistors (Figure 3) with improved performance due to full-channel gating (Figure 3b). The resulting transconductance is shown in Figure 3c with the maximum of 34 $\mu$S/\mu m, while the output curves are shown in Figure 3d. Above $V_{ds} = 1$ V, the current carrying capacity of the MoS$_2$ charge-carrying channel is saturating, making this the first observation of drain current saturation in single-layer MoS$_2$ FETs, with the drain–source conductance close to zero ($g_{ds} < 2 \mu$S/\mu m). These results show that MoS$_2$ could be interesting not only for applications in digital electronics\textsuperscript{13} but also for analogue applications where it could offer gain >10. Lembke and Kis also studied the breakdown current density in MoS$_2$, which is close to $5 \times 10^7$ A/cm$^2$. Considering that MoS$_2$ is a semiconductor, its extracted maximum current density is extremely high and 50 times larger than the limit set for metals by electromigration. This is due to strong intralayer Mo–S covalent bonds,\textsuperscript{14,15} which are much stronger than metallic bonds. The study of Lembke and Kis demonstrates that single-layer MoS$_2$ could be used in analogue circuit applications where it could provide power gain and support large current densities. High-frequency operation of MoS$_2$ FETs was also recently demonstrated by Krasnozhon et al. (Figure 4), who demonstrated current, voltage, and power gains in the gigahertz range in transistors based on 1–3 layer thick MoS$_2$ with a channel length of 240 nm.\textsuperscript{16} This shows that the mobility of MoS$_2$ is large enough to allow device operation in the technologically relevant gigahertz frequency range. These performance figures are currently limited rather by the contact resistance and relatively large gate lengths, and scaling is expected to result in further improvements of the operating speed.

### MoS$_2$ Logic Circuits and Amplifiers

After the demonstration of FETs based on single-layer MoS$_2$, Radisavljevic et al.\textsuperscript{13,17} realized electronic circuits based on the dual-gate geometry, including the first MoS$_2$ inverter demonstrated in 2011. Figure 5a shows an example of a circuit consisting of six FETs connected in series and fabricated on the same piece of single-layer MoS$_2$. A similar device was successfully operated as an inverter, the fundamental building block of digital electronics consisting of two FETs connected in series (Figure 5b) with the common electrode serving as the output and the gate of a transistor as the input (Inset Figure 5c). The circuit is characterized by a voltage gain larger than 4 (Figure 5c). This indicates that inverters based on single-layer MoS$_2$ FETs can be used for the fabrication of more complex circuits, consisting of cascaded inverters, where a gain larger than unity is required. Finally, single-layer MoS$_2$ transistors were also employed for the realization of analog circuits,\textsuperscript{17} such as a small signal amplifier, with gain up to 6.9 (Figure 5d).

#### Flash Memory

Combining the high conductivity of graphene with the semiconducting properties of single-layer MoS$_2$ represents a unique opportunity for creating high-performing devices with advanced functionalities, such as all-2D memory devices. Bertolazzi et al.\textsuperscript{18} demonstrated a nonvolatile memory cell based on MoS$_2$/graphene heterostructures. The device architecture, which resembles that of a floating gate FET, its electrical connections, and its
fabrication procedure are reported in Figure 6a–c. Single-layer MoS\textsubscript{2} acts as the transistor channel, and graphene electrodes are used to inject and collect the charge carriers. A piece of multilayer graphene, \( \sim 4\) nm thick, is employed as the charge trapping layer (floating gate) and is separated from the semiconducting channel by a 6 nm thick HfO\textsubscript{2} layer, serving as the tunnel oxide. The device is completed with the deposition of the control oxide and the control gate electrode.

For the first time, it was also shown that graphene can be used as a contact material to single-layer MoS\textsubscript{2}, resulting in ohmic-like \( I_{ds} \) vs \( V_{ds} \) output characteristics at room-temperature (Figure 6d). This result suggests the possibility to exploit the tunability of the graphene work function, e.g. by means of chemical functionalization or electrostatic doping, to develop efficient contacts and highly conducting graphene interconnects in advanced electronic circuits.

Due to the atomic scale thickness of single-layer MoS\textsubscript{2} and its large bandgap, the transistor operation was found to be very sensitive to the presence of charges in the floating gate. The proof-of-concept device by Bertolazzi et al. showed a program/erase current ratio as high as \( 10^6 \) (Figure 6e), which is considerably higher than what has been reported for memory devices employing graphene as the active material. The observed large memory window (\( > 8 \) V) revealed the possibility for engineering 2D memory cells with multilevel data storage capability.

**OPTOELECTRONIC DEVICES**

The direct bandgap\textsuperscript{19,20} of single-layer MoS\textsubscript{2} makes it suitable for optoelectronic applications. First single-layer MoS\textsubscript{2} phototransistors exhibited a photoresponsivity of 7.5 mA/W.\textsuperscript{21} Multilayer MoS\textsubscript{2} devices\textsuperscript{22,23} show higher photoresponsivity of \( \sim 100 \) mA/W, comparable to Si-based photodetectors. However, relatively low mobility (0.1 cm\textsuperscript{2}/(V s)) and high contact resistance in these devices offset the inherent advantage of using a material with direct bandgap. Lopez-Sanchez et al.\textsuperscript{24} demonstrated an ultrasensitive single-layer MoS\textsubscript{2} photodetector due to improved mobility and contact quality (Figure 7). The output curves depicted in Figure 7b show an increase of drain current by several orders of magnitude as the device is illuminated, proving photosresponse. As shown in Figure 7c, in the dark the device shows behavior typical of n-type MoS\textsubscript{2} FETs. When the device is illuminated, the device current increases for both the on and off states. This indicates that photocurrent dominates over thermionic and tunneling currents. Figure 7d shows the device photoresponsivity, reaching 880 A/W. Furthermore, the presence of a bandgap and high degree of electrostatic control allows the phototransistor to be efficiently turned off, resulting in a noise level lower than in commercial state-of-the-art Si avalanche photodiodes.

Sundaram et al.\textsuperscript{25} demonstrated that single-layer MoS\textsubscript{2} can be used as light emitter. While their study served as a proof-of-principle, the device was limited by a relatively high power...
threshold for light emission and only a small portion of the
device, restricted to the contact edge, was active in electro-
luminescence.25 A way to overcome these limiting factors is to
build light-emitting diodes based on vertical pn-junctions,
resulting in an increase of the junction’s area. Lopez-Sanchez
et al.26 demonstrated a new geometry in the form of a vertical
heterostructure composed of n-type MoS2 serving as electron
injection layer and p-type Si serving as the hole injection layer
(Figure 8). Figure 8a shows the current versus bias character-
istic of the MoS2/Si heterojunction diode (Inset Figure 8a),
exhibiting rectifying behavior. This shows that classical diodes
and related optoelectronic devices could be fabricated using a
combination of atomically thin 2D and classical 3D semi-
conductors. The electroluminescent emission intensity map is
shown in Figure 8b. Most of the heterojunction surface is active
in light emission. This is attractive from the practical point of
view because the total emitted light intensity can be scaled up
easily by increasing the device area. Figure 8c shows the
observed electroluminescence spectrum, closely matched to the
MoS2 PL spectrum. This shows that the relevant energy for the
radiative recombination process in the MoS2/Si heterojunction
is the direct bandgap in single-layer MoS2. The heterojunction
diode can also convert incident light into electrical current with
an external quantum efficiency of 4.4% and a broad spectral
response.26

MECHANICAL PROPERTIES

Measurements of mechanical properties of MoS2 showed that
these materials also have a strong potential in high-end flexible
and transparent electronics. Bertolazzi et al.15 reported the
first measurements of the in-plane stiffness and breaking strength
of single and bilayer sheets of MoS2 by means of nanoindentation
experiments performed with the tip of an atomic-force micro-
scope (AFM) on free-standing MoS2 membranes (Figure 9).
Such measurements provide simultaneous access to different
material properties, such as the 2D Young’s modulus ($E^{2D}$),
the pre-tension ($\sigma^0_{2D}$) of the membrane, and its breaking
strength ($\sigma_{\text{MAX}}$), which in the case of defect-free material is
directly related to the intrinsic strength of the interatomic
bonds.

Free-standing MoS2 membranes were fabricated on SiO2/Si
substrates by transferring atomically thin sheets of MoS2 over
an array of microfabricated circular holes, as shown in Figure
9a,b. During the indentation experiment, the AFM tip was
pushed against the center of the membrane, and the cantilever
deflection $\Delta z$ was monitored, resulting in force $F$ versus
membrane-deflection $\delta$ curves such as those shown in Figure 9c.
A semiempirical model was used to describe the curves, which
assumes that the suspended MoS2 sheets behave as a clamped
regular membrane, made of a linear isotropic elastic material.
Fitting the experimental curves with this model allows the
extraction of the 2D membrane pre-tension $\sigma^{2D}$ (in the 0.02–
0.2 N/m range) and of the Young’s modulus $E = ((E^{2D})/I)$ of

Figure 6. Flash memory device based on MoS2/graphene heterostructures. (a) Three dimensional representation of the MoS2/graphene floating
gate transistor. (b) Schematic of the device comprising a single-layer MoS2 semiconducting channel, graphene contacts, and multilayer graphene
floating gate. (c) Optical micrographs of two floating gate transistors at different stages of the fabrication process. (d) Output curves of the floating
gate transistor showing room-temperature ohmic-like behavior for graphene contacts to single-layer MoS2. (e) Time-dependent behavior of the
drain–source current for Program and Erase states. The magnitude of the Erase current at low bias ($V_{ds} = 50$ mV) is more than 3 orders of
magnitude higher than the Program current, due to the high electrostatic sensitivity of the 2D channel to the charges trapped in the floating gate.
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the ultrathin MoS$_2$ sheets. The membranes were further deformed up to the breaking point and the forces at which the rupture occurred were recorded (black and blue crosses in Figure 9c). After the occurrence of the breaking event, a hole was produced in the center of the membrane, at the location where the AFM tip punctured the MoS$_2$ sheet. Figure 9d,e shows intermittent-contact AFM topography images of the same free-standing MoS$_2$ membrane before and after performing the stretching and breaking experiments.

It was found that the Young’s modulus of single-layer MoS$_2$ is $\sim$270 GPa, higher than that of steel ($\sim$205 GPa). The Young’s modulus of single-layer MoS$_2$ was found to be higher than that of bulk MoS$_2$ ($\sim$240 GPa), most likely due to defects, interlayer sliding, or the absence of stacking faults in single-layers. Castellanos-Gomez et al.$^{27}$ performed similar experiments on multilayer MoS$_2$ with thickness ranging from 5 to 25 layers and obtained a Young’s modulus of 330 $\pm$ 70 GPa. Recent quantum-chemical molecular dynamics simulations performed by Lorenz et al.$^{28}$ show excellent agreement with the experimental observations of Bertolazzi et al. First-principles calculations by Li$^{29}$ also predicted that ultrathin sheets of MoS$_2$ are characterized by a biaxial elastic modulus of $\sim$250 GPa.

The breaking strength of single- and bilayer MoS$_2$ was between 6% and 11% of their Young’s modulus, which represents the upper theoretical limit of a material’s breaking strength and reflects the intrinsic properties of its interatomic bonds.$^{30}$ This suggests that MoS$_2$ can be suitable for integration with flexible plastic substrates, such as polyimide, whose maximum strain before failure is $\sim$7%.

All these results indicate that single-layer MoS$_2$ is at present the strongest and thinnest semiconducting material ever measured. Its ultimate thickness, low mass, and semiconducting energy gap could be exploited for building nanoelectromechanical systems (NEMS) with new functionalities. One first example was given by Castellanos-Gomez et al.$^{30}$ who demonstrated mechanical resonators based on single-layer MoS$_2$ membranes operating in the tension-dominated regime with resonance frequencies up to 30 MHz and a quality factor $\sim$55.

Further research is needed in order to explore the mechanical properties of other TMDs characterized by different electronic/optical properties that may be complementary to those of MoS$_2$.

Single-layer MoS$_2$, combined with conducting graphene and insulating BN, could be employed as semiconducting channel in all-2D FETs integrated on flexible substrates. Lee et al.$^{31}$ recently demonstrated MoS$_2$ FETs on flexible polymer substrates, which incorporate heterostructures of graphene and BN and show unaffected electronic mobility at strain levels of up to 1.5%. These results are a consequence of the excellent mechanical properties of 2D materials, which stem from their strong in-plane interatomic bonds.
The thermal and thermoelectric properties of single-layer MoS₂ have been investigated to assess its potential for applications in thermal nanodevices, such as on-chip power generators and nanosystems for waste thermal energy harvesting. Using temperature-dependent Raman spectroscopy on suspended sheets, Yan et al. extracted the thermal conductivity of single-layer MoS₂ (≈ 34.5 W/mK), which was found to be significantly lower than that of graphene. Buscema et al. measured the Seebeck coefficient in back-gated MoS₂ transistors and obtained a large tunable value, which could be varied between $-4 \times 10^2 \mu$V/K and $-1 \times 10^5 \mu$V/K by a gate voltage. This result shows that single-layer MoS₂ can be suitable for the aforementioned thermoelectric applications. Recent theoretical calculations by Huang et al. also confirm that 2D TMDs are promising thermoelectric materials.

## PROSPECTS AND CHALLENGES

In the coming years, several issues need to be addressed for MoS₂ devices to become mature. There is at this point no control over intrinsic doping levels in MoS₂. Possible doping strategies include the introduction of atoms (F, Cl, Re, etc.) substituting atoms of the crystal’s lattice (Mo, S) during the growth process as well as doping by adsorption of donor or acceptor molecules. Controlled doping will open up the route to achieve p-type or ambipolar behavior in single-layer MoS₂. Understanding how to make good electrical contacts to MoS₂ is also lagging. One strategy to reduce the contact resistance in 2D MoS₂/metal junctions could be based on heavily doping of the contact region by donor molecules. Another possibility is based on phase engineering of the MoS₂ region below the contact, for example, by converting the semiconductor into metallic 1T phase by treatment with n-butyl lithium. Another problem that needs to be addressed is related to atomic-layer deposition (ALD) of high-$\kappa$ dielectrics in extremely scaled devices. Since the presence of hydroxyl groups is crucial for the quality of the ALD deposition and pristine MoS₂ lacks out-of-plane covalent functional groups, surface functionalization might be required to fabricate extremely scaled devices. On the basis of theoretical predictions there is still much opportunity for improving the charge carrier mobility. Therefore, reducing charge impurity scattering centers at the dielectric/MoS₂ interfaces as well as within the material itself might result in devices of significantly enhanced performance.

## CONCLUSION

The discussed results indicate that MoS₂ is indeed a very promising material for practical electronic and optoelectronic applications. Figure 10 depicts a schematic comparison of the current on/off ratio, the bandgap, and the field-effect mobility of several materials suitable for electronic and optoelectronic applications, including traditional semiconductors and novel flexible materials. Each material class has its strength and weaknesses and which material to select for a given application depends on its specification. The particular advantages of single-layer MoS₂ include chemical stability, planar geometry, high bandgap, very high on/off ratio, competitive mobility, flexibility, and transparency. In combination with other 2D materials such as graphene and insulating layered materials (e.g., BN), MoS₂ adds to a repertoire suitable to construct electronic and optoelectronic heterostructures. Together with recent developments in large-scale production techniques, MoS₂ could potentially be incorporated in roll-to-roll processes to fabricate inexpensive, flexible, and transparent 2D devices.

![Figure 8. pn-Junction incorporating single-layer MoS₂. (a) Current vs bias curve of MoS₂/Si heterojunction diode studied by Lopez-Sanchez et al. Inset: Cross-sectional view of the device. Electrons are injected from n-type MoS₂, while holes are injected from the p-Si substrate. (b) Intensity map showing electroluminescent emission with superimposed outline of the device components. The entire surface of the heterojunction emits light (scale bar 5 μm). (c) Electroluminescence spectrum acquired under a forward bias $V = 15$ V ($I = 1.8 \mu$A). The spectrum is fitted with three Lorentzian lines that correspond to the A and B excitons of MoS₂ and the A⁻ trion resonance. Reproduced from ref 26. Copyright 2014 American Chemical Society.](image-url)
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Notes

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Figure 9. Nanoindentation of suspended MoS2 sheets. (a) Optical micrograph (left) of single-layer MoS2 deposited on a SiO2/Si substrate with holes. AFM image (right) showing the MoS2 flake forming a series of free-standing membranes. (b) Schematics of the AFM nanoindentation experiment. (c) Force F versus membrane-deflection, δ, curves for single- and bilayer MoS2. The breaking events occur at the points marked with the black (single-layer) and blue (bilayer) crosses. (d) AFM image of suspended single-layer MoS2 before the indentation process and (e) after the occurrence of the breaking event. Adapted with permission from ref 15. Copyright 2011 American Chemical Society.

Figure 10. Schematic comparison of single-layer MoS2 with other materials suitable for electronic and optoelectronic applications. Field-effect mobility as a function of (a) current on/off ratio and (b) electronic bandgap. For practical applications, an on/off ratio on the order of 10⁴−10⁷ is desirable. We compare the properties of graphene (including its bilayer), multilayered (10 nm thick) black phosphorus (P), single walled carbon nanotubes (SWNTs), few-layered transition metal dichalcogenides (TMDs), organic semiconductors, and several traditional semiconductors with single-layer MoS2 (red star).

Biographies

Dominik Lembke was born in 1984 in Munich, Germany. He graduated from Ludwig Maximilian University Munich with a Diplom in Physics (2010) and holds a Technology Management degree from Technical University Munich (2008). He is currently pursuing research towards his Ph.D. in Physics in the laboratory of Prof. Andras Kis at EPFL. His work focuses on electrical properties and applications of two-dimensional MoS2 transistors. His general research interests include novel materials suitable for flexible, transparent electronics.

Simone Bertolazzi was born in 1985 in Verona, Italy. He received the Laurea degree in 2007 from Politecnico di Milano (Engineering Physics) and M.Sc. degree in 2010 from both Politecnico di Milano and École
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Andras Kis was born in 1975 in Zagreb, Croatia. He graduated in Physics from the University of Zagreb in 1999 and obtained a Ph.D. in Physics from EPFL (École Polytechnique Fédérale de Lausanne) in 2003 after which he worked as a postdoctoral researcher at the University of California, Berkeley, in the group of Prof. Zettl. He has joined EPFL as faculty in 2008 when he formed a new research group with the aim of studying 2D materials beyond graphene, with a strong emphasis on transition metal dichalcogenide based electronics.

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