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Algebraic Nahm transform for parabolic Higgs bundles on $\mathbb{P}^1$

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We formulate the Nahm transform in the context of parabolic Higgs bundles on $\mathbb{P}^1$ and extend its scope in completely algebraic terms. This transform requires parabolic Higgs bundles to satisfy an admissibility condition and allows Higgs fields to have poles of arbitrary order and arbitrary behavior. Our methods are constructive in nature and examples are provided. The extended Nahm transform is established as an algebraic duality between moduli spaces of parabolic Higgs bundles. The guiding principle behind the construction is to investigate the behavior of spectral data near the poles of Higgs fields.

14H60; 14E05, 14J26

Introduction

This article brings a new geometric point of view to the Nahm transform for Higgs bundles. This new outlook is the analogue of the Fourier–Laplace transform for $\mathcal{D}$–modules (see Malgrange [7]) in the Dolbeault complex structure. For a survey of different aspects of the Nahm transform, see Jardim [6].

In this article, we limit ourselves to Higgs bundles $(\mathcal{E}, \theta)$ over the projective line consisting of:

- An effective divisor $\mathtt{D}$.
- A vector bundle $\mathcal{E}$.
- A morphism of coherent sheaves $\theta\colon \mathcal{E} \to \mathcal{E}(\mathtt{D})$, called *the Higgs field*.

The fundamental property of Higgs bundles is that they are reconstructible from the set of eigenvalues and eigenspaces of the endomorphism ϕ, which we call the *spectral data*. The standard construction formulates the spectral data corresponding to a Higgs bundle on a curve C as a divisor for the set of eigenvalues and a pure sheaf of homological dimension 1 for the eigenspaces in *a geometrically ruled surface over C*. This surface is *never* a trivial fibration unless the Higgs bundle is of no interest. When $C = \mathbb{P}^1$, the

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geometrically ruled surface is called a Hirzebruch surface, whose base C and fibers are all copies of $\mathbb{P}^1$.

In these circumstances, one is tempted to ask:

Question 1 Can the roles of the base and the fibers be reversed?

If the answer is affirmative, we will obtain a transform of Higgs bundles.

Nevertheless, the answer is *no* unless the Hirzebruch surface is a product surface.

In [10], the second author follows an analytical route to answer this question and describes the Nahm transform for parabolic Higgs bundles with a Hermitian–Einstein metric on the complex projective line $\mathbb{P}^1$ satisfying some semisimplicity and admissibility conditions. These parabolic Higgs bundles have at most regular singularities in points at finite distance and an irregular (Poincaré rank 1) singularity at infinity. Then the Nahm transform of a parabolic Higgs bundle $(\mathcal{E}, \theta)$ is defined by the following steps:

(1) Construct an eigensheaf $M^\flat$ on an open subset U of $\mathbb{P}^1 \times \widehat{\mathbb{P}}^1$.
(2) Push $M^\flat$ by the projection $\hat{\pi}\colon \mathbb{P}^1 \times \widehat{\mathbb{P}}^1 \to \widehat{\mathbb{P}}^1$.
(3) Choose the "right" extension $\widehat{\mathcal{E}}$ of $\hat{\pi}_*(M^\flat)$ to $\widehat{\mathbb{P}}^1$.

Here, $\widehat{\mathbb{P}}^1$ refers to the target projective line for the transform. For more details, we refer the reader to Section 2.

Since the answer to the above question is negative, we ask a different question in this paper:

Question 2 Can one define the spectral data corresponding to a Higgs bundle on $\mathbb{P}^1 \times \widehat{\mathbb{P}}^1$?

We will show that the answer to this question is *yes*.

In the spirit of nonabelian Hodge theory, it is important to understand whether this transform respects the parabolic structures on Higgs bundles. Although it doesn't for *arbitrary* parabolic Higgs bundles, this transform respects an important subclass of parabolic Higgs bundles satisfying a natural condition:

Admissibility condition At a parabolic puncture, either there is no jump corresponding to the weight zero, or otherwise, the jump is exactly prescribed by the residue of the Higgs field.

This condition will reappear in different, yet equivalent, forms throughout the paper (Admissibility conditions 1 and 2). For a more precise formulation of admissibility, we refer the reader to Section 1.

The main results of the present paper require only the admissibility condition on parabolic Higgs bundles; there is no further restriction on either

- the order of poles of the Higgs fields θ,
- the character of the Higgs fields θ.

For instance, *any* coefficient in the Laurent expansion of the Higgs field θ is allowed to have nilpotent part for any eigenvalue except for the eigenvalue 0. However if the eigenvalue of the most polar part of the Higgs field is 0 then our method needs to be refined both analytically and sheaf-theoretically to take into account the weight filtration corresponding to the action of the residue on the 0–eigenspace, along the lines of T Mochizuki's work in [9, Section 5.1].

Another attribute of the theory developed in the present paper is that despite being motivated by the Nahm transform in differential geometry, the results hold true over an arbitrary field and are not restricted to the complex number field, except for Theorem 8.5, which is used to compare our work with the main construction of [10].

Yet another advantage is that our methods are very explicit and allow one to calculate concrete examples; an important feature in a field where new examples are highly sought.

Along the way, we also introduce the notion of a *proper transform* to relate coherent sheaves on a scheme X to coherent sheaves on a blow-up of scheme X. We establish its basic properties required for the purposes of the present paper. We suspect that this new operation might be of independent interest in algebraic geometry.

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1 Outline of the paper

In Section 1.1 we explain the notation and the notions we use and the conditions under which our results hold. In Section 1.2 we then describe briefly the contents of the paper.

1.1 Notation

Let X be a projective scheme over a field $\boldsymbol{K}$. Given a birational morphism $\omega\colon X' \to X$, denote the total and proper transforms of a Cartier divisor $\mathtt{P}$ by $\omega^*\mathtt{P}$ and $\omega^\sigma\mathtt{P}$ respectively.

The vanishing locus of a global section t of a line bundle L on X is denoted by (t).

Let $\mathtt{P}$ be an effective Cartier divisor on X. Denote:

- $\mathbb{P}_X(\mathcal{O}\oplus\mathcal{O}(-\mathtt{P})) := \operatorname{Proj}(\operatorname{Sym}^\bullet(\mathcal{O}\oplus\mathcal{O}(-\mathtt{P}))^\vee)$ by $Z^{\mathtt{P}}$.
- The structure morphism by $\pi_{\mathtt{P}}\colon Z^{\mathtt{P}} \to X$.
- The relative hyperplane bundle by $\mathcal{O}_{Z^{\mathtt{P}}}(1)$.
- The canonical section of $\mathcal{O}_{Z^{\mathtt{P}}}(1)$ by $y_{\mathtt{P}}$.
- The canonical section of $\mathcal{O}_{Z^{\mathtt{P}}}(1)\otimes\mathcal{O}(\mathtt{P})$ by $x_{\mathtt{P}}$.
- The automorphism acting on $Z^{\mathtt{P}}$ by $(x_{\mathtt{P}}, y_{\mathtt{P}}) \mapsto (-x_{\mathtt{P}}, y_{\mathtt{P}})$ by $(-1)_{Z^{\mathtt{P}}}$.

We refer to the divisor $(y_{\mathtt{P}})$ as the infinity section and the divisor $(x_{\mathtt{P}})$ as the 0–section of $Z^{\mathtt{P}}$.

Definition 1.1 A *Higgs sheaf* $(\mathcal{E},\theta)$ on X consists of a coherent sheaf $\mathcal{E}$ on X and a homomorphism $\theta\colon \mathcal{E}\to\mathcal{E}(\mathtt{P})$ for an effective Cartier divisor $\mathtt{P}$.

As the divisor $\mathtt{P}$ controls where the Higgs field θ is allowed to have poles, we refer to it as the *polar divisor* and the points in its support as *polar points*.

A Higgs sheaf on a projective scheme X determines a unique coherent sheaf $M^{\mathtt{P}}$ on the surface $Z^{\mathtt{P}}$ so that $\dim M^{\mathtt{P}} = \dim\mathcal{E}$, $\operatorname{Supp} M^{\mathtt{P}}\cap(y_{\mathtt{P}}) = \varnothing$ and $\pi_{\mathtt{P}*}M^{\mathtt{P}} = \mathcal{E}(\mathtt{P})$. The sheaf $M^{\mathtt{P}}$ is called the *eigensheaf* corresponding to the Higgs sheaf $(\mathcal{E},\theta)$. The support of $M^{\mathtt{P}}$ is the *spectral scheme*. The sheaf $M^{\mathtt{P}}$ fits into an exact sequence

$$0 \longrightarrow \mathcal{E} \xrightarrow{x_{\mathtt{P}}-y_{\mathtt{P}}\theta} \mathcal{E}(\mathtt{P})\otimes\mathcal{O}_{Z^{\mathtt{P}}}(1) \longrightarrow M^{\mathtt{P}} \longrightarrow 0.$$

Let

$$\pi^H(\mathcal{E},\theta) := M^{\mathtt{P}}.$$

Conversely, let the Higgs bundle $\theta\colon \mathcal{E}\to\mathcal{E}(\mathtt{P})$ be the pushforward of the following sequence by $\pi_{\mathtt{P}}$:

$$\mathrm{Id}_M \otimes x_{\mathtt{P}}\colon M(-\mathtt{P}) \longrightarrow M\otimes\mathcal{O}_{Z^{\mathtt{P}}}(1).$$

Denote $(\mathcal{E},\theta)$ by $\pi_H(M, x_{\mathtt{P}})$, or simply by $\pi_H(M)$. It is clear that π^H and π_H are quasi-inverses. The correspondence extends to parabolic objects (see Definitions 3.5

and 3.15) in a straightforward manner. We refer to this construction as *the standard construction* and the resulting objects as the standard eigensheaf, the standard spectral cover etc.

Given a Higgs bundle $(\mathcal{E},\theta)$ on $\mathbb{P}^1$ (ie a Higgs sheaf with $\mathcal{E}$ the sheaf of holomorphic sections of a vector bundle) with polar divisor $\mathtt{P}=p_1+\cdots+p_n$ where $p_1,\ldots,p_n\in\mathbb{C}$, define the locally free sheaf

$$\mathcal{F}:=\ker\bigl(\mathcal{E}\to\operatorname{coker}(\theta(-\mathtt{P})_{\mathtt{P}})\bigr)$$

on $\mathbb{P}^1$. Then θ gives rise to well-defined residue maps

$$\operatorname{res}(\theta,p_j):\operatorname{End}(\mathcal{E}_{p_j}).$$

(Notice here that with respect to the usual terminology, we contract with the section $\partial/\partial z$ of the holomorphic tangent bundle; see also Remark 2.1.) The sheaf $\mathcal{F}$ can then be described as the sheaf of local sections of $\mathcal{E}$ whose evaluation at the polar divisor vanishes on the *generalized* 0–eigenspace of the residue. We suppose furthermore that a compatible *parabolic structure* is given at all p_j. Throughout the paper we will use the notion of parabolic structure in various degrees of generality. In this section, we will content ourselves by using the most classical such notion, first defined by Mehta and Seshadri [8]: this means the data of real numbers

$$0\le\alpha_0^j\le\cdots\le\alpha_{r-1}^j<1$$

(called *parabolic weights*) and a finite decreasing filtration

$$\{0\}=(\mathcal{E}_{p_j})_1\subset(\mathcal{E}_{p_j})_{\alpha_{r-1}^j}\subseteq\cdots\subseteq(\mathcal{E}_{p_j})_{\alpha_{r_j}^j}\subseteq(\mathcal{E}_{p_j})_0=\mathcal{E}_{p_j}\tag{1}$$

of the fiber of $\mathcal{E}$ at p_j, preserved by the residue $\operatorname{res}(\theta,p_j)$. Here, we denoted by r_j the smallest index k such that $\alpha_k^j>0$ (when such an index exists; $r_j=r$ otherwise). Let us introduce the notation

$$\operatorname{gr}_{\alpha_k^j}\mathcal{E}_{p_j}=(\mathcal{E}_{p_j})_{\alpha_k^j}\big/(\mathcal{E}_{p_j})_{\alpha_{k+1}^j}$$

for the graded vector spaces of the parabolic filtration, where by α_r^j we mean 1. For $r_j\le k<r$ we also require that $\operatorname{gr}_{\alpha_k^j}\mathcal{E}_{p_j}=0$ if and only if $\alpha_k^j<\alpha_{k+1}^j$. Also, let $\alpha_k(\mathcal{E}_p)$ stand for α_k^j for $p=p_j$.

The results of Sections 10 and 11 hold under the following admissibility condition for the parabolic structure:

Admissibility condition 1 (for Higgs bundles) For any polar point $p \in \mathtt{P}$, the Higgs bundle $(\mathcal{E}, \theta)$ satisfies one of the following conditions:

- $\alpha_0(\mathcal{E}_p) > 0$ and $\mathcal{E}_p = \mathcal{F}_p$.
- $\alpha_0(\mathcal{E}_p) = 0$ and $F_1\mathcal{E}_p = \mathrm{im}(\mathcal{F}_p \to \mathcal{E}_p)$ and the residue of θ acts on its 0–eigenspace with no nilpotent part.

1.2 Results

The paper is organized along the following lines: in Section 2, we give an overview of the conditions and results of [10] which are often referred to in the present paper.

In Section 3, we recall the notions which will be used throughout the paper: pure sheaves of dimension 1, parabolic sheaves, parabolic Euler characteristic, degree and stability of parabolic sheaves. We also prove some of their properties.

In Section 4, we define an *iterated* version of blow-up maps for nonreduced 0–dimensional subschemes. This will be essential for the generalization of the Nahm transform to Higgs bundles with higher-order poles.

In Section 5, analogous to the proper transform of a divisor with respect to a blow-up, we introduce the *proper transform of a coherent sheaf* with respect to a blow-up of a closed point. We study properties of the proper transform for 1–dimensional pure sheaves on surfaces. For such sheaves, the proper transform is related to the Hecke transform of locally free sheaves. In particular, for such sheaves, the proper transform is a quasi-inverse of the direct image (Lemma 5.13), and it preserves the Euler characteristic (Lemma 5.15). We also give a parabolic version of the proper transform, and prove that it preserves the parabolic Euler characteristic (Section 5.6).

In Section 6, we define two operations to modify the divisors of parabolic sheaves: *Deletion* along $\mathtt{E}$ removes an effective subdivisor $\mathtt{E}$ of the parabolic divisor, whereas *addition* along $\mathtt{E}$ appends an effective divisor $\mathtt{E}$ to the parabolic divisor. Under an assumption (which is equivalent to Admissibility condition 1), these operations are inverses of each other. Moreover, they preserve the parabolic Euler characteristic (Proposition 6.2).

In Section 7, we introduce what we call *spectral triples*, consisting of a smooth surface, flat over a base curve, an effective divisor on the surface, flat and finite over the base curve, and a coherent sheaf supported on the divisor, of homological dimension 1. Notice that the operation π^H associates to a given Higgs bundle a spectral triple $(Z^{\mathtt{P}}, \mathrm{Supp}(M^{\mathtt{P}}), M^{\mathtt{P}})$. We shall call this spectral triple the *standard spectral triple* associated to the Higgs bundle. However, there is another way of defining a spectral triple $(Z^0, \mathrm{Supp}(M^0), M^0)$ for the surface $Z^0 = \mathbb{P}^1 \times \widehat{\mathbb{P}}^1$: we call this the *naive spectral triple*. The surfaces Z^{P} and $\mathbb{P}^1 \times \widehat{\mathbb{P}}^1$ are related by a series of elementary transformations. Let Z be the resolution of indeterminacies of $Z^0 \dashrightarrow Z^{\mathrm{P}}$. We show that the proper transforms of M^{P} and M^0 agree on Z (Proposition 7.7).

In Section 8, we construct the Nahm transform of parabolic Higgs bundles on the projective line as a composition of the operations we have introduced thus far. The starting point is the diagram

$$
\begin{tikzcd}
 & & Z^{\mathrm{int}} \arrow[dll, "\boldsymbol{\rho}_{\mathrm{P}}"'] \arrow[d] \arrow[drr, "\widehat{\boldsymbol{\rho}}_{\widehat{\mathrm{P}}}"] & & \\
Z^{\mathrm{P}} \arrow[d] & & \mathbb{P}^1 \times \widehat{\mathbb{P}}^1 \arrow[dll] \arrow[drr] & & \widehat{Z}^{\widehat{\mathrm{P}}} \arrow[d] \\
\mathbb{P}^1 & & & & \widehat{\mathbb{P}}^1
\end{tikzcd}
\tag{2}
$$

(see (27)). Here the maps $\boldsymbol{\rho}_{\mathrm{P}}$ and $\widehat{\boldsymbol{\rho}}_{\widehat{\mathrm{P}}}$ are blow-up maps and Z^{int} is called the *intermediate spectral surface*. Starting from a 1–dimensional parabolic sheaf $M^{\mathrm{P}}_\bullet$ on Z^{P}, the Nahm transform produces a 1–dimensional parabolic sheaf $\widehat{M}^{\widehat{\mathrm{P}}}_\bullet$ on $\widehat{Z}^{\widehat{\mathrm{P}}}$ by the formula

$$M^{\mathrm{P}}_\bullet \mapsto (-1)^*_{\widehat{Z}^{\widehat{\mathrm{P}}}} \widehat{M}^{\widehat{\mathrm{P}}}_\bullet = (-1)^*_{\widehat{Z}^{\widehat{\mathrm{P}}}} (\widehat{\boldsymbol{\rho}}_{\widehat{\mathrm{P}}})_* \operatorname{Add}_{\widehat{\mathrm{E}}^+} \operatorname{Del}_{\mathrm{E}^+} (\boldsymbol{\rho}_{\mathrm{P}})^{\sigma} (M^{\mathrm{P}}_\bullet).$$

From right to left, this formula reads as a proper transform with respect to $\boldsymbol{\rho}_{\mathrm{P}}$, deletion along a divisor E^+, addition along a divisor $\widehat{\mathrm{E}}^+$, pushforward with $\widehat{\boldsymbol{\rho}}_{\widehat{\mathrm{P}}}$, and pullback with respect to the fiberwise (-1) multiplication. Here, E^+ and $\widehat{\mathrm{E}}^+$ are suitably chosen divisors related to the birational morphisms $\boldsymbol{\rho}_{\mathrm{P}}$ and $\widehat{\boldsymbol{\rho}}_{\widehat{\mathrm{P}}}$ respectively. Theorem 8.5 shows that our construction generalizes that of [10].

In Section 9, we describe two examples in which we use our method to compute the transformed Higgs bundle explicitly. These examples are beyond the scope of [10]. The first example features a Higgs field with a nilpotent residue. The second one is an example of a higher-order pole.

Section 10 provides a geometric proof of the fact that the transformation is involutive up to a sign.

In Section 11, we study the map induced by the Nahm transform on the moduli spaces of stable Higgs bundles of degree 0 with prescribed singularity behavior. First, we compute the dimension of these moduli spaces (Lemma 11.1). Then, we show that the Nahm transform preserves the parabolic degree, and for Higgs bundles of degree 0, it preserves stability (Lemma 11.3). Finally, in Theorem 11.4 we prove that the Nahm transformation induces a Kähler isometry between the corresponding Dolbeault moduli spaces. It is also true that the Nahm transform respects the de Rham complex structure (see Remark 11.5); however, since in this paper we are only concerned with the Dolbeault complex structure, we do not give a proof of this fact here.

2 An overview of the analytic Nahm transform

In this section, we give a summary of the results of [10] relevant for the present paper.

Let $\mathtt{P} = \{p_1, \dots, p_n\}$ be a finite set in $\mathbb{P}^1$ composed of distinct points at finite distance, $\mathcal{E}$ be a rank-r holomorphic vector bundle on $\mathbb{P}^1$ and

$$\theta\colon \mathcal{E} \to \mathcal{E} \otimes \mathcal{O}_{\mathbb{P}^1}(\mathtt{P})$$

be a holomorphic map (called the *Higgs field*), where $\mathcal{O}_{\mathbb{P}^1}(\mathtt{P})$ is the sheaf of meromorphic functions with at most simple poles in the points of $\mathtt{P}$ and no other poles.

Remark 2.1 This definition differs from the usual definition of the Higgs field, which is an endomorphism-valued 1–form. However, in the particular case of the projective line, tensor multiplication by the globally defined meromorphic 1–form dz establishes an isomorphism between the sheaf of holomorphic functions and the sheaf of 1–forms with a double pole at infinity. Therefore, the above definition is equivalent to that of a Higgs field in the usual sense, with logarithmic singularities at $\mathtt{P}$ and a double pole at infinity. Throughout the paper we tacitly use this isomorphism to simplify notation.

By definition, the residue of the Higgs field is *semisimple* if the endomorphism induced by $\operatorname{res}(\theta, p_j)$ on each graded vector space of the parabolic filtration is semisimple. In this section, in addition to admissibility we will assume that at any $p_j \in \mathtt{P}$ the Higgs field has semisimple residue satisfying the following additional properties:

(1) $\operatorname{res}(\theta, p_j)$ vanishes on $\operatorname{gr}_0 \mathcal{E}_{p_j}$,

(2) for all $k \geq r_j$, the space $\operatorname{gr}_{\alpha^j_k} \mathcal{E}_{p_j}$ is one-dimensional,

(3) for $l \neq k \geq r_j$, the eigenvalues of $\operatorname{res}(\theta, p_j)$ on $\operatorname{gr}_{\alpha^j_k} \mathcal{E}_{p_j}$ and on $\operatorname{gr}_{\alpha^j_l} \mathcal{E}_{p_j}$ are different constants (in particular, for $k \geq r_j$, the residue is nonvanishing on $\operatorname{gr}_{\alpha^j_k} \mathcal{E}_{p_j}$).

Remark 2.2 As it is easy to see, Admissibility condition 1 is weaker than the above Assumptions (1)–(3).

In more explicit terms, in the standard holomorphic coordinate z of $\mathbb{C}$ and in a convenient holomorphic trivialization $\{e_0^j, \dots, e_{r-1}^j\}$ of $\mathcal{E}$ in a neighborhood of p_j the Higgs field can be written

$$\theta = B_j \frac{1}{z - p_j} + O(1), \tag{3}$$

where $O(1)$ stands for holomorphic terms and

$$B_j = \begin{pmatrix} 0 & & & & \\ & \ddots & & & \\ & & 0 & & \\ & & & \lambda_{r_j}^j & & \\ & & & & \ddots & \\ & & & & & \lambda_{r-1}^j \end{pmatrix} \tag{4}$$

is a diagonal matrix with all the λ_k^j for $r_j \leq k < r$ nonvanishing and distinct. Here, the vectors $e_0^j, \dots, e_{r_j-1}^j$ span $\mathrm{gr}_0 \mathcal{E}_{p_j}$ and for $k \geq r_j$ the vector e_k^j spans $\mathrm{gr}_{\alpha_k^j} \mathcal{E}_{p_j}$.

At infinity, we suppose that θ is holomorphic, such that its Taylor series written in the local coordinate z^{-1} and some holomorphic trivialization of $\mathcal{E}$ near infinity

$$\theta = \frac{1}{2} A + B_\infty \frac{1}{z} + O\Big(\frac{1}{z^2}\Big) \tag{5}$$

satisfy that the constant term A is diagonal with eigenvalues $\xi_1, \dots, \xi_{\hat{n}}$ of multiplicity possibly higher than one:

$$A = \begin{pmatrix} \xi_1 & & & & & & \\ & \ddots & & & & & \\ & & \xi_1 & & & & \\ & & & \ddots & & & \\ & & & & \ddots & & \\ & & & & & \xi_{\hat{n}} & \\ & & & & & & \ddots & \\ & & & & & & & \xi_{\hat{n}} \end{pmatrix} \tag{6}$$

and the first-order term B_∞ is also diagonal (in the same trivialization)

$$B_\infty = \begin{pmatrix} \lambda_0^\infty & & & & & & \\ & \ddots & & & & & \\ & & \lambda_{-1+a_1}^\infty & & & & \\ & & & \ddots & & & \\ & & & & \ddots & & \\ & & & & & \lambda_{a_{\hat{n}}}^\infty & \\ & & & & & & \ddots & \\ & & & & & & & \lambda_{r-1}^\infty \end{pmatrix}. \tag{7}$$

Here the eigenvalues $\{\lambda^{\infty}_{a_l},\dots,\lambda^{\infty}_{-1+a_{l+1}}\}$ correspond to the basis vectors spanning the ξ_l–eigenspace of A (where we have put $a_0=0, a_{\hat{n}+1}=r$). We make the assumption that for a fixed $1\le l\le \hat{n}$, none of these eigenvalues $\{\lambda^{\infty}_{a_l},\dots,\lambda^{\infty}_{-1+a_{l+1}}\}$ vanish and they are all distinct. In light of Remark 2.1, we will call $A/2$ the *second-order term* and B_∞ the *residue* of θ at infinity. Furthermore, we suppose that a parabolic structure is given at this singularity as well: that is, we are given parabolic weights $0<\alpha^{\infty}_k<1$ for $k=0,\dots,r-1$, arranged in such a way that inside one block $a_l\le k<a_{l+1}$ they form an increasing sequence, and a corresponding filtration $(\mathcal{E}_\infty)_\alpha$ for $\alpha\in[0,1]$ of the fiber of $\mathcal{E}$ over infinity, spanned by the basis elements having parabolic weight greater than or equal to α. A sheaf $\mathcal{E}$ with a parabolic structure will often be denoted by $\mathcal{E}_\bullet$.

In the rest of this section, a holomorphic vector bundle with a parabolic structure in $\mathrm{P}\cup\{\infty\}$ will be called *parabolic vector bundle*; if moreover a compatible Higgs field is given, then we will call it a *parabolic Higgs bundle*.

Denote by $\deg(\mathcal{F})$ the usual algebraic geometric degree of a holomorphic vector bundle $\mathcal{F}$. It is clear that a *subbundle* $\mathcal{F}$ (or quotient bundle $\mathcal{Q}$) of a parabolic vector bundle $\mathcal{E}$ also admits an induced parabolic structure by intersecting with $\mathcal{F}$ the terms of the filtration of $\mathcal{E}$, and assigning the biggest of the weights to all filtered terms that become isomorphic after taking intersections with $\mathcal{F}$.

Definition 2.3 The *parabolic degree* of $\mathcal{E}_\bullet$ is the real number

$$\text{par-deg}(\mathcal{E}_\bullet)=\deg(\mathcal{E})+\sum_{j\in\{1,\dots,n,\infty\}}\sum_{k=0}^{r-1}\alpha^j_k.$$

The *parabolic slope* of $\mathcal{E}_\bullet$ is the real number $\text{par-}\mu(\mathcal{E}_\bullet)=\text{par-deg}(\mathcal{E}_\bullet)/\operatorname{rk}(\mathcal{E})$. Finally, $(\mathcal{E}_\bullet,\theta)$ is said to be *parabolically stable* if for any subbundle $\mathcal{F}_\bullet$ invariant with respect to θ with its induced parabolic structure, the inequality $\text{par-}\mu(\mathcal{F}_\bullet)<\text{par-}\mu(\mathcal{E}_\bullet)$ holds.

Suppose in all what follows that $(\mathcal{E},\theta)$ is not the trivial line bundle $\mathcal{O}_{\mathbb{P}^1}$ together with a constant multiplication map. Denote by $\widehat{\mathbb{C}}$ the *dual line* of $\mathbb{C}$ (another copy of $\mathbb{C}$), and by $\widehat{\mathbb{P}}^1$ the *dual sphere*, the compactification of $\widehat{\mathbb{C}}$ by the point $\widehat{\infty}$. By [10], the Nahm transform of a stable parabolic Higgs bundle $(\mathcal{E}_\bullet,\theta)$ of parabolic degree 0 is then a parabolic Higgs bundle $(\widehat{\mathcal{E}}_\bullet,\widehat{\theta})$ on $\widehat{\mathbb{P}}^1$, with regular singularities (ie $\widehat{\theta}\,d\xi$ having simple poles) in the set $\widehat{\mathrm{P}}=\{\xi_1,\dots,\xi_{\hat{n}}\}$ and an irregular singularity (ie $\widehat{\theta}\,d\xi$ having a double pole, therefore $\widehat{\theta}$ being holomorphic) at infinity. Also, the transform of a Hermitian–Einstein metric on $(\mathcal{E}_\bullet,\theta)$ is a Hermitian–Einstein metric on $(\widehat{\mathcal{E}}_\bullet,\widehat{\theta})$; in particular, this latter is polystable. We sketch the idea of the construction of the transform. First, introduce a twist of the Higgs field: for any $\xi \in \widehat{\mathbb{C}}$ set

$$\theta_\xi = \theta - \frac{\xi}{2}\,\mathrm{Id}_{\mathcal{E}},$$

where $\mathrm{Id}_{\mathcal{E}}$ is the identity bundle endomorphism of $\mathcal{E}$. Consider now the *open spectral curve* $\Sigma^\flat$ in $(\boldsymbol{C} \setminus \mathtt{P}) \times (\widehat{\mathbb{C}} \setminus \widehat{\mathtt{P}})$ defined by

$$\Sigma^\flat = \{(z,\xi) \mid \det(\theta_\xi)(z) = 0\}.$$

In other words, denoting by $\pi^\flat$ (respectively $\widehat{\pi}^\flat$) the projection on $\boldsymbol{C} \setminus \mathtt{P}$ (respectively $\widehat{\mathbb{C}} \setminus \widehat{\mathtt{P}}$) in the product $(\boldsymbol{C} \setminus \mathtt{P}) \times (\widehat{\mathbb{C}} \setminus \widehat{\mathtt{P}})$, this curve is the support of the cokernel sheaf $M^\flat$ of the map

$$\theta_\xi \colon (\pi^\flat)^*\mathcal{E} \to (\pi^\flat)^*\mathcal{E}.$$

Propositions 4.7, 4.15 and 4.24 together with Lemma 4.30 of [10] give an interpretation of the transform on an open set.

Theorem 2.4 *The Nahm transformed Higgs bundle restricted to* $\widehat{\mathbb{C}} \setminus \widehat{\mathtt{P}}$ *can be obtained as follows:*

- *The holomorphic bundle* $\widehat{\mathcal{E}}$ *is the pushdown* $\widehat{\pi}^\flat_* M^\flat$ *endowed with its induced holomorphic structure; we denote its rank by* $\widehat{r}$*.*
- *On the open set of* $\xi \in \widehat{\mathbb{C}} \setminus \widehat{\mathtt{P}}$ *over which the fiber of* $\Sigma^\flat$ *consists of distinct points* $\{z_1(\xi), \dots, z_{\widehat{r}}(\xi)\}$ *of multiplicity* 1*, the transformed Higgs field* $\widehat{\theta}$ *acts on the subspace* $\operatorname{coker}(\theta_\xi(z_k(\xi))) \subset \widehat{\mathcal{E}}|_\xi$ *as multiplication by* $-z_k(\xi)/2$*; this then admits a unique continuation into points where the fiber has multiple points.*

This description then gives an understanding of the behavior of the Higgs field near a point of $\widehat{\mathtt{P}}$ and near $\widehat{\infty}$: we only have to understand the behavior of the open spectral curve near these points. Because of the special form of θ in the singularities, we deduce that the eigenvalues of the transformed Higgs field have indeed simple poles in the points of $\widehat{\mathtt{P}}$, and are bounded near $\widehat{\infty}$. In different terms, this defines a natural compactification $\Sigma^0 \subset \mathbb{P}^1 \times \widehat{\mathbb{P}}^1$ of $\Sigma^\flat$. Moreover, we gain precise information about its asymptotic expansions near these points: namely, near a point $\xi_l \in \widehat{\mathtt{P}}$ the residue of the transformed Higgs field in a convenient trivialization of the transformed bundle is equal to

$$-\begin{pmatrix} 0 & & & & \\ & \ddots & & & \\ & & 0 & & \\ & & & \lambda^\infty_{a_l} & \\ & & & & \ddots \\ & & & & & \lambda^\infty_{-1+a_{l+1}} \end{pmatrix};$$

ie it is the direct sum of the opposite of the residue of the original Higgs field at infinity restricted to the ξ_l–eigenspace of the leading term and a 0–matrix (see [10, Theorem 4.32]), whereas its leading term at infinity in a convenient trivialization is

$$-\frac{1}{2}\begin{pmatrix} p_1 & & & & & & \\ & \ddots & & & & & \\ & & p_1 & & & & \\ & & & \ddots & & & \\ & & & & \ddots & & \\ & & & & & p_n & \\ & & & & & & \ddots \\ & & & & & & & p_n \end{pmatrix},$$

each p_j appearing with multiplicity $\operatorname{rk}(\operatorname{res}(\theta, p_j)) = r - r_j$, and the corresponding first-order term in the same trivialization is then

$$-\begin{pmatrix} \lambda^1_{r_1} & & & & & & \\ & \ddots & & & & & \\ & & \lambda^1_{r-1} & & & & \\ & & & \ddots & & & \\ & & & & \ddots & & \\ & & & & & \lambda^n_{r_n} & \\ & & & & & & \ddots \\ & & & & & & & \lambda^n_{r-1} \end{pmatrix}$$

(see [10, Theorem 4.33]). In particular, we deduce the formula

$$\hat{r} = \sum_{j=1}^{n} \operatorname{rk}(\operatorname{res}(\theta, p_j)). \tag{8}$$

Therefore, we see an intricate interplay between singularity behavior at the regular singularities and the one at the irregular singularity.

Afterwards, we use the extensions of $\mathcal{E}$ over the singularities to define an extension M^0 of $M^\flat$ to the compactified spectral curve Σ^0. These in turn induce an extension $\widehat{\mathcal{E}}^{\mathrm{ind}}_\bullet$ of $\widehat{\mathcal{E}}$ into a holomorphic bundle endowed with a parabolic structure in each point of $\widehat{\mathrm{P}} \cup \{\widehat{\infty}\}$, which we call the *induced extension* (cf [10, Section 4.4]). By definition, a local holomorphic section of this extension has a D''_ξ–harmonic representative obtained from a local section of the cokernel sheaf M^0 multiplied with a bump function of constant height concentrated near the spectral points of ξ, such that the diameter of their support converges to 0 up to first order near the points ξ_l and to ∞ near $\widehat{\infty}$ also up to first order. Next, we compute the parabolic weights of these extensions with respect to the transformed Hermitian–Einstein metric: for a point $\xi_l \in \widehat{\mathrm{P}}$ the nonzero weights are equal to $\alpha^\infty_{a_l} - 1, \dots, \alpha^\infty_{-1+a_{l+1}} - 1$; whereas the weights at $\widehat{\infty}$ are equal to $\alpha^1_{r_1} - 1, \dots, \alpha^1_{r-1} - 1, \dots, \alpha^n_{r_n} - 1, \dots, \alpha^n_{r-1} - 1$ (cf [10, Section 4.6]). In particular, we have that all the nonzero weights violate the requirement that they be between 0 and 1: they are actually shifted by -1. Therefore, in order to get a genuine parabolic Higgs bundle on $\widehat{\mathbb{P}}^1$, we have to change the induced extension at $\widehat{\infty}$ by a factor of ξ^{-1}, and the extension of the basis vectors corresponding to nonzero eigenvalues of the Higgs field at the logarithmic singularities ξ_l by factors of $(\xi - \xi_l)$. The result we obtain this way is called the *transformed extension*, and is denoted by $\widehat{\mathcal{E}}^{\mathrm{tr}}$ (cf [10, Section 4.7]). Finally, an application of the Grothendieck–Riemann–Roch theorem yields the degree of the transformed holomorphic bundle on $\widehat{\mathbb{P}}^1$ with respect to the transformed extensions:

$$\text{par-deg}(\widehat{\mathcal{E}}^{\mathrm{tr}}_\bullet) = \text{par-deg}(\mathcal{E}_\bullet). \tag{9}$$

3 Basic material

Fix a projective scheme X over a field $\boldsymbol{K}$ with an ample invertible sheaf $\mathcal{O}_X(1)$. For a given coherent $\mathcal{O}_X$–module $\mathcal{E}$, the *support* of $\mathcal{E}$ is the closed set $\mathrm{Supp}(\mathcal{E}) = \{x \in X \mid \mathcal{E}_x \neq 0\}$. Its dimension is called the *dimension* of the sheaf $\mathcal{E}$ and is denoted by $\dim(\mathcal{E})$.

Definition 3.1 For a given coherent sheaf $\mathcal{E}$ on X, the *Euler characteristic* is defined to be

$$\chi(\mathcal{E}) = \sum_{i=0}^{\dim(X)} (-1)^i \dim_k H^i(X, \mathcal{E}).$$

The *Hilbert polynomial* $P(\mathcal{E})$ of $\mathcal{E}$ is defined by $P(\mathcal{E}, m) := \chi(E \otimes \mathcal{O}_X(m))$.

Definition 3.2 A coherent sheaf $\mathcal{E}$ of $\mathcal{O}_X$–modules on X is *pure of dimension* d if $\dim(\mathcal{F}) = d$ for any nontrivial subsheaf $\mathcal{F}$ of $\mathcal{E}$.

Equivalently, $\mathcal{E}$ is pure if and only if the associated points of $\mathcal{E}$ are all of the same dimension.

Definition 3.3 A subsheaf $\mathcal{F}$ of a pure d–dimensional sheaf $\mathcal{E}$ is *saturated* if $\mathcal{E}/\mathcal{F}$ is either 0 or pure of dimension d.

Remark 3.4 In order to be able to define parabolic sheaves $\mathcal{E}$ with respect to an effective Cartier divisor $\mathtt{D}$ on X, we need to make the assumption,

$$\dim(\mathtt{D} \cap \mathrm{Supp}\,\mathcal{E}) < \dim \mathrm{Supp}\,\mathcal{E}.$$

This assumption with the purity of the sheaf $\mathcal{E}$ ensures that $\mathcal{E} \otimes \mathcal{O}_X(-\mathtt{D}) \to \mathcal{E}$ is injective. Hence, the sheaf $\mathcal{E} \otimes \mathcal{O}_X(-\mathtt{D})$ can be seen as a *subsheaf* of $\mathcal{E}$.

The condition $\dim(\mathtt{D} \cap \operatorname{Supp} \mathcal{E}) < \dim \operatorname{Supp} \mathcal{E}$ can be viewed as a very coarse transversality condition between the support of the sheaf $\mathcal{E}$ and the divisor $\mathtt{D}$, stating that no component of $\mathtt{D}$ contains a component of the support of the sheaf $\mathcal{E}$.

In the rest of this section, we assume $\mathcal{E}$ and $\mathtt{D}$ are as in Remark 3.4.

Definition 3.5 A triple $(\mathcal{E}, F_\bullet \mathcal{E}, \alpha_\bullet)$ is called a *parabolic sheaf* on X with parabolic divisor $\mathtt{D}$ and weights $\alpha_\bullet$ if $F_\bullet \mathcal{E}$ is a filtration of $\mathcal{E}$ by coherent subsheaves $F_i \mathcal{E}$ so that

$$\mathcal{E}(-\mathtt{D}) = F_l \mathcal{E} \subset F_{l-1} \mathcal{E} \subset \cdots \subset F_1 \mathcal{E} \subset F_0 \mathcal{E} = \mathcal{E}$$

and $\alpha_\bullet$ is a sequence of real numbers with $0 \le \alpha_0 < \alpha_1 < \cdots < \alpha_{l-1} < 1$. Set $\operatorname{gr}_i^F \mathcal{E} := F_i \mathcal{E} / F_{i+1} \mathcal{E}$.

One can view $\operatorname{gr}_i^F \mathcal{E}$ as coherent sheaves on $\mathtt{D} \cap \operatorname{Supp}(\mathcal{E})$.

Definition 3.6 Let $\mathtt{D}_1$ be an irreducible component of $\mathtt{D}$ such that for any other irreducible component $\mathtt{D}'$ one has $\mathtt{D}' \cap \mathtt{D}_1 \cap \operatorname{Supp} \mathcal{E} = \varnothing$. We say that the parabolic structure is trivial on $\mathtt{D}_1$ if $\operatorname{gr}_i^F \mathcal{E}|_{\mathtt{D}_1} = 0$ for all $i > 0$ and $\alpha_0 = 0$.

Definition 3.7 The pair $(\mathcal{E}, \mathcal{E}_\bullet)$ is an $\mathbb{R}$*–parabolic sheaf* on X if $\mathcal{E}_\bullet = \{\mathcal{E}_\alpha\}$ is a collection of coherent sheaves parametrized by $\alpha \in \mathbb{R}$ satisfying the following properties:

(1) $\mathcal{E}_0 = \mathcal{E}$.

(2) For all $\alpha < \beta$, $\mathcal{E}_\beta$ is a coherent subsheaf of $\mathcal{E}_\alpha$.

(3) For all α and small $\varepsilon > 0$, $\mathcal{E}_{\alpha-\varepsilon} = \mathcal{E}_\alpha$.

(4) For all α, $\mathcal{E}_{\alpha+1} = \mathcal{E}_\alpha(-\mathtt{D})$.

Set $\operatorname{gr}_\alpha \mathcal{E} := \mathcal{E}_\alpha / \mathcal{E}_{\alpha+\varepsilon}$ for sufficiently small $\varepsilon > 0$.

Parabolic sheaves and $\mathbb{R}$–parabolic sheaves are equivalent. To see this, set $\alpha_{l+1} = 1$ and $\alpha_0 := \alpha_l - 1$. For any real number α, let i be the unique integer such that $\alpha_{i-1} < \alpha - \lfloor \alpha \rfloor \le \alpha_i$, where $\lfloor \alpha \rfloor$ is the largest integer such that $\lfloor \alpha \rfloor \le \alpha$. Set $\mathcal{E}_\alpha := F_i \mathcal{E}(-\lfloor \alpha \rfloor \mathtt{D})$. Conversely, given an $\mathbb{R}$–parabolic sheaf $\mathcal{E}_\bullet$, inductively choose $0 \le \alpha_i < 1$ for $i = 1, \dots, l$ so that $\mathcal{E}_{\alpha_i}$ properly contains $\mathcal{E}_\beta$ for any $\beta > \alpha_i$. Set $F_i \mathcal{E} := \mathcal{E}_{\alpha_i}$ and $F_{l+1} = \mathcal{E}(-\mathtt{D})$. The resulting triple $(\mathcal{E}, F_\bullet, \alpha_\bullet)$ is a parabolic sheaf on X.

Therefore, from now on we will drop the adjective $\mathbb{R}-$ from the expression $\mathbb{R}$*–parabolic structure*. In addition, when the parabolic structure of a sheaf $\mathcal{E}$ is clear from the context, we will simply write $\mathcal{E}_\bullet$ for the pair $(\mathcal{E}, \mathcal{E}_\bullet)$.

Definition 3.8 Given two parabolic sheaves $\mathcal{E}'_\bullet$ and $\mathcal{E}_\bullet$, an $\mathcal{O}_X$–module homomorphism $\varphi\colon \mathcal{E}' \to \mathcal{E}$ is a *parabolic homomorphism* if $\varphi(\mathcal{E}'_\alpha) \subset \mathcal{E}_\alpha$ for all real numbers α. For $\mathcal{E}$ a parabolic sheaf and $\mathcal{E}'$ a subsheaf endowed with a parabolic structure, we say that $\mathcal{E}'_\bullet$ is a *parabolic subsheaf* if the inclusion is a parabolic homomorphism.

Definition 3.9 Given a saturated subsheaf $\mathcal{E}'$ of a parabolic sheaf $\mathcal{E}_\bullet$, the *induced parabolic structure* ${}^{\mathrm{ind}}\mathcal{E}'_\bullet$ on $\mathcal{E}'$ is defined as ${}^{\mathrm{ind}}\mathcal{E}'_\alpha := \mathcal{E}' \cap \mathcal{E}_\alpha$ for all $\alpha \in \mathbb{R}$.

Remark 3.10 The induced parabolic structure for a given saturated subsheaf $\mathcal{E}'$ is the largest among the parabolic structures $\mathcal{E}'_\bullet$ which makes $\mathcal{E}'$ into a parabolic subsheaf of $\mathcal{E}_\bullet$. As a consequence, it suffices to consider the saturated subsheaves of a parabolic sheaf $\mathcal{E}_\bullet$ with their induced parabolic structures to measure the stability of $\mathcal{E}_\bullet$.

Definition 3.11 Define the *parabolic Euler characteristic* of a parabolic sheaf $\mathcal{E}_\bullet$ by

$$\text{par-}\chi(\mathcal{E}_\bullet) := \chi(\mathcal{E}(-\mathtt{D})) + \sum_{i=0}^{l-1} \alpha_i \chi(\mathrm{gr}_i^F \mathcal{E}). \tag{10}$$

If X is a curve, then the *parabolic degree* of $\mathcal{E}_\bullet$ is defined as

$$\text{par-deg}(\mathcal{E}_\bullet) := \deg(\mathcal{E}(-\mathtt{D})) + \sum_{i=0}^{l-1} \alpha_i \dim(\mathrm{gr}_i^F \mathcal{E}). \tag{11}$$

One can check that $\text{par-}\chi(\mathcal{E}_\bullet) = \mathrm{int}_0^1 \chi(\mathcal{E}_\alpha)\, d\alpha$ (see Yokogawa [12]).

Proposition 3.12 *The parabolic Euler characteristic is additive: Given any short exact sequence*

$$0 \longrightarrow \mathcal{E}'_\bullet \longrightarrow \mathcal{E}_\bullet \longrightarrow \mathcal{E}''_\bullet \longrightarrow 0$$

of parabolic sheaves with the same parabolic divisor $\mathtt{D}$*, we have*

$$\text{par-}\chi(\mathcal{E}_\bullet) = \text{par-}\chi(\mathcal{E}'_\bullet) + \text{par-}\chi(\mathcal{E}''_\bullet).$$

Proof Recall that a sequence of parabolic sheaves is said to be exact if for all $\alpha \in \mathbb{R}$ the induced sequence on the α–filtered terms is exact. Taking $\alpha = -1$, we see that

$$0 \longrightarrow \mathcal{E}'(-\mathtt{D}) \longrightarrow \mathcal{E}(-\mathtt{D}) \longrightarrow \mathcal{E}''(-\mathtt{D}) \longrightarrow 0$$

is exact. By additivity of the usual Euler characteristic, $\chi(\mathcal{E}(-\mathtt{D})) = \chi(\mathcal{E}'(-\mathtt{D})) + \chi(\mathcal{E}''(-\mathtt{D}))$. On the other hand, the snake lemma implies that for any $\alpha \in \mathbb{R}$ the induced sequence on the α–graded pieces

$$0 \longrightarrow \operatorname{gr}_\alpha \mathcal{E}' \longrightarrow \operatorname{gr}_\alpha \mathcal{E} \longrightarrow \operatorname{gr}_\alpha \mathcal{E}'' \longrightarrow 0$$

is also exact. The statement follows by applying additivity of χ to these sequences. □

Definition 3.13 Given a parabolic sheaf $\mathcal{E}_\bullet$ and L a line bundle, define a parabolic structure on $\mathcal{E} \otimes L$ by setting $(\mathcal{E} \otimes L)_\alpha := \mathcal{E}_\alpha \otimes L$ for all $\alpha \in \mathbb{R}$.

On a smooth projective curve X, parabolic Euler characteristic and parabolic degree are related as follows:

Proposition 3.14 *If $\mathcal{E}$ is a parabolic sheaf on a smooth projective curve X, then*

$$\begin{aligned} \textit{par-deg}(\mathcal{E}(\mathtt{D})) &= \textit{par-deg}(\mathcal{E}) + r \deg \mathtt{D}, \\ \text{par-}\chi(\mathcal{E}(\mathtt{D})) &= \textit{par-deg}(\mathcal{E}) + r\,\chi(\mathcal{O}_X). \end{aligned}$$

Proof The first formula follows by definition, because the jumps α_i of the parabolic structures of $\mathcal{E}$ and $\mathcal{E}(\mathtt{D})$ are the same, and the graded pieces of the filtration corresponding to each α_i are isomorphic.

The second follows from Riemann–Roch and the isomorphism of the graded pieces. We have

$$\begin{aligned} \text{par-}\chi(\mathcal{E}(\mathtt{D})) = \chi(\mathcal{E}) + \sum_{i=1}^{l} \alpha_i \chi(\operatorname{gr}_i^F \mathcal{E}(\mathtt{D})) &= \deg(\mathcal{E}) + r\,\chi(\mathcal{O}_X) + \sum_{i=1}^{l} \alpha_i \chi(\operatorname{gr}_i^F \mathcal{E}) \\ &= \deg(\mathcal{E}) + r\,\chi(\mathcal{O}_X) + \sum_{i=1}^{l} \alpha_i \dim(\operatorname{gr}_i^F \mathcal{E}) \end{aligned}$$

because the $\operatorname{gr}_i^F \mathcal{E}$ are supported on the 0–dimensional subscheme $\mathtt{D}$. □

Definition 3.15 A *Higgs sheaf* $(\mathcal{E}, \theta)$ consists of a coherent sheaf $\mathcal{E}$ on X together with a $\mathcal{O}_X$–module homomorphism $\theta\colon \mathcal{E} \to \mathcal{E}(\mathtt{D})$. The resulting $\mathcal{O}(\mathtt{D})$–valued endomorphism θ is called a *Higgs field.* A *parabolic Higgs sheaf* $(\mathcal{E}_\bullet, \theta)$ with divisor $\mathtt{D}$ consists of a parabolic sheaf $\mathcal{E}_\bullet$ on X with divisor $\mathtt{D}$ and a parabolic homomorphism $\theta\colon \mathcal{E}_\bullet \to \mathcal{E}_\bullet(\mathtt{D})$. A *homomorphism of Higgs sheaves* $\psi\colon (\mathcal{E}^1, \theta^1) \to (\mathcal{E}^2, \theta^2)$ is a homomorphism of sheaves $\psi\colon \mathcal{E}^1 \to \mathcal{E}^2$ commuting with the Higgs fields: $(\psi \otimes 1) \circ \theta^1 = \theta^2 \circ \psi$. A *homomorphism of parabolic Higgs sheaves* is a homomorphism of Higgs sheaves respecting the parabolic structure.

A (parabolic) Higgs subsheaf of $\mathcal{E}$ is defined in the obvious way: it is a (parabolic) subsheaf preserved by the Higgs field.

Remark 3.16 Starting from Section 8, we will consider Higgs sheaves on $X = \mathbb{P}^1$ with polar divisor $\mathtt{P}$ and parabolic divisor $\mathtt{D} = \mathtt{P} + \infty$. In terms of Definition 3.15, these objects are defined as Higgs sheaves with polar divisor $\mathtt{D}$ with an apparent singularity at ∞. In other words, the Higgs field, as a rational section with values in $\mathtt{P}$, extends regularly at ∞. We take $Z^{\mathtt{P}}$ as the standard spectral surface for a Higgs sheaf with polar divisor $\mathtt{P}$ and parabolic divisor $\mathtt{D}$. It would also be possible to work with the surface $Z^{\mathtt{D}}$: these two surfaces are related by an elementary transformation over infinity. However, we work with $Z^{\mathtt{P}}$ because the poles of the Higgs field are already contained in $\mathtt{P}$.

Definition 3.17 A parabolic (Higgs) sheaf $\mathcal{E}_\bullet$ is said to be *semistable* if for any given proper parabolic (Higgs) sheaf $\mathcal{F}_\bullet \subset \mathcal{E}_\bullet$, par-$p(\mathcal{F}_\bullet, m) \leq$ par-$p(\mathcal{E}_\bullet, m)$ for large m. The (Higgs) sheaf $\mathcal{E}_\bullet$ is said to be *stable* if for all proper parabolic (Higgs) subsheaves $\mathcal{F}_\bullet \subset \mathcal{E}_\bullet$, par-$p(\mathcal{F}_\bullet, m) <$ par-$p(\mathcal{E}_\bullet, m)$ for large m.

The standard construction described in Section 1.1 adapts to the parabolic case as well. A parabolic Higgs sheaf $\theta\colon \mathcal{E}_\bullet \to \mathcal{E}_\bullet(\mathtt{P})$ with parabolic divisor $\mathtt{D}$ determines a parabolic sheaf $M_\bullet^{\mathtt{P}}$ on $Z^{\mathtt{P}}$ with parabolic divisor $\pi_{\mathtt{P}}^*(\mathtt{D})$, with $\pi_{\mathtt{P}*} M_\alpha^{\mathtt{P}} = \mathcal{E}(\mathtt{P})_\alpha$ for any $\alpha \in \mathbb{R}$ and $\operatorname{Supp} M_\bullet^{\mathtt{P}} \cap (y_{\mathtt{P}}) = \varnothing$. Write $\pi^H(\mathcal{E}_\bullet, \theta)$ for $M_\bullet^{\mathtt{P}}$ and $\pi_H(M_\bullet^{\mathtt{P}})$ for $(\mathcal{E}_\bullet, \theta)$.

Automorphism (-1) If M corresponds to the Higgs sheaf $\theta\colon \mathcal{E} \to \mathcal{E}(\mathtt{P})$, the pullback $(-1)^*_{Z^{\mathtt{P}}} M$ corresponds to $-\theta\colon \mathcal{E} \to \mathcal{E}(\mathtt{P})$. We formalize this for parabolic Higgs sheaves as well as Higgs sheaves:

Lemma 3.18 *We have*

$$\pi_H (-1)^*_{Z^{\mathtt{P}}} \pi^H(\mathcal{E}, \theta) = \pi^H(\mathcal{E}, -\theta),$$
$$\pi_H (-1)^*_{Z^{\mathtt{P}}} \pi^H(\mathcal{E}_\bullet, \theta) = \pi^H(\mathcal{E}_\bullet, -\theta).$$

4 Iterated blow-ups

A *sequence* of infinitesimally near points $(p_0, \dots, p_n)$ on X is defined recursively as follows. Let $X_0 := X$ and $p_0 \in X_0$. Let $\omega_j\colon X_j \to X_{j-1}$ denote the blow-up of X_{j-1} at p_{j-1}, let $\mathtt{E}_j$ denote the exceptional divisor $\omega_j^{-1}(p_{j-1})$ and let p_j be a point in $\mathtt{E}_j$ for $j = 1, \dots, n$. By abuse of notation, denote the total transform of the exceptional divisor $\mathtt{E}_j$ in X_n still by $\mathtt{E}_j$ for $j = 1, \dots, n$. For $1 < j < n$, set

$$\mathtt{C}_j := \mathtt{E}_j - \mathtt{E}_{j+1}. \tag{12}$$

The curve $\mathtt{C}_j$ is a (-2)–curve on X_n.

Definition 4.1 We call a 0–dimensional closed subscheme T of a smooth (projective) surface X *linear* if for each $p \in T_{\mathrm{red}}$, one can find $u, v \in \mathfrak{m}_{X,p}$ and a positive integer n so that

$$\mathfrak{m}_{X,p} = (u, v)\mathcal{O}_{X,p} \quad \text{and} \quad \mathcal{J}_{T,p} = (u^n, v)\mathcal{O}_{X,p}.$$

For $p \in T_{\mathrm{red}}$, the integer n is uniquely determined and is equal to the length of T at p, $\dim_k(\mathcal{O}_{X,p}/\mathcal{J}_{T,p})$. Denote this integer by n_p. The total length N of T equals the sum $\sum n_p$.

An *irreducible* linear subscheme T of local length $n+1$ with closed point p determines a sequence of infinitesimally near points $(p_0, \dots, p_n)$ with $p_0 = p$ as follows. Let:

- $p_0 := p$.
- $\mathtt{D}_0 := (u^n - v)$.
- $\mathtt{D}_j := \omega_j^{\sigma} \mathtt{D}_{j-1}$.
- p_j is the unique intersection point of $\mathtt{D}_j$ with $\mathtt{E}_j$ for $j = 1, \dots, n$.

The divisor $\mathtt{D}_0$ is a smooth curve, thus so are all $\mathtt{D}_j$ for $j > 0$. Because $\mathrm{mult}_{p_0} \mathtt{D}_0 = 1$, it follows that $\mathtt{D}_j \cdot \mathtt{E}_j = 1$, ie the intersection of $\mathtt{D}_j$ and $\mathtt{E}_j$ is a unique point, say p_j, for $j > 0$.

The surface X_n is the iterated blow-up of X at T and is denoted $\omega_T \colon \text{F-Bl}_T X \to X$.

Enumerate the components of a linear subscheme as $T_1, \dots, T_m$. Then define the iterated blow-up of X at T to be

$$\text{F-Bl}_T X := \text{F-Bl}_{T_1} X \times_X \cdots \times_X \text{F-Bl}_{T_m} X$$

and $\omega_T \colon \text{F-Bl}_T X \to X$ to be the corresponding morphism.

Clearly, $\text{F-Bl}_T X$ and ω_T do not depend on the enumeration chosen. However, we need the enumeration for better record keeping: denote the closed point of T corresponding to T_i by p_i and add the subscript i in front of previously written subscripts for the related data, thus making them $p_{ij}, \mathtt{E}_{ij}, \mathtt{C}_{ij}$ for appropriate values of j.

4.1 Formulas for exceptional divisors

Each leg of the following diagram is an iterated blow-up. To keep the notation simpler, assume that $\mathtt{D} = n \cdot \mathrm{pt}$ for some $n > 0$ and replace $\mathtt{D}$ with n in notation, making $Z^{\mathtt{D}}$ into Z^n etc:

$$\begin{tikzcd} & Z \arrow[dl, "\eta_0"'] \arrow[dr, "\eta_n"] & \\ Z^0 \arrow[rr, dashed] & & Z^n \end{tikzcd}$$

In order to construct this diagram, one has to fix a global section s of $\mathcal{O}(\mathtt{D})$ so that $\mathtt{D} = (s)$. All such sections differ by nonzero multiples. Let u be a global section of $\mathcal{O}(\mathrm{pt})$, and without loss generality assume that $s = u^n$.

Given a divisor in Z^0 or Z^n, denote its total transform in Z by the same letter. Attach a superscript $+/-$ to divisors related to η_n and η_0 respectively. Denote the fiber class by $\mathtt{F}$ on any of the surfaces Z^0, Z^n and Z. Moreover, set

$$\mathtt{C}_0^{\pm} := \mathtt{F} - \mathtt{E}_1^{\pm}, \qquad \mathtt{C}_n^{\pm} := \mathtt{E}_n^{\pm}. \tag{13}$$

Recall that

$$\mathtt{C}_j^{\pm} = \mathtt{E}_j^{\pm} - \mathtt{E}_{j+1}^{\pm} \quad \text{for } j = 1, \dots, n-1. \tag{14}$$

Then

$$\mathtt{F} = \mathtt{E}_j^{+} + \mathtt{E}_k^{-} \quad \text{for } j + k = n+1, \qquad \mathtt{F} = \sum_{i=0}^{n} \mathtt{C}_i^{+} = \sum_{i=0}^{n} \mathtt{C}_i^{-}. \tag{15}$$

Denoting the linear equivalence of divisors $\sim$, we see that

$$\mathtt{X}_0 \sim \mathtt{Y}_0, \qquad \mathtt{X}_n \sim \mathtt{Y}_n + n \cdot \mathtt{F}. \tag{16}$$

The formulae below relate various (exceptional) divisors.

Lemma 4.2 *We have*

$$\mathtt{X}_0 = \mathtt{X}_n - \sum_{j=1}^{n} \mathtt{E}_j^{+}, \qquad \mathtt{Y}_n = \mathtt{Y}_0 - \sum_{j=1}^{n} \mathtt{E}_j^{-}, \qquad \mathtt{C}_j^{-} = \mathtt{C}_{n-j}^{+} \quad \textit{for } j = 0, \dots, n.$$

The various relations are summarized by

$$\begin{array}{rcl} \mathtt{E}_n^{-} = \mathtt{C}_n^{-} & = & \mathtt{C}_0^{+} = \mathtt{F} - \mathtt{E}_1^{+} \\ \mathtt{E}_{n-1}^{-} - \mathtt{E}_n^{-} = \mathtt{C}_{n-1}^{-} & = & \mathtt{C}_1^{+} = \mathtt{E}_1^{+} - \mathtt{E}_2^{+} \\ \vdots \qquad \vdots & & \vdots \qquad \vdots \\ \mathtt{E}_1^{-} - \mathtt{E}_2^{-} = \mathtt{C}_1^{-} & = & \mathtt{C}_{n-1}^{+} = \mathtt{E}_{n-1}^{+} - \mathtt{E}_n^{+} \\ \mathtt{F} - \mathtt{E}_1^{-} = \mathtt{C}_0^{-} & = & \mathtt{C}_n^{+} = \mathtt{E}_n^{+}. \end{array} \tag{17}$$

We switch from the additive notation of divisors to multiplicative notation of line bundles and sections. Let $\mathtt{F}$ be the fiber above pt, ie it is cut out by the equation $u = 0$. Denote the section corresponding to a divisor by the same letter in lowercase, ie x_0 is the section which cuts the divisor $\mathtt{X}_0$. Set $\mathtt{C}_i = \mathtt{C}_i^{+}$. The equality $\mathtt{F} = \sum \mathtt{C}_i^{+}$ now becomes $u = \prod_0^n c_i$.

Given a Higgs bundle $\theta\colon \mathcal{E} \to \mathcal{E}(\mathtt{D})$, the eigensheaf $M^{\mathtt{D}}$ on $Z^{\mathtt{D}}$ equals $\operatorname{coker}(x_n - y_n\theta)$. We want to relate $M^{\mathtt{D}}$ to M^0 on Z^0. The sheaf M^0 is defined by using $sx_0 - y_0\theta$. We relate $sx_0 - y_0\theta$ to $x_n - y_n\theta$ by

$$sx_0 - y_0\theta = u^n x_0 - y_0\theta = \left(\prod_1^n e_j^-\right)(x_n - y_n\theta) = \left(\prod_0^{n-1} c_i^{n-i}\right)(x_n - y_n\theta). \tag{18}$$

For $0 \le k \le n$,

$$x_n = \left(\prod_1^n e_j^+\right)x_0 = \left(\prod_1^n c_j^{n+1-j}\right)x_0,$$

$$\gcd(x_n, u^k) = c_1^1 c_2^2 \cdots c_{k-1}^{k-1} c_k^k \cdots c_n^k,$$

$$\left(\prod_1^n e_j^-\right)\gcd(x_n, u^k) = c_0^n \cdots c_k^n c_{k+1}^{n-1} \cdots c_n^k = u^k\left(c_0^{n-k} \cdots c_k^{n-k} c_{k+1}^{n-(k+1)} \cdots c_{n-1}^1\right).$$

For $i = 1, \ldots, r$, let k_i be the largest power of u to divide all the elements of the i^{th} row of θ. Set P, Q, R to be diagonal matrices whose i^{th} diagonal entries are respectively

$$\begin{array}{ccc} P_{ii} & Q_{ii} & R_{ii} \\ \hline c_0^{n-k} \cdots c_k^{n-k} c_{k+1}^{n-(k+1)} \cdots c_{n-1}^1 & u^k & c_1^1 c_2^2 \cdots c_{k-1}^{k-1} c_k^k \cdots c_n^k \\ \parallel & \parallel & \parallel \\ e_1^- \cdots e_l^- & u^k & e_1^+ \cdots e_k^+ \end{array} \tag{19}$$

for $k = k_i$ and $l = n - k$. Then we have the following.

Lemma 4.3

$$P \cdot Q = \left(\prod_1^n e_j^-\right) R.$$

5 Proper transform of coherent sheaves

The goal of this section is to introduce a new tool, the *proper transform of coherent sheaves*, to compare eigensheaves constructed on different Hirzebruch surfaces. Hirzebruch surfaces are related to each other by a series of elementary transformations. An elementary transformation of a Hirzebruch surface $\boldsymbol{F}_n$ at a point of the intersection of the 0–section and a fiber F produces a Hirzebruch surface $\boldsymbol{F}_{n-1}$. First, blow up $\boldsymbol{F}_n$ at the intersection of its 0–section and the fiber F. Call this surface S. Then contract the proper transform of F in S to a point on the infinity section of $\boldsymbol{F}_{n-1}$. Similarly, one can revert the process to define an elementary transformation of $\boldsymbol{F}_{n-1}$ at a point of intersection of the infinity section and a fiber F.

The idea of defining a proper transform for coherent sheaves comes from the observation that under elementary transformations, spectral curves transform to each other.

For a given Higgs bundle, one can define eigensheaves defined on different Hirzebruch surfaces. We will find that these sheaves transform to each other under elementary transformations of Hirzebruch surfaces once we replace pullbacks of coherent sheaves by proper transforms of coherent sheaves.

5.1

We are interested in eigensheaves on rational surfaces, such as Hirzebruch surfaces. By definition, eigensheaves are of homological dimension 1 at all points in their support.

For a bit more generality and to emphasize the local nature of problem, we fix a smooth point x on a scheme X and concentrate on sheaves M which are of homological dimension 1 at this point. For such a coherent sheaf M, there exists an open neighborhood U of x on which M_U has a locally free resolution

$$0 \longrightarrow \mathcal{F}_1 \longrightarrow \mathcal{F}_0 \longrightarrow M_U \longrightarrow 0.$$

We denote the scheme corresponding to the single point x by T.

Recall that our interest is to analyze how eigensheaves on different Hirzebruch surfaces transform to each other under elementary transformations. For this reason, we need to introduce a new sheaf N_U using the short exact sequence

$$0 \longrightarrow N_U \longrightarrow M_U \longrightarrow M_T \longrightarrow 0.$$

We incorporate these two short exact sequences into an exact diagram which also defines the sheaf $\mathcal{H}$:

$$(\dagger) \qquad \begin{tikzcd}
 & & & 0 \arrow[d] & \\
 & 0 \arrow[d] & & N_U \arrow[d] & \\
0 \arrow[r] & \mathcal{F}_1 \arrow[r, "\Phi"] \arrow[d] & \mathcal{F}_0 \arrow[r] \arrow[d, equal] & M_U \arrow[r] \arrow[d] & 0 \\
0 \arrow[r] & \mathcal{H} \arrow[r] \arrow[d] & \mathcal{F}_0 \arrow[r] & M_T \arrow[r] \arrow[d] & 0 \\
 & N_U \arrow[d] & & 0 & \\
 & 0 & & &
\end{tikzcd}$$

Denote the blow-up of X at x by X^σ, the morphism to X by $\omega\colon X^\sigma \to X$ and the exceptional divisor by $\mathtt{E}$.

The pullback of diagram (†) via the blow-up morphism ω is not an exact diagram. The exactness fails at the pullback of the terms of $\mathcal{H}$ and N_U.

The notion of proper transform for coherent sheaves is introduced to eradicate the failing exactness of the pullback diagram. Once the terms $\omega^*\mathcal{H}$ and $\omega^* N_U$ are replaced with their proper transforms $\mathcal{H}^\sigma$ and N_V^σ respectively, the exactness of the diagram will be restored, where $V = \omega^{-1}(U)$. The resulting diagram

$$
(\ddagger)\qquad
\begin{tikzcd}
 & & & 0 \arrow[d] & \\
 & 0 \arrow[d] & & N_V^\sigma \arrow[d] & \\
0 \arrow[r] & \omega^*\mathcal{F}_1 \arrow[r, "\Phi"] \arrow[d] & \omega^*\mathcal{F}_0 \arrow[r] \arrow[d, equal] & \omega^* M_V \arrow[r] \arrow[d] & 0 \\
0 \arrow[r] & \mathcal{H}^\sigma \arrow[r] \arrow[d] & \omega^*\mathcal{F}_0 \arrow[r] & \omega^* M_T \arrow[r] \arrow[d] & 0 \\
 & N_V^\sigma \arrow[d] & & 0 & \\
 & 0 & & &
\end{tikzcd}
$$

is exact. More importantly, we have the following.

Proposition 5.1 *One can reconstruct diagram (†) from diagram (‡) by the pushforward ω_*. That is, $\omega_*(\ddagger) = (\dagger)$.*

The rest of this section is devoted to the proof of this statement, which is divided into smaller pieces as Lemmas 5.6, 5.11 and 5.12.

5.2

Given a smooth point x of a scheme X, denote the blow-up of X at x by X^σ, the morphism to X by $\omega\colon X^\sigma \to X$ and the exceptional divisor by $\mathtt{E}$. On X^σ, the pullback of any Cartier divisor $\mathtt{D}$ of X splits as $\omega^*\mathtt{D} = m\mathtt{E} + \mathtt{D}^\sigma$ where m is a nonnegative integer and $\mathtt{D}^\sigma$ is the proper transform of the divisor $\mathtt{D}$. One can imagine the proper transform $\mathtt{D}^\sigma$ as the total transform $\omega^*\mathtt{D}$ whose exceptional locus is trimmed.

For a divisor B, say a divisor A is a subdivisor of B and write $A \leq B$ if $B - A$ is effective. The relation $\leq$ is a partial order on the set of divisors of an arbitrary scheme. In this sense, the proper transform $\mathtt{D}^\sigma$ is the largest subdivisor of $\omega^*\mathtt{D}$ which does not contain the exceptional locus $\mathtt{E}$.

Keeping these properties of the proper transform of divisors in mind, we define the proper transform of coherent sheaves in complete analogy.

Let T denote the 0–dimensional subscheme corresponding to x. For a given coherent sheaf $\mathcal{F}$ on X, we call the pullback $\omega^*\mathcal{F}$ the *total* transform of $\mathcal{F}$. Let

$$\mathcal{F}^{\mathtt{E}} := \mathcal{T}or_1^{\mathcal{O}_{X^\sigma}}(\omega^*\mathcal{F}, \mathcal{O}_{X^\sigma}(\mathtt{E})_{\mathtt{E}}).$$

The sheaf $\mathcal{F}^{\mathtt{E}}$ coincides with the subsheaf of sections of $\omega^*\mathcal{F}$ supported along the exceptional divisor $\mathtt{E}$.

Definition 5.2 The *proper transform* of $\mathcal{F}$ is defined as the quotient $\omega^*\mathcal{F}/\mathcal{F}^{\mathtt{E}}$ and will be denoted by $\omega^\sigma\mathcal{F}$, or simply $\mathcal{F}^\sigma$ when suitable.

The following sequences are exact:

$$0 \longrightarrow \mathcal{O}_{X^\sigma} \longrightarrow \mathcal{O}_{X^\sigma}(\mathtt{E}) \longrightarrow \mathcal{O}_{X^\sigma}(\mathtt{E})_{\mathtt{E}} \longrightarrow 0,$$
$$0 \longrightarrow \mathcal{F}^{\mathtt{E}} \longrightarrow \omega^*\mathcal{F} \longrightarrow \mathcal{F}^\sigma \longrightarrow 0,$$
$$0 \longrightarrow \mathcal{F}^\sigma \longrightarrow \omega^*\mathcal{F}(\mathtt{E}) \longrightarrow \omega^*\mathcal{F}(\mathtt{E})_{\mathtt{E}} \longrightarrow 0.$$

To see the exactness of the latter two, tensor the first sequence with $\omega^*\mathcal{F}$ and split the resulting sequence into two short exact sequences.

Proposition 5.3 *Given a coherent sheaf* $\mathcal{F}$ *on* X, *we have the following.*

(1) *If* $\mathcal{F}$ *is torsion-free, then* $\mathcal{F}^{\mathtt{E}}$ *and* $\mathcal{F}^\sigma$ *coincide with the torsion and torsion-free parts of* $\omega^*\mathcal{F}$ *respectively.*

(2) *If* $\mathcal{F} = \mathcal{J}_x$ *is the ideal sheaf of the point* $x \in X$. *Then* $\mathcal{J}_x^\sigma = \mathcal{O}_{X^\sigma}(-\mathtt{E})$.

(3) *If* $\mathcal{F}$ *is locally free at* x, *then* $\omega^*\mathcal{F} = \mathcal{F}^\sigma$.

(4) *If* $x \notin \operatorname{Supp}\mathcal{F}$, *then* $\omega^*\mathcal{F} = \mathcal{F}^\sigma$ *and* $\omega_*\mathcal{F}^\sigma = \omega_*\omega^*\mathcal{F} = \mathcal{F}$.

(5) *Given a locally free sheaf* L *on* X, *then* $(\mathcal{F} \otimes L)^{\mathtt{E}} \cong \mathcal{F}^{\mathtt{E}} \otimes \omega^*(L)$ *and* $(\mathcal{F} \otimes L)^\sigma \cong \mathcal{F}^\sigma \otimes \omega^*(L)$.

Proof Part (1) follows as $\omega\colon X^\sigma \backslash \mathtt{E} \to X \backslash \{x\}$ is an isomorphism, $\omega^*\mathcal{F}$ is torsion-free over the open set $X^\sigma \backslash \mathtt{E}$ and the torsion locus is $\mathtt{E}$. Parts (2) and (3) follow from (1). Part (4) is clear. Part (5) follows from the locally freeness of the sheaf L. □

The third column of diagram (‡) is exact:

Lemma 5.4 *Given an exact sequence of sheaves on* X

$$0 \longrightarrow N \longrightarrow M \xrightarrow{\mathrm{ev}_T} M_T \longrightarrow 0,$$

where ev_T *is the evaluation map, then the sequence*

$$0 \longrightarrow N^\sigma \longrightarrow \omega^* M \longrightarrow \omega^* M_T \longrightarrow 0$$

is exact and $N^\sigma = M^\sigma(-\mathtt{E})$.

For any divisor $\mathtt{D}$ *with* $\mathrm{mult}_x \mathtt{D} = 1$, $M(-\mathtt{D})^\sigma = N^\sigma(-\mathtt{D}^\sigma)$.

Proof The map $N \to M$ is 0 at x and an isomorphism of the fibers away from x. Hence the map $\omega^* N \to \omega^* M$ vanishes along $\mathtt{E}$ and it is an isomorphism away from $\mathtt{E}$. As a result, the kernel and the image of this map are $N^{\mathtt{E}}$ and N^σ. This proves the exactness of the above sequence.

The proper transform M^σ fits into the exact sequence

$$0 \longrightarrow M^\sigma \longrightarrow \omega^* M(\mathtt{E}) \longrightarrow \omega^* M(\mathtt{E})_{\mathtt{E}} \longrightarrow 0.$$

Tensoring this sequence by $\mathcal{O}(-\mathtt{E})$ shows that $N^\sigma = M^\sigma(-\mathtt{E})$.

Given such a divisor $\mathtt{D}$, we see that $\omega^*\mathtt{D} = \mathtt{D}^\sigma + \mathtt{E}$. Starting from $N^\sigma = M^\sigma(-\mathtt{E})$, tensoring both sides by $\mathcal{O}(-\mathtt{D}^\sigma)$, we get $N^\sigma(-\mathtt{D}^\sigma) = M^\sigma(-\omega^*\mathtt{D}) = M(-\mathtt{D})^\sigma$. □

The second row of diagram (‡) is exact by the following observation:

Lemma 5.5 *Given an exact sequence of coherent sheaves on* U

$$0 \longrightarrow \mathcal{S} \longrightarrow \mathcal{F} \longrightarrow \mathcal{Q} \longrightarrow 0$$

with $\mathcal{F}$ *locally free and* $\mathcal{Q}$ *torsion, then*

$$0 \longrightarrow \mathcal{S}^\sigma \longrightarrow \omega^*\mathcal{F} \longrightarrow \omega^*\mathcal{Q} \longrightarrow 0$$

is exact on U^σ.

Proof The sequence $\omega^*\mathcal{S} \to \omega^*\mathcal{F} \to \omega^*\mathcal{Q} \to 0$ is exact. The first homomorphism factors through $\mathcal{S}^\sigma$ because $\mathcal{S}^\sigma$ is the torsion-free quotient of $\omega^*\mathcal{S}$. The homomorphism $\mathcal{S}^\sigma \to \omega^*\mathcal{F}$ is generically injective since $\omega^*\mathcal{Q}$ is torsion and injective since $\mathcal{S}^\sigma$ is torsion-free. Consequently, the image of $\omega^*\mathcal{S}$ in $\omega^*\mathcal{F}$ coincides with $\mathcal{S}^\sigma$. □

Lemma 5.6 *Given a coherent sheaf M on X, suppose that either*

(1) *M is a skyscraper sheaf supported at x, or*

(2) *M is of at most homological dimension 1 at x.*

Then

$$R^0\omega_*\omega^*M = M \quad \text{and} \quad R^i\omega_*\omega^*M = 0 \quad \text{for all } i > 0.$$

Proof (1) The pullback ω^*M is the constant sheaf $\mathcal{O}_{\mathtt{E}} \otimes M_x$ on the exceptional divisor $\mathtt{E}$. From this, the result follows.

(2) The result holds for $M = \mathcal{O}_X$ because the blown-up point x is a smooth point. It holds for arbitrary locally free sheaves by the projection formula. This takes care of homological dimension 0 at x.

If M is of homological dimension 1 at x, then there exists an open neighborhood U of x on which M_U has a locally free resolution

$$0 \longrightarrow \mathcal{F}_1 \longrightarrow \mathcal{F}_0 \longrightarrow M_U \longrightarrow 0.$$

The pullback of this sequence via ω stays as a short exact sequence since $\mathcal{F}_1$ is locally free. The pushforward of the pullback sequence is also a short exact sequence since the higher direct images of the locally free sheaves $\mathcal{F}_1$ and $\mathcal{F}_0$ vanish. We conclude $R^0\omega_*\omega^*M = M$ and $R^i\omega_*\omega^*M = 0$ for all $i > 0$. □

Lemma 5.7 *Suppose the following sequences are exact for coherent sheaves A, B and C on X and X^σ respectively:*

$$0 \longrightarrow A \longrightarrow B \longrightarrow C \longrightarrow 0,$$
$$0 \longrightarrow A^\sigma \longrightarrow \omega^*B \longrightarrow \omega^*C \longrightarrow 0.$$

Suppose the coherent sheaves B and C satisfy the conclusion of the previous lemma. Then

$$R^0\omega_*A^\sigma = A \quad \text{and} \quad R^i\omega_*A^\sigma = 0 \quad \text{for all } i > 0.$$

Proof The pushforward by ω of the second short exact sequence yields

$$\begin{aligned} 0 \longrightarrow R^0\omega_*A^\sigma &\longrightarrow R^0\omega_*\omega^*B \longrightarrow R^0\omega_*\omega^*C \\ \longrightarrow R^1\omega_*A^\sigma &\longrightarrow R^1\omega_*\omega^*B \longrightarrow R^1\omega_*\omega^*C \longrightarrow \cdots. \end{aligned}$$

From the first line of this long exact sequence and the assumptions $R^0\omega_*\omega^*B = B$, $R^0\omega_*\omega^*C = C$, it follows that $R^0\omega_*A^\sigma$ is the kernel of the natural map $B \to C$, which is isomorphic to A by assumption. Similarly, from the first line and the assumption $R^1\omega_*\omega^*B = 0$ we have that $R^1\omega_*A^\sigma$ is the cokernel of the natural map $B \to C$, that is to say, 0 again by assumption. The higher direct images $R^i\omega_*A^\sigma$ for $i > 1$ vanish by the five lemma and the assumption $R^i\omega_*\omega^*B = R^i\omega_*\omega^*C = 0$ for $i > 0$. □

From now on, assume that M is a coherent sheaf on X with $\mathrm{dh}(M_x) = 1$. Let $N := \ker(M \to M_x)$.

Lemma 5.8 *The sequence*

$$0 \longrightarrow M^{\mathrm{E}} \longrightarrow \omega^* M_{\mathrm{E}} \longrightarrow M^\sigma_{\mathrm{E}} \longrightarrow 0$$

is exact.

Proof Let $K := \ker(\omega^* M_{\mathrm{E}} \to M^\sigma_{\mathrm{E}})$. The following diagram is exact:

$$\begin{tikzcd}
 & & 0 \arrow[d] & 0 \arrow[d] & \\
 & & M^{\mathrm{E}} \arrow[r, equal] \arrow[d] & K \arrow[d] & \\
0 \arrow[r] & N^\sigma \arrow[r] \arrow[d, equal] & \omega^* M \arrow[r] \arrow[d] & \omega^* M_{\mathrm{E}} \arrow[r] \arrow[d] & 0 \\
0 \arrow[r] & N^\sigma \arrow[r] & M^\sigma \arrow[r] \arrow[d] & M^\sigma_{\mathrm{E}} \arrow[r] \arrow[d] & 0 \\
 & & 0 & 0 &
\end{tikzcd}$$

This completes the proof. □

5.3 The sheaf $\mathcal{H}^\sigma$ is a Hecke transform

We finalize our findings about diagram (‡).

Definition 5.9 Given a projective scheme X, a normal crossing divisor Σ, a locally free $\mathcal{O}_X$–module $\mathcal{F}$ and a locally free $\mathcal{O}_\Sigma$–module M together with a surjection $\phi\colon \mathcal{F} \to M$, the coherent sheaf $\ker\phi$ is called the Hecke transform of $\mathcal{F}$ with respect to M and ϕ.

Remark 5.10 Hecke transforms are locally free sheaves on X.

Lemma 5.11 (1) *If M_x is torsion $\mathcal{O}_{X,x}$–module, then $\mathcal{H}$ is a torsion-free $\mathcal{O}_X$–module of the same rank as $\mathcal{E}_1$.*

(2) *Diagram (‡) is exact.*

(3) *The proper transform $\mathcal{H}^\sigma$ of $\mathcal{H}$ is a Hecke transform of $\mathcal{F}_1$ along $\mathtt{E}$. In particular, $\mathcal{H}^\sigma$ is locally free.*

Proof (1) The sheaf $\mathcal{H}$ is torsion-free since any nontrivial subsheaf of a torsion-free sheaf is torsion-free. The sheaves $\mathcal{H}$ and $\mathcal{F}_0$ are of the same rank since they are isomorphic away from T.

(2) The exactness of the second row and the third column follows from Lemmas 5.5 and 5.4. The exactness of the first column is as a consequence.

(3) The locus $\omega^{-1}(T)$ is a normal crossing divisor and $\omega^* M_T$ a vector bundle on this divisor. The proper transform $\mathcal{H}^\sigma$ of $\mathcal{H}$ is the kernel of $\omega^*\mathcal{F}_0 \to \omega^* M_T$ which proves it is a Hecke transform. □

5.4 Purity

Lemma 5.12 *Given a coherent sheaf M on X with $\mathrm{dh}(M_x) = 1$, we have:*

(1) $M^{\mathtt{E}} \cong \mathcal{O}_{\mathtt{E}}(-1)^{\oplus m}$*, where $m = \dim_{k(x)} M_x \otimes k(x)$.*

(2) $R^0\omega_* M^\sigma = M$ *and* $R^i\omega_* M^\sigma = 0$ *for all $i > 0$.*

(3) *If X is of dimension 2 and M_x is torsion, then for all $y \in \mathtt{E}$, $\mathrm{dh}(M^\sigma_y) = 1$.*

(4) $\mathcal{T}or_i^{\mathcal{O}_{X^\sigma}}(M^\sigma, \mathcal{O}_{\mathtt{E}}) = 0$ *for all $i > 0$.*

(5) $\mathcal{E}xt^0_{\mathcal{O}_{X^\sigma}}(\mathcal{O}_{\mathtt{E}}, M^\sigma) = 0$.

(6) *If M is pure of dimension 1, then $\mathtt{E} \not\subseteq \operatorname{Supp} M^\sigma$.*

Proof of Lemma 5.12 (1) The stalk M_x has a two-step resolution by free $\mathcal{O}_{X,x}$–modules. Therefore, there exists an open neighborhood U of x on which M_U has a locally free resolution

$$0 \longrightarrow \mathcal{F}_1 \longrightarrow \mathcal{F}_0 \longrightarrow M_U \longrightarrow 0.$$

The locally free sheaves $\mathcal{F}_i$ may be assumed to be of the same rank, say r. Let m be the fiber dimension of M at x. Then $m \le r$. Denote $\omega^{-1}(U)$ by V. Because $\mathtt{E} \subset V$, $M^{\mathtt{E}}_V = \mathcal{T}or_1^{\mathcal{O}_V}(\omega^* M_U, \mathcal{O}_V(\mathtt{E})_{\mathtt{E}})$. Using the locally free resolution of M_U, we see that $M^{\mathtt{E}} \cong \mathcal{O}_{\mathtt{E}}(-1)^{\oplus m}$, where $m = \dim_{k(x)} M_x \otimes k(x)$.

(2) The sequence

$$0 \longrightarrow M^{\mathtt{E}} \longrightarrow \omega^* M \longrightarrow M^\sigma \longrightarrow 0$$

is exact. The first term is isomorphic to $\mathcal{O}_{\mathtt{E}}(-1)^{\oplus m}$. Therefore $R^i\omega_* M^{\mathtt{E}} = 0$ for all i, $R^0\omega_* M^\sigma = M$ and $R^i\omega_* M^\sigma = 0$ for all $i > 0$.

(3) For all $y \in \mathtt{E}$, $\mathrm{dh}(\omega^* M_y) = 1$ and $\omega^* M_y$ is a torsion $\mathcal{O}_{X^\sigma,y}$–module. By Lemma 5.4 the sheaf N^σ is a subsheaf of $\omega^* M$. By Corollary 5.19, N_y^σ is also torsion of homological dimension equal to 1. Again by Lemma 5.4 we have $M^\sigma = N^\sigma(\mathtt{E})$, so the same conclusion holds for M_y^σ too.

(4) We know that the third column of (‡) is exact; by the snake lemma so is its first column. Applying Lemma 5.4 shows that the sheaf M_V^σ has a two-step locally free resolution

$$0 \longrightarrow \mathcal{F}_1(\mathtt{E}) \longrightarrow \mathcal{H}^\sigma(\mathtt{E}) \longrightarrow M_V^\sigma \longrightarrow 0,$$

where $\mathcal{H}^\sigma$ is defined by a Hecke transform as $\ker(\mathcal{F}_0 \to \omega^* M_{\mathtt{E}})$ and hence locally free. Applying $\mathcal{O}_{\mathtt{E}} \otimes_{\mathcal{O}_{X^\sigma}} \bullet$ to this exact sequence and using the five lemma we get $\mathcal{T}or_i^{\mathcal{O}_{X^\sigma}}(M^\sigma, \mathcal{O}_{\mathtt{E}}) = 0$ for $i \geq 2$.

As for $\mathcal{T}or_1^{\mathcal{O}_{X^\sigma}}(M^\sigma, \mathcal{O}_{\mathtt{E}})$, by (1) we have $M^{\mathtt{E}} \cong \mathcal{O}_{\mathtt{E}}(-1)^{\oplus m}$. Therefore we have

$$\mathcal{T}or_1^{\mathcal{O}_{X^\sigma}}(M^{\mathtt{E}}, \mathcal{O}_{\mathtt{E}}) \cong \mathcal{O}_{\mathtt{E}}^{\oplus m}.$$

On the other hand, as we have seen in (1) we have

$$\begin{aligned}\mathcal{T}or_1^{\mathcal{O}_{X^\sigma}}(\omega^* M, \mathcal{O}_{\mathtt{E}}) &= \mathcal{T}or_1^{\mathcal{O}_{X^\sigma}}(\omega^* M, \mathcal{O}(\mathtt{E})_{\mathtt{E}}) \otimes \mathcal{O}(-\mathtt{E}) \\ &= M^{\mathtt{E}} \otimes \mathcal{O}(-\mathtt{E}) \cong \mathcal{O}(\mathtt{E})_{\mathtt{E}}^{\oplus m} \otimes \mathcal{O}(-\mathtt{E}) \cong \mathcal{O}_{\mathtt{E}}^{\oplus m}.\end{aligned}$$

From these equalities, the portion of the long exact sequence

$$\begin{aligned} &\mathcal{T}or_2^{\mathcal{O}_{X^\sigma}}(M^\sigma, \mathcal{O}_{\mathtt{E}}) \longrightarrow \\ \longrightarrow \mathcal{T}or_1^{\mathcal{O}_{X^\sigma}}(M^{\mathtt{E}}, \mathcal{O}_{\mathtt{E}}) \longrightarrow \mathcal{T}or_1^{\mathcal{O}_{X^\sigma}}(\omega^* M, \mathcal{O}_{\mathtt{E}}) \longrightarrow &\mathcal{T}or_1^{\mathcal{O}_{X^\sigma}}(M^\sigma, \mathcal{O}_{\mathtt{E}})\end{aligned}$$

and the already established vanishing of $\mathcal{T}or_2^{\mathcal{O}_{X^\sigma}}(M^\sigma, \mathcal{O}_{\mathtt{E}})$ we therefore deduce $\mathcal{T}or_1^{\mathcal{O}_{X^\sigma}}(M^\sigma, \mathcal{O}_{\mathtt{E}}) = \ker(M^{\mathtt{E}} \to \omega^* M_{\mathtt{E}})$; the latter is in turn 0 by Lemma 5.8.

(5) Apply $\mathcal{H}om(\mathcal{O}_{X^\sigma}, \bullet) M^\sigma$ to the exact sequence

$$0 \longrightarrow \mathcal{O}_{X^\sigma}(-\mathtt{E}) \longrightarrow \mathcal{O}_{X^\sigma} \longrightarrow \mathcal{O}_{\mathtt{E}} \longrightarrow 0.$$

The result is

$$0 \longrightarrow \mathcal{E}xt^0_{\mathcal{O}_{X^\sigma}}(\mathcal{O}_{\mathtt{E}}, M^\sigma) \longrightarrow M^\sigma \longrightarrow M^\sigma(\mathtt{E}) \longrightarrow M^\sigma(\mathtt{E})_{\mathtt{E}} \longrightarrow 0.$$

Then, $\mathcal{E}xt^0_{\mathcal{O}_{X^\sigma}}(\mathcal{O}_{\mathtt{E}}, M^\sigma) = \mathcal{T}or_1^{\mathcal{O}_{X^\sigma}}(M^\sigma, \mathcal{O}_{\mathtt{E}}) = 0$. The latter sheaf is trivial by (4).

(6) If M is pure of dimension 1, then M^σ is torsion and $\mathrm{dh}(M_y^\sigma) = 1$ for all $y \in X^\sigma$. Consequently, M^σ is pure of dimension 1. As $\mathtt{E}$ is not an associated point of M^σ by (4), $\mathtt{E} \not\subseteq \operatorname{Supp} M^\sigma$. □

5.5 Morphisms

Let M_1, M_2 be coherent sheaves on X. Then a homomorphism $\phi\colon M_1 \to M_2$ induces $\phi^*\colon \omega^* M_1 \to \omega^* M_2$ with $\omega^*(M_1^{\mathtt{E}}) \subset M_2^{\mathtt{E}}$, and we denote the induced morphism on the quotients $M_1^\sigma \to M_2^\sigma$ by ϕ^σ. For coherence, we also use $\omega^\sigma \phi$ for ϕ^σ.

From now on we assume that M_1, M_2 are pure of dimension 1.

Lemma 5.13 *For any homomorphism* $\phi\colon M_1 \to M_2$, $\omega_* \omega^\sigma \phi = \phi$. *For any homomorphism* $\psi\colon M_1^\sigma \to M_2^\sigma$, $(\omega_* \psi)^\sigma = \psi$.

Proof The image of $\phi - \omega_* \phi^\sigma$ is 0–dimensional. Because M_2 is pure of dimension 1, the image is trivial and hence $\phi = \omega_* \phi^\sigma$. The image of $\psi - (\omega_* \psi)^\sigma$ is contained in the 0–dimensional subscheme $\mathtt{E} \cap \operatorname{Supp} M_2^\sigma$. Hence, $\psi = (\omega_* \psi)^\sigma$. □

Lemma 5.14 *A homomorphism* $\phi\colon M_1 \to M_2$ *is injective if and only if* ϕ^σ *is injective. In this case,* $\omega_*(\operatorname{coker} \phi^\sigma) = \operatorname{coker} \phi$.

Proof $(\Leftarrow)$ As M_i are pure of dimension 1, $M_i = M_i^\sigma$ and $\phi = \omega_* \omega^\sigma \phi$. The injectivity follows from the left exactness of ω_*.

$(\Rightarrow)$ The homomorphism ϕ^σ is injective away from $\mathtt{E}$, thus $\ker \phi^\sigma \subset \mathtt{E} \cap \operatorname{Supp} M_1^\sigma$, hence trivial. □

Lemma 5.15 *Let* M *be a pure sheaf of dimension* 1. *Then*

$$\chi(\omega^* M) = \chi(M),$$
$$\chi(\omega^\sigma M) = \chi(M).$$

Proof Lemma 5.6 (respectively Lemma 5.12) shows that the sheaf cohomology of $\omega^* M$ (respectively $\omega^\sigma M$) matches the sheaf cohomology of M, and hence $\chi(\omega^* M) = \chi(M) = \chi(\omega^\sigma M)$. □

5.6 Parabolic case

Here, we introduce the proper transform for parabolic sheaves of dimension 1 on surfaces and investigate its basic properties.

Given a surface X, fix a smooth point x, an effective Cartier divisor $\mathtt{D}$ and parabolic sheaf $M_\bullet$ of dimension 1 with divisor $\mathtt{D}$.

No further transversality conditions other than Remark 3.4 are required of the parabolic sheaf $M_\bullet$, the divisor $\mathtt{D}$ and the point x: the divisor $\mathtt{D}$ or the support of $M_\bullet$ are allowed to contain the point x, or meet each other arbitrarily.

In Proposition 5.17, we show that the proper transform of a parabolic sheaf of dimension 1 with parabolic divisor $\mathtt{D}$ can be given the structure of a parabolic sheaf whose parabolic divisor is the *total transform* of $\mathtt{D}$.

In Section 6, we provide the details of how to alter the parabolic structure of the proper transform of a parabolic sheaf along the exceptional divisor $\mathtt{E}$ and obtain a parabolic sheaf whose associated divisor will be a proper transform of the sheaf $\mathtt{D}$.

Definition 5.16 Given a parabolic sheaf $M_\bullet$ of dimension 1 on X with parabolic divisor $\mathtt{D}$, the *proper transform* of $M_\bullet$ is defined by setting $\omega^\sigma(M)_\alpha := \omega^\sigma(M_\alpha)$ for $\alpha \in \mathbb{R}$.

Then $\omega^\sigma M_\bullet$ has the same weights as $M_\bullet$. We show that the definition of proper transform make sense for parabolic sheaves:

Proposition 5.17 *The proper transform $\omega^\sigma M_\bullet$ of $M_\bullet$ with divisor $\mathtt{D}$ is a parabolic sheaf on X^σ with divisor $\omega^*\mathtt{D}$, $\omega_* \operatorname{gr}_i^F \omega^\sigma M = \operatorname{gr}_i^F M$ for all i and* par-$\chi(\omega^\sigma M_\bullet) =$ par-$\chi(M_\bullet)$.

Proof We start by proving that the proper transform of a parabolic sheaf M is again a parabolic sheaf with divisor $\omega^*\mathtt{D}$.

First, $\omega^\sigma(M)_0 = \omega^\sigma M$. By assumption, $M_\bullet$ is parabolic. The sheaf M_β is a subsheaf of M_α for $\beta \geq \alpha$. Lemma 5.14 implies M_β^σ is a subsheaf of M_α^σ and $\omega_* \operatorname{gr}_i^F \omega^\sigma M = \operatorname{gr}_i^F M$ for all i. In addition, by Proposition 5.3,

$$\omega^\sigma(M)_{\alpha+1} = \omega^\sigma(M_\alpha(-\mathtt{D})) = [\omega^\sigma(M_\alpha)](-\omega^*\mathtt{D}).$$

The weights of $\omega^\sigma M_\bullet$ coincide with the weights of $M_\bullet$. These make $\omega^\sigma M_\bullet$ into a 1–dimensional parabolic sheaf with divisor $\omega^*\mathtt{D}$.

The parabolic Euler characteristic is preserved since

$$\omega^\sigma(M(-\mathtt{D})) = \omega^\sigma M(-\omega^*\mathtt{D})$$

and $\omega_* \operatorname{gr}_i^F M^\sigma = \operatorname{gr}_i^F M$ for all i:

$$\begin{aligned} \text{par-}\chi(\omega^\sigma M_\bullet) &= \chi(\omega^\sigma M(-\omega^*\mathtt{D})) + \sum \alpha_i \chi(\operatorname{gr}_i^F \omega^\sigma M) \\ &= \chi(M(-\mathtt{D})) + \sum \alpha_i \chi(\operatorname{gr}_i^F M) = \text{par-}\chi(M_\bullet). \end{aligned}$$

□

Similarly, we can define pushforwards of parabolic sheaves of dimension 1 whose divisor is $\omega^*\mathtt{D}$:

Definition 5.18 Given a parabolic sheaf $\widetilde{M}_\bullet$ of dimension 1 with parabolic divisor $\omega^*\mathtt{D}$, define the *pushforward* of $\widetilde{M}_\bullet$ by setting $(\omega_*\widetilde{M})_\alpha := \omega_*(\widetilde{M}_\alpha)$ for $\alpha \in \mathbb{R}$.

Then $\omega_*\widetilde{M}_\bullet$ has the same weights as $\widetilde{M}_\bullet$. In order to show that this definition makes sense we only need to check that $\omega_*(\widetilde{M}_{\alpha+1}) = \omega_*(\widetilde{M}_\alpha)(-\mathtt{D})$ for every $\alpha \in \mathbb{R}$. This follows from the projection formula:

$$\omega_*(\widetilde{M}_{\alpha+1}) = \omega_*(\widetilde{M}_\alpha(-\omega^*\mathtt{D})) = \omega_*(\widetilde{M}_\alpha) \otimes \omega_*\omega^*\mathcal{O}_X(-\mathtt{D}).$$

5.7 Appendix on commutative algebra

Let A be a local ring with maximal ideal $\mathfrak{m}$. For any A–module M, the depth of M is defined as

$$\operatorname{depth}(M) := \min\{i \mid \operatorname{Ext}^i_A(A/\mathfrak{m}, M) \neq 0\}$$

and the homological dimension $\dim(M)$ is defined as the minimal length of a projective resolution of M.

The Auslander–Buchsbaum formula relates the two invariants:

$$\operatorname{dh}(M) + \operatorname{depth}(M) = \operatorname{depth}(A).$$

If A is a regular local ring, then $\operatorname{depth}(A) = \dim(A)$.

Corollary 5.19 *Let A be a regular local ring of dimension 2, M a torsion A–module with* $\operatorname{dh}(M) = 1$. *Then, any submodule M' of M is a torsion A–module with* $\operatorname{dh}(M') = 1$.

Proof Any submodule M' of M is torsion, therefore not locally free and $\operatorname{dh}(M') \geq 1$. Since $\operatorname{depth}(M) = 1$,

$$\operatorname{Ext}^0_A(A/\mathfrak{m}, M') \subset \operatorname{Ext}^0_A(A/\mathfrak{m}, M) = 0.$$

Consequently, $\operatorname{depth}(M') \geq 1$. By the Auslander–Buchsbaum equality,

$$\operatorname{dh}(M') = \operatorname{depth}(M') = 1. \qquad \square$$

Lemma 5.20 *Assume X is a smooth projective surface. For a coherent sheaf M on X, the following are equivalent:*

(1) *M is pure of dimension 1.*

(2) *M is a torsion sheaf with $\dim(M_x) = 1$ for all $x \in X$.*

(3) *$\mathrm{dh}(M_x) = \mathrm{depth}(M_x) = 1$ for all $x \in X$.*

Moreover, any subsheaf of a given pure sheaf M of dimension 1 is also pure of dimension 1.

Proof Apply Corollary 5.19 and [5, Proposition 1.1.10] by Huybrechts and Lehn, which in this particular case, states that a coherent sheaf M of dimension 1 is pure if and only if $\mathrm{depth}(M_x) \geq 1$ for all $x \in X$. $\square$

6 Addition and deletion

In this section, we introduce two operations, addition and deletion, to modify the parabolic structure of a parabolic sheaf. Given a parabolic sheaf whose associated divisor is of the form $\mathtt{D}'$ and a different effective divisor $\mathtt{E}$, addition of $\mathtt{E}$ to the parabolic structure of the parabolic sheaf in question yields a new parabolic sheaf whose associated divisor is $\mathtt{D}' + \mathtt{E}$. Conversely, deletion of $\mathtt{E}$ removes the divisor $\mathtt{E}$ from the parabolic divisor $\mathtt{D}' + \mathtt{E}$ of some parabolic sheaf. In application, $\mathtt{D}'$ will be the proper transform of some divisor $\mathtt{D}$ with respect to a blow-up and $\mathtt{E}$ the exceptional divisor of the said blow-up.

Deletion Let $\mathbb{P}_\bullet$ be a parabolic sheaf on X with divisor $\mathtt{D}$ and $\mathtt{D}', \mathtt{E}$ be effective Cartier divisors in X such that $\mathtt{D} = \mathtt{D}' + \mathtt{E}$. Because $\dim(\mathrm{Supp}(\mathbb{P}) \cap \mathtt{D}) < \dim(\mathrm{Supp}(\mathbb{P}))$, the same holds for $\mathtt{E}$ and $\mathtt{D}'$ as well. We set $\mathbb{P}' := \mathbb{P}(-\mathtt{E})$. One can put a parabolic structure on the sheaf $\mathbb{P}'$ whose parabolic divisor is $\mathtt{D}'$: for $0 \leq \alpha < 1$, set $\mathbb{P}'_\alpha := \mathbb{P}' \cap \mathbb{P}_\alpha$. Extend this to a parabolic structure by setting $\mathbb{P}'_\alpha := \mathbb{P}'_{\alpha - \lfloor \alpha \rfloor}(\lfloor \alpha \rfloor \mathtt{D}')$. We call $\mathbb{P}'_\bullet$ the *deletion of $\mathtt{E}$ from the divisor of $\mathbb{P}_\bullet$*, and denote it by $\mathrm{Del}_{\mathtt{E}}(\mathbb{P}_\bullet)$.

Addition Conversely, given a parabolic sheaf $\mathbb{P}'_\bullet$ with divisor $\mathtt{D}'$ and an effective divisor $\mathtt{E}$ such that $\dim(\mathrm{Supp}(\mathbb{P}') \cap \mathtt{E}) < \dim(\mathrm{Supp}(\mathbb{P}'))$, one can put a parabolic structure with parabolic divisor $\mathtt{D} = \mathtt{D}' + \mathtt{E}$ on $\mathbb{P} = \mathbb{P}'(\mathtt{E})$ by setting $\mathbb{P}_0 = \mathbb{P}$ and $\mathbb{P}_\alpha = \mathbb{P}'_\alpha$ for $0 < \alpha < 1$ and extending this to $\mathbb{R}$ in the usual way. Denote $\mathbb{P}_\bullet$ by $\mathrm{Add}_{\mathtt{E}}(\mathbb{P}'_\bullet)$, and call it the *addition of $\mathtt{E}$ to the divisor of $\mathbb{P}'_\bullet$*.

Admissibility condition 2 (for parabolic sheaves) One has either

$$\alpha_0(\mathbb{P}) > 0 \quad \text{and} \quad \mathrm{Supp}(\mathbb{P}) \cap \mathtt{E} = \varnothing,$$
$$\alpha_0(\mathbb{P}) = 0 \quad \text{and} \quad F_1\mathbb{P} = \mathbb{P}(-\mathtt{E}).$$

Definition 6.1 A parabolic sheaf $\mathbb{P}_\bullet$ with associated divisor $\mathtt{D}$ will be called $\mathtt{E}$–admissible if for an effective Cartier divisor $\mathtt{E}$, the divisor $\mathtt{D}-\mathtt{E}$ is effective and $\mathbb{P}_\bullet$ obeys the dichotomy

$$\begin{aligned} &\alpha_0(\mathbb{P})>0 \quad \text{and} \quad \operatorname{Supp}(\mathbb{P})\cap\mathtt{E}=\varnothing, \quad \text{or} \\ &\alpha_0(\mathbb{P})=0 \quad \text{and} \quad F_1\mathbb{P}=\mathbb{P}(-\mathtt{E}). \end{aligned}$$

Proposition 6.2 *There is a one-to-one correspondence between the following two sets of objects induced by operations* $\mathrm{Del}_{\mathtt{E}}$ *and* $\mathrm{Add}_{\mathtt{E}}$*:*

$$\begin{aligned} &\mathtt{E}\text{–admissible parabolic sheaves } \mathbb{P} \text{ with associated divisor } \mathtt{D} \\ &\qquad\qquad \longleftrightarrow \text{Parabolic sheaves } \mathbb{P}' \text{ with divisor } \mathtt{D}-\mathtt{E} \text{ and } \alpha_0(\mathbb{P}')>0. \end{aligned}$$

Furthermore, $\text{par-}\chi(\mathrm{Del}_{\mathtt{E}}(\mathbb{P}_\bullet))=\text{par-}\chi(\mathbb{P}_\bullet)$*.*

Proof We show first that given an admissible parabolic sheaf $\mathbb{P}$, the parabolic sheaf $\mathrm{Del}_{\mathtt{E}}(\mathbb{P})$ obtained by deleting $\mathtt{E}$ from the parabolic divisor $\mathtt{D}$ produces a parabolic sheaf of the second kind.

Set $\mathbb{P}'_\bullet:=\mathrm{Del}_{\mathtt{E}}\mathbb{P}_\bullet$. Suppose first $\alpha_0(\mathbb{P})>0$. Then $\operatorname{Supp}(\mathbb{P})\cap\mathtt{E}=\varnothing$ implies $\mathbb{P}'_0=\mathbb{P}(-E)=\mathbb{P}$ and hence $\mathbb{P}'_\alpha=\mathbb{P}_\alpha\cap\mathbb{P}'_0=\mathbb{P}_\alpha\cap\mathbb{P}=\mathbb{P}_\alpha$ for $0<\alpha<1$. In short, in this case, $\mathbb{P}'_\bullet=\mathbb{P}_\bullet$ and $\alpha_0(\mathbb{P}')=\alpha_0(\mathbb{P})>0$.

Suppose now $\alpha_0(\mathbb{P})=0$. Therefore, $F_1\mathbb{P}=\mathbb{P}(-\mathtt{E})$ by definition. In other words, for small $\epsilon>0$, $\mathbb{P}_\epsilon=\mathbb{P}(-\mathtt{E})$. Recall that $\mathbb{P}'_0$ is set to be $\mathbb{P}(-\mathtt{E})$ and $\mathbb{P}'_\alpha:=\mathbb{P}_\alpha\cap\mathbb{P}(-\mathtt{E})$ for $0<\alpha<1$. Consequently, we see that at $\alpha=0$, there is no jump, ie for small $\epsilon\geq 0$, we have $\mathbb{P}'_0=\mathbb{P}(-\mathtt{E})=\mathbb{P}'_\alpha$. Hence, $\alpha_0(\mathbb{P}')>0$.

Traversing the steps back one observes that $\mathrm{Add}_{\mathtt{E}}$ is the inverse of $\mathrm{Del}_{\mathtt{E}}$. These prove that $\mathrm{Del}_{\mathtt{E}}$ and $\mathrm{Add}_{\mathtt{E}}$ form a one-to-one correspondence between the two sets of parabolic sheaves above.

One of the conclusions of the above paragraph is that $\mathrm{Del}_{\mathtt{E}}$ does not modify the sheaves $\mathbb{P}_\alpha$ for any $0<\alpha\leq 1$, ie for α in this range, $\mathbb{P}_\alpha=\mathrm{Del}_{\mathtt{E}}(\mathbb{P}_\bullet)_\alpha$. This also forces $\mathrm{gr}_\alpha(\mathbb{P}_\bullet)=\mathrm{gr}_\alpha(\mathrm{Del}_{\mathtt{E}}(\mathbb{P}_\bullet))$ for α in the same range. Finally, we conclude that $\text{par-}\chi(\mathrm{Del}_{\mathtt{E}}(\mathbb{P}_\bullet))=\text{par-}\chi(\mathbb{P}_\bullet)$. □

Lemma 6.3 *Let* $\mathbb{P}_\bullet$ *be an* $\mathtt{E}$*–admissible parabolic sheaf on* X *and* $\mathcal{S}_\bullet$ *be a saturated subsheaf endowed with the induced parabolic structure. Then* $\mathcal{S}_\bullet$ *is also* $\mathtt{E}$*–admissible.*

Proof The induced parabolic structure is defined by the formula $\mathcal{S}_\alpha=\mathcal{S}\cap\mathbb{P}_\alpha$ for all $0\leq\alpha<1$. For small enough $\varepsilon>0$, one then has $\mathcal{S}_\varepsilon=\mathcal{S}\cap\mathbb{P}_\varepsilon=\mathcal{S}\cap\mathbb{P}(-\mathtt{E})$. Because $\mathcal{S}\subset\mathbb{P}$, this implies $\mathcal{S}_\varepsilon=\mathcal{S}(-\mathtt{E})$. Assuming $\operatorname{Supp}(\mathcal{S})\cap\mathtt{E}\neq\varnothing$, one has $\mathcal{S}_\varepsilon\neq\mathcal{S}$. Letting $\varepsilon\to 0$, this implies $\alpha_0(\mathcal{S})=0$, and $F_1\mathcal{S}=\mathcal{S}(-\mathtt{E})$ as claimed. □

7 Spectral triples

From now on, let C be a smooth projective curve over the field $\boldsymbol{K}$ and $\mathtt{P}$ an effective Cartier divisor on C.

Definition 7.1 A spectral triple (Z, Σ, M) consists of a smooth surface Z together with a morphism $\omega\colon Z \to C$, an effective divisor $\Sigma \subset Z$ and an $\mathcal{O}_Z$–module M of homological dimension 1 with support Σ so that $\pi\colon Z \to C$ is flat and $\pi_{|\Sigma}\colon \Sigma \to C$ is finite and flat.

Definition 7.2 A parabolic spectral triple $(Z, \Sigma, M_\bullet)$ is a spectral triple so that $M_\bullet$ is a parabolic sheaf whose parabolic divisor is $\omega^*(\mathtt{P})$.

Remark 7.3 In fact, it would be enough to consider the pair (Z, M) as Σ is determined by $\Sigma = \operatorname{Supp}(M)$. We include Σ in the definition for expository purposes.

Starting from a Higgs bundle $\theta\colon \mathcal{E} \to \mathcal{E}(\mathtt{P})$ one can define a spectral triple $(Z^\flat, \Sigma^\flat, M^\flat)$: set $C^\flat := C - \mathtt{P}$ and $\boldsymbol{A}^1 := \operatorname{Spec}(k[\lambda])$. Here:

(20) $$Z^\flat := C^\flat \times \boldsymbol{A}^1 \text{ with the obvious choice for } \omega_\flat\colon Z^\flat \longrightarrow C.$$

(21) $$\Sigma^\flat := (\det(\lambda \operatorname{Id}_{\mathcal{E}} - \theta)).$$

(22) $$M^\flat := \operatorname{coker}(\lambda \operatorname{Id}_{\mathcal{E}} - \theta).$$

Definition 7.4 A compactification of the triple $(Z^\flat, \Sigma^\flat, M^\flat)$ is a spectral triple (Z, Σ, M) and an open immersion $i\colon Z^\flat \to Z$ so that

(1) Z is a connected smooth projective surface,

(2) $i(\Sigma^\flat)$ is contained in Σ as a dense open subset,

(3) $i^*M = M^\flat$,

so that the following diagram commutes:

$$\begin{tikzcd} Z^\flat \arrow[rr, "i"] \arrow[dr, "\pi_\flat"'] & & Z \arrow[dl, "\pi"] \\ & C & \end{tikzcd}$$

By (1) and (2), $Z^\flat$ is dense inside Z and Σ is the closure of $i(\Sigma^\flat)$ in Z.

There are many compactifications of the triple $(Z^\flat, \Sigma^\flat, M^\flat)$: given one compactification, one can obtain other via suitable blow-ups.

The standard choice for a compactification of $(Z^\flat, \Sigma^\flat, M^\flat)$ is

$$Z^{\mathtt{P}} := \mathbb{P}_C(\mathcal{O} \oplus \mathcal{O}(-\mathtt{P})), \quad \Sigma^{\mathtt{P}} := \big(\det(x_{\mathtt{P}} \operatorname{Id}_{\mathcal{E}} - y_{\mathtt{P}}\theta)\big), \quad M^{\mathtt{P}} := \operatorname{coker}(x_{\mathtt{P}} \operatorname{Id}_{\mathcal{E}} - y_{\mathtt{P}}\theta).$$

Remark 7.5 The divisor $\Sigma^{\mathtt{P}}$ does not meet the line at infinity $(y_{\mathtt{P}})$.

7.1 Naive compactification: The surface and the curve

We describe the compactification $Z^0 = C \times \mathbb{P}^1$. The spectral curve Σ^0 is determined by $\Sigma^\flat$ as its closure. The sheaf M^0 is the cokernel of a morphism between locally free sheaves as was $M^{\mathtt{P}}$.

Remark 7.6 The spectral curve Σ^0 meets the line at infinity (y_0) over the points $p \in \mathtt{P}$ for which θ_p is not a nilpotent endomorphism.

7.2 Naive compactification: The sheaf

Fix a section s so that $\mathtt{P} = (s)$ and let $\Theta_0 := sx_0 \operatorname{Id}_{\mathcal{E}} - y_0\theta \colon \mathcal{E} \to \mathcal{E}(\mathtt{P}) \otimes \mathcal{O}_{Z^0}(1)$. Then

$$(\det \Theta_0) = \Sigma^0 + (\det \operatorname{coker}(\theta)).$$

We construct another map Φ_0 related to $\Theta_0 \colon \mathcal{E} \to \mathcal{E}(\mathtt{P}) \otimes \mathcal{O}_{Z^0}(1)$ so that

$$\Sigma^0 = (\det \Phi_0).$$

Set $\mathcal{F} := \ker(\mathcal{E}(\mathtt{P}) \to \operatorname{coker}\theta(-\mathtt{P})_{\mathtt{P}})$ and denote the inclusion $\mathcal{F} \hookrightarrow \mathcal{E}$ by Q. The map $\theta \colon \mathcal{E} \to \mathcal{E}(\mathtt{P})$ naturally factors through $\mathcal{F}(\mathtt{P})$. Similarly, Θ_0 factors through $\mathcal{F}(\mathtt{P}) \otimes \mathcal{O}_{Z^0}(1)$. Denote the map $\mathcal{E} \to \mathcal{F}(\mathtt{P}) \otimes \mathcal{O}_{Z^0}(1)$ by Φ_0. The following diagram is exact:

$$\begin{tikzcd}
 & & & 0 \arrow[d] & \\
 & 0 \arrow[d] & & M^0 \arrow[d] & \\
0 \arrow[r] & \mathcal{E} \arrow[r, "\Theta_0"] \arrow[d, "\Phi_0"'] & \mathcal{E}(\mathtt{P}) \otimes \mathcal{O}_{Z^0}(1) \arrow[r] \arrow[d, equal] & \eta_{T^-}^* \operatorname{coker}\Theta_0 \arrow[r] \arrow[d] & 0 \\
0 \arrow[r] & \mathcal{F}(\mathtt{P}) \otimes \mathcal{O}_{Z^0}(1) \arrow[r, "Q\otimes 1"] \arrow[d] & \mathcal{E}(\mathtt{P}) \otimes \mathcal{O}_{Z^0}(1) \arrow[r] & \eta_{T^-}^* \operatorname{coker}\theta_D \arrow[r] \arrow[d] & 0 \\
 & M^0 \arrow[d] & & 0 & \\
 & 0 & & &
\end{tikzcd} \tag{23}$$

From this diagram, we see that

$$\operatorname{Supp} M^0 = (\det \Phi_0) = \Sigma^0.$$

7.3

We use the ideas we have developed in Section 5 to compare the sheaves M^0 and M^{P} constructed on the Hirzebruch surfaces Z^0 and Z^{P} respectively. The Hirzebruch surfaces transform to each under other arbitrary elementary transformations. However, in our description, elementary transformations are prescribed by the data derived from the Higgs bundle $(\mathcal{E}, \theta)$. Moreover, instead of analyzing a single elementary transformation, we analyze all elementary transformations which will be needed in a single diagram:

$$\begin{tikzcd} & Z \arrow[dl, "\eta_{T^-}"'] \arrow[dr, "\eta_{T^+}"] & \\ Z^0 \arrow[rr, dashed] & & Z^{\mathrm{P}} \end{tikzcd}$$

Here:

- $T^- := \Sigma^- \cap (y_0)$ and η_{T^-} is the iterated blow-up map $Z \to Z^0$.
- $T^+ := \Sigma^+ \cap (x_{\mathrm{P}})$ and η_{T^+} is the iterated blow-up map $Z \to Z^{\mathrm{P}}$.

Also, set:

- $\mathcal{F}_1^0 := \mathcal{E}$ and $\mathcal{F}_0^0 := \mathcal{F}(\mathrm{P}) \otimes \mathcal{O}_{Z^0}(1)$.
- $\mathcal{F}_1^{\mathrm{P}} := \mathcal{E}$ and $\mathcal{F}_0^{\mathrm{P}} := \mathcal{E}(\mathrm{P}) \otimes \mathcal{O}_{Z^{\mathrm{P}}}(1)$.

View $\mathcal{F}_1^{\mathrm{B}} \to \mathcal{F}_0^{\mathrm{B}}$ as a two-step locally free resolution for the sheaf M^{B} on $U = Z^{\mathrm{B}}$. Set

$$\mathcal{H}_{\mathrm{B}} := \ker\big(\mathcal{F}_0^{\mathrm{B}} \longrightarrow M^{\mathrm{B}}_{T_{\mathrm{B}}}\big).$$

Proposition 7.7 *For $\mathtt{B} = 0, \mathtt{P}$, T^{B} is linear, $Z = \text{F-Bl}_{T^{\mathrm{B}}} Z^{\mathrm{B}}$ and $\omega_{\mathrm{B}} = \omega_{T^{\mathrm{B}}}$. Also:*

(24) $$\eta_{T^+}^{\sigma} \mathcal{H}^{\mathrm{P}} = \eta_{T^-}^{\sigma} \mathcal{H}^0.$$

(25) $$\eta_{T^+}^{\sigma} N^{\mathrm{P}} = \eta_{T^-}^{\sigma} N^0.$$

(26) $$\Sigma^{0^{\sigma}} = \Sigma^{\mathrm{P}^{\sigma}}.$$

We denote the proper transformed spectral curve (26) by Σ.

Proof For simplicity, assume $\mathtt{P} = n \cdot \mathrm{pt}$ with local coordinate u at pt. The subscheme $T^{\mathtt{P}}$ is linear because $T^{\mathtt{P}} = (u^n) \cap (t) = (u^n, t)$, where t is $x_{\mathtt{P}}$ or y_D. The pair of u and t is clearly transversal. Lemma 4.3 says that the following is diagram commutative and $\eta_{T^+}{}^{\sigma}\mathcal{H}^{\mathtt{P}} = \eta_{T^-}{}^{\sigma}\mathcal{H}^0$:

$$\begin{tikzcd}
{\mathcal{H}_{\mathtt{P}}}^{\sigma} \arrow[rr, dashed, "(\cong)"] \arrow[d, "P"'] & & {\mathcal{H}_0}^{\sigma} \arrow[d, "R"] \\
\mathcal{F}(\mathtt{P}) \otimes \mathcal{O}_{Z^0}(1) \arrow[dr, "Q"'] & & \mathcal{E}(\mathtt{P}) \otimes \mathcal{O}_{Z^{\mathtt{P}}}(1) \arrow[dl, "\left(\prod_1^n e_i^-\right)"] \\
& \mathcal{E}(\mathtt{P}) \otimes \mathcal{O}_{Z^{\mathtt{P}}}(1) &
\end{tikzcd}$$

The rest follows. $\square$

8 Algebraic Nahm transform

Let $(Z^{\mathtt{P}}, \Sigma^{\mathtt{P}}, M^{\mathtt{P}}_{\bullet})$ be the standard parabolic spectral triple of the parabolic Higgs bundle $(\mathcal{E}_{\bullet}, \theta)$ whose parabolic divisor $\mathtt{D}$ is $\mathtt{P}+\infty$. Define the following 0–dimensional subschemes of $Z^{\mathtt{P}}$:

$$T^+ := \pi^*(\mathtt{P}) \cap (x_{\mathtt{P}}),$$

which is the intersection of the pullback divisor $\pi^*(\mathtt{P})$ and the 0–section in $Z^{\mathtt{P}}$, and

$$\hat{T}^- := \pi^*(\infty) \cap \Sigma^{\mathtt{P}},$$

which is the fiber of $\Sigma^{\mathtt{P}}$ over $\infty \in Z^{\mathtt{P}}$. Applying the ideas of Section 5.2 to the 0–dimensional subscheme T^+ produces a new spectral triple $(Z^{\mathtt{P}}, \Sigma^{\mathtt{P}}, N^{\mathtt{P}})$ out of $M^{\mathtt{P}}_{\bullet}$:

$$N^{\mathtt{P}} := \ker\bigl(M^{\mathtt{P}} \to M^{\mathtt{P}}_{T^+}\bigr).$$

By definition, $N^{\mathtt{P}}$ consists of the local sections of $M^{\mathtt{P}}$ vanishing in T^+.

Next, define another spectral triple $(Z, \Sigma, N_{\bullet})$: the surface Z is the blow-up η_{T^+} of $Z^{\mathtt{P}}$ at T^+, the divisor $\mathtt{E}^+$ is the exceptional divisor of η_{T^+} and the coherent sheaf N is defined as the proper transform $\eta_{\mathtt{P}^{\sigma}}(N^{\mathtt{P}})$ of $N^{\mathtt{P}}$. The support Σ of N is the proper transform $\eta^{\sigma}_{T^+}(\Sigma^{\mathtt{P}})$ of $\Sigma^{\mathtt{P}}$. Let

$$N_{\bullet} := \mathrm{Del}_{\mathtt{E}^+} \eta^{\sigma}_{T^+} M^{\mathtt{P}}_{\bullet}.$$

The sheaf N now has a parabolic structure whose divisor is $\eta^{\sigma}_{T^+}\mathtt{D} = \eta^*_{T^+}\mathtt{D} - \mathtt{E}^+$. By our convention, the parabolic structure of $N^{\mathtt{P}}_{\bullet}$ (as the ones of all the other sheaves involved in the construction) has weights between 0 and 1 in all parabolic points.

Proposition 8.1 *If Admissibility condition 1 holds for $(\mathcal{E},\theta)$, then Admissibility condition 2 holds for the parabolic sheaf $N_\bullet$ and the divisor $\mathtt{E}^+$.*

Proof It is sufficient to work in a neighborhood of a parabolic point $p \in \mathtt{P}$. The Admissibility condition 2 says that if the residue of θ at p has a 0 eigenvalue, the relation $\mathcal{F} = F_1\mathcal{E}$ holds, and the smallest parabolic weight is 0. This then implies that $\mathrm{gr}_0^F(\mathcal{E}) = \mathrm{coker}(\theta(-p)_p)$, and that the parabolic weight associated to this graded piece is 0. By the definition of the parabolic structure of $M^{\mathtt{P}}$, this is equivalent to saying that $\mathrm{gr}_0^F(M^{\mathtt{P}}) = \mathrm{coker}(\theta(-p)_p)$ on the fiber $\mathtt{F} = \pi^{-1}(p)$, with 0 parabolic weight associated to this graded piece. Since the sheaf $\mathrm{coker}(\theta(-p)_p)$ is supported in the point $t = \mathtt{F} \cap (x_{\mathtt{P}})$, we see that the 0–weight subspace of $M^{\mathtt{P}}|_{\mathtt{F}}$ is precisely $M^{\mathtt{P}}|_t$. Therefore, we have $F_1 M^{\mathtt{P}} = \ker(M^{\mathtt{P}} \to M^{\mathtt{P}}|_t)$. Now let us blow up the point t, and call $\mathtt{E}$ the exceptional divisor and M the proper transform of $M^{\mathtt{P}}$. Then for the parabolic structure of M the relation $F_1 M = \ker(M \to M|_{\mathtt{E}}) = M(-\mathtt{E})$ holds, and the corresponding parabolic weight is 0. This shows that Admissibility condition 2 is true for M. It then follows for N as well because of our convention of keeping the same weights for kernel and cokernel sheaves. $\square$

For simplicity, write π for the projection $\pi \circ \boldsymbol{\eta}_{T^+}\colon Z^{\mathtt{P}} \to \mathbb{P}^1$, whenever this does not create any ambiguity, and do the same for all other projections to $\mathbb{P}^1$ composed with blow-up maps. Similarly, identify the 0–dimensional subscheme $T \subset Z^2$ with $\omega^{-1}(T)$ if $\omega\colon Z^1 \to Z^2$ does not affect T. For example, view $\widehat{T}^- = \pi^*(\infty) \cap \Sigma^{\mathtt{P}}$ both as a 0–dimensional subscheme of $Z^{\mathtt{P}}$ and Z.

Now, apply the blow-up construction of Section 5.2 to the 0–dimensional scheme $\widehat{T}^-$ in Z in order to obtain a parabolic spectral triple $(Z^{\mathrm{int}}, \Sigma^{\mathrm{int}}, N_\bullet^{\mathrm{int}})$, called the *absolute parabolic spectral triple* of $(\mathcal{E}_\bullet,\theta)$: the surface Z^{int} is the blow-up $\boldsymbol{\rho}_{\widehat{T}^-}$ of Z along $\widehat{T}^-$, and the coherent sheaf N^{int} is defined as the proper transform $\boldsymbol{\rho}^{\sigma}_{\widehat{T}^-}\boldsymbol{\eta}^{\sigma}_{T^+} N$ of N with respect to $\boldsymbol{\eta}_{T^+}\boldsymbol{\rho}_{\widehat{T}^-}$. It is supported on the proper transform

$$\Sigma^{\mathrm{int}} = \boldsymbol{\rho}^{\sigma}_{\widehat{T}^-}\Sigma$$

of Σ with respect to $\boldsymbol{\rho}_{\widehat{T}^-}$. Set

$$N_\bullet^{\mathrm{int}} := \mathrm{Del}_{\mathtt{E}^+}(\boldsymbol{\eta}_{T^+} \circ \boldsymbol{\rho}_{\widehat{T}^-})^{\sigma}(M_\bullet^{\mathtt{P}}).$$

The parabolic divisor of $N_\bullet^{\mathrm{int}}$ is

$$\mathtt{D}^{\mathrm{int}} = \boldsymbol{\rho}^*_{\widehat{T}^-}\boldsymbol{\eta}^{\sigma}_{T^+}\mathtt{D}.$$

We call $\widehat{\mathtt{E}}^-$ the exceptional divisor of $\boldsymbol{\rho}_{\widehat{T}^-}$. Now, set $\mathtt{P}^{\mathrm{int}} = \mathtt{D}^{\mathrm{int}} \setminus \pi^{-1}\infty$. We call Z^{int} the *absolute surface*, Σ^{int} the *absolute spectral curve* and N^{int} the *absolute spectral sheaf*.

It is possible to reconstruct the original parabolic Higgs bundle from the absolute parabolic spectral triple: by Lemma 5.12, we have

$$N^{\mathrm{P}} = (\boldsymbol{\eta}_{T^+} \circ \boldsymbol{\rho}_{\hat{T}^-})_* N^{\mathrm{int}} \quad \text{and} \quad M^{\mathrm{P}}(-\mathrm{P})_\bullet = (\boldsymbol{\eta}_{T^+} \circ \boldsymbol{\rho}_{\hat{T}^-})_* N^{\mathrm{int}}_\bullet(-\mathrm{P}^{\mathrm{int}}).$$

The parabolic vector bundle $\mathcal{E}_\bullet$ is $\pi_* M^{\mathrm{P}}(-\mathrm{P})_\bullet$, and the Higgs field θ is the direct image of multiplication map by the global section x_{P} on Z^{P}.

The dual divisor $\widehat{\mathrm{P}}$ is defined as the image of $\widehat{T}^-$ under $\hat{\pi}\colon \mathbb{P}^1 \times \widehat{\mathbb{P}}^1 \to \widehat{\mathbb{P}}^1$, ie

$$\widehat{\mathrm{P}} := \hat{\pi}(\widehat{T}^-).$$

We equally set $\widehat{\mathrm{D}} = \widehat{\mathrm{P}} + \widehat{\infty}$. All the maps we have constructed so far fit into the commutative diagram

$$\begin{tikzcd}[column sep=small]
 & & Z^{\mathrm{int}} \arrow[dl, "\boldsymbol{\rho}_{\hat{T}^-}"'] \arrow[dr, "\boldsymbol{\rho}_{T^-}"] & & \\
 & Z \arrow[dl, "\boldsymbol{\eta}_{T^+}"'] \arrow[dr, "\boldsymbol{\eta}_{T^-}"] & & \widehat{Z} \arrow[dl, "\hat{\boldsymbol{\eta}}_{\hat{T}^-}"'] \arrow[dr, "\hat{\boldsymbol{\eta}}_{\hat{T}^+}"] & \\
Z^{\mathrm{P}} & & \mathbb{P}^1 \times \widehat{\mathbb{P}}^1 & & \widehat{Z}^{\widehat{\mathrm{P}}}.
\end{tikzcd} \tag{27}$$

Here, the surfaces $\widehat{Z}^{\widehat{\mathrm{P}}}$, $\widehat{Z}$ and the related maps are defined in an analogous manner to above. More precisely, recall that $\boldsymbol{\eta}_{T^+}$ is the blow-up of Z^{P} at the points T^+, and $\boldsymbol{\eta}_{T^-}$ is the blow-down of Z along the proper transforms $\mathrm{E}^- = \boldsymbol{\eta}^{\sigma}_{T^+}(\pi^*(\mathrm{P}))$ in Z of the fibers of π in Z^{P} over the points of P. As usual, call π and $\hat{\pi}$ the two projections of $\mathbb{P}^1 \times \widehat{\mathbb{P}}^1$. For any $p \in \mathrm{P}$ the proper transform $\boldsymbol{\eta}^{\sigma}_{T^+}(\pi^*(p))$ of the fiber over p contracts to the point in $\mathbb{P}^1 \times \widehat{\mathbb{P}}^1$ which is the intersection of the ∞–fiber of $\hat{\pi}$ with the fiber of π over p: denote by $T^- \subset \hat{\pi}^*(\widehat{\infty})$ the union of these points for all $p \in \mathrm{P}$; this is a finite set. Furthermore, recall that $\widehat{T}^-$ is the intersection of Σ^0 and the ∞–fiber of π, or said differently, the intersection of the fibers of $\hat{\pi}$ over $\widehat{\mathrm{P}}$ in $\mathbb{P}^1 \times \widehat{\mathbb{P}}^1$ and the ∞–fiber of π. Then, the map $\hat{\boldsymbol{\eta}}_{\hat{T}^-}$ is the blow-up of $\mathbb{P}^1 \times \widehat{\mathbb{P}}^1$ in the 0–dimensional subscheme $\widehat{T}^-$, $\boldsymbol{\rho}_{T^-}$ is the blow-up of $\widehat{Z}$ in the 0–dimensional subscheme T^-, and finally $\hat{\boldsymbol{\eta}}_{\hat{T}^+}$ is the blow-down of the proper transform $\widehat{\mathrm{E}}^+ = \hat{\boldsymbol{\eta}}^{\sigma}_{\hat{T}^-}(\hat{\pi}^*(\widehat{\mathrm{P}}))$ of the fibers of $\hat{\pi}$ over $\widehat{\mathrm{P}}$ with respect to $\hat{\boldsymbol{\eta}}_{\hat{T}^-}$. Call $\widehat{T}^+$ the finite set where these fibers contract in the 0–section of $\hat{\pi}$ in $\widehat{Z}^{\widehat{\mathrm{P}}}$. In other words, the relation between $\mathbb{P}^1 \times \widehat{\mathbb{P}}^1$, $\widehat{Z}^{\widehat{\mathrm{P}}}$ and $\widehat{Z}$ with respect to the points $\widehat{\mathrm{P}}$ and the projection $\hat{\pi}$ is the same as the relation between $\mathbb{P}^1 \times \widehat{\mathbb{P}}^1$, Z^{P} and Z with respect to the points P and the projection π; whereas Z^{int} is the fibered product of Z and $\widehat{Z}$ over $\mathbb{P}^1 \times \widehat{\mathbb{P}}^1$. Therefore, Z^{int} has two projections to projective lines: $\pi = \pi \circ \boldsymbol{\eta}_{T^+} \circ \boldsymbol{\rho}_{\hat{T}^-}$ to the base $\mathbb{P}^1$ of the geometrically ruled surface Z^{P}, and $\hat{\pi} = \hat{\pi} \circ \hat{\boldsymbol{\eta}}_{\hat{T}^+} \circ \boldsymbol{\rho}_{T^-}$ to the base $\widehat{\mathbb{P}}^1$ of $\widehat{Z}^{\widehat{\mathrm{P}}}$. Let $\widehat{M}^{\widehat{\mathrm{P}}}$ be the direct image sheaf $(\hat{\boldsymbol{\eta}}_{\hat{T}^+} \circ \boldsymbol{\rho}_{T^-})_*(\mathrm{Add}_{\widehat{\mathrm{E}}^+} N^{\mathrm{int}})$ and denote by $\widehat{\Sigma}^{\widehat{\mathrm{P}}}$ its support.

Definition 8.2 The direct image parabolic sheaf $\widehat{\pi}_* \widehat{M}^{\widehat{\mathrm{P}}}_\bullet(-\widehat{\mathrm{P}})$ on $\widehat{\mathbb{P}}^1$ with parabolic points $\widehat{\mathrm{D}}$ will be called $\widehat{\mathcal{E}}'_\bullet$.

By the definition of $\widehat{M}^{\widehat{\mathrm{P}}}_\bullet$, $\widehat{\mathcal{E}}'_\bullet$ is isomorphic to

$$\big(\widehat{\pi}_* \circ (\widehat{\boldsymbol{\eta}}_{\widehat{T}^+})_* \circ (\boldsymbol{\rho}_{T^-})_* \circ \mathrm{Add}_{\widehat{E}^+} \mathrm{Del}_{E^+} \circ \boldsymbol{\rho}^{\sigma}_{\widehat{T}^-} \circ \boldsymbol{\eta}^{\sigma}_{T^+}(M^{\mathrm{P}}_\bullet)\big)(-\widehat{\mathrm{P}}).$$

Remark that by construction, the parabolic weights of $\widehat{\mathcal{E}}'_\bullet$ are between 0 and 1. Furthermore, similarly to $(x_{\mathrm{P}}, y_{\mathrm{P}})$ on Z^{P}, there exists a pair of globally well-defined parameters $(\widehat{x}_{\widehat{\mathrm{P}}}, \widehat{y}_{\widehat{\mathrm{P}}})$ on $\widehat{Z}^{\widehat{\mathrm{P}}}$.

Definition 8.3 Denote the direct image of multiplication by the global section $-\widehat{x}_{\widehat{\mathrm{P}}}$ on $\widehat{M}^{\widehat{\mathrm{P}}}_\bullet(-\widehat{\mathrm{P}})$ by $\widehat{\theta}'$.

Remark 8.4 One checks without difficulty that $\widehat{\theta}'$ respects the parabolic filtration of $\widehat{\mathcal{E}}'_\bullet$; hence $(\widehat{\mathcal{E}}'_\bullet, \widehat{\theta}')$ is a parabolic Higgs bundle. Using the notation introduced in Section 1, the definitions above can be written as

$$(\widehat{\mathcal{E}}'_\bullet, \widehat{\theta}') = \pi_H(\widehat{M}^{\widehat{\mathrm{P}}}_\bullet, -\widehat{x}_{\widehat{\mathrm{P}}}),$$

or equivalently by Lemma 3.18 as

$$(\widehat{\mathcal{E}}'_\bullet, \widehat{\theta}') = \pi_H\big((-1)^*_{\widehat{Z}^{\widehat{\mathrm{P}}}} \widehat{M}^{\widehat{\mathrm{P}}}_\bullet, \widehat{x}_{\widehat{\mathrm{P}}}\big).$$

Notice that the construction of $(\widehat{\mathcal{E}}'_\bullet, \widehat{\theta}')$ only assumes that θ has first-order poles at finite distance and no singularity at infinity, but no assumption is made on the residues of θ in these singularities, nor about stability or the degree of $\mathcal{E}$. However, the reason why we introduced this construction is that under the assumptions of [10], the two definitions of the Nahm transform agree:

Theorem 8.5 *Assume $(\mathcal{E}_\bullet, \theta)$ satisfies the conditions of Section 2: the residue of θ is semisimple at all points of P, with no multiple eigenvalues except possibly for 0, and the additional properties (1)–(3) of Section 2 are fulfilled. Furthermore the Higgs field at infinity has a second-order pole of the form ((5), (6), (7)) with all $\{\lambda^\infty_{1+a_l}, \dots, \lambda^\infty_{a_l+1}\}$ nonvanishing and distinct for a fixed $1 \le l \le \widehat{n}$ and a compatible parabolic structure with positive weights. Then the parabolic Higgs bundles $(\widehat{\mathcal{E}}'_\bullet, \widehat{\theta}')$ and $(\widehat{\mathcal{E}}^{\mathrm{tr}}_\bullet, \widehat{\theta})$ are isomorphic.*

Definition 8.6 In what follows of this paper, this common object will be referred to as $(\widehat{\mathcal{E}}_\bullet, \widehat{\theta})$.

Proof It follows from the results discussed in Theorem 2.4 that on the open set $\widehat{\mathbb{C}} \setminus \widehat{\mathrm{P}}$ the two Higgs bundles $(\widehat{\mathcal{E}}', \widehat{\theta}')$ and $(\widehat{\mathcal{E}}, \widehat{\theta})$ agree. Indeed, over $\widehat{\mathbb{C}} \setminus \widehat{\mathrm{P}}$ the two surfaces $\mathbb{P}^1 \times \widehat{\mathbb{P}}^1$ and $\widehat{Z}^{\widehat{\mathrm{P}}}$ are isomorphic, the same holds for the sheaves M^0 and $\widehat{M}^{\widehat{\mathrm{P}}}$ and the coordinates ξ and $\hat{x}_{\widehat{\mathrm{P}}}$. Finally, the two factors of $\frac{1}{2}$ in the definition of $(\widehat{\mathcal{E}}, \widehat{\theta})$ (namely, that of $\theta_\xi = \theta - \xi/2$ and $\widehat{\theta} = -z(\xi)/2$) cancel each other. Hence, we only need to check that the extensions to the singularities agree as well.

Lemma 5.12 implies

$$(\boldsymbol{\eta}_{T^-})_* \circ (\boldsymbol{\rho}_{\widehat{T}^-})_*(N^{\mathrm{int}}) = N^0(\pi^{-1}(\infty))$$

as sheaves, and because we have the equality $\widehat{\mathrm{E}}^- \cap \Sigma^{\mathrm{int}} = (\pi \circ \boldsymbol{\eta}_{T^-} \circ \boldsymbol{\rho}_{\widehat{T}^-})^{-1}(\infty) \cap \Sigma^{\mathrm{int}}$ and $(\hat{\pi} \circ \boldsymbol{\eta}_{T^-} \circ \boldsymbol{\rho}_{\widehat{T}^-})^{-1}(\widehat{\mathrm{P}}) = \widehat{\mathrm{E}}^- \cup \widehat{\mathrm{E}}^+$, this implies

$$(\boldsymbol{\eta}_{T^-})_* \circ (\boldsymbol{\rho}_{\widehat{T}^-})_* \mathrm{Add}_{\widehat{\mathrm{E}}^+}(N^{\mathrm{int}}) = N^0(\hat{\pi}^{-1}(\widehat{\mathrm{P}})).$$

Since the projections $\hat{\pi} \circ \boldsymbol{\eta}_{T^-} \circ \boldsymbol{\rho}_{\widehat{T}^-}$ and $\hat{\pi} \circ \hat{\boldsymbol{\eta}}_{\widehat{T}^+} \circ \boldsymbol{\rho}_{T^-}$ from Z^{int} to $\widehat{\mathbb{P}}^1$ are the same, we have the isomorphism of sheaves

$$\hat{\pi}_* N^0(\widehat{\mathrm{P}}) = \hat{\pi}_* \widehat{M}^{\widehat{\mathrm{P}}},$$

because both are equal to the direct image of $\mathrm{Add}_{\widehat{\mathrm{E}}^+} N^{\mathrm{int}}$ with respect to the same projection. However, the direct image of the parabolic structure of N^0 is not the correct one: indeed, on one hand the set of parabolic points of N^0 contains the points of $\mathrm{E}^+ \cap \Sigma^0$ with trivial parabolic structure, so these will induce extra parabolic points on $\widehat{\mathbb{P}}^1$ with trivial structure; and on the other hand it does not contain the points $\widehat{\mathrm{E}}^+ \cap \Sigma^0 \subset \hat{\pi}^{-1}\widehat{\mathrm{P}}$ so that the direct image with respect to $\hat{\pi}$ of the parabolic structure on N^0 does not really make sense. On the other hand, we modified the parabolic divisor of N^{int} so that these problems do not occur when we push it down. Hence, in order to prove equality of the bundles $\widehat{\mathcal{E}}^{\mathrm{tr}}$ and $\widehat{\mathcal{E}}'$ it is sufficient to prove that $\pi_* N^0 = \widehat{\mathcal{E}}^{\mathrm{tr}}$; whereas for their parabolic structure, we will work directly with $\widehat{M}^{\widehat{\mathrm{P}}}(-\hat{\pi}^{-1}(\widehat{\mathrm{P}}))$.

Now, since local sections of N^0 are sections of M^0 which vanish in the points $T^- = (\{\widehat{\infty}\}) \cap \Sigma^0$ and $\widehat{T}^- = (\widehat{\mathrm{P}} \times \{\infty\}) \cap \Sigma^0$, this means precisely that if ζ is a local coordinate of $\widehat{\mathbb{P}}^1$ at $\widehat{\infty}$ then the local sections near $\widehat{\infty}$ of the direct image $\hat{\pi}_* N^0$ can be represented by a local section of the sheaf M^0 multiplied by bump functions concentrated at the spectral points of ζ whose heights converge to 0 up to first order as $\zeta \to 0$. On the other hand, the induced extension at infinity is defined precisely by admitting a representation by bump functions of constant height as $\zeta \to 0$, and the local sections of the transformed extension are obtained from those of the induced extension upon multiplication of the latter with $\xi^{-1} = \zeta$ (cf the discussion before (9)). It is proved in [10, Proposition 4.24] that a D''_ξ–harmonic 1–form $\varphi = \varphi_1 dz + \varphi_{\bar{1}} d\bar{z}$ represents the element of M^0_ξ which is the class of $\{\varphi_1(q(\xi))dz\}$ modulo the image of θ, where $q(\xi)$ runs over the finite set of spectral points of ξ. It follows that multiplying the harmonic representatives by bump functions of height converging to 0 instead of constant height amounts to taking sections of M^0 that vanish at $\widehat{\infty}$. Therefore, locally near the dual infinity the isomorphism of holomorphic bundles $\hat{\pi}_* N^0 = \widehat{\mathcal{E}}^{\mathrm{tr}}$ holds. Similarly, near a logarithmic singularity ξ_l the change of trivializations to obtain $\widehat{\mathcal{E}}^{\mathrm{tr}}$ from $\widehat{\mathcal{E}}^{\mathrm{ind}}$ is to take bump functions of height decaying as $|\xi - \xi_l|$ near the spectral points converging to $\infty \in \mathbb{P}^1$. This amounts precisely to taking local sections of M^0 vanishing up to first order on the divisor $\{\infty\} \times \widehat{\mathtt{P}}$. Such local sections of M^0 are by definition the local sections of N^0, therefore the direct image of N^0 in the logarithmic singularities $\widehat{\mathtt{P}}$ of $\hat{\theta}$ is also equal to $\widehat{\mathcal{E}}^{\mathrm{tr}}$; this implies an isomorphism of the bundles and Higgs fields.

It remains to identify the parabolic structures of the direct image. First of all, the set of parabolic points of $\widehat{\mathcal{E}}^{\mathrm{tr}}_\bullet$ in $\widehat{\mathbb{P}}^1$ is $\widehat{\mathtt{D}} = \widehat{\mathtt{P}} \cup \{\widehat{\infty}\}$, and the same holds for $\widehat{\mathcal{E}}'_\bullet$ because the deletion procedure removes extra parabolic points with trivial structure from $\widehat{M}^{\widehat{\mathtt{P}}}_\bullet$. Second, by [10, Section 4.6], near the punctures the local bases of the transformed bundle defined by the representatives given in the previous paragraph are adapted to the transformed harmonic metric. Hence, the direct images of the parabolic filtrations of $\widehat{M}^{\widehat{\mathtt{P}}}(-\hat{\pi}^{-1}(\widehat{\mathtt{P}}))_\bullet$ are the filtrations of $\widehat{\mathcal{E}}^{\mathrm{tr}}_\bullet$. Furthermore, the same thing holds for the weights as well. Indeed, the weights of $\widehat{M}^{\widehat{\mathtt{P}}}(-\hat{\pi}^{-1}(\widehat{\mathtt{P}}))_\bullet$ in the points of $\widehat{\mathtt{E}}^- \cap \widehat{\Sigma}^{\widehat{\mathtt{P}}}$ above ξ_l are equal to the parabolic weights α^∞_k of the original bundle on the ξ_l–eigenspace of A at infinity. This follows because the weights of $M^{\mathtt{P}}_\bullet$ at infinity are α^∞_k, since near infinity the isomorphism of sheaves $N^{\mathtt{P}} \cong M^{\mathtt{P}}$ holds, and because our convention is to keep the same parabolic weights for sheaves isomorphic near a parabolic divisor. The weights of $\widehat{M}^{\widehat{\mathtt{P}}}(-\hat{\pi}^{-1}(\widehat{\mathtt{P}}))_\bullet$ in the points $\widehat{T}^+$ are in turn equal to 0 by the definition of adding a new divisor to the parabolic divisor of a parabolic structure. On the other hand, by [10, Theorem 4.37] the nonzero weights of $\widehat{\mathcal{E}}^{\mathrm{tr}}_\bullet$ in $\widehat{\mathtt{P}}$ corresponding to the sections defined above are equal to α^∞_k. This proves equality of the parabolic structures of the two extensions in the logarithmic singularities. Similarly, the weights of $\widehat{M}^{\widehat{\mathtt{P}}}(-\hat{\pi}^{-1}(\widehat{\mathtt{P}}))_\bullet$ at $\widehat{\infty}$ are the nonzero weights α^j_k of $\mathcal{E}_\bullet$ in points of $\mathtt{P}$: again, this follows from the fact that the weights of $M^{\mathtt{P}}_\bullet$ in the points of $\pi^{-1}(\mathtt{P}) \cap \Sigma^{\mathtt{P}} \setminus T^+$ are α^j_k, combined with the local isomorphism of sheaves $N^{\mathtt{P}} \cong M^{\mathtt{P}}$ away from the 0–section of π. By [10, Theorem 4.34] the corresponding weights of $\widehat{\mathcal{E}}^{\mathrm{tr}}_\bullet$ in $\widehat{\infty}$ are also equal to α^j_k; whence the theorem. □

9 Examples

In this section, we illustrate by two examples how our approach allows us to increase the degree of generality of the setup of the transform defined in [10].

9.1 Nilpotent residues

In this first example we show that the transform can be defined for a Higgs field whose polar parts are not necessarily semisimple. The conclusion is that a zero residue matrix at infinity can induce two nilpotent residues of rank one in points of finite distance of the transformed object; hence, there is no analog to nilpotent parts of the preservation of the sum of the ranks of the semisimple parts of the residues by the transform. Parallel to this, the multiplicity of the parabolic weights is also not preserved by the transform. In concrete terms, this means that a parabolic weight α of multiplicity 1 splits up to a multiplicity 2 weight $\alpha/2$; in particular, the total parabolic degree is preserved. Notice finally that in this example we start out with a Higgs field with a rank-two singularity at infinity, and we arrive at one with a logarithmic pole at infinity.

Let u_0 and v_0 be the standard coordinates on $\mathbb{P}^1$ in a neighborhood of 0 and ∞ respectively, let $\mathtt{P} = \{0\} \subset \mathbb{P}^1$, let $\mathcal{E}$ be the rank-two trivial holomorphic bundle $\mathcal{O}_{\mathbb{P}^1} \oplus \mathcal{O}_{\mathbb{P}^1}$, and let θ on the open affine $v_0 = 1$ containing 0 be given in matrix form

$$\begin{pmatrix} \frac{1}{u_0} & 1 \\ -1 & -\frac{1}{u_0} \end{pmatrix}, \tag{28}$$

or, in homogeneous coordinates,

$$\begin{pmatrix} v_0 & u_0 \\ -u_0 & -v_0 \end{pmatrix}. \tag{29}$$

The residue in 0 has two distinct eigenvalues ± 1; whereas the limit of the field at infinity is

$$\begin{pmatrix} 0 & 1 \\ -1 & 0 \end{pmatrix}$$

with eigenvalues $\pm i$. Furthermore, setting $u_0 = 1$ in (29), an easy computation shows that the eigenvalues $\xi(v_0)$ of the matrix can be written

$$\xi(v_0) = \pm i\sqrt{1 - v_0^2} = \pm i\left(1 - \frac{v_0^2}{2} + O(v_0^4)\right). \tag{30}$$

In particular, since the eigenvalues are distinct for $v_0 = 0$, in a neighborhood of ∞ there exists a trivialization of $\mathcal{E}$ in which the matrix of this endomorphism is diagonal. This trivialization then clearly satisfies the properties required by (5)–(7) with first-order term $B_\infty = 0$, since the eigenvalues are functions of v_0^2. However, the assumption that the eigenvalues of B_∞ are all nonvanishing and distinct obviously fails. Finally, let $\alpha_+^0, \alpha_-^0 \in [0, 1[$ be arbitrary weights at the singularity 0 corresponding to the $1-$ and

(-1)–eigenspaces respectively, and let $\alpha_+^\infty, \alpha_-^\infty \in [0,1[$ be arbitrary weights at infinity corresponding to the i– and $(-i)$–eigenspaces.

Consider the standard spectral surface $Z^1 = \mathbb{P}_{\mathbb{P}^1}(\mathcal{O}_{\mathbb{P}^1} \oplus \mathcal{O}_{\mathbb{P}^1}(\mathtt{P}))$, and call the preferred sections of $\mathcal{O}_{Z^1}(1) \otimes \mathcal{O}_{\mathbb{P}^1}(\mathtt{P})$ and $\mathcal{O}_{Z^1}(1)$ x_1 and y_1 respectively. The standard spectral curve Σ^1 in Z^1 is defined by the polynomial

$$\det(x_1 - y_1\theta).$$

Since the standard spectral curve does not intersect the ∞–section, we may assume $y_1 = 1$. Then this polynomial becomes

$$x_1^2 - v_0^2 + u_0^2.$$

The solution of this homogeneous equation is

$$y_1 = 1, \quad x_1 = s^2 - t^2, \quad u_0 = 2st, \quad v_0 = s^2 + t^2, \tag{31}$$

where $[s:t]$ stands for homogeneous coordinates on a smooth rational curve. In other words, the mapping $[s:t] \mapsto u_0, v_0, x_1, y_1$ defined by (31) is a closed embedding whose image is the smooth spectral curve Σ^1. In particular, it has two branches over $[u_0 = 0 : v_0 = 1] = 0 \in \mathbb{C}$ which do not intersect: one of them through $s = 0, t = 1$, the other one through $s = 1, t = 0$. The first corresponds to $x_1 = -1$, that is, the eigenvalue -1 of the residue of θ in 0, the second $x_1 = 1$ to the eigenvalue 1. Similarly, the two branches of Σ^1 over $\infty \in \mathbb{P}^1$ pass through $x_1 = i$ and $x_1 = -i$. The pullback of $\mathcal{E}$ to Σ^1 is the rank-two trivial holomorphic bundle $\mathcal{O}_{\Sigma^1} \oplus \mathcal{O}_{\Sigma^1}$ and the map Θ is then by definition

$$x_1 - y_1\theta \colon \mathcal{O}_{\Sigma^1} \oplus \mathcal{O}_{\Sigma^1} \longrightarrow \mathcal{O}_{\Sigma^1}(2) \oplus \mathcal{O}_{\Sigma^1}(2).$$

Using the above formulae, we obtain for this map the matrix form

$$\begin{pmatrix} -2t^2 & -2st \\ 2st & 2s^2 \end{pmatrix}.$$

A cokernel map for this is left matrix multiplication

$$(s,t) \colon \mathcal{O}_{\Sigma^1}(2) \oplus \mathcal{O}_{\Sigma^1}(2) \longrightarrow \mathcal{O}_{\Sigma^1}(3);$$

in particular, the sheaf $M^{\mathtt{P}}$ is $\mathcal{O}_{\Sigma^1}(3)$. Now, since the residue of θ in 0 has two distinct nonzero eigenvalues $\{\pm 1\}$ and the rank of $\mathcal{E}$ is equal to 2, the set $t^{=0}$ is empty. Therefore, the sheaf $N^{\mathtt{P}}$ is the kernel of evaluation of $\mathcal{O}_{\Sigma^1}(3)$ in the points of the spectral curve Σ^1 over the point at infinity $[u = 0 : v = 1]$. Because Σ^1 is a double cover of $\mathbb{P}^1$ and smooth over infinity, we deduce that $N^{\mathtt{P}} = \mathcal{O}_{\Sigma^1}(1)$. Furthermore, it has four parabolic points: the two points over $0 \in \mathbb{C}$ and the two points over $\infty \in \mathbb{P}^1$

discussed above. The weights are as follows: the one in $x_1=-1$ over $0\in\mathbb{C}$ is α^0_-; the one in $x_1=1$ over $0\in\mathbb{C}$ is α^0_+; the one in $x_1=-i$ over $\infty\in\mathbb{P}^1$ is α^∞_-; the one in $x_1=i$ over $\infty\in\mathbb{P}^1$ is α^∞_+.

Let us now consider the compactification Σ^0 of the open spectral curve $\Sigma^\flat$ in the surface $\mathbb{P}^1\times\widehat{\mathbb{P}}^1$. We have seen that it is the proper transform of Σ^1 with respect to the elementary transformation linking Z^1 to $\mathbb{P}^1\times\widehat{\mathbb{P}}^1$. Let x_0, y_0 be homogeneous coordinates of $\widehat{\mathbb{P}}^1$: they can be thought of as sections of $\mathcal{O}_{\widehat{\mathbb{P}}^1}(1)$ vanishing in 0 and $\widehat{\infty}$ respectively. Since the elementary transformation in question blows up the point $u_0=0$, the relation between these coordinates and the sections of $\mathcal{O}_{Z^1}(1)$ is

$$x_0=x_1,\quad y_0=y_1u_0. \tag{32}$$

Therefore, the parametrization (31) transforms into

$$y_0=2st,\quad x_0=s^2-t^2,\quad u_0=2st,\quad v_0=s^2+t^2, \tag{33}$$

and the equation defining the curve becomes

$$x_0^2u_0^2-y_0^2v_0^2+y_0^2u_0^2. \tag{34}$$

This curve in $\mathbb{P}^1\times\widehat{\mathbb{P}}^1$ is not smooth. Indeed, it is straightforward to check that it has a node in the point $(0,\widehat{\infty})\in\mathbb{P}^1\times\widehat{\mathbb{P}}^1$. On the other hand, it has no other singularity, because on the complement of the fiber of π over 0 it is isomorphic to Σ^1.

The first thing we need to identify in order to perform the elementary transformations for the projection $\hat{\pi}$ is the polar divisor $\widehat{\mathrm{P}}\subset\widehat{\mathbb{P}}^1$. Recall that it is given as the intersection points of Σ^1 with the ∞–fiber of π. Plugging $u_0=0, v_0=1$ into the equation of Σ^1 we get $x_1^2=-1$. We deduce $\widehat{\mathrm{P}}=\{[i:1],[-i:1]\}$; in other words, the points $\{\pm i\}$ of $\widehat{\mathbb{C}}$. We now proceed in two steps: first, compute the coordinates of the surface $\widehat{Z}^i$ obtained by performing an elementary transformation on the point $[i:1]$; then perform another elementary transformation, this time on $[-i:1]$, to obtain the surface $\widehat{Z}^2=\widehat{Z}^{\{-i,i\}}$.

The transformation in i is

$$u_1=u_0(x_0-iy_0)=u_0(s-it)^2,\quad v_1=v_0;$$

since these are projective variables, this is equivalent to

$$u_1=u_0(s-it),\quad v_1=\frac{v_0}{s-it}=s+it.$$

The transformation in $-i$ is

$$u_2=u_1(x_0+iy_0)=u_1(s+it)^2,\quad v_2=v_1,$$

which is again equivalent to

$$u_2 = u_1(s+it) = u_0(s-it)(s+it) = 2st(s^2+t^2), \quad v_2 = \frac{v_1}{s+it} = 1. \tag{35}$$

Meanwhile, since over the points $\pm i \in \widehat{\mathbb{C}}$ the spectral curve has only one branch through $\infty \in \mathbb{P}^1$, these modifications do not introduce new singularities. In other words, the spectral curve is transformed by these modifications into a rational curve $\widehat{\Sigma}^2 = \widehat{\Sigma}^{-i,i}$ with one node, in the 0–section over the point $\widehat{\infty}$.

Next, we need to compute the transformed bundle $\widehat{\mathcal{E}}$. First, because Σ^1 does not pass through the intersection of the 0–section and the 0–fiber of Z^1, the curve $\Sigma \subset Z$ is isomorphic to Σ^1, and Σ is disjoint from the proper transform $\mathtt{E}^+$ of the blow-up $Z \to Z^1$. Furthermore, because the multiplicity of each intersection point of Σ^1 and the ∞–fiber of Z^1 is 1, the curve $\Sigma^{\mathrm{int}} \subset Z^{\mathrm{int}}$ is also isomorphic to Σ^1, which is a projective line. It follows that the proper transform of M^1 in Z^{int} is $\mathcal{O}_{\Sigma^{\mathrm{int}}}(3)$, and the intermediate spectral sheaf is by definition $N^{\mathrm{int}} = \mathrm{Del}_{\mathtt{E}^+}(\mathcal{O}_{\Sigma^{\mathrm{int}}}(3))$, which is $\mathcal{O}_{\Sigma^{\mathrm{int}}}(3)$ because $\mathtt{E}^+$ is disjoint from Σ^{int}. The curve $\widehat{\Sigma}^2$ intersects the 0–section of $\widehat{Z}^2$ over each of the points $\pm i$ with multiplicity 1. Hence, both components of the exceptional divisor $\widehat{\mathtt{E}}^+$ of the blow-up of $\widehat{Z}^2$ in these points intersects the curve Σ^{int} in one point. Therefore, the sheaf $\mathrm{Add}_{\widehat{\mathtt{E}}^+}(N^{\mathrm{int}}) = (\mathcal{O}_{\Sigma^{\mathrm{int}}}(3))(\widehat{\mathtt{E}}^+)$ is isomorphic to $\mathcal{O}_{\Sigma^{\mathrm{int}}}(5)$. By definition, $\widehat{\mathcal{E}}(\{-i,i\})$ is the direct image of $\mathrm{Add}_{\widehat{\mathtt{E}}^+}(N^{\mathrm{int}})$ with respect to the projection of Z^{int} to $\widehat{\mathbb{P}}^1$. By the computations above and because Σ^{int} is a double cover of $\widehat{\mathbb{P}}^1$, this means that $\widehat{\mathcal{E}}$ is the direct image of $\mathcal{O}_{\Sigma^{\mathrm{int}}}(1)$. We claim that

$$\widehat{\mathcal{E}} = \mathcal{O}_{\widehat{\mathbb{P}}^1} \oplus \mathcal{O}_{\widehat{\mathbb{P}}^1}.$$

Indeed, since Σ^{int} is a double cover of $\widehat{\mathbb{P}}^1$, clearly $\widehat{\mathcal{E}}$ is of rank 2. Now $\mathcal{O}_{\Sigma^{\mathrm{int}}}(1)$ has exactly two independent global sections: s and t. They induce global sections of $\widehat{\mathcal{E}}$. Conversely, any global section of $\widehat{\mathcal{E}}$ induces a global section of $\mathcal{O}_{\Sigma^{\mathrm{int}}}(1)$, and is therefore a linear combination of s and t. In different terms, s and t give a global trivialization of $\widehat{\mathcal{E}}$ on $\widehat{\mathbb{P}}^1$.

The last thing to compute is the transformed Higgs field $\widehat{\theta}$. Here, the section $\hat{x}_{\widehat{\mathtt{P}}}$ is called u_2, and $\widehat{\theta}$ is the direct image of multiplication by $-u_2$ on $\mathcal{O}_{\widehat{\Sigma}^{-i,i}}(1)$. Notice that using (33) and (35) we obtain

$$u_2 \cdot s = 2st(s^2+t^2) \cdot s = x_0 y_0 \cdot s + y_0^2 \cdot t, \tag{36}$$

$$u_2 \cdot t = 2st(s^2+t^2) \cdot t = y_0^2 \cdot s - x_0 y_0 \cdot t, \tag{37}$$

therefore the matrix of $\widehat{\theta}$ in the above trivialization is

$$-\begin{pmatrix} x_0 y_0 & y_0^2 \\ y_0^2 & -x_0 y_0 \end{pmatrix}.$$

Here x_0, y_0 are standard projective coordinates for $\widehat{\mathbb{P}}^1$, vanishing at the points 0 and $\widehat{\infty}$ respectively. The matrix form of this map on the affine $\widehat{\mathbb{C}}$ with poles in the points $\{\pm i\}$ can be obtained by setting $y_0 = 1$, and dividing each entry by $(x_0 - i)(x_0 + i)$. The result is

$$\widehat{\theta} = -\frac{1}{x_0^2 + 1}\begin{pmatrix} x_0 & 1 \\ 1 & -x_0 \end{pmatrix}.$$

This matrix clearly has poles in $x_0 = \pm i$, with residues

$$-\begin{pmatrix} \pm i & 1 \\ 1 & \mp i \end{pmatrix},$$

both of which are nilpotent. Furthermore, the spectral curve is ramified over both of these points, as it can be seen for example from the defining equation (34) of Σ^0 upon putting $x_0 = 1$, $v_0 = 1$, and using the observation that taking the proper transform with respect to an elementary transformation does not change the property of the projection to $\widehat{\mathbb{P}}^1$ being ramified or not. Another way of seeing the same thing is as follows. Express v_0 in terms of $x_0 = \xi$ in (30) for example near the value $x_0 = i$; we get

$$v_0^2 = x_0 - i, \tag{38}$$

which is clearly an index–2 ramification for $\widehat{\pi}$. It follows that over the logarithmic poles $\pm i$ of the transformed field, both branches of the spectral curve pass through the ∞–section of $\widehat{\pi}$.

We deduce that the parabolic filtration of $\widehat{\mathcal{E}}$ in these points has to be trivial, hence with only one weight in i (resp. $-i$). Moreover, the norm squared with respect to h of the cokernel vector is equivalent to $|v_0|^{2\alpha_+^\infty}$, which is equal to $|x_0 - i|^{\alpha_+^\infty}$ because of (38); therefore, this unique parabolic weight in the point i (resp. $-i$) is $\alpha_+^\infty/2$ (resp. $\alpha_-^\infty/2$). On the other hand, near $\widehat{\infty}$ the matrix for $\widehat{\theta}$ looks up to higher-order terms like

$$-\begin{pmatrix} \frac{1}{x_0} & \frac{1}{x_0^2} \\ \frac{1}{x_0^2} & -\frac{1}{x_0} \end{pmatrix};$$

this converges to 0 as x_0 goes to infinity, and the first-order term in its Taylor series is

$$\begin{pmatrix} -1 & 0 \\ 0 & 1 \end{pmatrix},$$

with eigenvalues $\{\pm 1\}$. The parabolic weight of the ± 1–eigenspace is $\alpha_\pm^0$.

9.2 Higher order pole

Although so far we assumed that the Higgs field has at most logarithmic singularities in the singular points at finite distance, it is relatively clear that iterating the construction several times according to the order of the poles, one can get a transform for Higgs bundles with higher-order poles; the transformed Higgs field will then have a ramification at infinity. Here we describe the archetype of this phenomenon: the original Higgs bundle has a maximal ramification at infinity, and the transformed field has a pole in the origin whose order equals the index of this ramification.

Let $r \geq 2$ and take $\mathcal{E}$ to be the rank-r trivial holomorphic bundle $\mathcal{O}_{\mathbb{P}^1}^{\oplus r}$, with $\mathtt{P} = \{0\}$ the only regular parabolic point. Define the Higgs field as the map

$$\theta\colon \mathcal{O}_{\mathbb{P}^1}^{\oplus r} \longrightarrow \mathcal{O}_{\mathbb{P}^1}(1)^{\oplus r}$$

defined in some global trivialization $\{\tau_1, \dots, \tau_r\}$ by the matrix

$$\begin{pmatrix} 0 & u_0 & 0 & \cdots & 0 \\ 0 & 0 & u_0 & \cdots & 0 \\ \vdots & \vdots & \ddots & \ddots & \vdots \\ 0 & 0 & \cdots & 0 & u_0 \\ v_0 & 0 & \cdots & 0 & 0 \end{pmatrix}, \tag{39}$$

where u_0, v_0 are the global sections of $\mathcal{O}_{\mathbb{P}^1}(1)$ vanishing in 0 and ∞ respectively. Here and in the rest of this section we identify $\mathcal{O}_{\mathbb{P}^1}(\mathtt{P})$ with $\mathcal{O}_{\mathbb{P}^1}(1)$, and correspondingly replace $\mathtt{P}$ by 1 in all upper and lower indices.

We immediately see that at infinity the matrix (39) has only 0 eigenvalues, so we deduce that the singular set of the transform will be $\widehat{\mathtt{P}} = \{0\}$. Furthermore, in a local affine coordinate v centered at infinity, the matrix becomes

$$\begin{pmatrix} 0 & 1 & 0 & \cdots & 0 \\ 0 & 0 & 1 & \cdots & 0 \\ \vdots & \vdots & \ddots & \ddots & \vdots \\ 0 & 0 & \cdots & 0 & 1 \\ v & 0 & \cdots & 0 & 0 \end{pmatrix}.$$

It is clear that there exists no nontrivial subspace invariant by both the constant and the first-order term of this matrix. Since the parabolic filtration has to be preserved by the polar part of the field, this then implies that the only possible filtration in this point is the trivial filtration

$$\mathcal{E}|_\infty = F_0\mathcal{E}|_\infty \supset F_1\mathcal{E}|_\infty = \{0\}.$$

Let us call the corresponding weight α^∞. By the general hypotheses made on the weights, α^∞ is in $]0,1[$. For the sake of simplicity, let us also suppose that $\alpha^\infty < 1/r$.

On the other hand, the residue of the Higgs field (39) in the only regular singular point $0 \in \mathbb{C}$ is of rank 1, so the transformed bundle will be of rank 1. Moreover, since this residue has only 0 eigenvalues, the standard spectral curve passes through the 0–section over the polar point 0 with maximal multiplicity r. The Admissibility condition 1 then forces the parabolic structure in this polar point to be trivial, ie the filtration is trivial and the only weight is 0.

The spectral surface is $Z^1 = \mathbb{P}_{\mathbb{P}^1}(\mathcal{O}_{\mathbb{P}^1} \oplus \mathcal{O}_{\mathbb{P}^1}(1))$, and

$$\Theta_1 = x_1 - y_1\theta = \begin{pmatrix} x_1 & -y_1u_0 & 0 & \cdots & 0 \\ 0 & x_1 & -y_1u_0 & \cdots & 0 \\ \vdots & \vdots & \ddots & \ddots & \vdots \\ 0 & 0 & \cdots & x_1 & -y_1u_0 \\ -y_1v_0 & 0 & \cdots & 0 & x_1 \end{pmatrix}$$

implies that the spectral curve is

$$\Sigma^1 = (\det(\Theta_1)) = (x_1^r - (-y_1)^r u_0^{r-1} v_0).$$

This curve is singular in the point $x_1 = 0, u_0 = 0$ if $r > 2$. Therefore, instead of working on this surface, we choose first to perform an elementary transformation of the point 0 and reduce our problem to the case of a smooth curve in $\mathbb{P}^1 \times \widehat{\mathbb{P}}^1$. The elementary transformation to apply is given by the coordinate changes $x_0u_0 = x_1$, $y_0 = y_1$, and now we consider $Q^{-1}\theta$ as a map $\mathcal{E} \to \mathcal{F}(1) = \mathcal{O}_{\mathbb{P}^1}^{\bigoplus(r-1)} \oplus \mathcal{O}_{\mathbb{P}^1}(1)$. In concrete terms, denoting by $\boxtimes$ external tensor product of sheaves on a product space, the map Q in the Diagram (23) has the form $\operatorname{diag}(u_0, \dots, u_0, 1)$ in the same basis as above, and this means that

$$\Theta_0 = Q^{-1}(x_0u_0 - y_0\theta)\colon \mathcal{O}_{\mathbb{P}^1}^{\bigoplus r} \longrightarrow \big(\mathcal{O}_{\mathbb{P}^1}^{\bigoplus(r-1)} \oplus \mathcal{O}_{\mathbb{P}^1}(1)\big) \boxtimes \mathcal{O}_{\widehat{\mathbb{P}}^1}(1)$$

is given by

$$\begin{pmatrix} x_0 & -y_0 & 0 & \cdots & 0 \\ 0 & x_0 & -y_0 & \cdots & 0 \\ \vdots & \vdots & \ddots & \ddots & \vdots \\ 0 & 0 & \cdots & x_0 & -y_0 \\ -y_0v_0 & 0 & \cdots & 0 & x_0u_0 \end{pmatrix}.$$

Therefore, the spectral curve Σ^0 is defined by the equation

$$x_0^r u_0 - (-y_0)^r v_0 = 0,$$

and this is clearly a nonsingular rational curve smoothly parametrized by (x_0, y_0): namely, one has $u_0 = (-y_0)^r$, $v_0 = x_0^r$. As an effect of passing to the product surface we therefore desingularize the curve, and applying the map Q^{-1} we get rid of the extra fiber of multiplicity $(r-1)$ over 0 of the total transform of the curve. Over the origin in $\mathbb{P}^1$ one has $u_0 = 0, v_0 = 1$, so necessarily $y_0 = 0$, which means that the only point of the spectral curve over the origin is the point $(0, \widehat{\infty}) \in \mathbb{P}^1 \times \widehat{\mathbb{P}}^1$. Moreover, putting $x_0 = 1$ the equation for the curve near this point becomes

$$u_0 = (-1)^r y_0^r,$$

which means that the projection to $\mathbb{P}^1$ has a ramification of index r in this point. Similarly, the curve passes through $(\infty, 0)$ and projection to $\mathbb{P}^1$ has a ramification of index r over infinity. In particular, it intersects the 0– and ∞–fibers of π in these points with multiplicity r. A simple computation shows that the map

$$A\colon \big(\mathcal{O}_{\mathbb{P}^1}^{\oplus(r-1)} \oplus \mathcal{O}_{\mathbb{P}^1}(1)\big) \boxtimes \mathcal{O}_{\widehat{\mathbb{P}}^1}(1) \longrightarrow \mathcal{O}_{\mathbb{P}^1}(1) \boxtimes \mathcal{O}_{\widehat{\mathbb{P}}^1}(r),$$
$$A = \big(y_0 x_0^{r-2} v_0, y_0^2 x_0^{r-3} v_0, \dots, y_0^{r-1} v_0, x_0^{r-1}\big),$$

is a cokernel for Θ_0. In particular, the cokernel sheaf M^0 of Θ_0 is the restriction of $\mathcal{O}_{\mathbb{P}^1}(1) \boxtimes \mathcal{O}_{\widehat{\mathbb{P}}^1}(r)$ to Σ^0, which is equal to $\mathcal{O}_{\Sigma^0}(2r)$ because Σ^0 is an r-to-1 cover of $\mathbb{P}^1$ and a 1-to-1 cover of $\widehat{\mathbb{P}}^1$. By definition, the sheaf N^0 is the kernel of evaluation of M^0 in the intersection points $(\infty, 0) \in \mathbb{P}^1 \times \widehat{\mathbb{P}}^1$ and $(0, \widehat{\infty}) \in \mathbb{P}^1 \times \widehat{\mathbb{P}}^1$ of the spectral curve with the ∞–fiber of π and the 0–fiber in the ∞–section. We have seen above that these intersections are of multiplicity r. It follows that $N^0 = M^0(-2r) = \mathcal{O}_{\Sigma^0}$. Since $\hat{\pi}\colon \Sigma^0 \to \widehat{\mathbb{P}}^1$ is an isomorphism, we deduce that $\widehat{\mathcal{E}} = \mathcal{O}_{\widehat{\mathbb{P}}^1}$.

Let us now identify the transformed Higgs field $\widehat{\theta}$; for this purpose, we need to perform additional elementary transformations on $\mathbb{P}^1 \times \widehat{\mathbb{P}}^1$, but this time with respect to the projection $\hat{\pi}$. Namely, since the spectral curve Σ^0 intersects the ∞–section of $\hat{\pi}$ in its 0–fiber, we need to introduce $u_1 = u_0 x_0, v_1 = v_0$. The equation of the proper transformed curve $\widehat{\Sigma}^1$ is then given by $x_0^{r-1} u_1 - (-y_0)^r v_1 = 0$. However, this still intersects the ∞–section of $\hat{\pi}$, so we need to do another elementary transformation $u_2 = u_1 x_0, v_2 = v_1$, and continue this procedure until the proper transformed curve no longer intersects the ∞–section, that is, $u_r = u_0 x_0^r, v_r = v_0$. The equation of the proper transformed curve $\widehat{\Sigma}^r$ is now $u_r - (-y_0)^r v_r = 0$. Since this curve does not intersect the ∞–section of $\hat{\pi}$, we may set $v_r = 1$. Then the curve is given by $u_r = (-y_0)^r$. Now, one has by definition $\widehat{\theta} = \hat{\pi}_*(-u_r \cdot)$, hence we see that the transformed Higgs field has the form

$$-(-y_0)^r\colon \mathcal{O}_{\widehat{\mathbb{P}}^1} \longrightarrow \mathcal{O}_{\widehat{\mathbb{P}}^1}(r),$$

where $\mathcal{O}_{\widehat{\mathbb{P}}^1}(r)$ stands for $\mathcal{O}_{\widehat{\mathbb{P}}^1}(r\{0\})$. This map therefore has an order r pole at 0 (on the affine $\widehat{\mathbb{C}}$ it can be written as $\pm 1/x_0^r$), and on the other hand it clearly has an order r zero at infinity. Since the fibers are of dimension one, the parabolic filtrations are trivial in both of these points. Similar arguments as in Section 9.1 show that the corresponding parabolic weights in 0 and $\widehat{\infty}$ are $r\alpha^{\infty}$ and 0 respectively.

10 Quasi-involutibility

As the second author has proved in [10, Chapter 5], one of the features of the Nahm transform is its involutibility up to a sign. The proof there is done in the framework of integrable connections, and relies on the analysis of a spectral sequence. Our aim in this section is to give a new, more geometric proof of the same result in terms of the techniques developed in this paper (Theorem 10.2).

Define the $\boldsymbol{K}$–linear map

$$-1\colon \mathbb{P}^1 \longrightarrow \mathbb{P}^1$$

to be the extension to $\mathbb{P}^1$ of the ($\boldsymbol{K}$–linear) involution taking an element of $\boldsymbol{K}$ to its additive inverse. It has two fixed points: 0 and ∞. We will denote by $(-1)_{\mathrm{rel}}$ the relative version of this map on any $\mathbb{P}^1$–fibration over a curve. It then induces a map on any blow-up of points in the 0–section or the ∞–section, that we will still denote with the same symbol.

Remark 10.1 In what follows, it will be important to distinguish the divisors $(-1)^*\mathtt{P}$ and $-\mathtt{P}$: the first is the set of points $-p$ where $p \in \mathtt{P}$, whereas the second is the inverse of the divisor $\mathtt{P}$ in the divisor group.

Recall that given a parabolic Higgs bundle $(\mathcal{E}, \theta)$ on $\mathbb{P}^1$ with singularities on $D \cup \{\infty\}$ we have constructed in Section 8 its Nahm transformed Higgs bundle $(\widehat{\mathcal{E}}, \widehat{\theta})$: it is a parabolic Higgs bundle on $\widehat{\mathbb{P}}^1$, the "dual" projective line, with singularities on $\widehat{\mathtt{D}} = \widehat{\mathtt{P}} \cup \{\widehat{\infty}\}$. Here $\widehat{\mathtt{P}}$ is the set of eigenvalues of the leading term of θ at ∞. We have also computed the eigenvalues of $\widehat{\theta}$ at $\widehat{\infty}$, and we realized that they agree with the image $-1(\mathtt{P})$ of $\mathtt{P}$ under the involution. The bidual $\widehat{\widehat{\mathbb{P}}}{}^1$ of $\mathbb{P}^1$ identifies naturally with $\mathbb{P}^1$ itself, hence applying the Nahm transform to $(\widehat{\mathcal{E}}, \widehat{\theta})$, we obtain a parabolic Higgs bundle $(\widehat{\widehat{\mathcal{E}}}, \widehat{\widehat{\theta}})$ on $\mathbb{P}^1$ with singularities on the set $-1(\mathtt{P}) \cup \{\infty\}$. The main result of this section can now be formulated.

Theorem 10.2 *If $(\mathcal{E}, \theta)$ satisfies Admissibility condition 1, then there is a natural isomorphism of parabolic Higgs bundles between $(\widehat{\widehat{\mathcal{E}}}, \widehat{\widehat{\theta}})$ and $(-1)^*(\mathcal{E}, -\theta)$.*

Remark 10.3 There is a sign change of θ between this and the corresponding formula in [10]. This is because there we considered the Higgs field as a 1–form valued endomorphism, and $d(-z) = -dz$.

Proof Starting with the parabolic Higgs bundle $(\widehat{\mathcal{E}}, \widehat{\theta})$ on $\widehat{\mathbb{P}}^1$, we wish to construct its transform. The first object we need to understand is the standard spectral triple

$(W^{\widehat{\mathtt{P}}}, \Xi^{\widehat{\mathtt{P}}}, Q^{\widehat{\mathtt{P}}})$ of $(\widehat{\mathcal{E}}, \widehat{\theta})$. Remember that in Section 8 we constructed the spectral triple $(\widehat{Z}^{\widehat{\mathtt{P}}}, \widehat{\Sigma}^{\widehat{\mathtt{P}}}, \widehat{M}_*^{\widehat{\mathtt{P}}})$ out of $(\mathcal{E}, \theta)$. First, because of the definition $W^{\widehat{\mathtt{P}}} = \boldsymbol{P}(\mathcal{O}_{\widehat{\mathbb{P}}^1} \oplus \mathcal{O}_{\widehat{\mathbb{P}}^1}(\widehat{\mathtt{P}}))$, we see immediately that $W^{\widehat{\mathtt{P}}}$ is naturally isomorphic to the surface $\widehat{Z}^{\widehat{\mathtt{P}}}$. Since $\widehat{\mathcal{E}}$ is the direct image under $\widehat{\pi}\colon \widehat{Z}^{\widehat{\mathtt{P}}} \to \widehat{\mathbb{P}}^1$ of $\widehat{M}^{\widehat{\mathtt{P}}}(-\widehat{\mathtt{P}})$ and $\widehat{\theta}$ is that of multiplication by $-\widehat{x}^{\widehat{\mathtt{P}}}$, it follows also that $\Xi^{\widehat{\mathtt{P}}} = (-1)_{\mathrm{rel}}\widehat{\Sigma}^{\widehat{\mathtt{P}}}$. Finally, by the results of Beauville, Narasimhan and Ramanan [1, Proposition 3.6 and Remark 3.7], we obtain $Q^{\widehat{\mathtt{P}}} = (-1)^*_{\mathrm{rel}}\widehat{M}^{\widehat{\mathtt{P}}}(\widehat{\mathtt{P}})$ as a sheaf. We know furthermore that the parabolic structure of $\widehat{\mathcal{E}}$ induces a parabolic structure on $Q^{\widehat{\mathtt{P}}}$. On the other hand, because $\widehat{M}^{\widehat{\mathtt{P}}}$ has a parabolic structure, the above isomorphism makes $Q^{\widehat{\mathtt{P}}}$ into a parabolic sheaf as well. These two parabolic structures on $Q^{\widehat{\mathtt{P}}}$ agree: indeed, the parabolic structure of $\widehat{\mathcal{E}}$ is the direct image of the one of $\widehat{M}^{\widehat{\mathtt{P}}}$, so the filtration comes from the restriction of $\widehat{M}^{\widehat{\mathtt{P}}}$ to some branches of the spectral curve over the parabolic points of $\widehat{\mathcal{E}}$, hence the two filtrations of $Q^{\widehat{\mathtt{P}}}$ are the same; a similar argument works for the weights.

The next ingredient in the construction is the analog of diagram (27). The surface Z was obtained from $Z^{\mathtt{P}}$ by blowing up the points t^+ of the 0–section mapping to $\mathtt{P}$ under π. Therefore, its analog W for $W^{\widehat{\mathtt{P}}}$ is the blow-up in the points T^+ of the 0–section of $\widehat{\pi}$ over $\widehat{\mathtt{P}}$: clearly, this is the surface $\widehat{Z}$. Now, Z^{int} was obtained from $Z^{\mathtt{P}}$ by blowing up the points $\widehat{T}^-$ in the intersection of the ∞–fiber of π and the spectral curve. Because of $\Xi^{\widehat{\mathtt{P}}} = (-1)_{\mathrm{rel}}\widehat{\Sigma}^{\widehat{\mathtt{P}}}$, the intersection points of the ∞–fiber of $\widehat{\pi}$ and the spectral curve $\Xi^{\widehat{\mathtt{P}}}$ are the points $(-1)^*T^-$ of $\widehat{Z}$. It follows that (-1) induces a natural isomorphism between the absolute surface W^{int} of $(\widehat{\mathcal{E}}, \widehat{\theta})$ and Z^{int}; hence, we will simply write $W^{\mathrm{int}} = (-1)^*_{\mathrm{rel}}Z^{\mathrm{int}}$. Notice that this surface has a projection to both projective lines $\mathbb{P}^1$ and $\widehat{\mathbb{P}}^1$ (although these projections are only rational, and not every fiber is a single line), and the map $(-1)_{\mathrm{rel}}$ above is induced by inversion on the fibers when it is considered as a fibration over $\widehat{\mathbb{P}}^1$. However, it is possible to interpret the same map as induced by inversion of the basis of the other fibration; we will simply write (-1) for this map in the sequel, for any fibration over $\mathbb{P}^1$. Therefore, W^{int} can equally be written as $(-1)^*Z^{\mathrm{int}}$. We now come to an analog of $\widehat{Z}$: this surface was the blow-down in Z^{int} of the proper transforms $\mathtt{E}^-$ of the fibers of π over the points $\mathtt{P}$. Applying this to W^{int} we obtain the result $\widehat{W} = (-1)^*Z$. Finally, arguments similar to the above yield that the analog of $\widehat{Z}^{\widehat{\mathtt{P}}}$ for $(\widehat{\mathcal{E}}, \widehat{\theta})$ is $\widehat{W}^{\mathtt{P}} = (-1)^*Z^{\mathtt{P}}$, that is, the surface $\boldsymbol{P}(\mathcal{O}_{\mathbb{P}^1} \oplus \mathcal{O}_{\mathbb{P}^1}((-1)^*\mathtt{P}))$. We deduce that diagram (27) corresponding to $(\widehat{\mathcal{E}}, \widehat{\theta})$ is

$$\begin{tikzcd}
 & & (-1)^*Z^{\mathrm{int}} \arrow[dl, "(-1)^*\boldsymbol{\rho}_{T^-}"'] \arrow[dr, "(-1)^*\boldsymbol{\rho}_{\widehat{T}^-}"] & & \\
 & \widehat{Z} \arrow[dl, "\widehat{\boldsymbol{\eta}}_{\widehat{T}^+}"'] \arrow[dr, "\widehat{\boldsymbol{\eta}}_{\widehat{T}^-}"] & & (-1)^*Z \arrow[dl, "(-1)^*\boldsymbol{\eta}_{T^-}"'] \arrow[dr, "(-1)^*\boldsymbol{\eta}_{T^+}"] & \\
\widehat{Z}^{\widehat{\mathtt{P}}} & & \mathbb{P}^1 \times \widehat{\mathbb{P}}^1 & & (-1)^*Z^{\mathtt{P}}
\end{tikzcd} \tag{40}$$

The surface in the lower-left corner is the standard spectral surface of $(\widehat{\mathcal{E}}, \widehat{\theta})$, it has a projection to $\widehat{\mathbb{P}}^1$, and the transformed bundle is obtained by taking the proper transform of $Q^{\widehat{\mathtt{P}}}$ in $(-1)^* Z^{\mathrm{int}}$ with respect to $\widehat{\boldsymbol{\eta}}_{\widehat{T}^+} \circ (-1)^* \boldsymbol{\rho}_{T^-}$, deleting the exceptional divisor of $\widehat{\boldsymbol{\eta}}_{\widehat{T}^+}$ from the parabolic divisor, adding the exceptional divisor of $(-1)^* \boldsymbol{\eta}_{T^+}$ to the parabolic divisor, pushing down the result to $(-1)^* Z^{\mathtt{P}}$, then pushing down the result to $\mathbb{P}^1$ by the projection map π of $(-1)^* Z^{\mathtt{P}}$, and finally tensoring by $(-1)^* \mathtt{P}$.

We have seen that the sheaf $Q^{\widehat{\mathtt{P}}}$ is isomorphic to $(-1)^*_{\mathrm{rel}} \widehat{M}^{\widehat{\mathtt{P}}}(\widehat{\mathtt{P}})$. It follows from the property $(\widehat{\boldsymbol{\eta}}_{\widehat{T}^+} \circ \boldsymbol{\rho}_{T^-})^{\sigma} \circ (\widehat{\boldsymbol{\eta}}_{\widehat{T}^+} \circ \boldsymbol{\rho}_{T^-})_* = \mathrm{Id}$ for pure sheaves of dimension 1 on a smooth surface (see Lemma 5.12) that

$$(\widehat{\boldsymbol{\eta}}_{\widehat{T}^+} \circ \boldsymbol{\rho}_{T^-})^{\sigma} Q^{\widehat{\mathtt{P}}} = (-1)^* \mathrm{Add}_{\widehat{\mathtt{E}}^+} N^{\mathrm{int}}.$$

To obtain the absolute spectral sheaf of $(\widehat{\mathcal{E}}, \widehat{\theta})$ we need to delete from the parabolic divisor of $(-1)^* \mathrm{Add}_{\widehat{\mathtt{E}}^+} N^{\mathrm{int}}$ the exceptional divisor $\widehat{\mathtt{E}}^+$ of the blow-up map $\widehat{\boldsymbol{\eta}}_{\widehat{T}^+}$. We deduce from Proposition 6.2 that the absolute spectral sheaf of $(\widehat{\mathcal{E}}, \widehat{\theta})$ is $(-1)^* N^{\mathrm{int}}$ on $(-1)^* Z^{\mathrm{int}}$. The next step in the construction is to add the exceptional divisor $(-1)^* \mathtt{E}^+$ of $(-1)^* \boldsymbol{\eta}_{T^+}$ to the parabolic divisor of $(-1)^* N^{\mathrm{int}}$. By Proposition 8.1, Admissibility condition 1 for $(\mathcal{E}, \theta)$ implies Admissibility condition 2 for N^{int} and $\mathtt{E}^+$. Again by Proposition 6.2, addition and deletion of a divisor are inverses to each other under Admissibility condition 2. We obtain that

$$\mathrm{Add}_{(-1)^* \mathtt{E}^+} (-1)^* N^{\mathrm{int}} = (-1)^* (\mathrm{Add}_{\mathtt{E}^+} N^{\mathrm{int}}) = (-1)^* (\widehat{\boldsymbol{\eta}}_{\widehat{T}^+} \circ \boldsymbol{\rho}_{T^-})^{\sigma} M^{P}.$$

We then consider the direct image of this parabolic sheaf with respect to the blow-up map $(-1)^* (\widehat{\boldsymbol{\eta}}_{\widehat{T}^+} \circ \boldsymbol{\rho}_{T^-})$: by Lemma 5.12, the direct image is $(-1)^* M^{\mathtt{P}}$. The push-down of this to $\mathbb{P}^1$ by the projection $(-1)^* \pi$ of $(-1)^* Z^{\mathtt{P}}$ is $(-1)^* (\mathcal{E}(\mathtt{P})) = ((-1)^* \mathcal{E})((-1)^* \mathtt{P})$; see Lemma 3.18. The final step is to tensor this sheaf by the inverse (in the divisor group) of the effective divisor corresponding to the parabolic set. Here this effective divisor is $(-1)^* \mathtt{P}$. Therefore tensoring $((-1)^* \mathcal{E})((-1)^* \mathtt{P})$ by its inverse, we get precisely $(-1)^* \mathcal{E}$. This proves equality of the bundles $\widehat{\widehat{\mathcal{E}}}$ and $(-1)^* \mathcal{E}$. Clearly, the modifications of the sheaves involved so far transform the parabolic structure of $\widehat{M}^{\widehat{\mathtt{P}}}$ into the parabolic structure of $(-1)^* M^{\mathtt{P}}$ induced via pullback by (-1) of the original structure of $M^{\mathtt{P}}$. Since the direct image by π of the parabolic structure of $M^{\mathtt{P}}(-\mathtt{P})$ is the parabolic structure of $\mathcal{E}$, we also see that the direct image by π of the parabolic structure of $(-1)^* (M^{\mathtt{P}}(-\mathtt{P}))$ is the parabolic structure of $(-1)^* \mathcal{E}$ induced by pullback under (-1) from the parabolic structure of $\mathcal{E}$. Finally, the canonical section $\widehat{\widehat{x}}_{(-1)^* \mathtt{P}}$ of $(-1)^* Z^{\mathtt{P}}$ is $(-1)^* x_{\mathtt{P}}$, where $x_{\mathtt{P}}$ is the canonical section of $Z^{\mathtt{P}}$. By definition, the double transformed Higgs field $\widehat{\widehat{\theta}}$ is the direct image with respect to π of multiplication by $-\widehat{\widehat{x}}_{(-1)^* \mathtt{P}}$. On the other hand, the Higgs field θ is equal to the direct image of multiplication by $x_{\mathtt{P}}$. It follows that $-\widehat{\widehat{\theta}} = (-1)^* \theta$. □

11 The map on moduli spaces

In this section, we show that the Nahm transform is a Kähler isometry between Dolbeault moduli spaces (Theorem 11.4).

As shown by Biquard and Boalch in [2, Theorem 0.2], the moduli space of stable Higgs bundles of parabolic degree 0 with fixed simultaneously diagonalizable polar parts of arbitrary order and fixed parabolic structures is a hyper-Kähler manifold. The two anticommuting complex structures J and I are by definition given by the local holomorphic variations of Higgs bundles (inducing the *Dolbeault* complex structure) and integrable connections (inducing the *de Rham* complex structure) respectively, and the Kähler metric on the moduli space is defined as the L^2–inner product of the harmonic representatives of tangent vectors with respect to the harmonic metric on the endomorphism-bundle and the usual Euclidean metric on the base $\mathbb{C}$. Let now $\mathcal{M}$ stand for the moduli space of Higgs bundles on $\mathbb{P}^1$ with logarithmic singularities in the points of $\mathtt{D}$ with fixed equivalence class of polar parts (3)–(4) and fixed parabolic filtration (1) and with weights α_k^j, and with an irregular singularity of rank one at infinity with fixed equivalence class of polar parts (5)–(7) and fixed parabolic structure with weights α_k^∞, up to complex gauge transformations preserving the parabolic structures.

Lemma 11.1 *The complex dimension of the Zariski tangent space of $\mathcal{M}$ in any point is*

$$2r\hat{r}+2-r-\hat{r}-\sum_{j=1}^{n}\mathrm{rk}(\mathrm{res}(\theta,p_j))^2-\sum_{l=1}^{\hat{n}}\mathrm{rk}(\mathrm{res}(\hat{\theta},\xi_l))^2,$$

where the last two sums are taken for all logarithmic singularities of θ and $\hat{\theta}$ respectively.

Remark 11.2 This formula is in fact invariant under exchanging r with $\hat{r}$ and θ with $\hat{\theta}$, as it should be because of invertibility of the transform.

Proof The computation is done by Boden and Yokogawa in [3] for the case of parabolic Higgs bundles of rank r with only logarithmic singularities on a curve of arbitrary genus g. Notice that the authors there fix the residues of the Higgs field to be block nilpotent, but in fact the same proof works for any other fixed block-diagonal parts as well. The result obtained there is

$$2(g-1)r^2+2+\sum_{j=1}^{n}2f_{p_j}, \tag{41}$$

where $2f_{p_j}$ is the dimension of the adjoint orbit of $\mathrm{res}(\theta, p_j)$ in $\mathfrak{g} = \mathfrak{gl}(r, \boldsymbol{K})$ (see also the count of the dimension of the moduli space of parabolic vector bundles in [8]). We can understand this as coming from excision: the term $2(g-1)r^2+2$ is the degree of $\boldsymbol{\Omega}^1 \otimes \mathrm{End}(\mathcal{E})$ plus the constant 2 coming from global endomorphisms of $\mathcal{E}$, and in the last sum we add up terms arising in a neighborhood of each of the singular points. Explicitly, because we only consider deformations of the Higgs field whose residues in any singular point can be taken into the initial residue by a holomorphic change of basis, the residue of an infinitesimal deformation corresponding to a one-parameter family of such deformations has to be in the adjoint orbit of the residue of θ in $\mathfrak{g}$; the dimension of such choices for the residue in p_j is by definition $2f_{p_j}$. In the case where irregular singularities occur, by the same excision argument we need to define the quantity $2f_p$ in the last sum of (41) as the dimension of the adjoint orbit of the polar part of the Higgs field in $\mathfrak{g} \otimes_{\boldsymbol{K}} \boldsymbol{K}[\varepsilon]/(\varepsilon^{n_p+1})$, where n_p is the Poincaré rank of the singularity in p. Indeed, the local contribution to the dimension formula comes from two distinct sources: first, the change in the parabolic structure of the underlying bundle; second, the modification of the Higgs field. The first amounts to a parameter in the flag manifold associated to the parabolic subgroup of the filtration of $\mathcal{E}_p$. Once the parabolic structure is fixed, the highest-order term of the change in the Higgs field is further restricted to lie in the parabolic subalgebra. Clearly, the sum of the dimensions of these two contributions is equal to the dimension of the orbit of the polar part of the Higgs field.

In our case, the only irregular singularity is infinity, of Poincaré rank 1; let us compute the dimension of its orbit in $\mathfrak{g} \otimes_{\boldsymbol{K}} \boldsymbol{K}[\varepsilon]/(\varepsilon^2)$. For the sake of simplicity, we only do this in the special case $\hat{n} = 2$; the generalization to higher $\hat{n}$ is immediate. So we suppose that at infinity the Higgs field can be written in the block form

$$\frac{1}{2}\begin{pmatrix} \Xi_1 & 0 \\ 0 & \Xi_2 \end{pmatrix} + \varepsilon \begin{pmatrix} \Lambda_1 & 0 \\ 0 & \Lambda_2 \end{pmatrix}, \tag{42}$$

where $\Xi_1 = \xi_1 \,\mathrm{Id}_a$ for some $1 \le a < r$, $\Xi_2 = \xi_2 \,\mathrm{Id}_{r-a}$, $\Lambda_1 = \mathrm{diag}(\lambda_1^\infty, \dots, \lambda_a^\infty)$, and $\Lambda_2 = \mathrm{diag}(\lambda_{1+a}^\infty, \dots, \lambda_r^\infty)$ (see (6)–(7)). We also assume $\xi_1 \neq \xi_2$, and that all the $\lambda_1^\infty, \dots, \lambda_a^\infty$ are nonzero and pairwise distinct, and the same condition for $\lambda_{1+a}^\infty, \dots, \lambda_r^\infty$. Then the stabilizer of the adjoint action is by definition the block matrices

$$\begin{pmatrix} A & B \\ C & D \end{pmatrix} + \varepsilon \begin{pmatrix} \alpha & \beta \\ \gamma & \delta \end{pmatrix}, \tag{43}$$

which commute with (42) modulo ε^2. It is straightforward to check that this holds if and only if $B = 0$, $C = 0$, $\beta = 0$, $\gamma = 0$ and A and $\mathtt{D}$ are diagonal; under these assumptions, α and δ can be arbitrary. Therefore, the dimension of the stabilizer of the polar form (42) is $a + (r-a) + a^2 + (r-a)^2$, and so the dimension of its orbit is $2r^2 - r - a^2 - (r-a)^2$. For general $\hat{n}$, the same argument gives for this dimension $2r^2 - r - \sum_{l=1}^{\hat{n}} (a_l - a_{l-1})^2$. Now, since the transformed Higgs field $\hat{\theta}$ has residue of rank $(a_l - a_{l-1})$ in ξ_l, we can rewrite this as

$$2r^2 - r - \sum_{l=1}^{\hat{n}} \operatorname{rk}(\operatorname{res}(\hat{\theta}, \xi_l))^2.$$

Similarly, it is easy to check that under the assumptions made in (4), for all logarithmic singularity p_j the formula

$$\begin{aligned}
2f_{p_j} &= r^2 - r_j^2 - (r - r_j) \\
&= r^2 - (r - \operatorname{rk}(\operatorname{res}(\theta, p_j)))^2 - \operatorname{rk}(\operatorname{res}(\theta, p_j)) \\
&= 2r \cdot \operatorname{rk}(\operatorname{res}(\theta, p_j)) - \operatorname{rk}(\operatorname{res}(\theta, p_j))^2 - \operatorname{rk}(\operatorname{res}(\theta, p_j)) \\
&= (2r-1) \operatorname{rk}(\operatorname{res}(\theta, p_j)) - \operatorname{rk}(\operatorname{res}(\theta, p_j))^2
\end{aligned}$$

holds, where $r - r_j$ is the rank of $\operatorname{res}(\theta, p_j)$. Plugging these into (41) and using (8) one obtains the dimension of the Zariski tangent as claimed. □

Similarly to $\mathcal{M}$, let us denote by $\widehat{\mathcal{M}}$ the moduli space of stable Higgs bundles of parabolic degree 0 on $\widehat{\mathbb{P}}^1$ with logarithmic singularities in the points of $\hat{\mathtt{D}}$ with fixed equivalence class of polar parts and parabolic structures induced by the transform from the corresponding structures of $(\mathcal{E}, \theta)$ at infinity, and with an irregular singularity of rank one with fixed equivalence class of polar parts and fixed parabolic structure induced by the transform from the corresponding structures of $(\mathcal{E}, \theta)$ in the points of $\mathtt{D}$ — as explained in Section 2 — again up to complex gauge transformations preserving the parabolic structures.

Lemma 11.3 *Let $(\mathcal{E}_\bullet, \theta)$ be a Higgs bundle on $\mathbb{P}^1$ which satisfies Admissibility condition 1. Then the parabolic degrees of $\mathcal{E}_\bullet$ and of its Nahm transform $\hat{\mathcal{E}}_\bullet$ are the same. Furthermore, if $(\mathcal{E}_\bullet, \theta)$ is stable of degree 0, then the same is true for $(\hat{\mathcal{E}}_\bullet, \hat{\theta})$.*

Proof The claim on parabolic degrees follows from Grothendieck–Riemann–Roch, as explained in [10, Section 4.7]. It is also possible to deduce it using the fact that under Admissibility condition 1 all operations involved in passing from $M_\bullet^{\mathtt{P}}$ to $\widehat{M}_\bullet^{\hat{\mathtt{P}}}$ preserve the parabolic Euler characteristic. Indeed, by the parabolic Riemann–Roch theorem (Proposition 3.14) applied to the curve $\mathbb{P}^1$, one has $\text{par-}\chi(\mathcal{E}_\bullet(\mathtt{P}+\infty)) = \text{par-deg}(\mathcal{E}_\bullet) + r$, or equivalently, $\text{par-}\chi(\mathcal{E}_\bullet(P)) = \text{par-deg}(\mathcal{E}_\bullet)$. Of course, a similar relation holds for $\hat{\mathcal{E}}_\bullet$ as well. Finally, we obtain the result using the fact that $\text{par-}\chi(\mathcal{E}_\bullet(\mathtt{P})) = \text{par-}\chi(M_\bullet^{\mathtt{P}})$ because $\pi_* M_\bullet^{\mathtt{P}} = \mathcal{E}(\mathtt{P})_\bullet$, and the analogous statement for $\hat{\mathcal{E}}_\bullet$.

Suppose now $(\mathcal{E}'_\bullet, \theta')$ is a parabolic Higgs subbundle of $(\mathcal{E}_\bullet, \theta)$. By Remark 3.10, we may suppose that the parabolic structure of $\mathcal{E}'_\bullet$ is the structure induced by $\mathcal{E}_\bullet$. By Lemma 6.3, the standard spectral sheaf $(M')^{\mathrm{P}}$ of $\mathcal{E}'_\bullet$ and the divisor $\mathtt{E}^+$ also satisfy Admissibility condition 2, because $(M')^{\mathrm{P}}_\bullet$ is a parabolic subsheaf of $M^{\mathrm{P}}_\bullet$ with the induced parabolic structure. By Lemma 5.14, the proper transform preserves injective maps of sheaves. The same thing holds clearly for the direct image by a blow-up map because it is the inverse of the proper transform, and for addition and deletion, since on the level of sheaves the latter are simply tensoring operations. We conclude that $(\widehat{M}')^{\widehat{\mathrm{P}}}_\bullet$ is a parabolic subsheaf of $\widehat{M}^{\widehat{\mathrm{P}}}_\bullet$, hence $(\widehat{\mathcal{E}}'_\bullet, \widehat{\theta}')$ is a parabolic Higgs subbundle of $(\widehat{\mathcal{E}}_\bullet, \widehat{\theta})$. On the other hand, by the first part of the lemma, the parabolic degree of $\widehat{\mathcal{E}}'_\bullet$ is equal to that of $\mathcal{E}'_\bullet$. In particular, the parabolic degree of $\widehat{\mathcal{E}}'_\bullet$ is positive if and only if the parabolic degree of $\widehat{\mathcal{E}}_\bullet$ is positive. In different terms, if par-deg$(\mathcal{E}_\bullet) = 0$, then $(\mathcal{E}'_\bullet, \theta')$ is destabilizing for $(\mathcal{E}_\bullet, \theta)$ if and only if $(\widehat{\mathcal{E}}'_\bullet, \widehat{\theta}')$ is destabilizing for $(\widehat{\mathcal{E}}_\bullet, \widehat{\theta})$. This proves preservation of stability. □

The lemma allows us to introduce the map

$$\mathcal{N}\colon \mathcal{M} \longrightarrow \widehat{\mathcal{M}} \tag{44}$$

defined by mapping the gauge equivalence class of the Higgs bundle $(\mathcal{E}, \theta)$ to the gauge equivalence class of the Higgs bundle $(\widehat{\mathcal{E}}, \widehat{\theta})$; by an easy adaptation of [6, Lemma 1] to the parabolic case over a curve, this map is well defined. It is a bijective map between hyper-Kähler manifolds. We have the following result.

Theorem 11.4 *The map $\mathcal{N}$ is an algebraic Kähler isometry for the Dolbeault complex structures.*

Proof By Theorem 10.2, the map $\mathcal{N}$ is invertible. For the fact that $\mathcal{N}$ preserves the L^2–metric, we refer to Braam and van Baal [4]: the computations there carry through to this case, because they only make use of the invertibility of the transform and general algebraic properties of Green's operator that are satisfied in this case as well.

Therefore, all that remains is to check that it preserves the complex structure J, and is moreover algebraic. By [10, Proposition 4.15 and Equation (4.14)], the restriction of the holomorphic bundle $\widehat{\mathcal{E}}$ to the affine $\widehat{\mathbb{C}}$ is the first hypercohomology

$$\boldsymbol{H}^1\Big(\mathcal{E} \xrightarrow{\theta_\xi} \mathcal{F}(\mathrm{P})\Big),$$

together with the holomorphic structure induced by the trivial holomorphic structure of $\mathcal{F}(\mathrm{P})$ relative to $\widehat{\mathbb{P}}^1$. (Notice that the sheaf we denoted by $\mathcal{F}$ in [10] is called $\mathcal{F}(\mathrm{P})$ in the present paper.) Furthermore, by the extension [10, Equation (4.35)] of this holomorphic bundle to infinity, we have the isomorphism of holomorphic bundles over $\widehat{\mathbb{P}}^1$,

$$\widehat{\mathcal{E}}^{\mathrm{ind}} = \boldsymbol{H}^1\Big(\mathcal{E} \xrightarrow{y_0\theta - x_0} \mathcal{F}(\mathrm{P}) \otimes \mathcal{O}_{\widehat{\mathbb{P}}^1}(1)\Big), \tag{45}$$

where the right-hand side is endowed with the holomorphic structure induced from the holomorphic structure of the sheaf $\mathcal{F}(\mathrm{P}) \otimes \mathcal{O}_{\widehat{\mathbb{P}}^1}(1)$ relative to $\widehat{\mathbb{P}}^1$. Let T be an open set in an affine complex line, and $(\mathcal{E}(t), \theta(t))$ with $t \in T$ be a local 1–parameter algebraic family of Higgs bundles with fixed singularity data. The transform maps each Higgs bundle in this family to a Higgs bundle $(\widehat{\mathcal{E}}(t), \widehat{\theta}(t))$; we need to show that this is an algebraic family of Higgs bundles on $\widehat{\mathbb{P}}^1$ over T. Because of (45), the induced extensions of $\widehat{\mathcal{E}}$ vary algebraically over T: indeed, the sheaves $\mathcal{E}(t)$ and $\mathcal{F}(t)$ depend algebraically on t, as well as the map $y_0\theta(t) - x_0$, and the first hypercohomology spaces of an algebraic family of sheaf complexes such that all the other hypercohomologies vanish, form again an algebraic family. In order to obtain the transformed extensions from the induced ones we need to modify the local holomorphic sections with nonzero weights. In terms of the spectral sheaf, this amounts to a Hecke-modification at the intersection points of the spectral curve with the fibers of $\widehat{\pi}$ over the set $\widehat{\mathrm{P}}$. Now, since $\theta(t)$ varies algebraically with t, it follows that so does the standard spectral curve $\Sigma^{\mathrm{P}}(t)$ and standard spectral sheaf $M^{\mathrm{P}}(t)$ corresponding to $(\mathcal{E}(t), \theta(t))$. The same then holds for the intermediate spectral data as well as for the spectral data $(\widehat{\Sigma}^{\widehat{\mathrm{P}}}(t), \widehat{M}^{\widehat{\mathrm{P}}}(t))$ because they arise from the standard data by proper transform along divisors which do not depend on t. Since the spectral curve changes algebraically and the fibers are fixed, the intersection scheme of the spectral curve with the fibers of $\widehat{\pi}$ also varies algebraically. Because the transformed extension of $\widehat{\mathcal{E}}(t)$ is the result of a Hecke modification at this scheme, it follows that the bundles $\widehat{\mathcal{E}}(t)$ depend algebraically on t as well.

On the other hand, by Theorem 8.5 the map $\widehat{\theta}(t)$ is the direct image by $\widehat{\pi}(t)$ of multiplication by $-\widehat{x}^{\widehat{\mathrm{P}}}$; here $\widehat{\pi}(t)$ is the projection map from $\widehat{\Sigma}^{\widehat{\mathrm{P}}}(t)$ to $\widehat{\mathbb{P}}^1$. Clearly, since $\widehat{\Sigma}^{\widehat{\mathrm{P}}}(t)$ depends algebraically on t, so does $\widehat{\pi}(t)$. Therefore, since $\widehat{\theta}(t)$ is the direct image of multiplication by a coordinate that does not depend on t, with respect to a projection depending algebraically on t, the resulting map depends also algebraically on t. This proves that $\mathcal{N}$ preserves the complex structure J, and is an algebraic map. □

Remark 11.5 Actually, it is also true that the Nahm transform preserves the complex structure I. Indeed, in the proof of Theorem 11.4 we only made use of an algebraic interpretation of the Nahm transform for the Dolbeault complex structure established in earlier sections of this paper. Therefore, once we have an algebraic interpretation of the Nahm transform for the de Rham complex structure, the analogous statement for I can be treated along the same lines as J. Such an interpretation is provided in [11] by the second author, as the Laplace transform of $\mathcal{D}$–modules.

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