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# Simulation of SOL turbulence in tokamak plasmas

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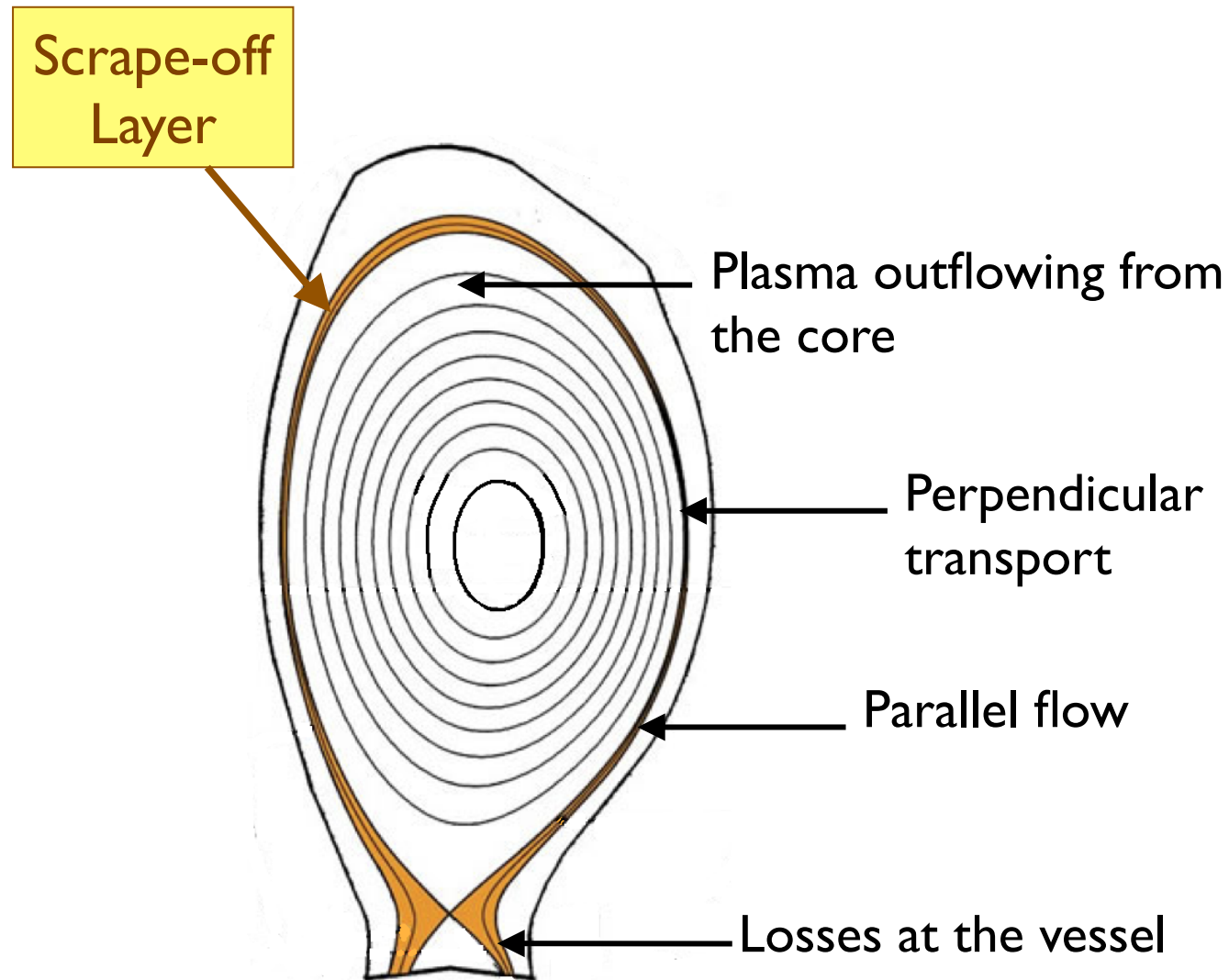
The reduced model to study SOL turbulence

The GBS code and its path towards SOL simulations

Anatomy of SOL turbulence: from linear instabilities to SOL width and  
intrinsic toroidal rotation

# SOL channels particles and heat to the wall

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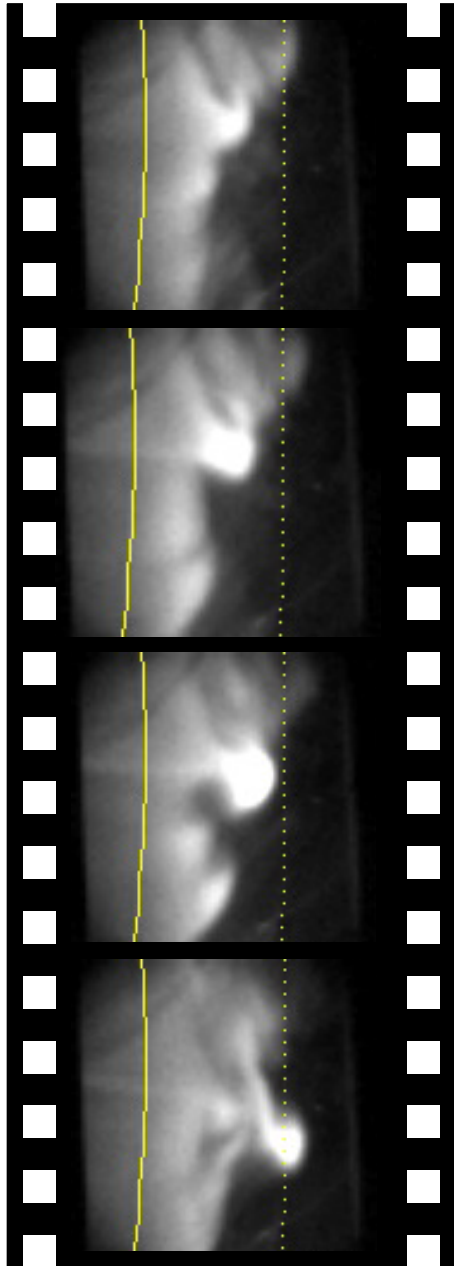
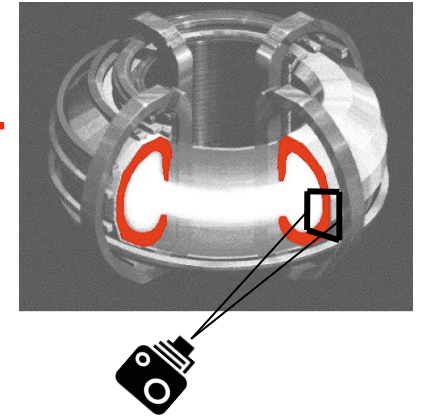
# The key questions

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- What is the mechanism setting the SOL turbulent level and the perpendicular transport?
- How is the SOL width established?
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# Properties of SOL turbulence

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Courtesy of R. Maqueda

- $n_{fluc} \sim n_{eq}$
- $L_{fluc} \sim L_{eq}$
- Fairly cold magnetized plasma

# A reduced model for the SOL

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- Delta-n vs full-n?

- $n_{fluc} \sim n_{eq}$ , need **full-n**

- Local vs global?

- Flux tube valid for  $k_r L_{eq} \gg 1$ , but  $k_r L_{eq} \gtrsim 1$ , need **global**

- Gradient-driven vs flux-driven?

- Evolution equilibrium profile needed, need **flux-driven**

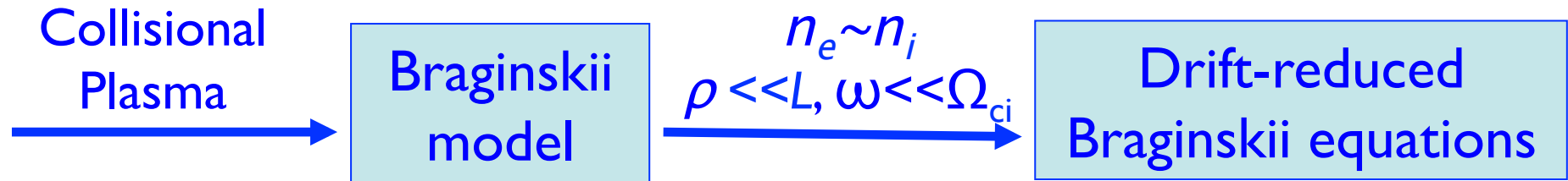
- Kinetic vs fluid?

- $\lambda_{ei} \ll L_{\parallel}$ ,  $\nu^* \gg 1$ , **fluid** is good starting point

- Full v and FLR vs drift-reduced?

- $\omega \ll \omega_{ci}$  and  $k_{\perp} \rho \sim 0.1$ , **drift-reduced** is reasonable

# The GBS code, a tool to simulate SOL turbulence



$$\frac{\partial n}{\partial t} + [\phi, n] = \hat{C}(nT_e) - n\hat{C}(\phi) - \nabla_{\parallel}(nV_{\parallel e}) + S$$

Convection      Magnetic curvature      Parallel dynamics      Source

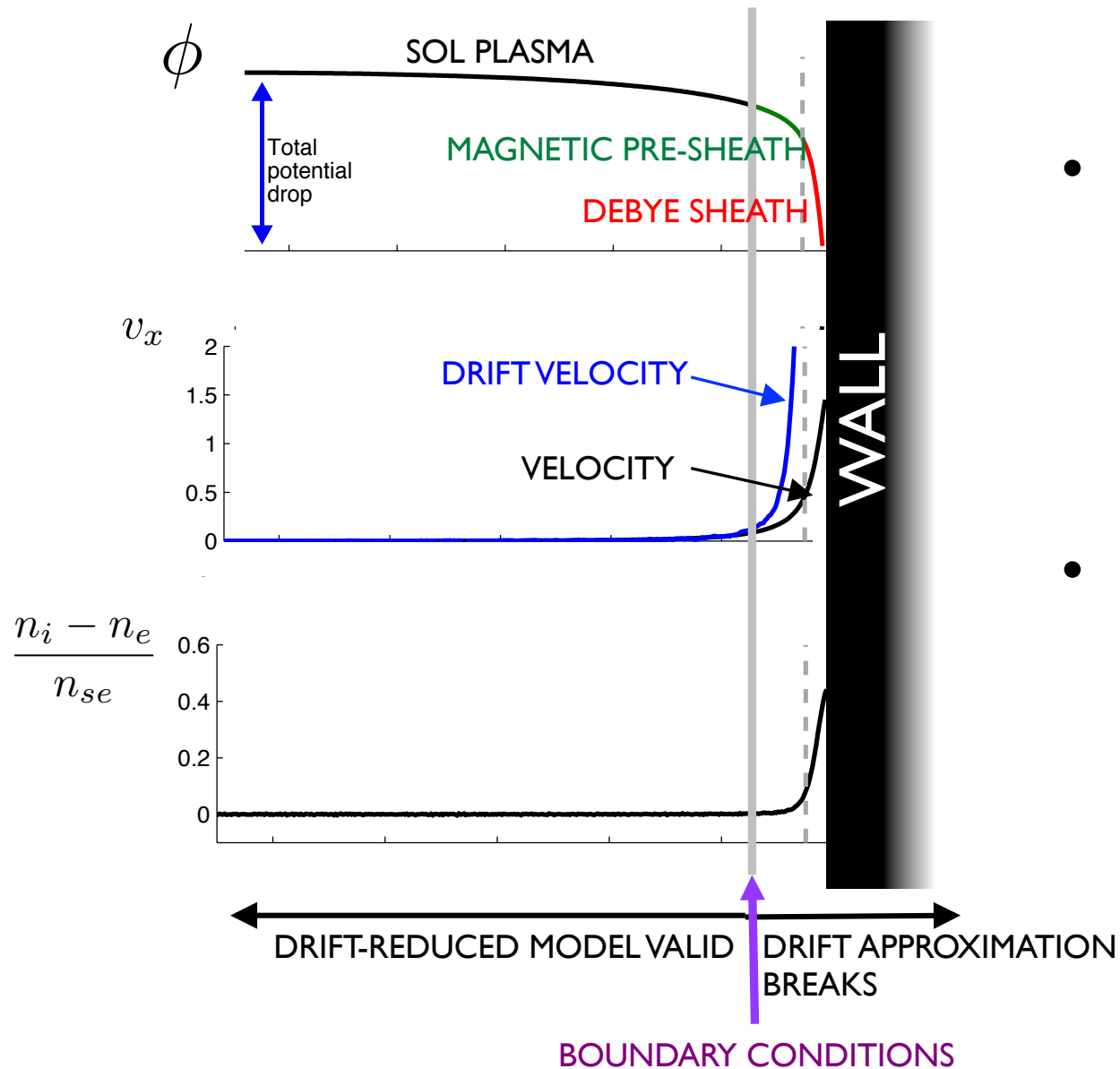
$T_e, \Omega$  (vorticity) → similar equations ( $T_i \ll T_e$ )

$V_{\parallel e}, V_{\parallel i}$  → parallel momentum balance

$$\nabla_{\perp}^2 \phi = \Omega$$

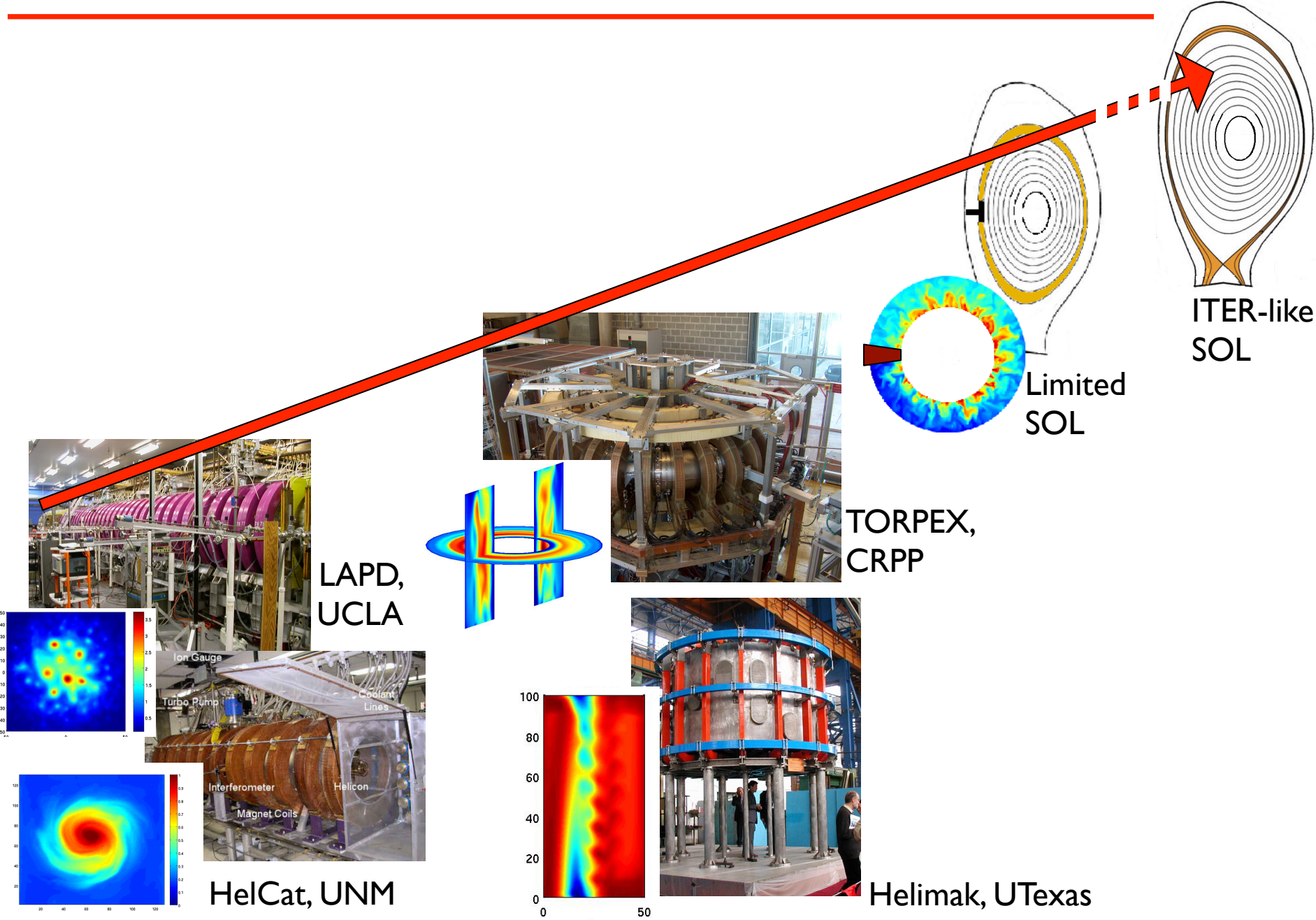
Solved in 3D geometry, taking into account plasma outflow from the core, turbulent transport, and losses at the vessel

# Boundary conditions at the plasma-wall interface



- Set of b.c. for all quantities, generalizing Bohm-Chodura
- Checked agreement with PIC simulations

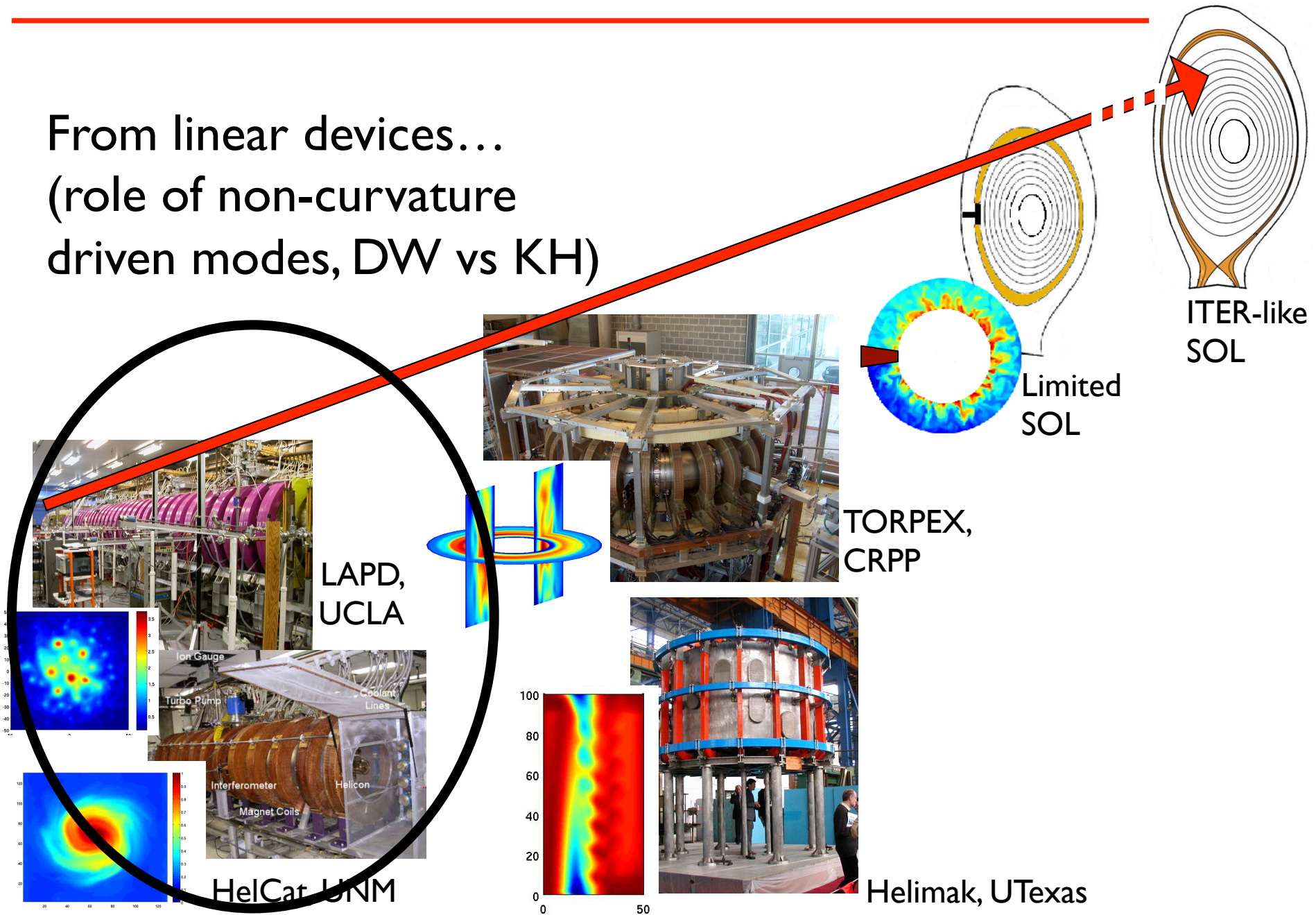
# GBS analysis of configurations of increasing complexity





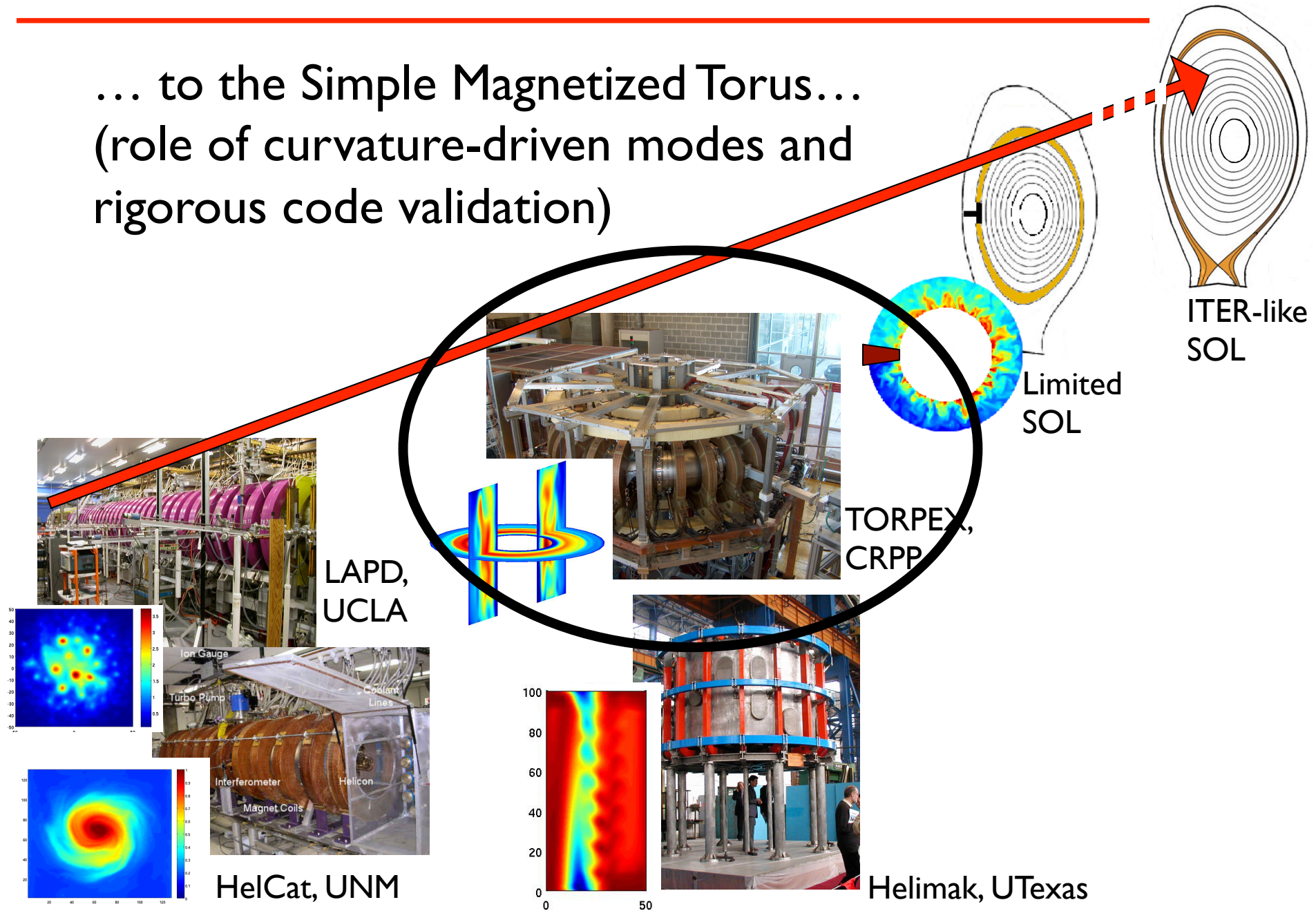
# GBS analysis of configurations of increasing complexity

From linear devices...  
(role of non-curvature  
driven modes, DW vs KH)



# GBS analysis of configurations of increasing complexity

... to the Simple Magnetized Torus...  
(role of curvature-driven modes and rigorous code validation)



ITER-like SOL

Limited SOL

TORPEX, CRPP

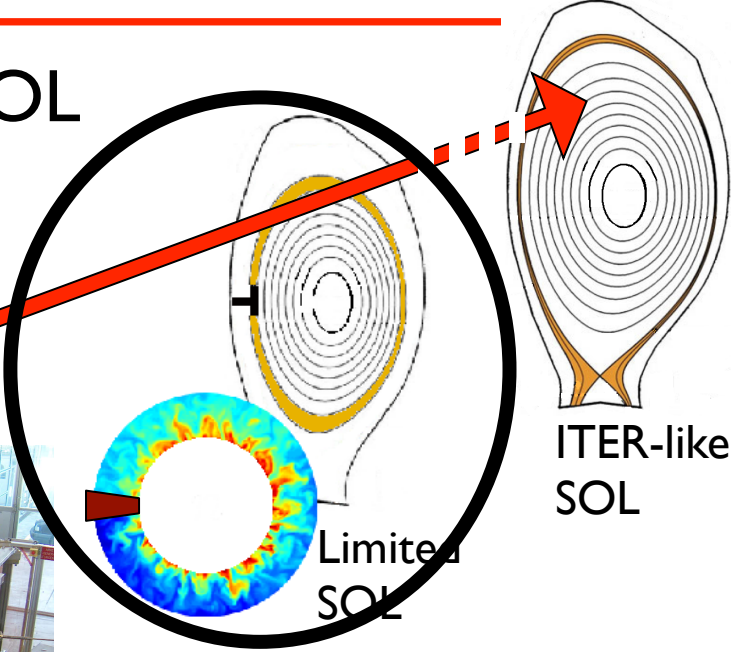
LAPD, UCLA

HelCat, UNM

Helimak, UTexas

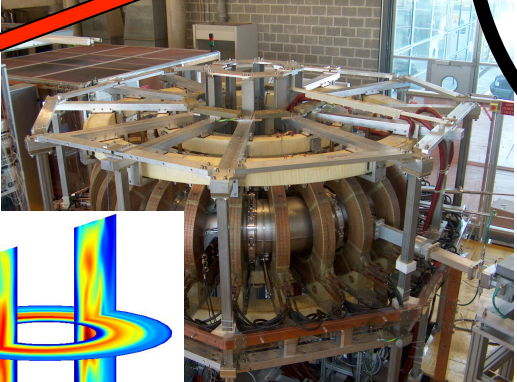
# GBS analysis of configurations of increasing complexity

...to limited SOL

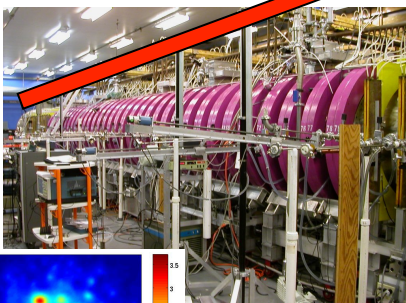
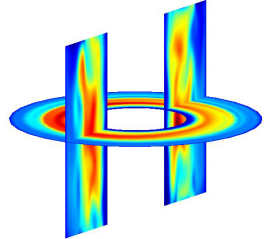


ITER-like SOL

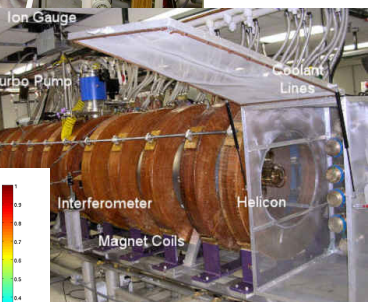
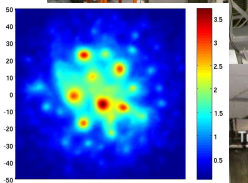
Limited SOL



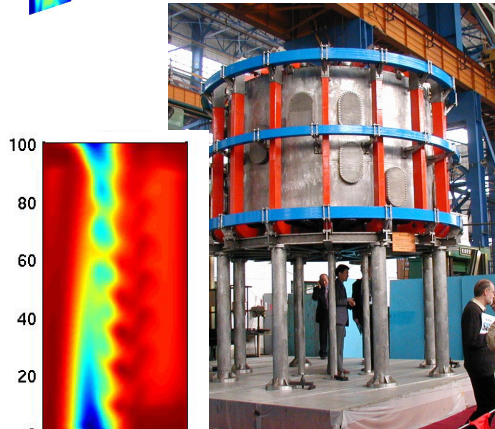
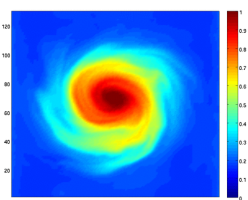
TORPEX, CRPP



LAPD, UCLA



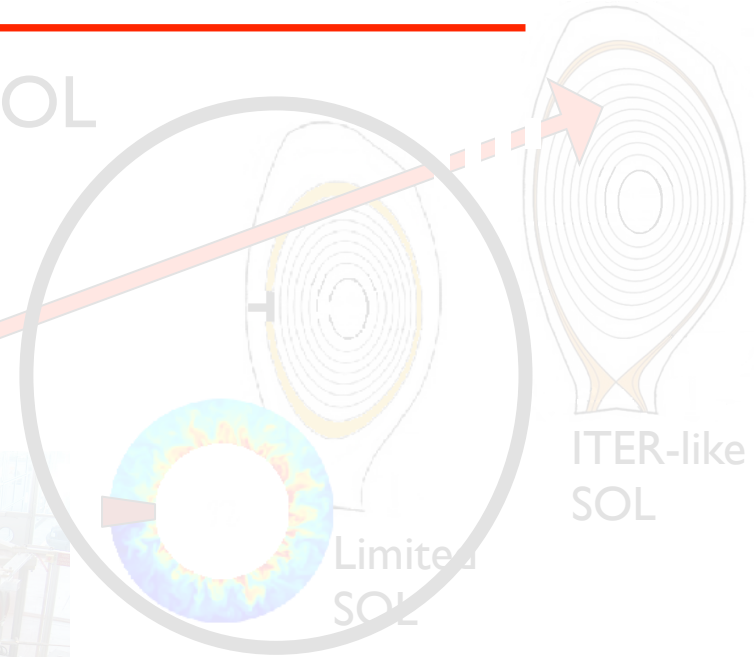
HelCat, UNM



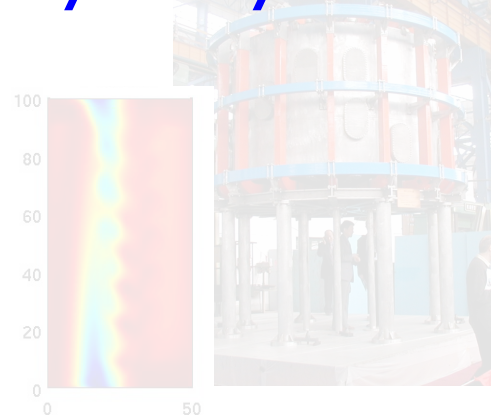
Helimak, UTexas

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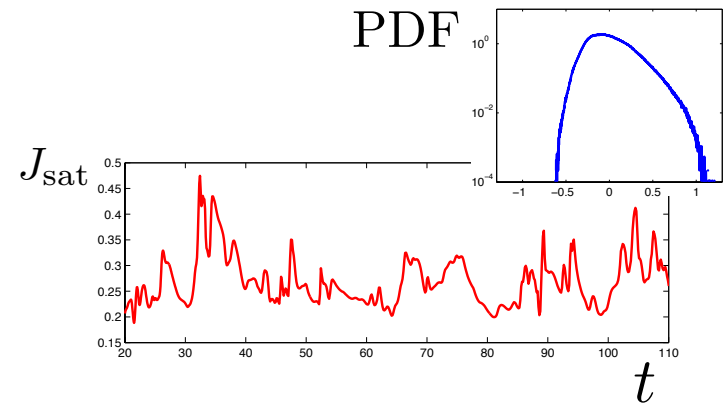
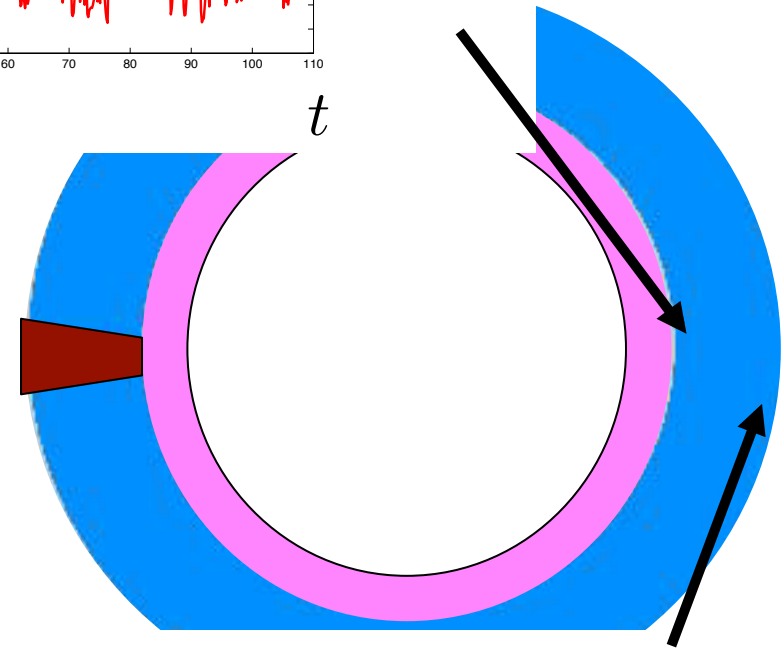
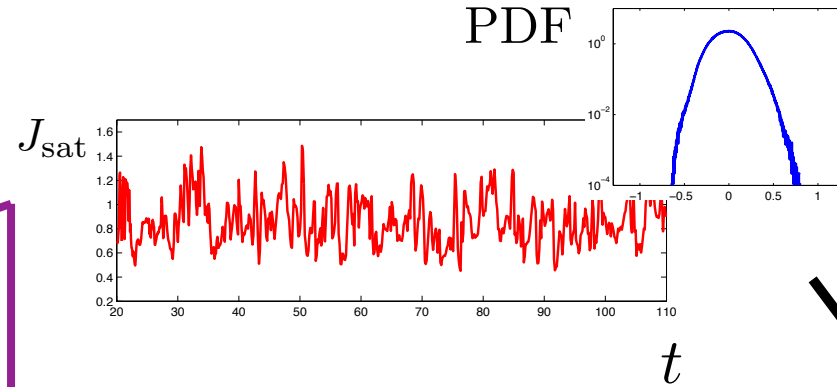
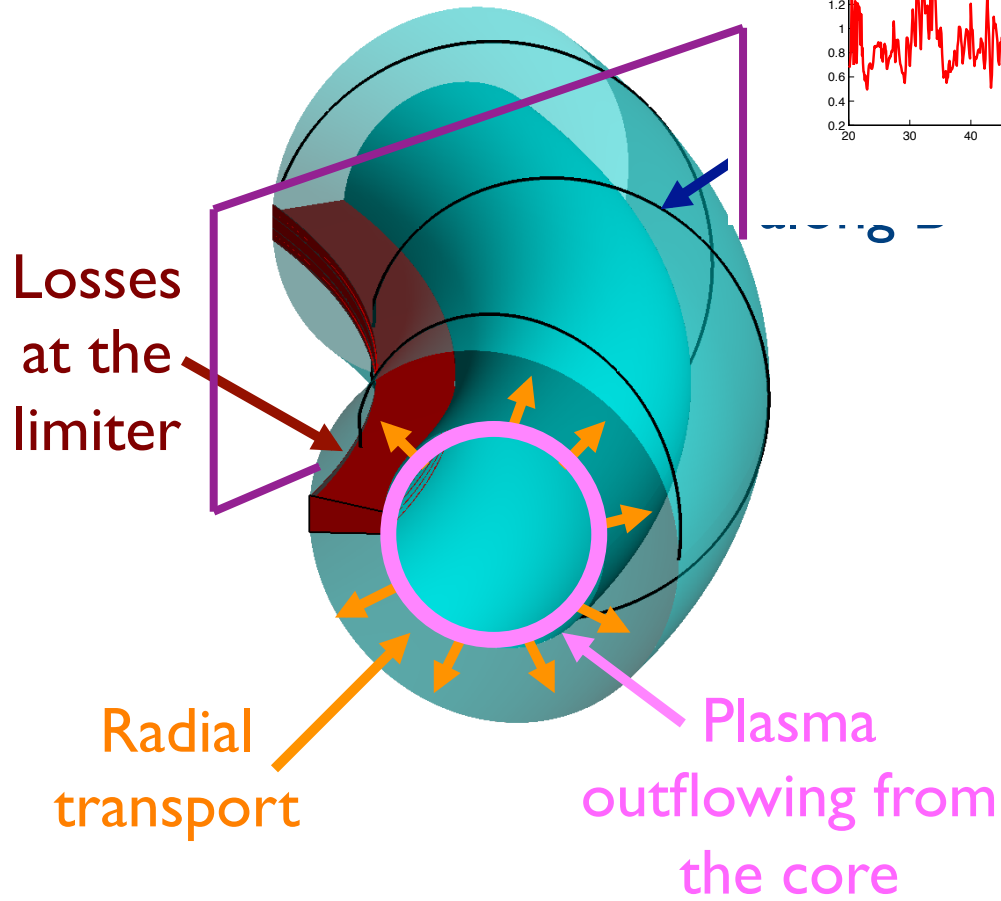
...to limited SOL



... supported by analytical investigations



# Tokamak



S,

Simulations contain physics of ballooning modes, Kelvin-Helmholtz, blobs, parallel transport

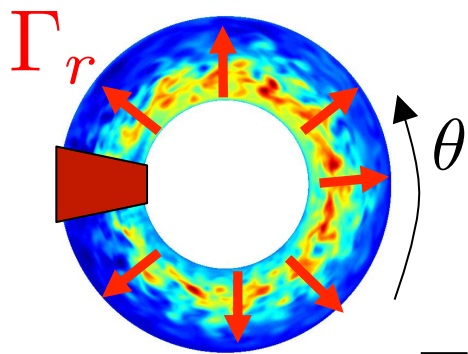
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# Turbulent transport with gradient removal (GR) saturation

Turbulence saturates when it removes its drive  $\rightarrow \frac{\partial p_{e1}}{\partial r} \sim \frac{\partial p_{e0}}{\partial r} \rightarrow k_r p_{e1} \sim p_{e0}/L_p$



$$\frac{\partial p_e}{\partial t} \simeq [p_e, \phi]$$

$$\Gamma_r = \left\langle \left( p_{e1} \frac{\partial \phi_1}{\partial \theta} \right) \right\rangle_t$$

GR hypothesis

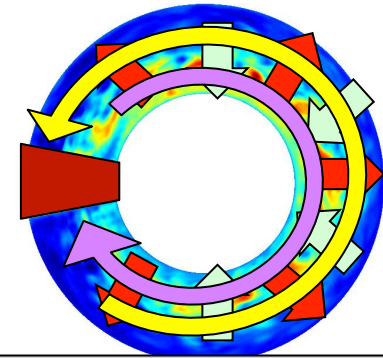
Nonlocal linear theory,  $k_r \sim \sqrt{k_\theta/L_p}$



$$D_{GR} = \frac{\Gamma_r}{p_{e0}/L_p} \sim \frac{\gamma L_p}{k_\theta}$$

# Turbulence saturation due to Kelvin-Helmholtz instability (KH)

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Primary instability grows until it causes KH unstable shear flow

$$\rightarrow \frac{\partial \omega}{\partial t} \sim [\phi, \omega] \rightarrow \phi_1 \sim \frac{\gamma}{k_\theta^2}$$

$$\Gamma_r = \left\langle p_{e1} \frac{\partial \phi_1}{\partial \theta} \right\rangle_t \sim \frac{\gamma p_{e0}}{L_p k_\theta^2}$$



$$D_{KH} \sim \frac{\gamma}{k_\theta^2}$$

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## KH vs GR mechanism:

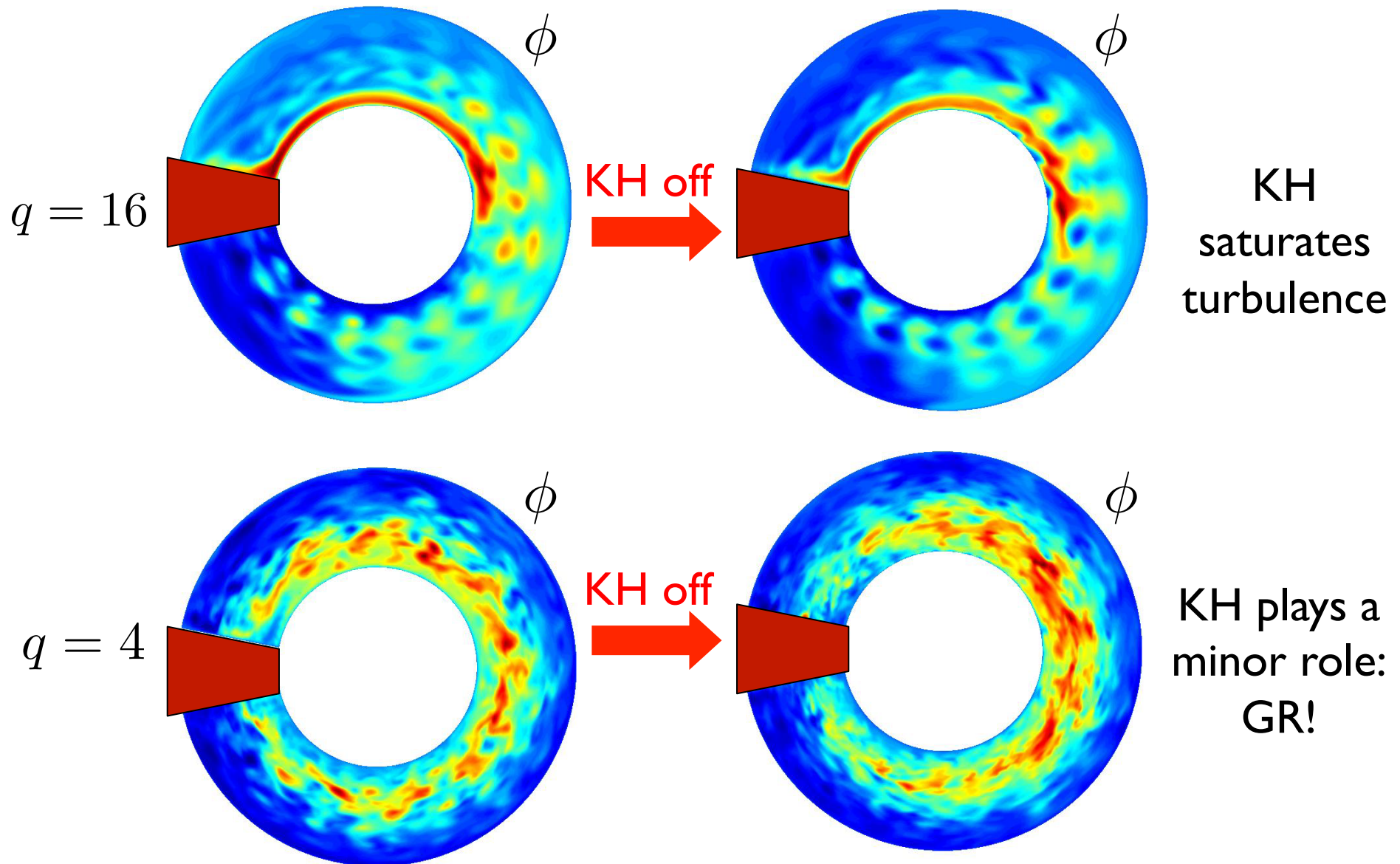
$$\frac{D_{KH}}{D_{GR}} \sim \frac{1}{k_\theta L_p} < 1$$

We expect KH to limit the transport, provided that KH is unstable!



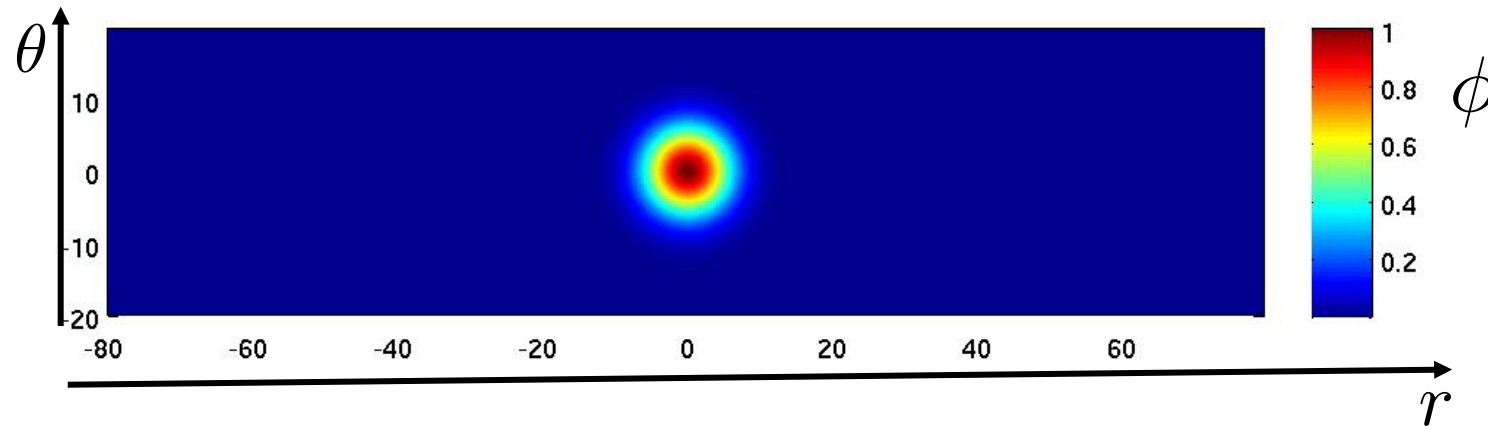
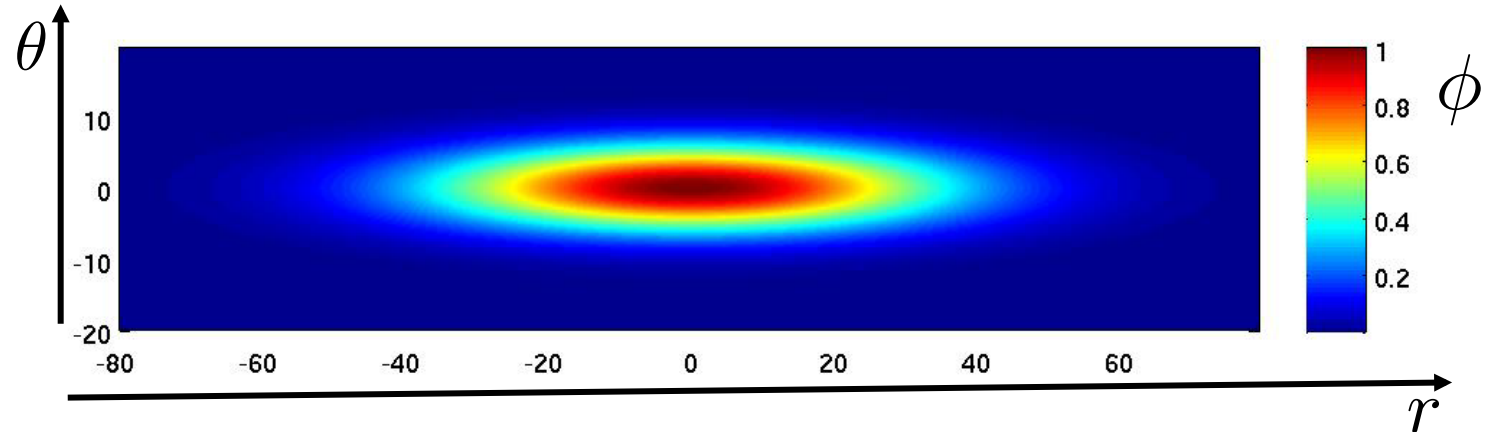
# Is KH really setting transport?

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# Why is KH stable at low $q$ but not higher $q$ ?

Only  
elongated  
eddies  
are KH  
unstable



By comparing eddy turn over time and KH growth rate,  
KH unstable if:  $\sqrt{k_\theta L_p} > 3$  (as in the  $q = 16$  case)

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# Transport and profile scaling for KH stable cases

Balance of perpendicular transport and parallel losses

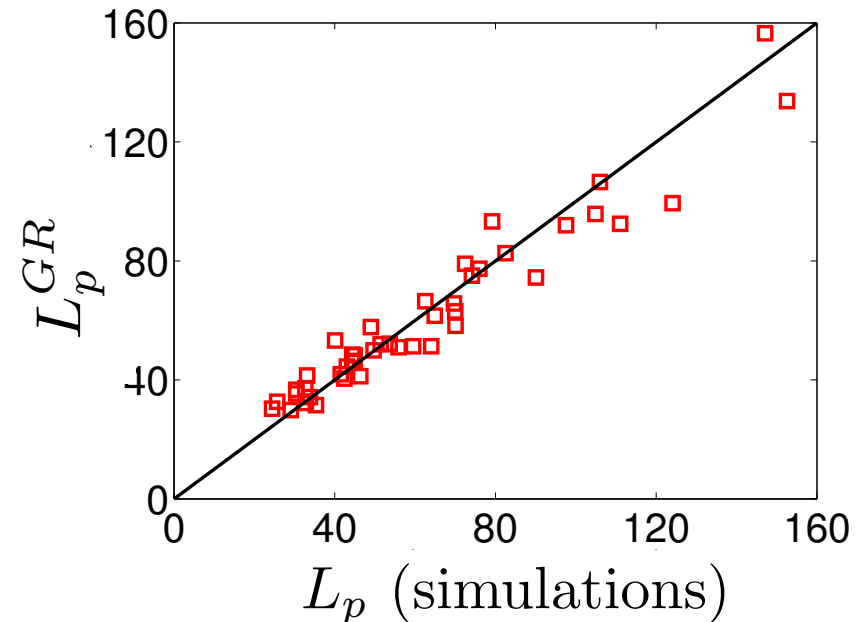
$$\frac{d\Gamma_r}{dr} \sim L_{\parallel} \sim \frac{n_0 c_s}{qR}$$

↑  
Bohm's

→  $L_p^{GR} \simeq q \left( \frac{\gamma}{k_{\theta}} \right)_{\max}$

→  $L_p^{GR} = L_p^{GR}(R, q, \hat{s}, \nu)$

Simulations show expected scaling



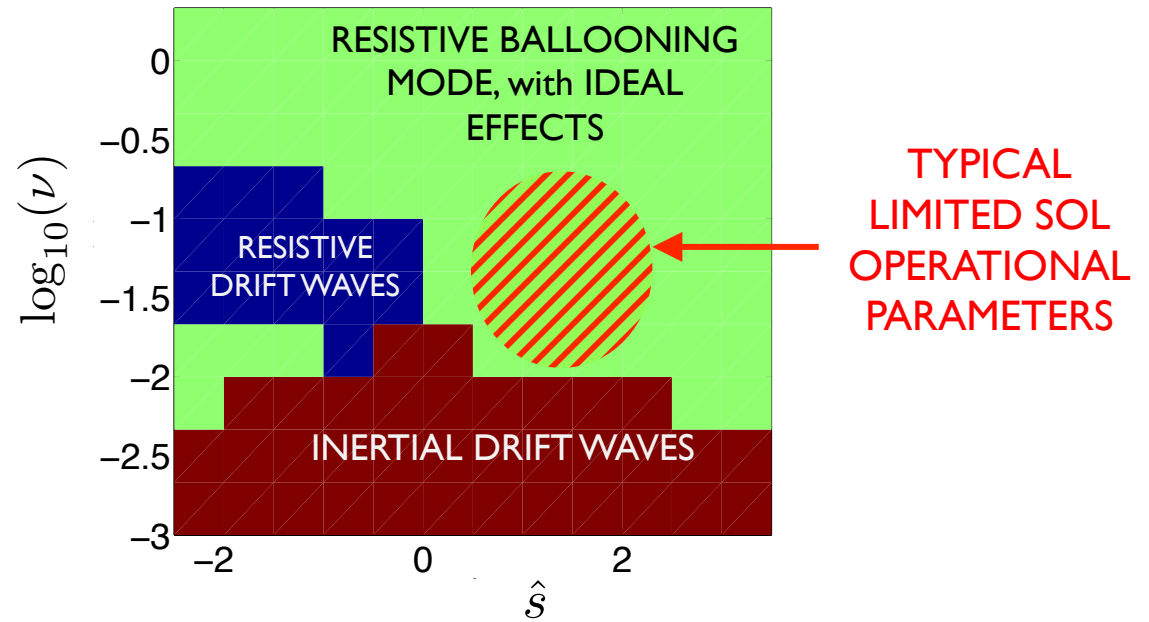
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# SOL Turbulent regimes

Instabilities driving turbulence depends mainly on  $q$ ,  $\nu$ ,  $\hat{s}$ .



LIMITED SOL:

$$L_p^{GR} \simeq q \begin{pmatrix} \gamma \\ k_{\theta} \end{pmatrix}$$

$\gamma \sim \gamma_b = \sqrt{2R/L_p}$  (RBM)  
 $k_{\theta} \sim k_b = \sqrt{\frac{1 - \alpha_{MHD}}{\nu q^2 \gamma_b}}$  (RBM<sub>max</sub>)

$$L_p = R^{1/2} [2\pi (1 - \alpha_{MHD}) \alpha_d / q]$$

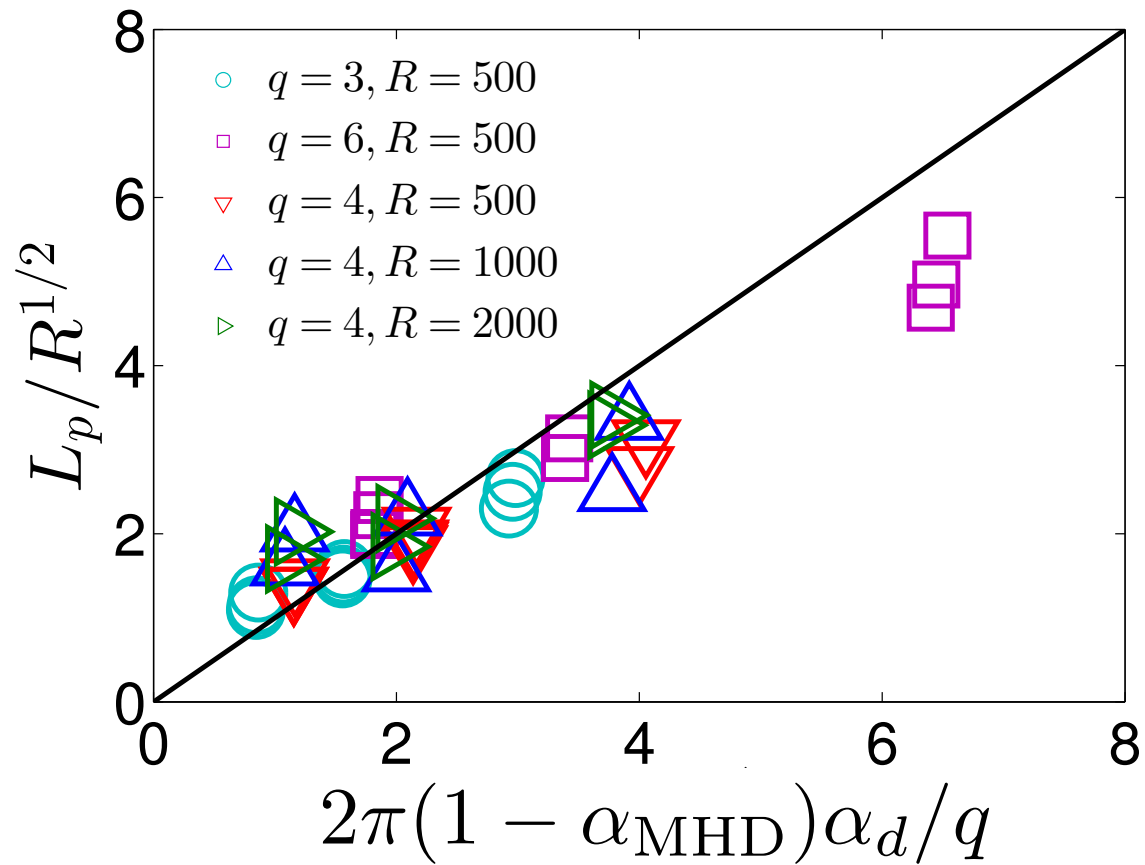
MAJOR RADIUS

$$\alpha_{MHD} \sim q^2 \beta R / L_p$$

$$\alpha_d \sim (R/L_p)^{1/4} \nu^{-1/2} / q$$

# Simulations agree with ballooning estimates

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# The key questions

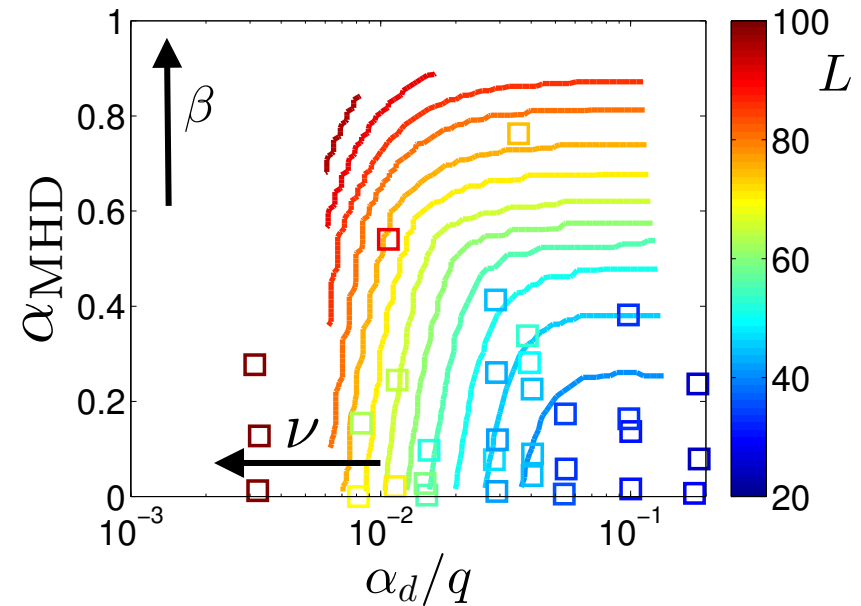
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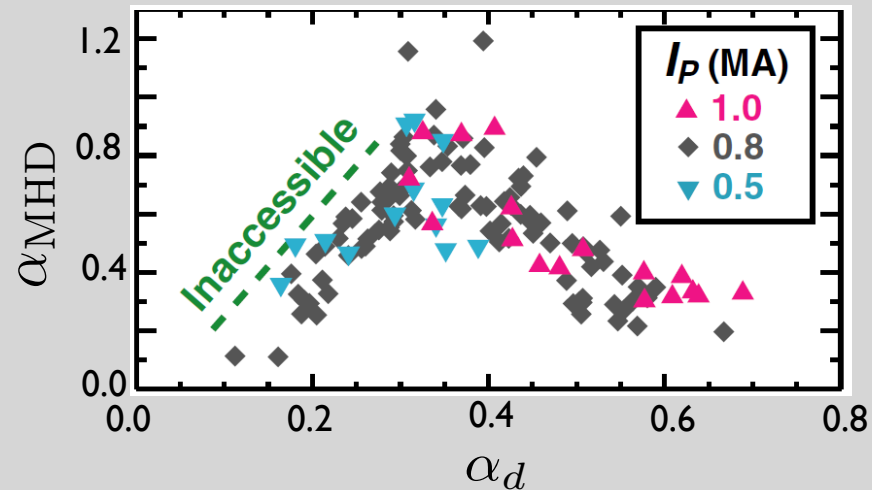
# Limited SOL transport increases with $\beta$ and $\nu$

$$L_p = R^{1/2} [2\pi(1 - \alpha_{\text{MHD}})\alpha_d/q]$$



Maybe related to the density limit?

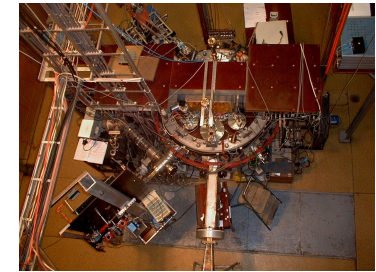
Coupling with core physics needs be addressed...



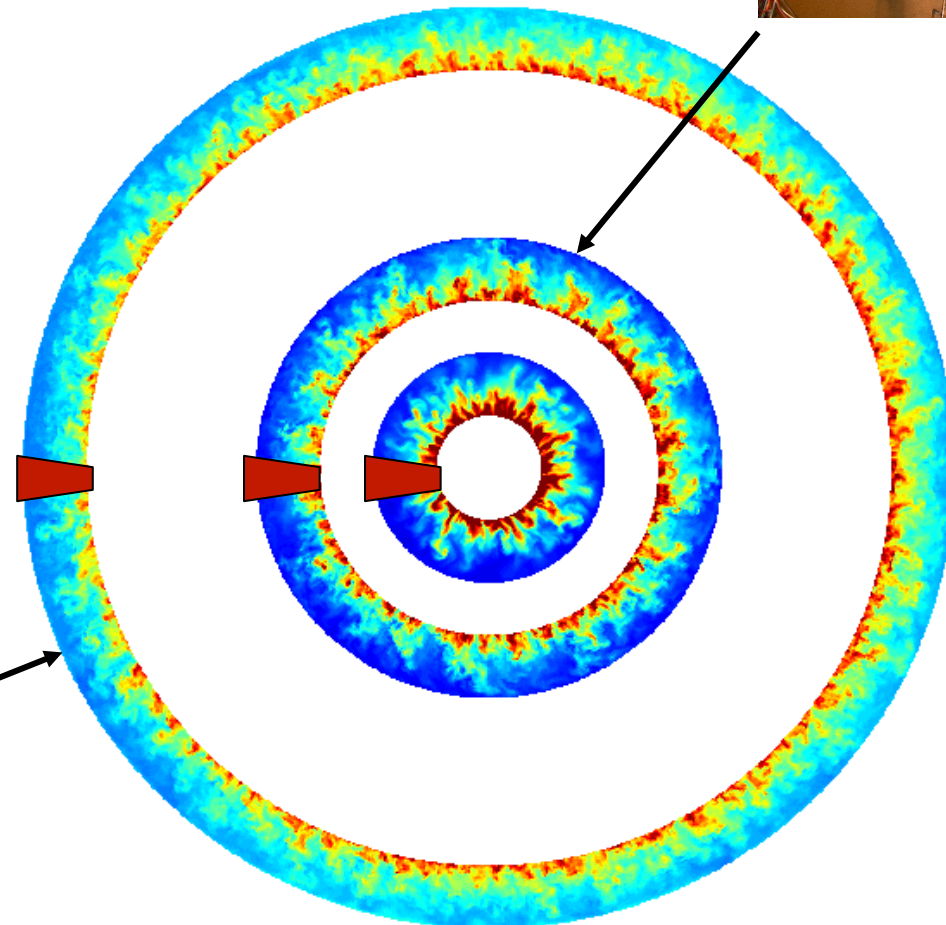
LaBombard, NF 2005

# Limited SOL width widens with $R$

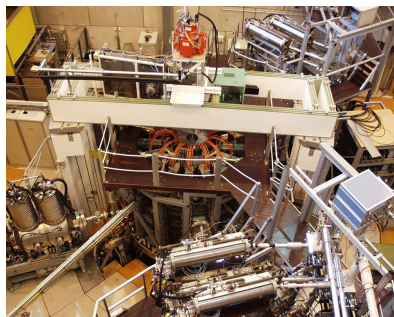
$$L_p = R^{1/2} [2\pi(1 - \alpha_{\text{MHD}})\alpha_d/q]$$



CASTOR



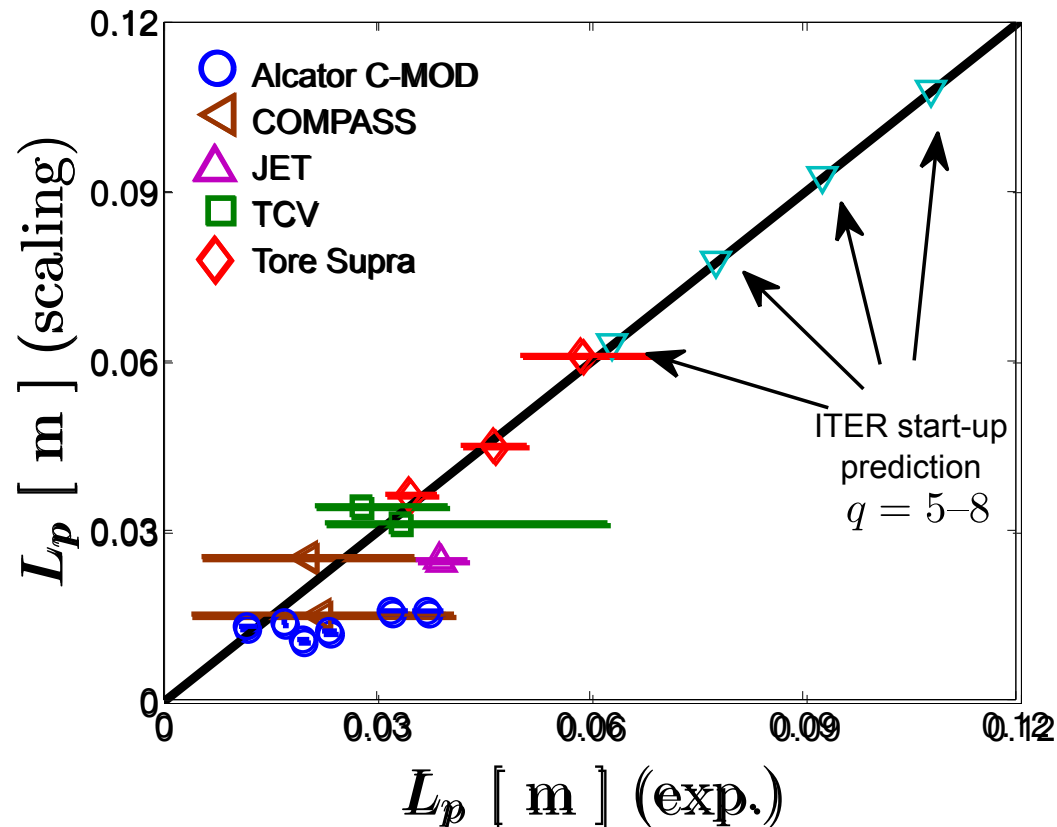
TCV



# Good agreement with multi-machine measurements

The ballooning scaling, in SI units:

$$L_p \simeq 7.97 \times 10^{-8} q^{8/7} R^{5/7} B^{-4/7} T_e^{-2/7} n_e^{2/7}$$



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# Potential in the SOL set by sheath and electron adiabaticity

Typical estimate: at the sheath

$$v_{\parallel i} = c_s \quad v_{\parallel e} = c_s \exp(\Lambda - e\phi/T_e^{\text{sh}})$$

to have ambipolar flows,  $v_{\parallel i} = v_{\parallel e}$

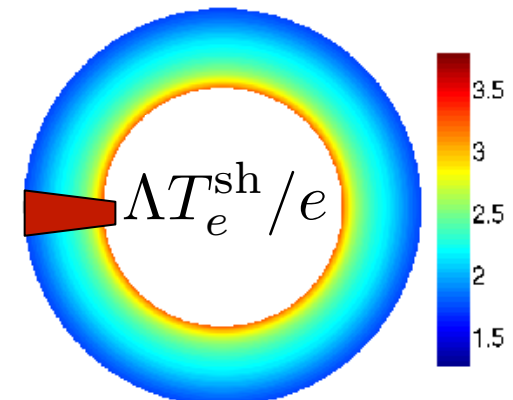
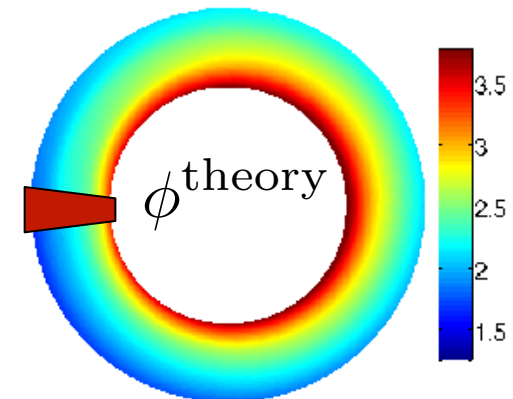
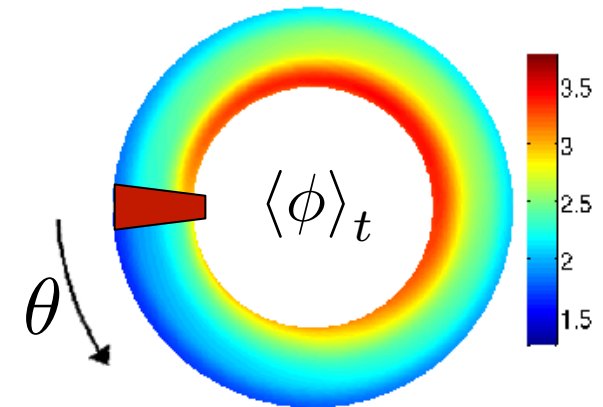
$$\phi = \Lambda T_e^{\text{sh}}/e \simeq 3T_e^{\text{sh}}/e$$

Our more rigorous treatment, from Ohm's law

$$\phi = \underbrace{\Lambda T_e^{\text{sh}}/e}_{\text{Sheath}} + \underbrace{2.71(T_e - T_e^{\text{sh}})/e}_{\text{Adiabaticity}}$$

Sheath

Adiabaticity



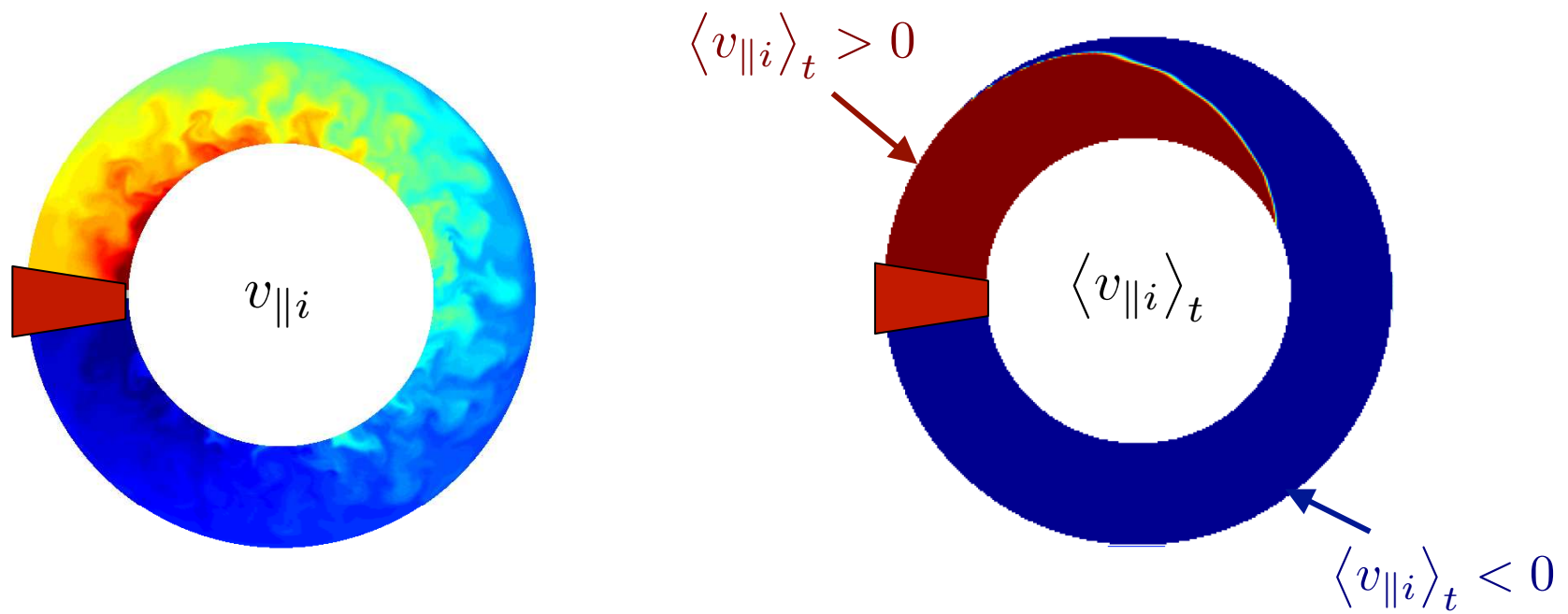
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# GBS simulations show intrinsic toroidal rotation

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# A model for the SOL intrinsic toroidal rotation

Time-averaging the momentum equation:

$$\frac{\partial}{\partial r} D \frac{\partial v_{\parallel i}}{\partial r} + \frac{\partial \phi}{\partial r} \frac{\partial v_{\parallel i}}{\partial \theta} + \epsilon \frac{v_{\parallel i}}{q} \frac{\partial v_{\parallel i}}{\partial \theta} + \frac{\epsilon}{nq} \frac{\partial p}{\partial \theta} = 0$$

Turbulent driven  
radial transport,  
gradient-removal  
estimate

Poloidal  
convection

Parallel  
convection

Pressure poloidal  
asymmetry

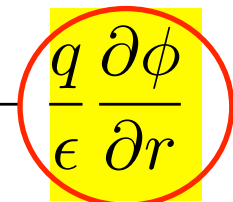
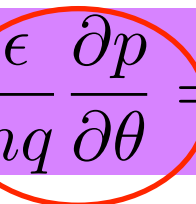
solved with boundary conditions:

$$v_{\parallel i} \Big|_{se} = c_s - \frac{q}{\epsilon} \frac{\partial \phi}{\partial r}$$

Bohm's  
criterion

E<sub>x</sub>B  
correction

Sources of toroidal  
rotation

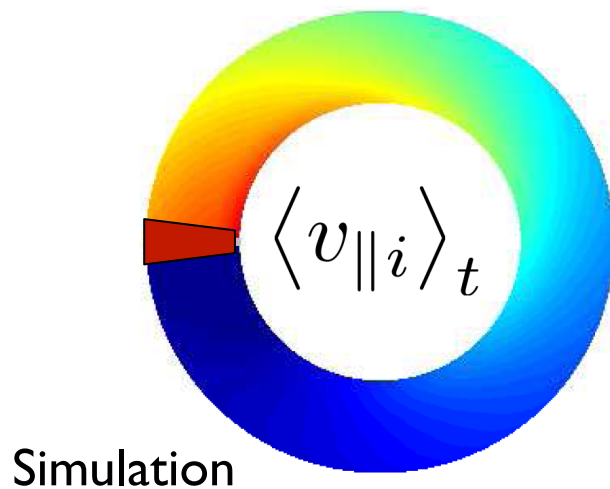
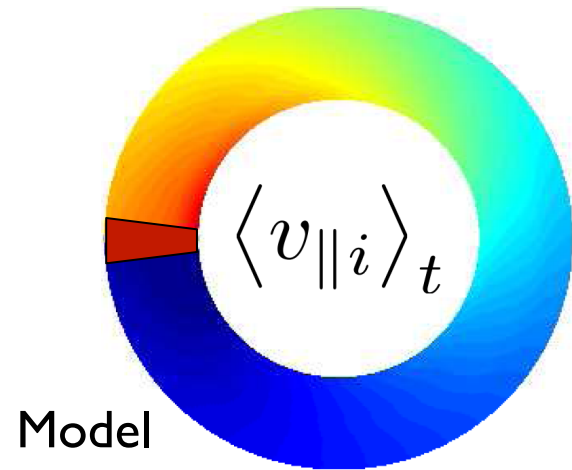




# Our model explains experimental and simulation rotation

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Good agreement between model and simulations:



Able to explain the experimental trends:

- $M_{||} \lesssim 1$
- Typically co-current
- Can become counter-current by reversing  $\mathbf{B}$  or divertor position

Incidentally, a Rice Scaling is observed,

$$v_{\varphi} \sim T_e / I_p$$

# What are we learning from GBS simulations?

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- The use of a progressive simulation approach to investigate plasma turbulence, supported by analytical investigations
- SOL turbulence:
  - Saturation mechanism given by gradient removal or Kelvin-Helmholtz instability
  - Turbulent regimes: in limited plasmas, resistive ballooning modes
  - Good agreement of the scaling of the pressure scale length with multi-machine measurements
  - Sheath dynamics and electron adiabaticity set the electrostatic potential in the SOL
  - Toroidal rotation generated by sheath dynamics and pressure poloidal asymmetry