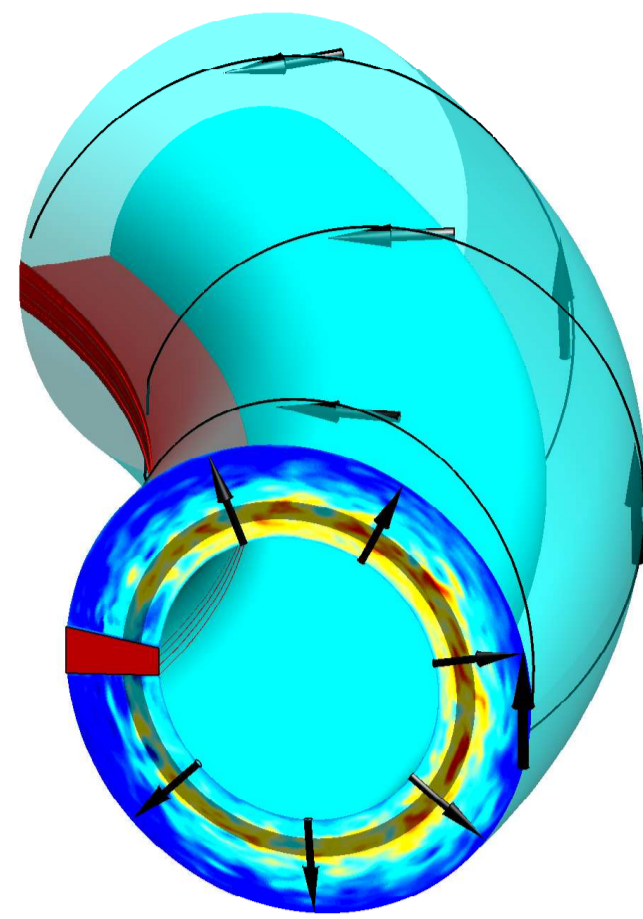
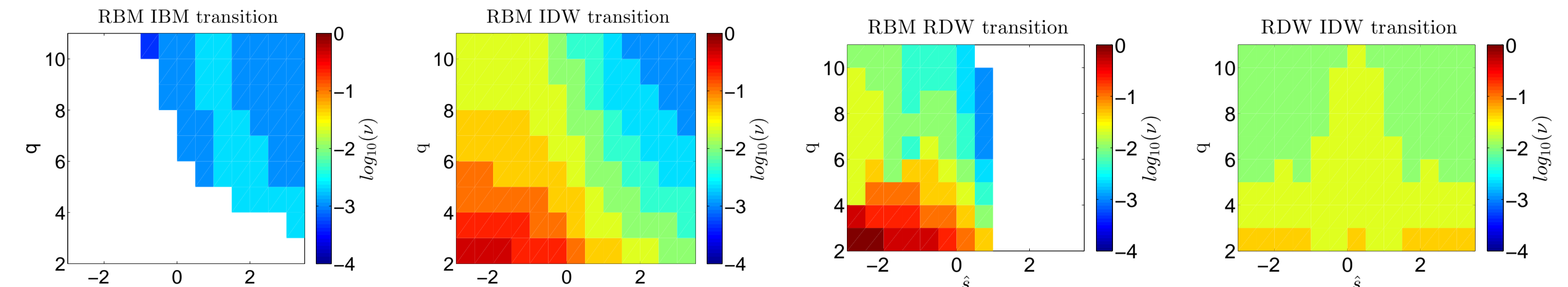


## 1- Introduction

- We study Scrape-Off Layer (SOL) turbulence through simulations that evolve the plasma dynamics as the interplay of plasma source from the core, perpendicular transport, and losses at the limiter plates
- We identify the SOL turbulence regimes, defining the regions of existence of the Ballooning Modes [resistive (RBM) and inertial (IBM)] and the Drift Waves [resistive (RDW) and inertial (IDW)] instabilities, focusing on the role of magnetic shear

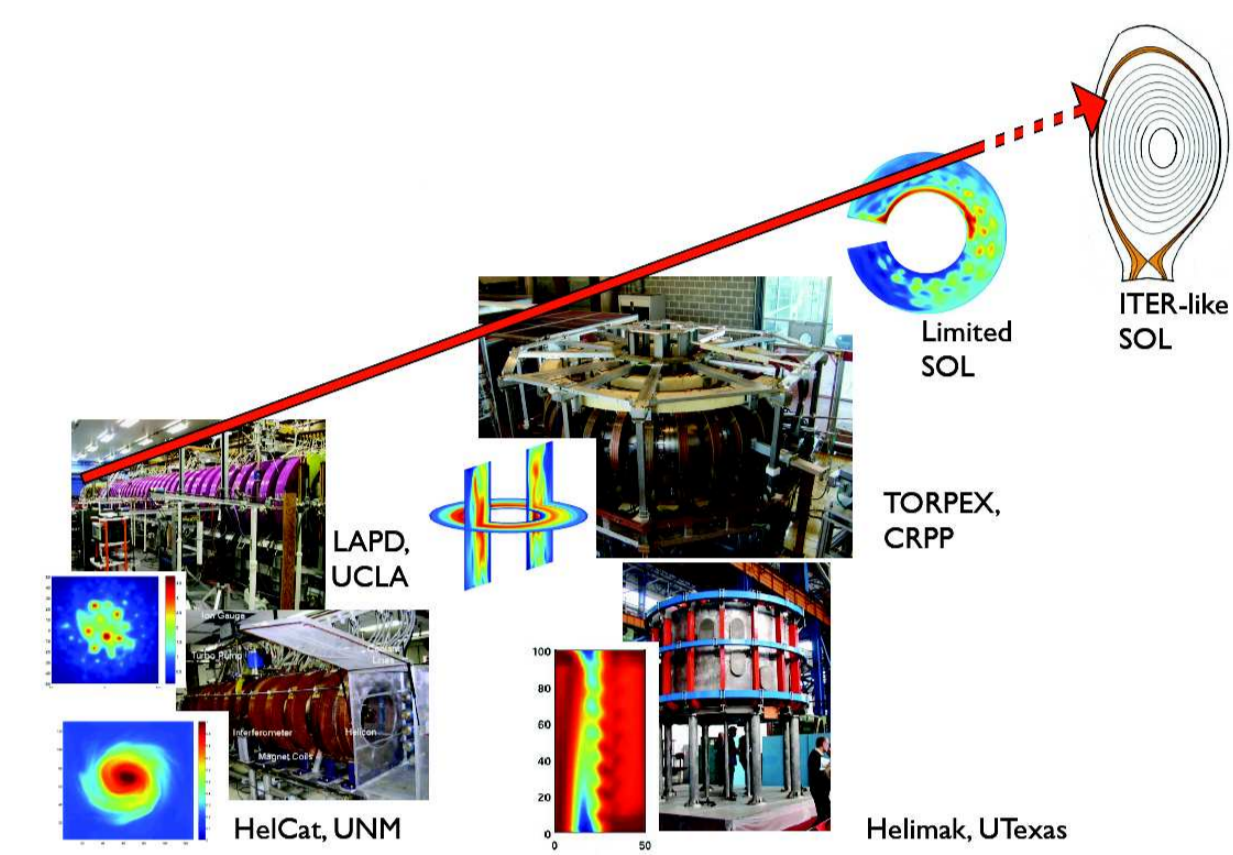


### Transitions:



- For each graph: value of  $\nu$  at which the transition between the first and the second instability takes place (white region: the first instability prevails on the second one for any value of  $\nu$ )
- The transition between IBM and IDW is independent of  $\nu$ ; for  $\hat{s} \lesssim 1$  IDW prevails over IBM; for  $\hat{s} \gtrsim 1$  IBM prevails over IDW for  $q > -3/2\hat{s} + 23/2$

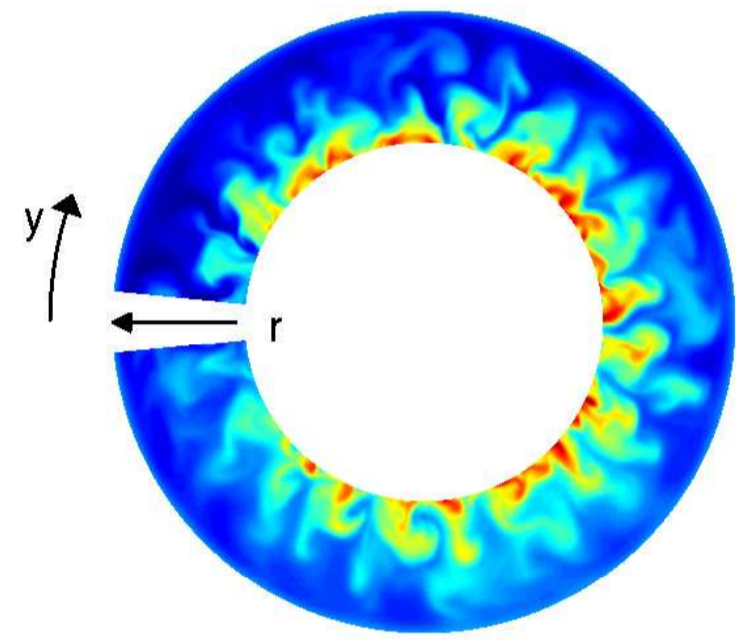
## 2- The Global Braginskii Solver (GBS) code



- The code is based on the non-linear, drift-reduced two-fluid Braginskii equations ([1],[2] and [3])
- Self-consistent global evolution of equilibrium and fluctuations
- The understanding of SOL plasma turbulence has been approached by studying systems of increasing complexity

Topics currently under investigation:

- turbulent saturation mechanism
- identification of the main instabilities
- magnetic shear
- size scaling
- intrinsic rotation (J. Loizu talk I3.409 on Wednesday)
- toroidicity effects (finite aspect ratio, Shafranov shift, ...)
- impurity transport



Continuity:  $\frac{\partial n}{\partial t} = \frac{c}{B} [\Phi, n] + \frac{c}{eRB} (\hat{C}p_e - n\hat{C}\Phi) - \nabla_{\parallel} (nV_{\parallel e})$

Vorticity:  $\frac{\partial \nabla_{\perp}^2 \Phi}{\partial t} = \frac{c}{B} [\Phi, \nabla_{\perp}^2 \Phi] + \frac{B}{m_i c n R} \hat{C}p_e - V_{\parallel i} \nabla_{\parallel} \nabla_{\perp}^2 \Phi + \frac{m_i \Omega_e^2}{e^2 n} \nabla_{\parallel} |j_{\parallel}|$

Ohm's law:  $m_e n \frac{\partial V_{\parallel e}}{\partial t} + \frac{en}{c} \frac{\partial \psi}{\partial t} = m_e n \frac{c}{B} [\Phi, V_{\parallel e}] - m_e n V_{\parallel e} \nabla_{\parallel} V_{\parallel e} - T_e \nabla_{\parallel} n + en \nabla_{\parallel} \Phi - 1.71 n \nabla_{\parallel} T_e + \frac{\nu m_i}{e} j_{\parallel}$

Parallel ion velocity:  $\frac{\partial V_{\parallel i}}{\partial t} = \frac{c}{B} [\Phi, V_{\parallel i}] - V_{\parallel i} \nabla_{\parallel} V_{\parallel i} - \frac{1}{nm_i} \nabla_{\parallel} p_e$

Electron temperature:  $\frac{\partial T_e}{\partial t} = \frac{c}{B} [\Phi, T_e] + \frac{2c}{3eRB} \left( \frac{7}{2} T_e \hat{C}T_e + \frac{T_e^2}{n} \hat{C}n - T_e \hat{C}\Phi \right) + \frac{2T_e}{3en} 0.71 \nabla_{\parallel} |j_{\parallel}| - \frac{2}{3} T_e \nabla_{\parallel} V_{\parallel e} - V_{\parallel e} \nabla_{\parallel} T_e$

Parallel gradient:  $\nabla_{\parallel} = \frac{\partial}{\partial z} + \frac{B}{B^2} \times \nabla \psi \cdot \nabla$

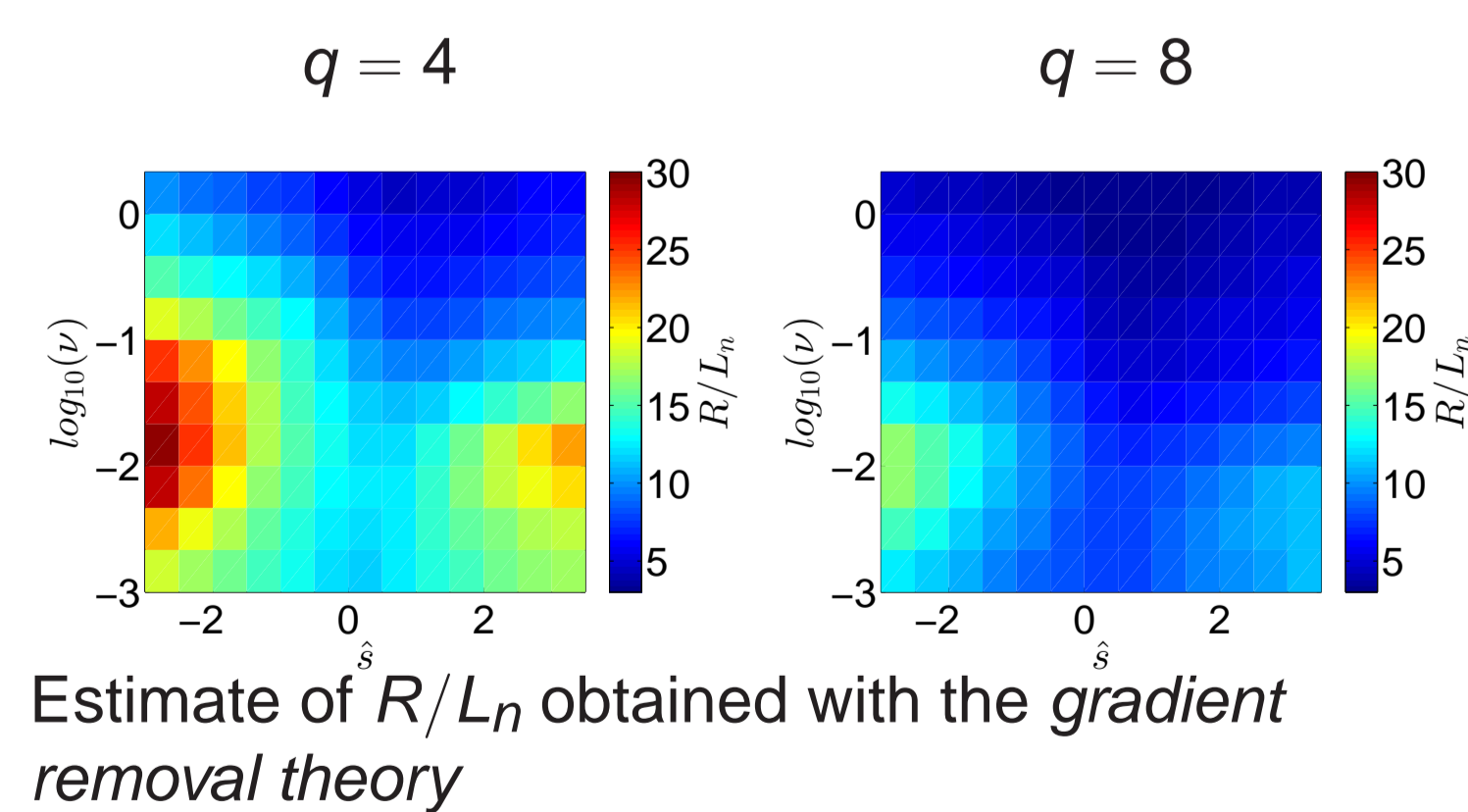
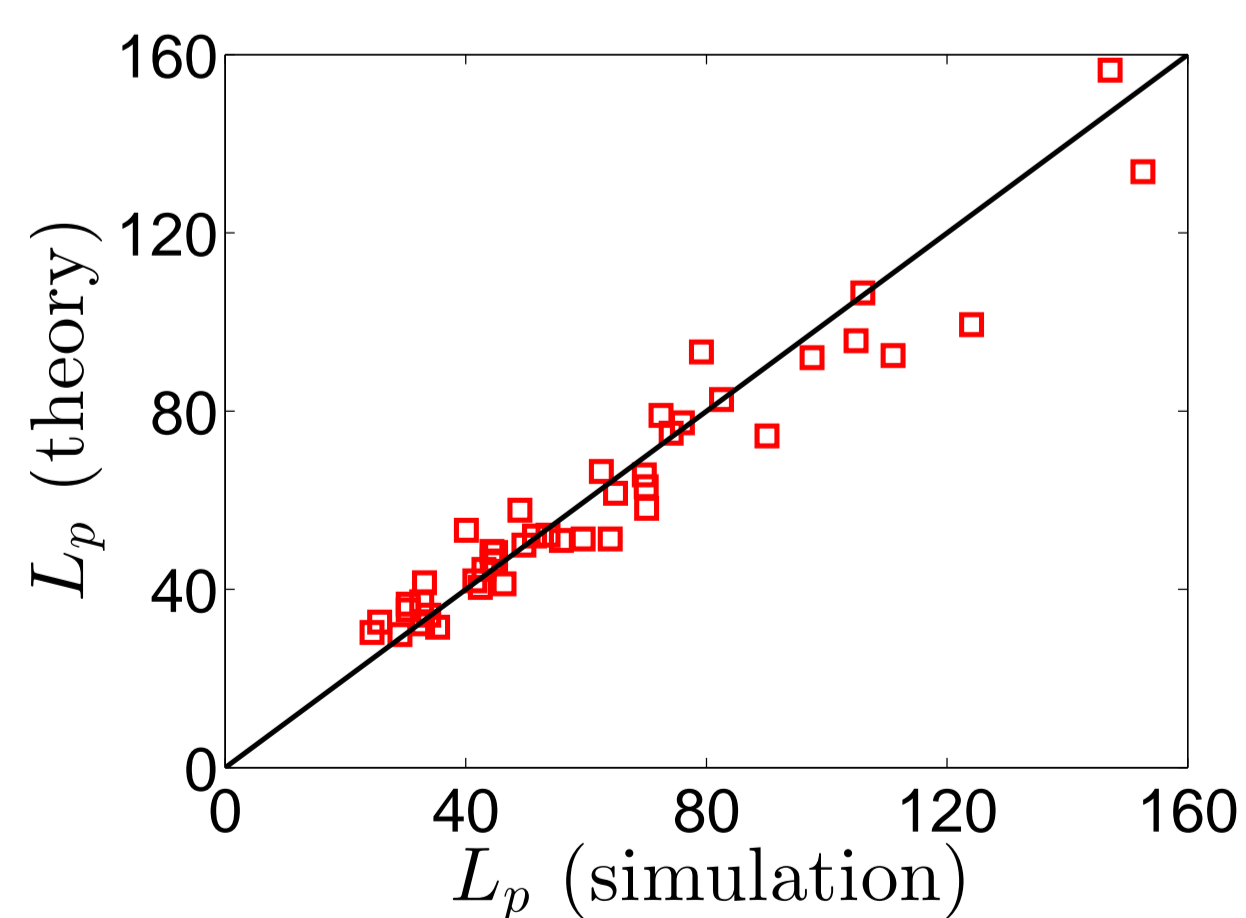
Curvature operator:  $\hat{C} = -2 \left[ \sin \theta \frac{\partial}{\partial x} + \left( \sin \theta \frac{y \hat{s}}{a} + \cos \theta \right) \frac{\partial}{\partial y} \right]$

Boundary conditions: see Ref. [4]

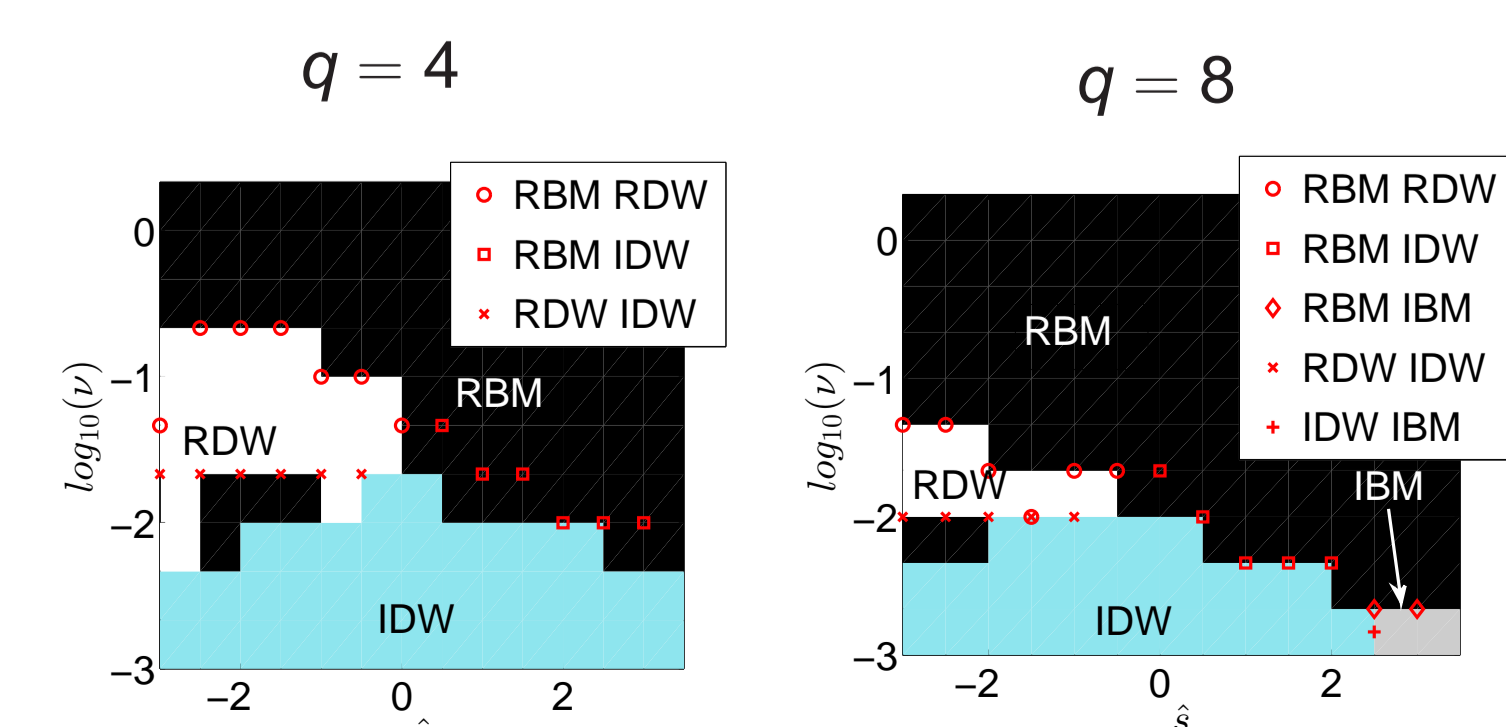
## 3- Gradient removal saturation mechanism

$\partial_r \tilde{n} \sim \partial_r \bar{n}$   
 $\partial_r \Gamma_r \sim \nabla_{\parallel} (\bar{n} V_{\parallel i})$   
 $\Gamma_r = R k_y \bar{\phi} \tilde{n} + \Gamma_r / L_n \sim \bar{n} c_s / q \rightarrow L_n \sim \frac{q}{c_s} \left( \frac{\gamma}{k_y} \right)_{MAX}$   
 $\tilde{\phi} \sim \gamma \tilde{n} L_n / (\bar{n} R k_y)$   
 $\partial_t n = -R [\phi, n]$

- red: saturation occurs when the radial gradient of the perturbed density becomes comparable to the radial gradient of the background density
- blue: potential estimate from the leading order term of the density equation
- magenta: balance between radial particle flux and parallel losses



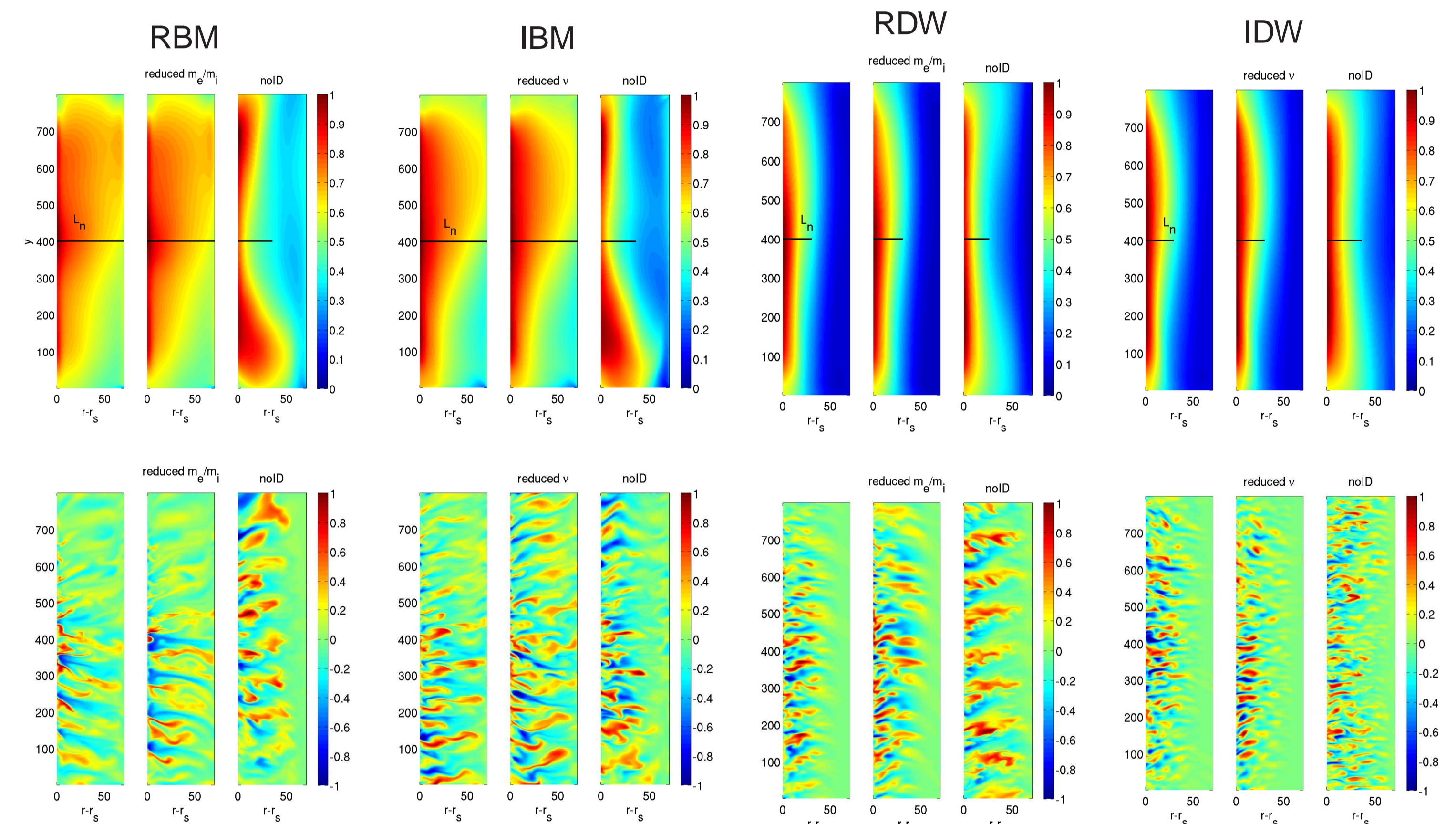
## 4- Identification of the SOL turbulent regimes



- The turbulence regime is identified as the one having the maximum  $\gamma/k_y$  for the  $R/L_n$  predicted by the *gradient removal theory*
- The transitions among different regimes are identified by comparing  $\gamma/k_y$  for each pair of instabilities (red symbols)

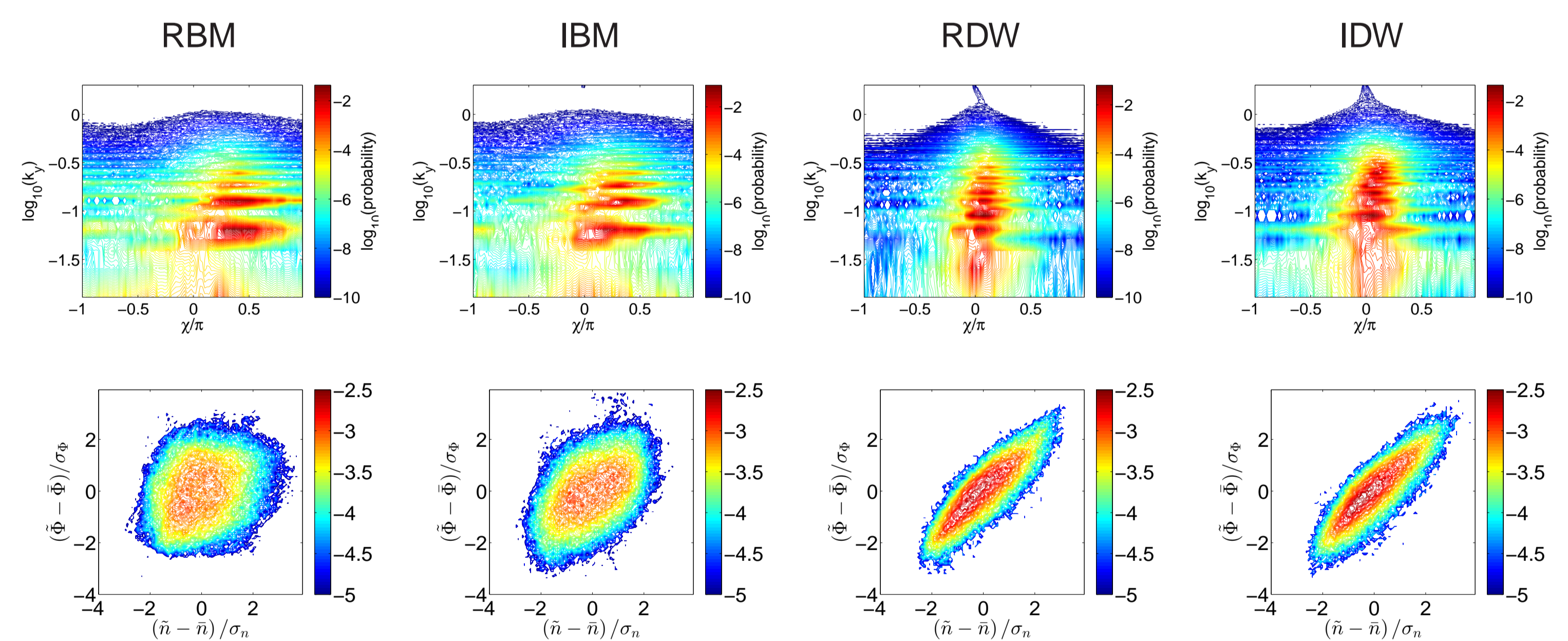
## 5- Verification of the non-linear turbulent regimes with GBS

### Equilibrium and fluctuation profiles of the turbulent regimes



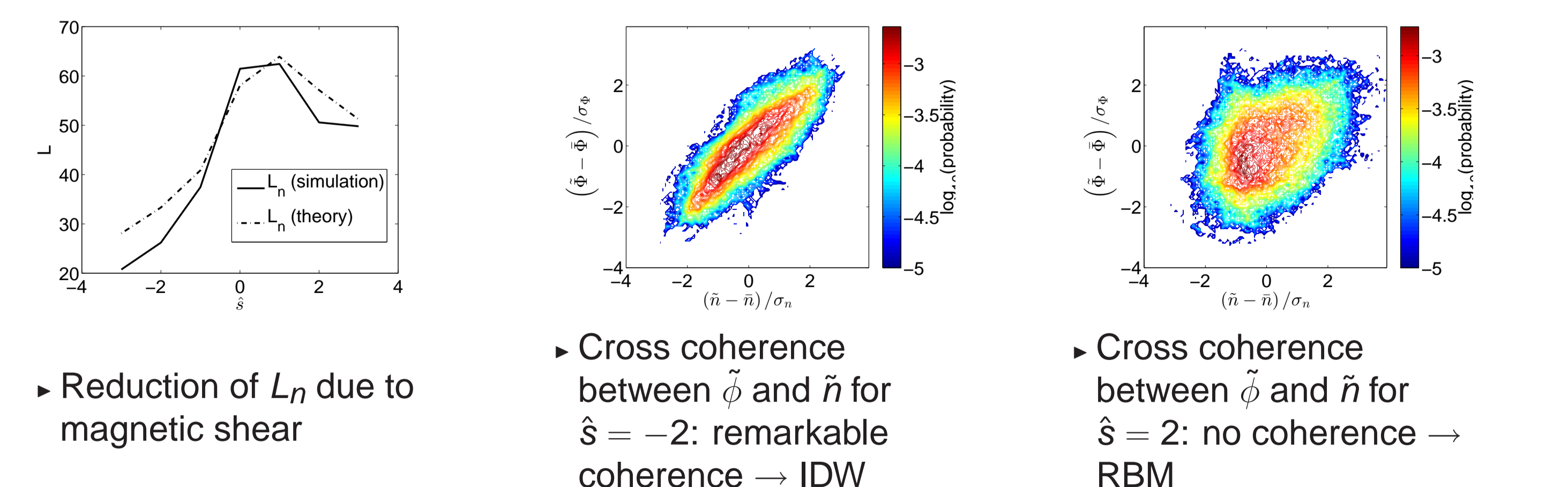
- No changes if  $\nu$  ( $m_e/m_i$ ) is reduced  $\rightarrow$  identification of an inertial (resistive) regime
- Turning off of the Interchange Drive (ID): no effect  $\rightarrow$  DW, significant changes  $\rightarrow$  BM

### Phase shift probability and cross coherence between $\tilde{n}$ and $\tilde{\phi}$



- Phase shift probability: maximum at  $\chi/\pi \approx 0 \rightarrow$  DW, maximum at  $\chi/\pi \approx 0.5 \rightarrow$  BM
- Cross coherence: remarkable coherence  $\rightarrow$  DW, no coherence  $\rightarrow$  BM

## 6- Influence of the magnetic shear



- Reduction of  $L_n$  due to magnetic shear
- Cross coherence between  $\tilde{\phi}$  and  $\tilde{n}$  for  $\hat{s} = -2$ : remarkable coherence  $\rightarrow$  IDW
- Cross coherence between  $\tilde{\phi}$  and  $\tilde{n}$  for  $\hat{s} = 2$ : no coherence  $\rightarrow$  RBM

## 7- Conclusions

- Estimate of the gradient length by means of the *gradient removal theory*
- Identification of the SOL turbulent regimes
- Study of the shear induced steepening of the pressure profile and turbulence regime transition

### References :

- [1] A. Zeiler et al., *Phys. Plasmas*, Vol. 4, Issue 6, 1997
- [2] B. N. Rogers and P. Ricci, *Physical Review Letters*, 104, 225002 (2010)
- [3] P. Ricci and B. N. Rogers, *Physical Review Letters*, 104, 145001 (2010)
- [4] J. Loizu, F. D. Halpern, S. Jolliet and P. Ricci., *Phys. Plasmas*, Vol. 19, Num 12, 2012
- [5] A. Masetto, F. D. Halpern, S. Jolliet, J. Loizu and P. Ricci., *submitted to Phys. Plasmas*