SUSTAINED NANO-MECHANICAL OSCILLATION OF A RESONANT-BODY TRANSISTOR BY FREQUENCY-MODULATED HETERODYNE PHASE-LOCKED-LOOP

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Many applications based on resonant nanoelectromechanical systems (NEMS) require monitoring their natural frequency of oscillation over time with high precision, e.g. for gas sensing or nanomechanical mass spectrometry [1]. In this study, we integrated for the first time a very-high frequency, nanomechanical resonant-body field-effect transistor (RB-FET) into a frequency-modulated phase-locked loop (FM-PLL) which operates analog, requires only one frequency source, and simultaneously exploits the low-noise motion detection based on FM-demodulation with resonant transistors [2]. We demonstrate sustained mechanical oscillation of a 120 MHz doubly-clamped nano-resonator (54 nm thick, 158 nm wide, 2.65 µm long) by using the FM-PLL in vacuum and in air, reaching a frequency stability in the low *ppm*-range at room temperature.

Fig.1 shows a SEM top view of a representative device. The device fabrication was based on a large-scale, silicon-on-insulator (SOI) surface nano-machining process. The initial 37 nm thin device layer was patterned using e-beam lithography. Details on fabrication and co-integration with SOI-CMOS were reported in [3]. The fundamental, flexural mode at $\omega_0/2\pi=121$ MHz was actuated through lateral gate electrodes and 65 nm capacitive air-gaps. We used the RB-FET as efficient FM-demodulator [2] to detect the nanomechanical oscillation on-chip; this technique provides a large signal-to-noise ratio (SNR) and ease of implementation. However, the first-order response does not provide a signal suitable to build a negative feedback loop, due to its symmetry around the resonator centre frequency (see Fig.2). Therefore, most NEMS-PLLs have relied on two-source mixing techniques [4] or digital implementations [5], which required algorithms and computer interface and can limit the measurement bandwidth.

Here, we used a nonlinear effect to create a linear, negative loop-feedback signal. In FM, the bandwidth required to transmit the signal is $\sim 2(\omega_{ref} + \Delta \omega)$, where ω_{ref} is the modulation frequency and $\Delta \omega$ the frequency deviation. When increasing $\Delta \omega$, the FM-bandwidth increased and approached the resonator natural line-width $\omega_0 / Q \approx 120$ kHz. The response became then increasingly nonlinear due to higher-order terms, which arise from the $(\Delta \omega)^n \cos^n(\omega_{ref}t)$ terms (n is the nth order harmonic). The mathematical origin of this nonlinearity was shown in [5]. The odd-terms n=1,3,... showed an amplitude-frequency relation that is asymmetric with respect to the resonator center frequency, and we could use this property to generate a feedback signal (see Fig.3).

The resonator was measured in a vacuum-probe station (Süss Microtech) using a lock-in amplifier (Stanford Research) and an RF signal source with FM-capability. The FM-PLL circuitry is shown in Fig.4. When closing the loop, we observed that a small phase error remained in the loop (parasitic cable capacitances); we used a DC voltage (V_{offset}) in combination with an adder in order to compensate for this error. Interestingly, both the feedback gain and the loop capture range could be controlled via $\Delta\omega$ (Fig.5), which implies an additional degree of freedom for the design of integrated NEMS-PLLs.

When using NEMS as gravimetric sensors, e.g. for measuring a gas concentration, it is useful to introduce the surface mass resolution:

$$\delta \mathcal{M}_s = 2 \, \mathcal{M}_{eff} / \mathcal{A}_{eff} \, \delta f / f_0 \approx 1.47 \, \delta f / f_0 \, t \, \rho \tag{1}$$

where $\delta f/f_0$ is the fractional frequency stability, ρ the material density, t the thickness and \mathcal{A}_{eff} the effective surface area. The modal mass M_{eff} was estimated $\sim 34x10^{-15}g$ for the fundamental mode (uniform mass loading). Eq.1 underlines that the design of RB-FETs must be carefully considered depending on their final application. We implemented different cc-beam resonator geometries, with the widest resonator showing the best areal mass sensitivity, owing to the largest $Q \sim 1000$ and output SNR achieved. Fig.6 shows the experimental stability of 2 ppm of the RB-FET resonator integrated in the FM-PLL, with the equivalent mass sensitivity of ~ 380 zg/ μ m² (1 zepto-gram= 10^{-24} kg). In air ($Q \sim 150$), 12 ppm stability was achieved.

We have demonstrated a novel heterodyne feedback loop based on a nonlinear effect in FM-demodulation with a resonant-body transistor. Importantly, the FM-PLL is compatible with the parallel actuation and readout of a large resonator arrays, which remains a crucial aspect for the design of real-world nano-sensors with enhanced the output SNR and capture cross-section [1, 9]. As such, the presented FM-NEMS-PLL could emerge a valuable scheme to realize low-power, low-noise, ultra-sensitive NEMS systems hybridized with CMOS on a single chip.

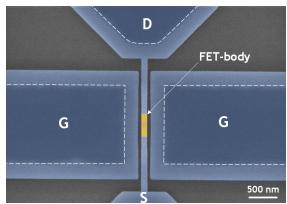


Fig.1: Top-view SEM image (false color) of the n^+pn^+ RB-FET with lateral drive electrodes. The dashed lines indicate the buried ox ide beneath the SOI device layer.

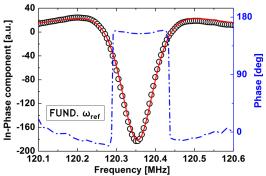


Fig.2: The mechanical resonance at f_0 =120.35MHz is shown as in-phase component of the drain current detected at ω_{ref} . The right axis shows the phase response.

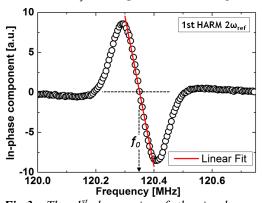


Fig.3: The 1^{st} harmonic of the in-phase component detected at $2\omega_{ref}$ is used to provide a negative loop feedback signal and lock onto the resonance at $f_0 = \omega_0/2\pi$.

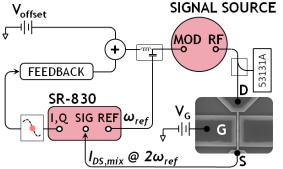


Fig.4: The circuit schematic of the single-source FM-NEMS-PLL (ω_{ref} =8kHz, V_G =10V). The signal source generates the FM-signal $\cos(\omega t + \Delta \omega/\omega_{ref} \sin \omega_{ref} t)$.

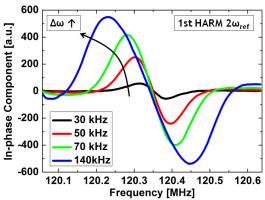


Fig.5: The amplitude response of the detected 1^{st} harmonic as function of the frequency deviation $\Delta \omega / 2\pi$. The loop capture range is reached at the point where the slope of the signal falls to zero.

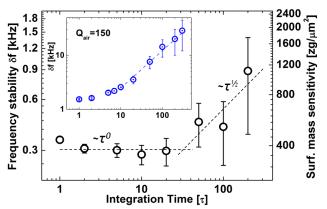


Fig.6: Experimental frequency stability and corresponding surface mass sensitivity vs. the integration time (295K, in high vacuum) for a cc-beam measuring 158nm x 2.65 μ m (W x L). The drive power is -26 dBm and the deviation $\Delta\omega/2\pi=70$ kHz. The inset shows the frequency stability of ~12 ppm measured at ambient conditions ($P_{in}=-18$ dBm).

ACKNOWLEDGMENT: This work was partially funded by the FP7 project NEMSIC.

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