

1 Solute transport in **nearly saturated** porous media under
2 landfill clay liners: A finite deformation approach

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13 **Abstract**

14 For solute transport in a deformable clay liner, the importance of consolidation
15 in the presence of sorption and consolidation-induced advection are well known.
16 Here a one-dimensional coupled consolidation and solute transport model for a
17 partially saturated porous medium, including the new features of finite strain and
18 geometric and material nonlinearity, is proposed. A new boundary condition at the
19 compacted clay liner (CCL) base is also introduced. A comprehensive compar-
20 ison demonstrates the significance of finite strain, compressibility of pore water
21 (CPW), longitudinal dispersion (LD) and the degree of saturation on the solute
22 transport in an unsaturated porous medium.

23 Consolidation in the presence of sorption and consolidation-induced advection
24 both affect solute transport in a deformable clay liner. Here, a one-dimensional
25 coupled consolidation and solute transport model for a **nearly saturated** porous
26 medium, including finite strain and geometric and material nonlinearity, was pro-

27 posed. A new boundary condition at the compacted clay liner base was also in-
28 troduced. The model demonstrates the significance of finite strain, pore water
29 compressibility, dispersion and the degree of saturation on solute transport in an
30 unsaturated, consolidating porous medium.

31 *Keywords:* material coordinates; finite deformation; degree of saturation; linear
32 equilibrium sorption; porous flow

33 **1. Introduction**

34 Land-based containment facilities are commonly used for the disposal of mu-
35 nicipal solid waste and contaminated dredged material (Liu, 2007). In mod-
36 ern landfills, liner systems are designed to isolate the landfill contents from the
37 surrounding environment to protect the groundwater from pollution. For well-
38 constructed composite liners, the geo-membrane typically has few defects, so re-
39 stricting advection through it (Giroud and Bonaparte, 1989; Foose et al., 2002).
40 However, volatile organic compounds (VOCs) can diffuse through membranes
41 with magnitude four to six orders greater than the possible advection. Therefore,
42 diffusion of VOCs in composite liners is viewed as a critical issue in the design of
43 landfill liners (Foose, 2002).

44 The VOC transit time was traditionally estimated using the diffusion equation
45 (Rowe and Badv, 1996; Fityus et al., 1999; Foose, 2002). However, several field
46 tests have reported that the transit of VOCs is much earlier than theoretical pre-
47 dictions (Workman, 1993; Othman et al., 1997). Many researchers attribute this to
48 consolidation and associated advective transport. Several theoretical models cou-
49 pling mechanical consolidation with solute transport were constructed in recent
50 years (Smith, 2000; Fox, 2007; Lewis et al., 2009).

51 There are opposing opinions regarding the importance of consolidation-induced
52 advection. Based on a model coupling finite deformation consolidation with so-
53 lute transport, Lewis et al. (2009) claimed that consolidation is essentially com-
54 plete before the VOC breaks through the clay liner, and its influence is further
55 minimized if sorption occurs. In their illustrative example, with linear sorption
56 at the level of $K_d = 0.001$ l/g, the consolidation made no discernible difference
57 to the concentration at the compacted clay liner (CCL) drainage base. Conse-
58 quently, they concluded that the advective transport flux has less influence on
59 solute migration than the combination of geometric and void ratio variation. With
60 this assumption, Lewis et al. (2009) proposed several simplified models, such as
61 the *instant deformation-diffusion only* model (calculates the final layer thickness
62 and void ratio before performing a diffusion-only analysis), and the *no advection*
63 model (ignores the advective transport component in the coupled model), to ap-
64 proximate the coupled consolidation and transport model. It is noted that, in their
65 model (Lewis et al., 2009), the boundary condition for the void ratio at the CCL
66 base is constant, which is not consistent with Smith (2000). On the other hand,
67 Fox (2007) presented contrary simulation results and stated that the advective flux
68 caused by consolidation has a lasting effect on transport even after the consolida-
69 tion has completed, and that its relative importance does not diminish for a VOC
70 sorption level up to 0.001 l/g.

71 In real environments, the clay barrier below the waste content is not fully sat-
72 urated (Fityus et al., 1999). Furthermore, when the liner materials are compacted,
73 the required optimal water content will cause the engineered clay to be partially
74 saturated. **The optimal water content in compacted clay is close to saturation**
75 **(Vaughan, 2003). Within a nearly saturated soil, the air phase is not continuous**

76 and exists in the form of occluded bubbles (Wang et al., 1997). Soil parameters,
77 such as hydraulic conductivity and effective diffusion, depend on the degree of
78 saturation. Effective diffusion decreases with consolidation. Consequently, the
79 relative importance of the mechanical dispersion component to effective diffusion
80 may reach a level at which it cannot be neglected. Moreover, the compressibility
81 of pore water, which is more pronounced in partially saturated soil (Fredlund and
82 Rahardjo, 1993), has been reported to reduce the rate of porous flows and con-
83 solidation (Booker and Carter, 1987; Vaziri and Christian, 1994). Since advective
84 solute transport is induced by pore-water flow, the compressibility of pore water
85 is expected to affect solute migration.

86 Based on the one-dimensional Biot consolidation theory, Zhang et al. (2012)
87 proposed an advection-diffusion equation that incorporates the degree of satura-
88 tion, compressibility of the pore fluid (CPW) and dispersivity of the solute trans-
89 port in a nearly saturated deforming porous medium. Both CPW and dispersivity
90 were found to influence solute migration within the CCL, significantly so in some
91 circumstances. However, Zhang et al. (2012) considered an infinitesimal strain,
92 (i.e., small deformation) model. Additionally, they did not consider the material
93 and geometric nonlinearity, factors that could be important in some circumstances
94 (Lewis et al., 2009). Financial constraints sometimes limit deployment of the
95 relatively costly CCLs. Natural clay deposits (sometimes with relatively high
96 compressibility) are used as substitutes. Since the soft clayey soil generally pro-
97 vides a good contact adhesion with a geomembrane, high effectiveness is a priori
98 expected. However, the finite deformation caused by the emplacement of waste
99 cannot be neglected.

100 The objective of this study is to extend the small deformation model for solute

101 transport in a **nearly saturated** medium (Zhang et al., 2012) to finite deforma-
 102 tions. This allows us to clarify the influence of consolidation in the progress of
 103 solute transport (using a time-dependent boundary in terms of void ratio at the
 104 CCL base). The influence of **the degree of saturation** on the VOC transit time in
 105 clay barriers will also be examined. To account for the geometric nonlinearity, a
 106 material coordinate system is used. Both CPW and dispersivity are considered in
 107 the new model. Further, our approach incorporates nonlinearity of the constitutive
 108 properties related to soil compressibility, the hydraulic conductivity and decreas-
 109 ing effective diffusion coefficient. A parametric study is carried out to examine
 110 the influence of several dominant parameters on the process of solute transport in
 111 porous medium.

112 **2. Model Formulation**

113 Recently, Lewis et al. (2009) and Peters and Smith (2002) developed a model
 114 coupling finite strain consolidation and solute transport in a fully saturated soil.
 115 Below, the CPW and dispersion in a nearly saturated soil is included.

116 *2.1. Coordinates systems*

117 A Lagrangian coordinate system (z, t) is employed to derive the flow and trans-
 118 port equations. We define $\xi(z, t)$ as the particle displacement with $\xi(z, 0) = z$. The
 119 relationship between Lagrangian and Eulerian (ξ, t) coordinate systems then im-
 120 plies that for any variable $F(z, t) = f(\xi(z, t), t)$:

$$\frac{\partial F}{\partial z} = \frac{\partial f}{\partial \xi} \frac{\partial \xi}{\partial z}, \quad \frac{\partial F}{\partial t} = \frac{\partial f}{\partial \xi} \frac{\partial \xi}{\partial t} + \frac{\partial f}{\partial t} = \frac{\partial f}{\partial \xi} v_s + \frac{\partial f}{\partial t}, \quad (1)$$

121 where $v_s = \partial \xi / \partial t$ is the solid velocity.

122 *2.2. Consolidation equations*

123 The equation describing changes in void ratio, $e(z, t)$, are derived from the
 124 continuity equations for the solid and fluid phases together with Darcy's law. The
 125 mass balance equation of the solid phase in differential form is::

$$\frac{\partial}{\partial t} \left[\rho_s (1 - n) \frac{\partial \xi}{\partial z} \right] = 0, \quad (2)$$

126 where ρ_s is the soil grain density, $n = e/(1 + e)$ is the current porosity, and $n_0 =$
 127 $n(z, 0)$ is the initial porosity. Note that, for constant ρ_s , the Jacobian, M , for the
 128 coordinate transformation is:

$$M = \frac{\partial \xi}{\partial z} = \frac{1 - n_0}{1 - n} = \frac{1 + e}{1 + e_0}, \quad (3)$$

129 where e_0 is the initial void ratio.

130 The continuity equation for the fluid phase (i.e., pore water) is

$$\frac{\partial}{\partial t} \left(n S_r \rho_f \frac{\partial \xi}{\partial z} \right) = - \frac{\partial}{\partial z} (\rho_f q), \quad (4)$$

131 where ρ_f is the pore fluid density.

132 According to Darcy's Law, the fluid flux is given by

$$q = - \frac{k_v}{\rho_f g} \frac{\partial p}{\partial \xi}, \quad (5)$$

133 where k_v is hydraulic conductivity and p is excess pore pressure. If the hydraulic
 134 gradient is constant, the Darcy equation in terms of total pressure can be trans-
 135 formed to this form (Peters and Smith, 2002).

136 Assuming ρ_f varies with pore pressure as $\partial\rho_f/\partial p = \beta\rho_f$ (Barry et al., 2007),
 137 substituting Eq. (5) into Eq. (4), then the continuity equation for the fluid phase
 138 becomes:

$$nS_r\beta\frac{\partial\xi}{\partial z}\frac{\partial p}{\partial t} + \frac{\partial}{\partial t}\left(S_r\frac{\partial\xi}{\partial z}\right) = \frac{1}{\rho_f g}\frac{\partial}{\partial z}\left(k_v\frac{\partial p}{\partial z}\frac{\partial\xi}{\partial z}\right), \quad (6)$$

139 where the compressibility of pore fluid (β) can be estimated by (Fredlund and
 140 Rahardjo, 1993):

$$\beta = \frac{S_r}{K_{w0}} + \frac{1 - S_r + r_h S_r}{P_a + P_0}, \quad (7)$$

141 in which K_{w0} is the pore water bulk modulus, r_h denotes volumetric fraction of
 142 dissolved air within pore water, P_a denotes gauge air pressure and P_0 represents
 143 the atmospheric pressure. In a nearly saturated soil, for example, $r_h = 0.02$,
 144 $S_r = 0.8 \sim 1.0$, β falls into the range of $2 \times 10^{-6} \sim 2 \times 10^{-7} \text{ Pa}^{-1}$.

145 Because n and n_0 (implicitly embedded in $\partial\xi/\partial z$) appear simultaneously, and
 146 n is unknown, Eq. (6) can not be directly solved in terms of p . In the following
 147 derivation, it turns out that once the relationship between the derivative of p (with
 148 respect to t and a) and the corresponding derivative of e is known, it is straightfor-
 149 ward to convert Eq. (6) to an equation in terms of e .

150 Assuming self-weight is negligible due to the relatively small thickness of the
 151 CCL (Zhang et al., 2012), the vertical force equilibrium is:

$$\frac{\partial\sigma}{\partial z} = 0, \quad (8)$$

152 where σ (now a function of t only) is the total normal stress of the soil and the
 153 z coordinate is vertically upwards. Assuming the compressive normal stress is
 154 positive, i.e., $\sigma = \sigma' + p$ (σ' is the effective normal stress), Eq. (8) leads to:

$$\frac{\partial p}{\partial \xi} = \frac{\partial}{\partial z} (-\sigma' + \sigma) \frac{\partial z}{\partial \xi} = \frac{1 + e_0}{1 + e} \frac{1}{\alpha_v} \frac{\partial e}{\partial z}, \quad (9)$$

155 where $\alpha_v = -de/d\sigma'$ is the coefficient of soil compressibility.

156 In the absence of self-weight, the rate of change of total stress at an arbitrary
 157 location equals that of the external top loading,

$$\frac{\partial \sigma}{\partial t} = \frac{\partial Q}{\partial t}, \quad (10)$$

158 where Q is the external load. The rate of change of the excess pore water pressure
 159 in the time domain is:

$$\frac{\partial p}{\partial t} = \frac{\partial}{\partial t} (\sigma - \sigma') = \frac{\partial Q}{\partial t} + \frac{1}{\alpha_v} \frac{\partial e}{\partial t}. \quad (11)$$

160 Substituting Eq. (3, 9, 11) into Eq. (6) yields:

$$\left(\frac{e S_r \beta}{(1 + e_0) \alpha_v} + \frac{S_r}{1 + e_0} \right) \frac{\partial e}{\partial t} - \frac{1 + e_0}{\rho_f g} \frac{\partial}{\partial z} \left(\frac{k_v}{\alpha_v (1 + e)} \frac{\partial e}{\partial z} \right) = - \frac{S_r \beta e}{1 + e_0} \frac{\partial Q}{\partial t}. \quad (12)$$

161 For the fully saturated case and when the CPW is neglected, i.e., $\beta = 0$, Eq. (12)
 162 reduces to:

$$\frac{1}{1 + e_0} \frac{\partial e}{\partial t} = \frac{1 + e_0}{\rho_f g} \frac{\partial}{\partial z} \left(\frac{k_v}{\alpha_v (1 + e)} \frac{\partial e}{\partial z} \right), \quad (13)$$

163 which is identical to Eq. (1) of Lewis et al. (2009).

164 2.3. Solute transport equations

165 Solute transport occurs in both solid and fluid phases. Here, for the nearly
 166 saturated soil, the mixture of pore water and entrapped air is taken as a homoge-
 167 neous fluid. Due to the discrete air bubbles, VOC transport by gas diffusion can
 168 be neglected in a nearly-saturated soil. Therefore, the mass conservation equation
 169 for the solute in the solid phase is:

$$\frac{\partial}{\partial t} \left[(1 - n) \rho_s S \frac{\partial \xi}{\partial z} \right] = f'_{a \rightarrow s}, \quad (14)$$

170 where S is the mass of solute sorbed on or within the solid phase per unit mass
 171 of the solid phase and $f'_{a \rightarrow s}$ denotes rate of solute loss in the water phase by solid
 172 phase sorption.

173 The mass conservation equation for solute in the fluid phase is:

$$\frac{\partial}{\partial t} \left(n S_r c_f \frac{\partial \xi}{\partial z} \right) = - \frac{\partial J_f}{\partial z} - f'_{a \rightarrow s}, \quad (15)$$

174 where c_f is the concentration of the solute in the pore fluid. In Eq. (15), the
 175 term $\partial \xi / \partial z$ comes from the volumetric change (Peters and Smith, 2002) and J_f
 176 represents solute flux in the fluid phase, which is described by (Peters and Smith,
 177 2002):

$$J_f(z, t) = n S_r (v_f - v_s) c_f - \frac{n S_r D}{M} \frac{\partial c_f}{\partial z}, \quad (16)$$

178 where D is the hydrodynamic dispersion coefficient. It is given by the sum of the
 179 effective diffusion coefficient (D_e) and the coefficient of mechanical dispersion
 180 (D_m):

$$D_m = \alpha_L (v_f - v_s), \quad (17)$$

181 where α_L is dispersion coefficient, v_f is the pore fluid velocity and $v_f - v_s$ denotes
182 the relative velocity of the pore fluid.

183 Based on Eq. (14-16), we have:

$$\frac{\partial}{\partial t} \left\{ [nS_r c_f + (1-n)\rho_s S] \frac{\partial \xi}{\partial z} \right\} = \frac{\partial}{\partial z} \left(\frac{nS_r D}{M} \frac{\partial c_f}{\partial z} \right) - \frac{\partial}{\partial z} [nS_r (v_f - v_s) c_f]. \quad (18)$$

184 The above equation can be further simplified with Darcy's Law, Eq. (5), and
185 the mass balance equations for both solid and fluid phases, Eqs. (2) and (4),
186 respectively. Equation (18) can then be expressed as:

$$\begin{aligned} nS_r \frac{\partial \xi}{\partial z} \frac{\partial c_f}{\partial t} + (1-n)\rho_s \frac{\partial \xi}{\partial z} \frac{\partial S}{\partial t} = \frac{\partial}{\partial z} \left(\frac{nS_r D}{M} \frac{\partial c_f}{\partial z} \right) + \frac{k_v}{\rho_f g} \frac{\partial p}{\partial \xi} \frac{\partial c_f}{\partial z} \\ + \left(nS_r \beta \frac{\partial \xi}{\partial z} \frac{\partial p}{\partial t} - \frac{\beta k_v}{\rho_f g} \frac{\partial p}{\partial \xi} \frac{\partial p}{\partial z} \right) c_f. \end{aligned} \quad (19)$$

187 Substituting Eq. (9) and Eq. (11) into Eq. (19) results in:

$$\begin{aligned} \left(S_r \frac{e}{1+e_0} + \frac{\rho_s K_d}{1+e_0} \right) \frac{\partial c_f}{\partial t} = S_r \frac{\partial}{\partial z} \left(\frac{e(1+e_0)}{(1+e)^2} D \frac{\partial c_f}{\partial z} \right) + \frac{k_v}{\rho_f g} \frac{1+e_0}{\alpha_v (1+e)} \frac{\partial e}{\partial z} \frac{\partial c_f}{\partial z} \\ + \beta \left[S_r \frac{e}{1+e_0} \left(\frac{\partial Q}{\partial t} + \frac{1}{\alpha_v} \frac{\partial e}{\partial t} \right) - \frac{k_v}{\rho_f g \alpha_v^2} \frac{1+e_0}{1+e} \left(\frac{\partial e}{\partial z} \right)^2 \right] c_f, \end{aligned} \quad (20)$$

188 where K_d describes the partitioning coefficient.

189 2.4. Special cases

190 In this section, three special cases of the present model are outlined.

191 *A. Saturated soil with finite deformation*

192 For a saturated soil, where $S_r = 1$, and incompressible pore fluid, i.e., $\beta = 0$,
 193 Eq. (20) reduces to:

$$\left(\frac{e}{1+e_0} + \frac{\rho_s K_d}{1+e_0} \right) \frac{\partial c_f}{\partial t} = \frac{\partial}{\partial z} \left(\frac{e(1+e_0)}{(1+e)^2} D \frac{\partial c_f}{\partial z} \right) + \frac{k_v}{\rho_f g} \frac{1+e_0}{\alpha_v(1+e)} \frac{\partial e}{\partial z} \frac{\partial c_f}{\partial z}, \quad (21)$$

194 which is identical to Eq. (4) of Lewis et al. (2009) and Eq. (44) in Peters and
 195 Smith (2002).

196 *B. Small deformation model*

197 Under the assumptions of negligible self-weight and small deformation (con-
 198 stant porosity, i.e., $n = n_0$), the coupled small deformation model is (Zhang et al.,
 199 2012):

$$S_r n_0 \beta \frac{\partial p}{\partial t} + S_r \frac{\partial^2 u}{\partial t \partial \xi^2} = \frac{1}{\rho_w g} \frac{\partial}{\partial \xi} \left(k_v \frac{\partial p}{\partial \xi} \right), \quad (22)$$

$$G \frac{2(1-\nu)}{(1-2\nu)} \frac{\partial^2 u}{\partial \xi^2} = \frac{\partial p}{\partial \xi} \quad (23)$$

200 and:

$$\begin{aligned} [S_r n_0 + (1-n_0)\rho_s K_d] \frac{\partial c_f}{\partial t} &= S_r n_0 D_e \frac{\partial^2 c_f}{\partial \xi^2} - \alpha_L \frac{k_v}{\rho_w g} \frac{\partial p}{\partial \xi} \frac{\partial^2 c_f}{\partial \xi^2} \\ &+ \frac{\partial c_f}{\partial \xi} \left\{ -\alpha_L S_r n_0 \beta \frac{\partial p}{\partial t} - \alpha_L S_r \frac{\partial^2 u}{\partial \xi \partial t} \right. \\ &+ \frac{\alpha_L \beta k_v}{\rho_w g} \left(\frac{\partial p}{\partial \xi} \right)^2 + S_r D_e (1-n_0) \frac{\partial^2 u}{\partial \xi^2} \\ &\left. + \frac{k_v}{\rho_w g} \frac{\partial p}{\partial \xi} - [S_r n_0 + (1-n_0)\rho_s K_d] \frac{\partial u}{\partial t} \right\} \\ &+ S_r n_0 \beta \frac{\partial p}{\partial t} c_f - \beta \frac{k_v}{\rho_w g} \left(\frac{\partial p}{\partial \xi} \right)^2 c_f + S_r n_0 \beta \frac{\partial u}{\partial t} \frac{\partial p}{\partial \xi} c_f, \end{aligned} \quad (24)$$

201 where u is the soil displacement, G the shear modulus and ν Poisson's ratio. The
 202 constant material coefficients can be described as:

$$G = \frac{c_v \rho_f g (1 - 2\nu)}{2k_v (1 - \nu)} = \frac{(1 + e_p)(1 - 2\nu)}{2(1 - \nu)\alpha_{vp}}, \quad (25)$$

$$k_v = k_p, \quad D_e = D_{e0},$$

203 where c_v is the consolidation coefficient; k_s and k_p the saturated hydraulic con-
 204 ductivity and hydraulic conductivity of the soil corresponding to e_p (the void ratio
 205 corresponding to pre-consolidation stress), respectively.

206 *C. Nearly saturated soil with no deformation*

207 For the partially saturated no deformation model, i.e., $e = e_0$, $\xi = z$, the
 208 overloading, Q , does not affect solute transport. In the spatial coordinate system
 209 (ξ, t) , Eq. (20) reduces to the linear diffusion equation:

$$\frac{\partial c_f}{\partial t} = D \left(1 + \frac{\rho_s K_d}{S_r e_0} \right)^{-1} \frac{\partial^2 c_f}{\partial \xi^2}. \quad (26)$$

210 **3. Variations of parameters in consolidation and solute transport processes**

211 *The finite deformation model allows consideration of the effects of variations*
 212 *in the coefficients of consolidation and transport (such as the coefficient of com-*
 213 *pressibility, α_v , hydraulic conductivity, k_v and hydrodynamic dispersion, D) on*
 214 *solute transport process.* Lewis et al. (2009) utilized void ratio-dependent func-
 215 tions for the related coefficients while Li and Liu (2006) used a fractal pore-space
 216 theory to develop fractal models of water flow and solute diffusion in rigid un-
 217 saturated soils. Their approach allowed comparison of these coefficients between

218 the fully saturated and unsaturated cases. Here, a combination of both models is
219 employed so that the hydraulic conductivity and the effective diffusion depend on
220 both the void ratio and the degree of saturation. Linear, reversible solute sorption
221 is assumed in this study; however, the approach can be adapted for other sorption
222 models.

223 3.1. Soil compressibility

224 The soil layer is assumed to be over-consolidated, and compression of the soil
225 layer commences when the applied stress exceeds the pre-consolidation stress,
226 i.e., deformation due to re-compression is neglected. In this case, the void ratio is
227 idealized as a linear function of the logarithm of the effective stress (Means and
228 Parcher, 1964):

$$e = e_p - C_c \log\left(\frac{\sigma'}{\sigma'_p}\right), \quad (27)$$

229 where σ' is effective stress, σ'_p denotes the pre-consolidation stress and C_c is the
230 compression index of the soil (defined by the absolute value of the slope of the
231 idealized virgin compression line). For a nearly saturated soil, the degree of sat-
232 uration is sufficiently high so that the air phase exists in the form of occluded
233 bubbles. Vaughan (2003) claimed that the presence of occluded air bubbles is un-
234 likely to affect soil effective stresses. Therefore, Eq. (27) is employed to describe
235 the volumetric change of a nearly saturated soil.

236 The coefficient of compressibility in terms of void ratio can be obtained by
237 differentiation of Eq. (27) with respect to effective normal stress (Lewis et al.,
238 2009):

$$\alpha_v = \alpha_{vp} \exp \left[\ln 10 \left(\frac{e - e_p}{C_c} \right) \right], \quad (28)$$

239 where α_{vp} is the coefficient of compressibility corresponding to σ'_p , i.e.,

$$\alpha_{vp} = \frac{C_c}{\sigma'_p \ln 10}. \quad (29)$$

240 3.2. Hydraulic characteristic

241 For hydraulic conductivity, an empirical relationship describing its variation
242 with void ratio in saturated clay soils is given as (Mitchell, 1993)::

$$k_s = k_p \exp \left[\ln(10) \left(\frac{e - e_p}{C_k} \right) \right], \quad (30)$$

243 where C_k is the hydraulic conductivity index.

244 The power law relationship equation for hydraulic conductivity versus water
245 content θ ($= S_r n$) is (Li and Liu, 2006):

$$k_v = k_s \left(\frac{\theta}{\theta_s} \right)^\alpha, \quad (31)$$

246 where θ_s is saturated water content, and α falls in the range of 2.68 to 2.78 for clay
247 loam.

248 3.3. Dispersion coefficient

249 In a saturated soil, the effective solute diffusion coefficient is defined as the
250 product of the free diffusion coefficient of the solute in the pore fluid (D_f) and the
251 tortuosity factor (t_f), which accounts for the irregular path that diffusing molecules

252 must take through the pore space (Acar and Haider, 1990). Lewis et al. (2009)
 253 claimed that it is rational to take D_e as constant, because uncertainty of the range
 254 of τ_f can be the same order of consolidation-induced change of D_e . Alternatively,
 255 the reduction of D_e can be expressed with a hypothetical relationship associated
 256 with the overall void ratio change as (Lewis et al., 2009; Morel-Seytour et al. ,
 257 1996):

$$D_e = \left(\frac{e_0 - e}{3(e_0 - e_f)} + \frac{e - e_f}{e_0 - e_f} \right) D_{e_0}, \quad (32)$$

258 where e_f denotes the final void ratio, and D_{e_0} is the initial effective dispersion
 259 coefficient.

260 In variably saturated soils, the effective diffusion coefficient, D_e , depends on
 261 soil water content, bulk density, and soil type for soils with different textures. Re-
 262 garding the water content, there is a threshold value under which solute diffusivity
 263 vanishes (Hunt and Ewing, 2003; Hamamoto et al., 2009). The impedance factor
 264 (Porter et al., 1960) (i.e., the ratio of solute diffusion coefficient in soil to prod-
 265 uct of solute diffusion coefficient in free water and volumetric soil water content),
 266 decreased with increasing bulk density for each soil type, but the effect of the
 267 overall bulk density on the impedance factor is minor compared with the effect of
 268 soil water content and soil type (Hamamoto et al., 2009). The effective diffusion
 269 coefficient was found to decrease with decreasing saturation in laboratory exper-
 270 iments (Barbour et al., 1996). The decrease was found to be quite rapid initially,
 271 followed by a near-linear decline for degree of saturation below 60%. Here, the
 272 soil diffusion coefficient is expressed as (Li and Liu, 2006)::

$$D_e = 1.1D_f\theta(\theta - \theta_t), \quad (33)$$

273 where θ_t denotes threshold water content, which was observed to become higher
274 with increasing clay content and varies between 3% and 20% for clay soil.

275 3.4. Sorption

276 It has been reported that the effect of the degree of saturation on the adsorp-
277 tion coefficient is insignificant from full saturation to a degree of saturation of
278 10% (Barbour et al., 1996). A significant decrease in the adsorption coefficient
279 only occurs in cases with a low degree of saturation. In this study the degree
280 of saturation varies from 1 to 0.8, i.e., the effect on sorption can be neglected.
281 Therefore, the concentration of solute in the solid phase, S , is expressed as:

$$S = K_d c_f. \quad (34)$$

282 This assumption of a linear sorption is valid at the relatively low concentrations
283 that are usually found in the municipal waste disposal sites (Mathur and Jayawar-
284 dena, 2008).

285 4. Application to a landfill liner

286 4.1. Problem description

287 As the schematic in Fig. 1 shows, the composite landfill liner beneath a pri-
288 mary leachate collect system (PLCS) consists of an impermeable (to diffusion of
289 inorganic solute) geomembrane, an underlying engineered compacted clay layer
290 (CCL), and a second leachate collecting system (SLCS).

291 The model parameters employed in the following analyses are based on those
 292 used in recent studies of solute transport in composite liners (Foose, 2002; Lewis
 293 et al., 2009). Because of the unavailability of consolidation data in the literature,
 294 hypothetical values of the applied stress, pre-consolidation stress, compression
 295 index, hydraulic conductivity index, threshold moisture content and other param-
 296 eters in calculating the D_e and k_v are used. As a primary parameter, the com-
 297 pression index covers a large range to account for the high-compressibility soil
 298 considered (Lewis et al., 2009). However, the related applied stress was selected
 299 to avoid negative and unrealistically low void ratios. The parameters used are
 300 given in Table 1.

301 4.2. *Boundary conditions for consolidation*

302 The following boundary conditions are introduced. Assuming there are no de-
 303 fects in the geomembrane, the top boundary ($z = 0$) is assumed to be impermeable,
 304 i.e., $q = 0$. Therefore, from Eq. (5) and Eq. (9),

$$\frac{\partial e}{\partial z} = 0 \text{ at } z = 0. \quad (35)$$

305 At the bottom drainage boundary ($z = L$), the excess pore pressure is zero
 306 and a Dirichlet-type boundary condition for void ratio (e) can be derived from the
 307 effective stress–void ratio equilibrium relationship, Eq. (27):

$$e = e_p - C_c \log \left(\frac{\sigma'_L}{\sigma'_p} \right), \quad (36)$$

308 where σ'_L denotes the effective stress at bottom.

309 The excess pore pressure vanishes at the bottom boundary, so $\sigma'_L = \sigma_a$, where
 310 σ_a is a time-varying stress due to the external overburden. Note that σ_a is the
 311 maximum loading in the model of Lewis et al. (2009). The void ratio rapidly ap-
 312 proaches a steady value, which consequently leads to a spurious higher fluid veloc-
 313 ity and faster solute transportation. To distinguish the cases, we label the present
 314 boundary condition at the CCL bottom as ‘BCC’ and ‘BCL’, i.e., the boundary
 315 conditions used by Lewis et al. (2009).

316 4.3. Boundary conditions for solute transport

317 At the top of the CCL, VOC diffusion through the geo-membrane is described
 318 by Fick’s law (Booker et al., 1997), so the concentration gradient is proportional
 319 to the difference in concentrations on each side of the (sufficiently thin) geomem-
 320 brane. In the material coordinate system, the boundary condition is (Lewis et al.,
 321 2009):

$$\frac{\partial c_f}{\partial z}(0, t) = \frac{(1 + e(0, t))^2}{e_0(1 + e_0)} \frac{P_G}{hD_e} (c_f(0, t) - C_{f0}), \quad (37)$$

322 where C_{f0} is the (constant) solute concentration at the top surface of the geo-
 323 membrane with the assumption that the landfill waste volume is large (Peters and
 324 Smith, 2002); h and P_G are, respectively, the thickness and the permeation coeffi-
 325 cient for the solute in the geo-membrane.

326 The lower boundary condition for the solute concentration (c_f) is (Peters and
 327 Smith, 2001):

$$\frac{\partial c_f}{\partial z} = 0, \text{ at } z = L, \quad (38)$$

328 which assumes negligible diffusion below the CCL base (Barry and Sposito, 1988).

329 5. Numerical results and discussion

330 A numerical solution was constructed using COMSOL 3.5a (Comsol, 2010).
331 It discretized the domain into unstructured Lagrange-linear elements with a max-
332 imum global element size of 10^{-2} m, and maximum local element size at the end
333 boundaries (where the most rapid changes occur) of 10^{-4} m. Temporally, the sub-
334 time step was 10^{-2} y. To be easily interpreted, solution curves were plotted in the
335 spatial coordinate x :

$$x = z + \int_z^L \frac{e_0 - e(\zeta)}{1 + e_0} d\zeta. \quad (39)$$

336 Thus, the first-order PDE,

$$\frac{\partial x}{\partial z} = 1 - \frac{e_0 - e(z)}{1 + e_0}, \quad (40)$$

337 with boundary conditions $x(0, t) = S_{mt}$ and $x(L, t) = L$ was constructed to find x ,
338 where the settlement S_{mt} is given by:

$$S_{mt} = \int_0^L \frac{e_0 - e(\zeta)}{1 + e_0} d\zeta. \quad (41)$$

339 5.1. Model verification

340 Since there are no experimental data available in the literature, the present
341 model was reduced to the full-saturation case using the same boundary condition
342 at the CCL bottom for e as used by Lewis et al. (2009), i.e., σ_a is taken as the
343 maximum loading; and $K_d = 0$, $\alpha_L = 0$, $C_c = 0.8$, $k_p = 10^{-9}$ m/s. A com-
344 parison between the present and previous models is illustrated in Fig. 2. In the

345 figure, the results of the finite deformation with constant and decreasing hydro-
346 dynamic dispersion, Eq. (32), small deformation model (Zhang et al., 2012) and
347 the pure diffusion model (i.e., no deformation model) are included. Both consol-
348 idation (i.e., void ratio, e , distribution) and relative concentration obtained from
349 the present model are in excellent agreement with results of Lewis et al. (2009).
350 As shown in Fig. 2, with the constant effective diffusion coefficient, the small
351 deformation model (Zhang et al., 2012) predicts a slower solute migration than
352 the corresponding finite deformation model.

353 *5.2. Correctness of the boundary condition at CCL base*

354 The differences due to the different boundary conditions, ‘BCL’ (used by
355 Lewis et al. (2009)) and ‘BCC’ (used in the present model), are presented in Fig.
356 3, where $C_c = 0.8$ and $k_p = 10^{-9}$ m/s. A comparison of Fig. 3(a) (BCC) and 2(a)
357 (BCL) shows that taking σ_a as the maximum loading leads to a greater void ratio
358 gradient and a faster consolidation process, although the final value of e is very
359 close. This initially speeds up the solute transit slightly, and then slows it down
360 in the long-term (Fig. 3(b)). The reason the trend reverses after the consolidation
361 completes for the ‘BCL’ case is that the higher solute concentration level during
362 the consolidation phase of ‘BCC’ occurs later resulting in an increased advective
363 flux. The separation is more obvious for the relatively soft and higher permeabil-
364 ity cases. In the following sections all numerical results are based on the boundary
365 condition ‘BCC’.

366 *5.3. Effect of consolidation*

367 On basis of the ‘BCL’ boundary condition, Lewis et al. (2009) observed that
368 there is no noticeable solute concentration at the CCL base when consolidation of

369 the liner is completed even for the case of very high compressibility ($C_c = 0.8$).
370 They thus concluded that transport can be simulated using the pure diffusion
371 model with the final void ratio value. However, during consolidation the dis-
372 tribution of solute concentration changes, which is the initial condition of what
373 follows. Thus, advective transport due to consolidation may not be negligible.

374 Figures 4 and 5 illustrate the consolidation processes and solute transport in a
375 saturated soil for two cases with different compression indices (C_c) and hydraulic
376 conductivities (k_v). Consolidation lasts 2.2 and 34.5 y for $C_c = 0.2$ and $C_c = 0.8$,
377 respectively. For the ‘soft’ case, a noticeable concentration difference from the
378 no deformation model appears at the CCL base during consolidation, as shown in
379 Fig. 5. The difference decreases with higher levels of sorption (Fig. 5(b)). The
380 effect of consolidation on transport exists during both the consolidation and post-
381 consolidation stages, which is consistent with Fox (2007). Since the advection
382 results in a notable concentration level at the CCL base, simplifying assumptions
383 such as instant deformation, pure diffusion and finite deformation without advec-
384 tion modelling are not appropriate. **The magnitude of solute concentration C_f in**
385 **Fig. 5(a) is an order greater than that in Fig. 5(b).** Here, the influence of sorption
386 is noticeable as it drastically retards the solute transport.

387 Figures 6 and 7 present the results for a **nearly saturated** soil. We see again
388 that soft clay consolidation has a noticeable effect on solute transport (Fig. 6).
389 However, since the effective diffusion (D_e) reduces with deformation, concentra-
390 tions for the pure diffusion model surpass those of coupled models, as is obvious
391 for the case of $K_d = 1$ ml/g.

392 Consolidation effects are composed of the variation of void ratio and the oc-
393 currence of pore water flow, which in turn causes the advective transport flux. As

394 mentioned previously, Lewis et al. (2009) claimed the advection component can
395 be ignored as long as the variation of void ratio is considered. Here, we included
396 in Fig. 8 the case of finite deformation without advection, i.e., advection is re-
397 moved from Eq. (20). Exclusion of advection underestimates the concentration
398 level and consequently leads to a longer transit time. In the absence of sorption,
399 at the nominal 10% breakthrough, a nearly twofold change occurs in the transit
400 time; this change increases when sorption is included.

401 5.4. *Effect of degree of saturation*

402 Fig. 9 demonstrates that the higher saturation of the no-deformation (ND)
403 model results in faster solute transport due to the saturation (S_r)-dependent ef-
404 fective diffusion; the gap is larger in the presence of sorption. Concentrations
405 predicted by the coupled finite deformation and solute transport model are shown
406 in Figures 10 and 11. For cases with parameters $C_c = 0.8$ and $k_p = 10^{-9}$ m/s,
407 consolidation lasts for approximately 12.8 y. Higher saturation results in faster
408 solute transport because of greater effective diffusion, regardless of the sorption.
409 For decreasing D_e , the transit time increases. With sorption, finite deformation
410 with $S_r = 0.8$ and constant D_e leads to almost the same concentration as for the
411 ND model (Fig. 11(b)). Again, this demonstrates that the effect of unsaturation
412 is more apparent in the presence of sorption. Interestingly, with both sorption
413 and decreasing D_e taken into account, finite deformation (FD) models will not
414 always produce faster solute transport (Fig. 10(b)). During consolidation and in
415 the early post-consolidation stage, the FD models have a faster transit, but then
416 are surpassed by the ND model because the effective diffusion is reduced due to
417 compaction. However, the decreasing D_e with compaction is inevitable. In the
418 field, VOC has been shown to appear earlier than predicted by the pure diffusion

419 model has been observed (Peters and Smith, 2002). Possible explanations are: (1)
420 the constitutive relationships for soil parameters are not accurate enough; or (2)
421 other factors, such as heat transfer, should be also included in the model.

422 5.5. *Effects of compressibility of pore water (CPW)*

423 As shown in Fig. 12, the effect of compressibility of pore water (CPW) is
424 related to the soil consolidation coefficient. The influence of CPW on the relative
425 concentration at the CCL becomes more significant for the cases with smaller con-
426 solidation coefficients. When the soil is relatively soft ($C_c = 0.8$ and $k_p = 2 \times 10^{-10}$
427 m/s), CPW causes twofold longer transit times for the nominal 10% breakthrough.
428 However, at the early consolidation stage, the retarding effect of CPW is more
429 pronounced for ‘stiffer’ soils and then the trend reverses (Fig. 12) after consol-
430 idation completes. These graphs are not shown as the numerical values are too
431 small to present in the same figure. This can be explained by the slowing fluid
432 flow and longer consolidation time due to CPW. Since the separation of curves at
433 a relatively higher concentration level, i.e., absolute concentration difference, is
434 of interest, it follows that the influence of CPW is more significant in softer soil.

435 To investigate further the influence of CPW, three models examining the three
436 terms containing β are considered here.

- 437 • Model A: eliminate $\frac{eS_r\beta}{(1+e_0)\alpha_v} \frac{\partial e}{\partial t}$ from Eq. (12);
- 438 • Model B: eliminate $-\frac{S_r\beta e}{1+e_0} \frac{\partial Q}{\partial t}$ from Eq. (12);
- 439 • Model C: eliminate the term involving β from Eq. (20).

440 As shown in Fig. 13, each of the missing terms leads to a large deviation from
441 the full model, so all terms involving β should be retained for the cases considered.

442 *5.6. Effect of dispersion*

443 Lewis et al. (2009) neglected mechanical dispersion on the assumption that
444 the pore fluid velocity in fine-grain soil is less than 10^{-6} m²/s. However, as shown
445 in Fig. 14, its influence cannot be neglected when the clay is relatively soft, even
446 when the maximum fluid average linear velocity is approximately 4.5×10^{-9} m/s
447 for the case $C_c = 0.8$ and $k_p = 2 \times 10^{-10}$ m/s. Its influence becomes more signif-
448 icant as the hydraulic conductivity increases with the same soil compressibility,
449 C_c . This is because decreasing D_e increases the Péclet number (ratio of the rate
450 of advection to the rate of diffusion). Therefore, a rough estimate using pore fluid
451 velocity alone as proposed by Lewis et al. (2009) is not always definitive.

452 Figure 15 illustrates the individual influence of decreasing D_e , dispersion and
453 CPW. The effect of reducing D_e causes slower transport, while dispersion a faster
454 transit. Although the influence of CPW is not as significant as decreasing D_e and
455 dispersion, it is not negligible, as shown in Fig. 15.

456 *5.7. Effect of finite deformation*

457 For the soil without sorption (see Fig. 2b, 10a, 11a, 15a), the ND model always
458 leads to a longer transit time than the finite deformation model. In the presence
459 of sorption (as shown in Fig. 11b), the difference between the ND model and the
460 finite deformation model is negligibly small. However, when the decrease of the
461 effective diffusion coefficient due to deformation is also considered (Fig. 10b and
462 15b), the results of the two models differ.

463 Compared with the finite deformation model, the small deformation model can
464 overestimate the contaminant transit time in a liner undergoing large consolida-
465 tion (Fig. 2b). This demonstrates that the significance of geometric nonlinearity
466 is noticeable for relatively soft soil. This finding is consistent with that of Peters

467 and Smith (2002) and Lewis et al. (2009). Regarding the consolidation, the small
468 deformation model can predict settlement that is non-physical for soft soil (i.e.,
469 larger than the total soil thickness). Therefore, for a relatively compressible soil,
470 where the consolidation effect is more significant, a finite deformation consolida-
471 tion is necessary when being coupled with the solute transport.

472 **6. Conclusion**

473 In this paper, a finite deformation model for coupling consolidation and solute
474 transport processes in partially saturated soil has been presented. It was applied to
475 predict the VOC breakthrough in a landfill clay liner. CPW, dispersion, the non-
476 linear variation of soil compaction, hydraulic conductivity and effective diffusion
477 are included in the model. Based on the numerical simulation results, we conclude
478 that:

- 479 1. Consolidation-induced advection has a lasting effect on solute transport dur-
480 ing and after the deformation for relatively compressible soil regardless
481 of the sorption level, though the sorption can dramatically slow the solute
482 transport process rate.
- 483 2. After an initial acceleration effect on transport, the finite-deformation cou-
484 pled model with decreasing effective diffusion and sorption produces a lower
485 concentration at the CCL base than the pure diffusion model.
- 486 3. A lower degree of saturation leads to a slower pore fluid flow and solute
487 transport(since larger pores drain preferentially with decreasing saturation).
488 The CPW associated with unsaturated conditions cannot be ignored when
489 the consolidation is required to be coupled with solute transport. In the

490 model, CPW terms exist in both the consolidation and transport equations,
491 none of which can be neglected for simplification. Effective diffusion de-
492 creases during consolidation and consequently the relative importance of
493 mechanical dispersion becomes profound. For a long-term prediction, me-
494 chanical dispersion could cause significant solute transport. Therefore, it
495 should be included in modelling efforts.

496 4. Generally speaking, reducing soil compressibility and improving sorption
497 levels of clay are the most effective ways to retard contaminant migration.
498 At the same level of stiffness and sorption, the lower hydraulic conductivity
499 and lower degree of saturation can lengthen the time for contaminants to
500 break through the protective liner.

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597 **List of symbols**

z , material coordinate, L

C_c , compression index of the soil

C_k , hydraulic conductivity index
 c_{f0} , solute mass concentration at top of geo-membrane, ML^{-3}
 c_f , concentration of the solute in the fluid phase, ML^{-3}
 c_s , concentration of the solute in the solid phase, ML^{-3}
 D , hydrodynamic dispersion coefficient, L^2T^{-1}
 e , void ratio
 e_0 , initial void ratio
 e_p , void ratio corresponding to the pre-consolidation stress
 P_G , mass transfer coefficient of geomembrane, L^2T^{-1}
 D_e , effective diffusion coefficient, L^2T^{-1}
 D_{e0} , initial effective dispersion coefficient, L^2T^{-1}
 D_f , free diffusion coefficient of the solute in the pore fluid, L^2T^{-1}
 D_m , coefficient of mechanical dispersion, L^2T^{-1}
 $f_{a \rightarrow s}$, rate of solute loss in aquatic phase by sorption onto solid phase, $ML^{-3}T^{-1}$
 G , shear modulus of soil, $ML^{-1}T^{-2}$
 g , gravity acceleration, LT^{-2}
 h , thickness of geomembrane, L
 J_f , solute flux in fluid phase, $M^2L^{-3}T^{-1}$
 k_p , hydraulic conductivity corresponding to e_p , LT^{-1}
 k_s , saturated hydraulic conductivity, LT^{-1}
 k_v , hydraulic conductivity, LT^{-1}
 K_d , contaminant partitioning coefficient, L^3M^{-1}
 K_{w0} , pore water bulk modulus, $ML^{-1}T^{-2}$
 L , thickness of CCL, L
 M , Jacobian of coordinate transformation

n , current soil porosity
 n_0 , initial soil porosity
 P_a , atmospheric pressure, $ML^{-1}T^{-2}$
 P_0 , atmosphere air pressure, $ML^{-1}T^{-2}$
 p , excess pore pressure, $ML^{-1}T^{-2}$
 r_h , volumetric fraction of dissolved air
 q , Darcy flow velocity, LT^{-1}
 Q , external load, $ML^{-1}T^{-2}$
 S , mass of contaminant sorbed onto the solid phase per unit mass of solid phase
 S_r , degree of saturation
 t , time, T
 u , soil displacement, L
 u' , arbitrary variable
 U , arbitrary variable
 v_f , average fluid velocity, LT^{-1}
 v_s , solid velocity, LT^{-1}
 x , spatial coordinate, L

598 **Greek symbols**

ξ , spatial coordinate, L
 τ_f , the tortuosity factor
 σ , total soil stress, $ML^{-1}T^{-2}$
 σ' , effective soil stress, $ML^{-1}T^{-2}$
 σ_a , the time varying stress due to external overburden, $ML^{-1}T^{-2}$

σ'_L , the effective stress at bottom, $\text{ML}^{-1}\text{T}^{-2}$

σ'_p , effective soil stress corresponding to the pre-consolidation stress

ρ_f , density of pore water, ML^{-3}

ρ_s , density of soil gain, ML^{-3}

β , compressibility of pore water, LT^2M^{-1}

ν , Poisson's ratio

α , coefficient in calculating k_v

α_L , longitudinal dispersion, L

α_v , coefficient of compressibility, LT^2M^{-1}

α_{vp} , coefficient of compressibility corresponding to σ'_p , LT^2M^{-1}

θ , water content

θ_s , saturated water content

θ_t , threshold water content

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600 1 Values of input parameters 35

Table 1: Values of input parameters

Parameter	Value
Maximum applied stress (ramp loading for 2 years), σ_a	450 kPa
Preconsolidation stress, σ'_p	50 kPa
Compression index, C_c	0.2, 0.8
Preconsolidation hydraulic conductivity, k_p	10^{-9} , 2×10^{-10} m/s
Constant, α	2.7
Hydraulic conductivity index, C_k	0.585
Thickness of geomembrane, h	0.0015 m
Thickness of CCL, L	1.22 m
Mass transfer coefficient of geomembrane, P_G	4×10^{-11} m ² /s
Initial effective diffusion coefficient, D_{e0}	2×10^{-10} m ² /s
Free diffusion coefficient in the pore fluid, D_f	10^{-9} m ² /s
Threshold moisture content, θ_t	0.05
Partitioning coefficient, K_d	0, 0.2, 1 ml/g
Dispersion, α_L	0, 0.1 m
Initial void ratio, e_0 ($= e_p$)	1.17
Acceleration due to gravity, g	9.81 m/s ²
Initial density of pore water, ρ_f	10^3 kg/m ³
Density of the solid phase, ρ_s	2.7×10^3 kg/m ³
Degree of saturation of clay, S_r	1, 0.9, 0.8

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603 2 Comparison of (a) void ratio evolution and (b) breakthrough curves
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606 formation model, ND: no deformation model. 40

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610 and dash-dot line for ‘BCL’. Case 1: $k_p = 2 \times 10^{-10}$ m/s, $C_c = 0.8$;
611 Case 2: $k_p = 10^{-9}$ m/s, $C_c = 0.8$; and Case 3: $k_p = 10^{-9}$ m/s, $C_c = 0.2$. 41

612 4 Consolidation settlements in a saturated soil ($S_r = 1$). 42

613 5 Effect of consolidation on relative concentration C_f/C_{f0} in a sat-
614 urated soil (a) $K_d = 0$ and (b) $K_d \neq 0$ ($S_r = 1$, without CPW,
615 $\alpha_L = 0$, constant D_e). Notations: solid line (FD, finite deforma-
616 tion model): $C_c = 0.8, k_p = 2 \times 10^{-10}$ m/s; dash-dot line (FD,
617 finite deformation model): $C_c = 0.2, k_p = 10^{-9}$ m/s; and dashed
618 line: no deformation model (ND). 43

619 6 Consolidation settlement in partially saturated soils ($S_r = 0.8$). . . 44

620 7 Effect of consolidation on relative concentration C_f/C_{f0} (a) $K_d =$
621 0 and (b) $K_d \neq 0$ in partially saturated soils ($S_r = 0.8$, with CPW,
622 $\alpha_L = 0.1$ m, varying D_e as in Equation Eq. (33)). Notations: solid
623 line (FD, finite deformation model): $C_c = 0.8, k_p = 2 \times 10^{-10}$ m/s;
624 dash-dot line (FD, finite deformation model): $C_c = 0.2, k_p = 10^{-9}$
625 m/s; and dashed line: no deformation model (ND). 45

626	8	Effect of advection flux on concentration level at CCL base for	
627		partially saturated cases ($S_r = 0.8$, with CPW, $\alpha_L = 0.1$ m,	
628		varying D_e as in (33)). For finite deformation model, solid line:	
629		$C_c = 0.8$, $k_p = 2 \times 10^{-10}$ m/s; dash-dot line: without advection	
630		flux in transport, (20); dashed line: No deformation model.	46
631	9	Effect of saturation S_r on transport for no-deformation model . . .	46
632	10	Concentration level at CCL base for partially saturated cases with	
633		decreasing D_e . ($C_c = 0.8$, $k_p = 10^{-9}$ m/s). Notation: FD: finite	
634		deformation model and ND: no deformation model.	47
635	11	Concentration level at CCL base for partially saturated cases with	
636		a constant D_e ($\theta = S_r n_0$ in (33)). ($C_c = 0.8$ and $k_p = 10^{-9}$ m/s).	
637		Notation: FD: finite deformation model and ND: no deformation	
638		model.	48
639	12	Effect of CPW on concentration level at CCL base for partially	
640		saturated cases ($S_r = 0.8$) with varying D_e and without sorption	
641		($K_d = 0$). Solid lines: $C_c = 0.8$, $k_p = 2 \times 10^{-10}$ m/s; Dashdot lines:	
642		$C_c = 0.8$, $k_p = 10^{-9}$ m/s; Dotted lines: $C_c = 0.2$, $k_p = 10^{-9}$ m/s.	
643		Cross symbol: with CPW; circle symbol: without CPW ($\beta = 0$). .	49
644	13	Significance of each term involving β on concentration level at	
645		CCL base for partially saturated cases ($S_r = 0.8$, $C_c = 0.8$, $k_p =$	
646		2×10^{-10} m/s) with varying D_e and without sorption ($K_d = 0$). . .	49

647	14	Effect of dispersion on concentration level at CCL base for partially saturated cases ($S_r = 0.8$) with varying D_e and without sorption ($K_d = 0$). Solid lines: $C_c = 0.8$, $k_p = 2 \times 10^{-10}$ m/s; Dashdot lines: $C_c = 0.8$, $k_p = 10^{-9}$ m/s; Dotted lines: $C_c = 0.2$, $k_p = 10^{-9}$ m/s. Cross symbol: $\alpha_L = 0.1$ m; circle symbol: $\alpha_L = 0$ (no dispersion).	50
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653	15	Comparison of the concentration level at CCL base for various variable associative in partially saturation soils ($S_r = 0.8$, $C_c = 0.8$, $k_p = 10^{-9}$ m/s). Notation: FD: finite deformation model; CD: constant D_e ; NLGD: excluding the dispersion; NCPW: excluding the CPW; ND: no deformation model.	51
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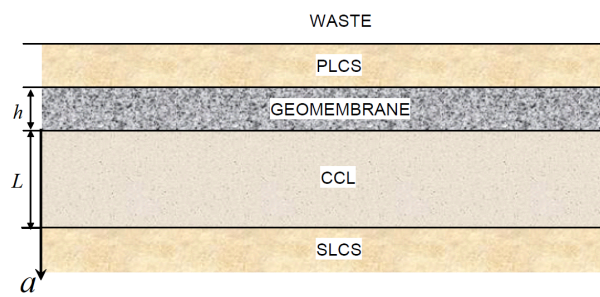
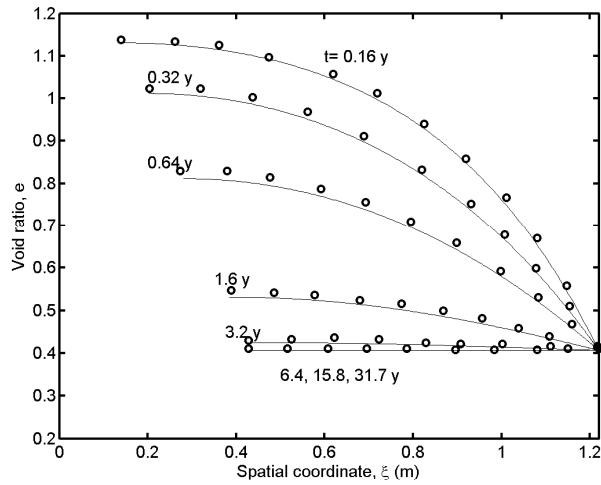
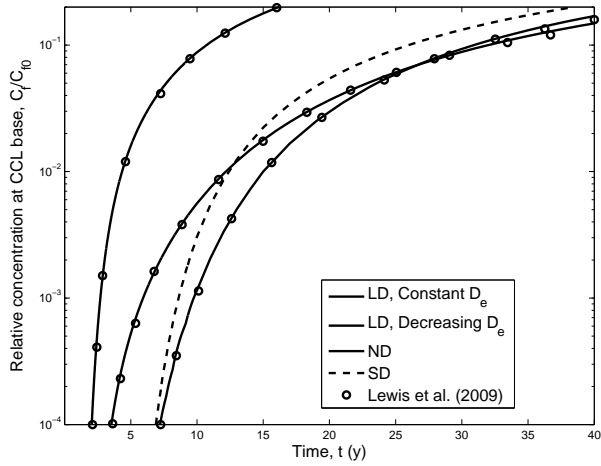


Figure 1: A Schematic of a composite landfill liner

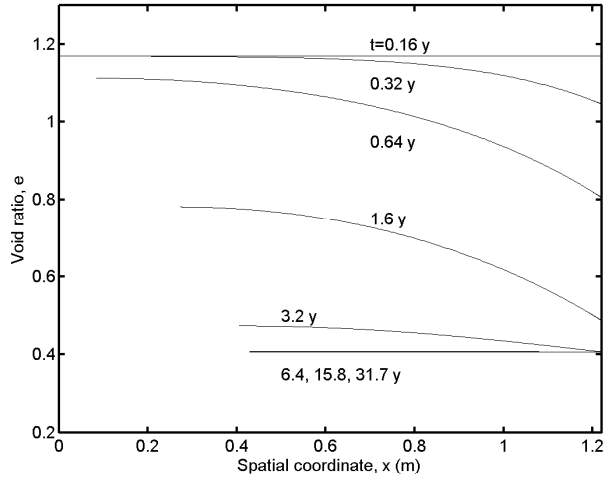


(a) Void ratio evolution

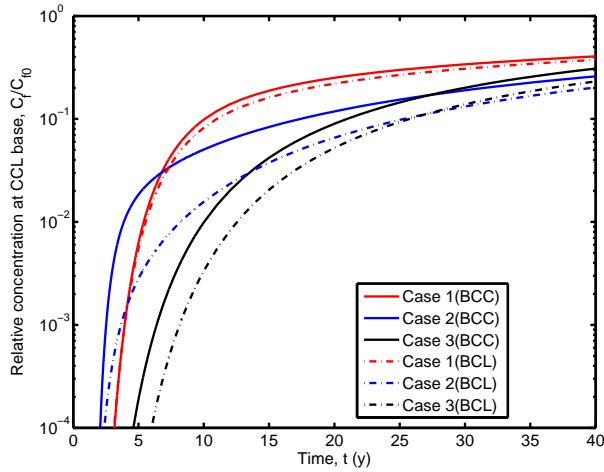


(b) Breakthrough curves

Figure 2: Comparison of (a) void ratio evolution and (b) breakthrough curves between the present model (solid line) and Lewis et al. (2009) (circle). Notations: FD: finite deformation model, SD: small deformation model, ND: no deformation model.



(a) Void ratio evolution (BCC only)



(b) Breakthrough curves (BCC and BCL)

Figure 3: Influence of Boundary condition of void ratio (e) at CCL base (a) void ratio evolution (BCC only) and (b) breakthrough curves ($S_r = 1, \beta = 0, \alpha_L = 0, \text{constant } D_e$). In (b), solid line for 'BCC', and dash-dot line for 'BCL'. Case 1: $k_p = 2 \times 10^{-10}$ m/s, $C_c = 0.8$; Case 2: $k_p = 10^{-9}$ m/s, $C_c = 0.8$; and Case 3: $k_p = 10^{-9}$ m/s, $C_c = 0.2$.

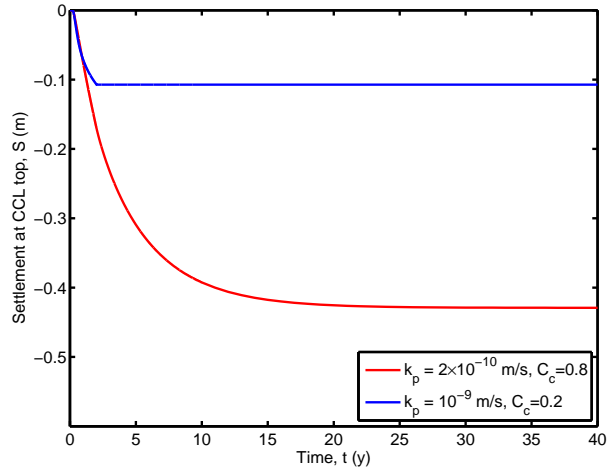
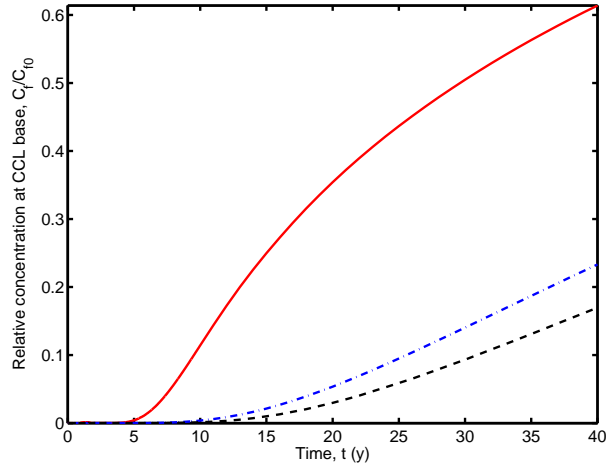
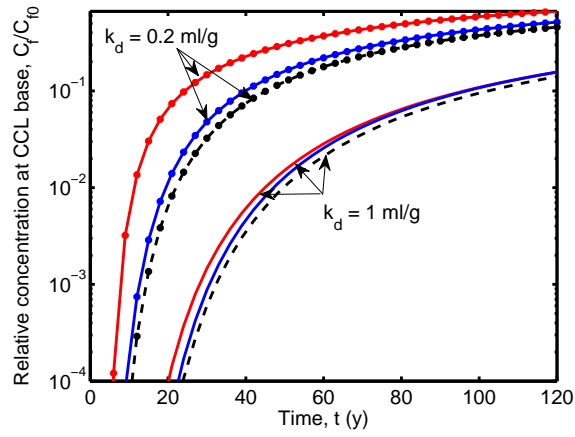


Figure 4: Consolidation settlements in a saturated soil ($S_r = 1$).



(a) $K_d = 0$



(b) $K_d \neq 0$

Figure 5: Effect of consolidation on relative concentration C_f/C_{f0} in a saturated soil (a) $K_d = 0$ and (b) $K_d \neq 0$ ($S_r = 1$, without CPW, $\alpha_L = 0$, constant D_e). Notations: solid line (FD, finite deformation model): $C_c = 0.8$, $k_p = 2 \times 10^{-10}$ m/s; dash-dot line (FD, finite deformation model): $C_c = 0.2$, $k_p = 10^{-9}$ m/s; and dashed line: no deformation model (ND).

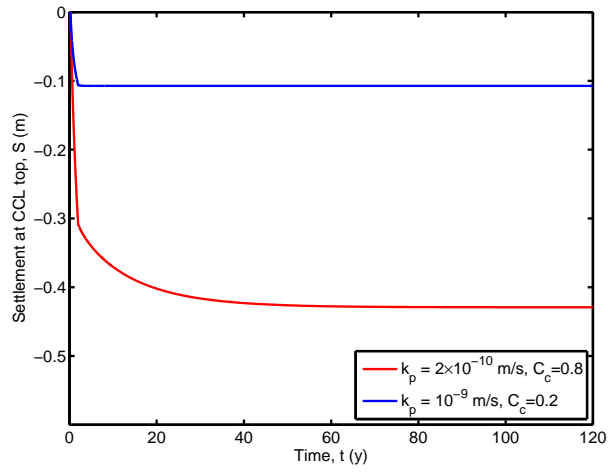
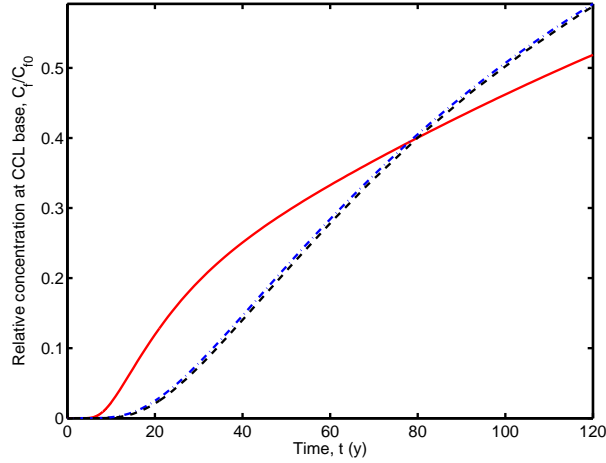
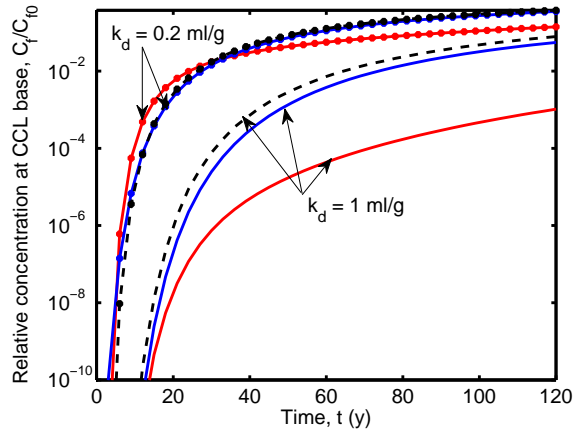


Figure 6: Consolidation settlement in partially saturated soils ($S_r = 0.8$).



(a) $K_d = 0$



(b) $K_d \neq 0$

Figure 7: Effect of consolidation on relative concentration C_f/C_{f0} (a) $K_d = 0$ and (b) $K_d \neq 0$ in partially saturated soils ($S_r = 0.8$, with CPW, $\alpha_L = 0.1$ m, varying D_e as in Equation Eq. (33)). Notations: solid line (FD, finite deformation model): $C_c = 0.8$, $k_p = 2 \times 10^{-10}$ m/s; dash-dot line (FD, finite deformation model): $C_c = 0.2$, $k_p = 10^{-9}$ m/s; and dashed line: no deformation model (ND).

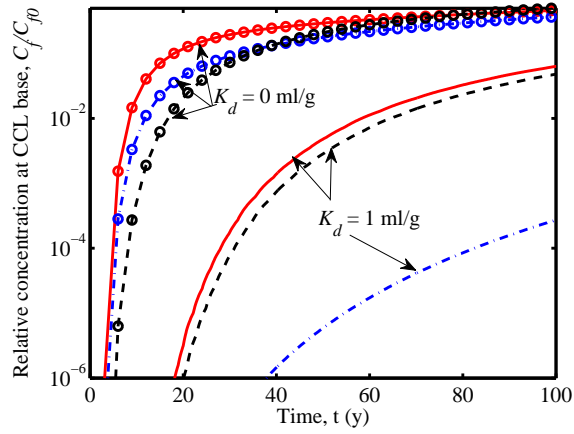


Figure 8: Effect of advection flux on concentration level at CCL base for partially saturated cases ($S_r = 0.8$, with CPW, $\alpha_L = 0.1$ m, varying D_e as in (33)). For finite deformation model, solid line: $C_c = 0.8$, $k_p = 2 \times 10^{-10}$ m/s; dash-dot line: without advection flux in transport, (20); dashed line: No deformation model.

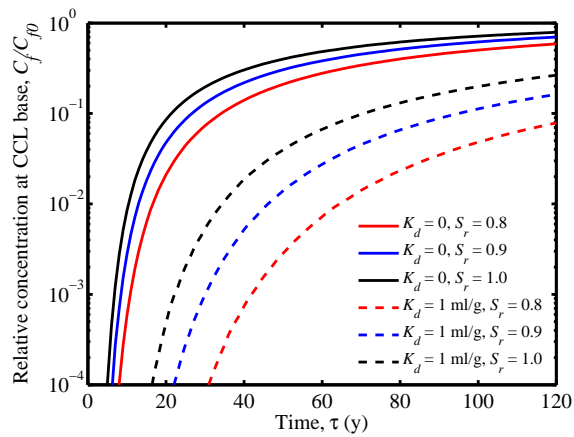
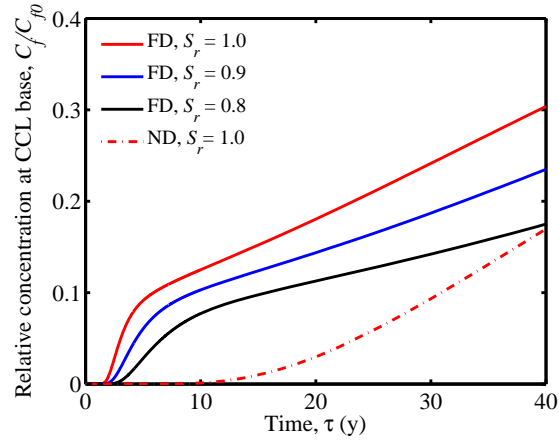
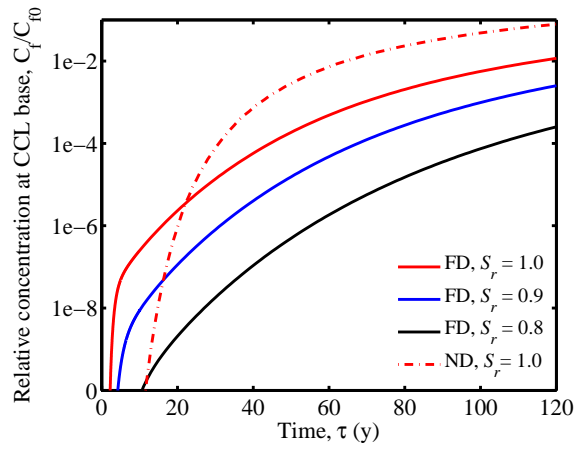


Figure 9: Effect of saturation S_r on transport for no-deformation model

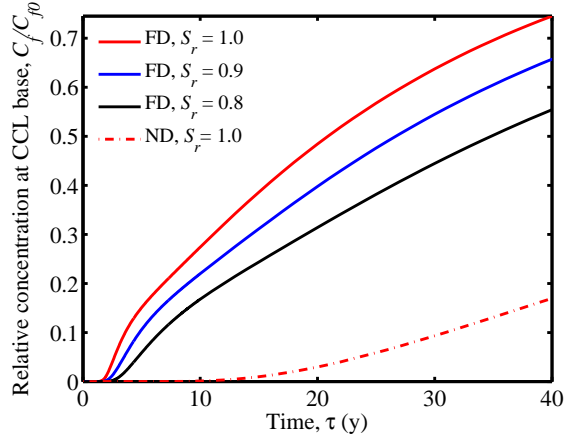


(a) $K_d = 0$

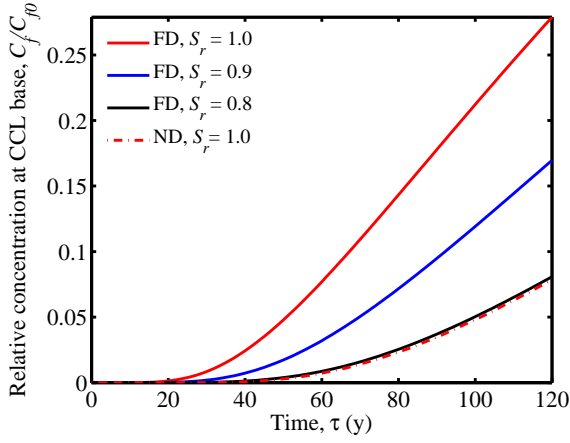


(b) $K_d = 1 \text{ ml/g}$

Figure 10: Concentration level at CCL base for partially saturated cases with decreasing D_e . ($C_c = 0.8$, $k_p = 10^{-9} \text{ m/s}$). Notation: FD: finite deformation model and ND: no deformation model.



(a) $K_d = 0$



(b) $K_d = 1 \text{ ml/g}$

Figure 11: Concentration level at CCL base for partially saturated cases with a constant D_e ($\theta = S_r n_0$ in (33)). ($C_c = 0.8$ and $k_p = 10^{-9}$ m/s). Notation: FD: finite deformation model and ND: no deformation model.

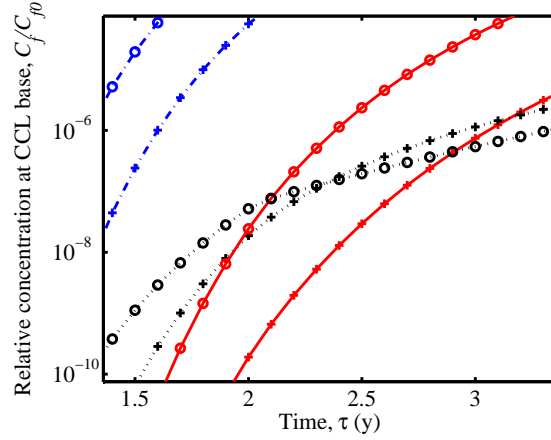


Figure 12: Effect of CPW on concentration level at CCL base for partially saturated cases ($S_r = 0.8$) with varying D_e and without sorption ($K_d = 0$). Solid lines: $C_c = 0.8$, $k_p = 2 \times 10^{-10}$ m/s; Dashdot lines: $C_c = 0.8$, $k_p = 10^{-9}$ m/s; Dotted lines: $C_c = 0.2$, $k_p = 10^{-9}$ m/s. Cross symbol: with CPW; circle symbol: without CPW ($\beta = 0$).

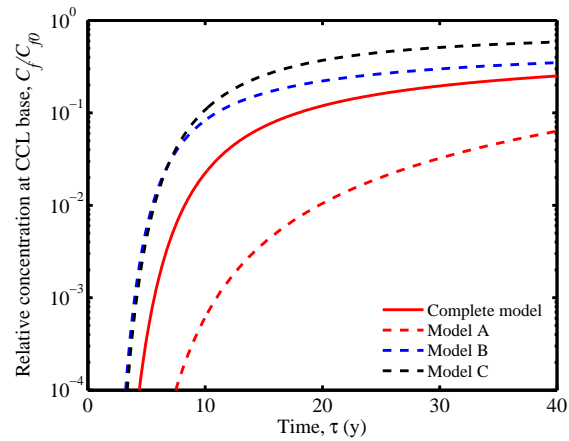


Figure 13: Significance of each term involving β on concentration level at CCL base for partially saturated cases ($S_r = 0.8$, $C_c = 0.8$, $k_p = 2 \times 10^{-10}$ m/s) with varying D_e and without sorption ($K_d = 0$).

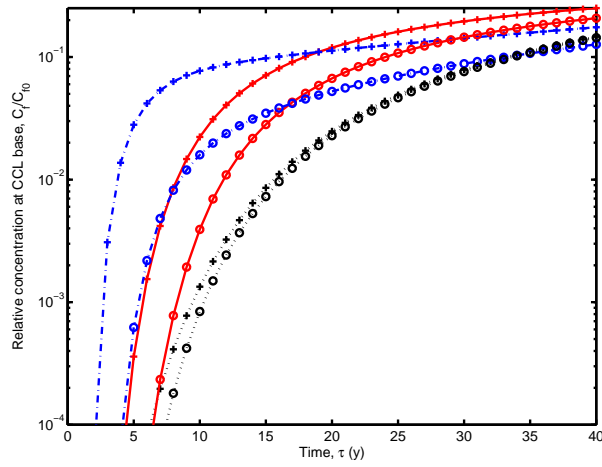
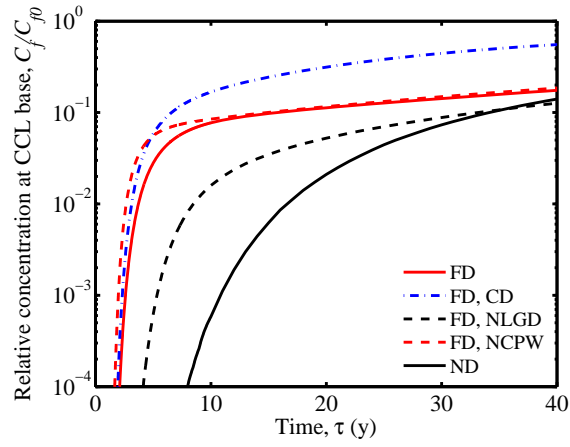
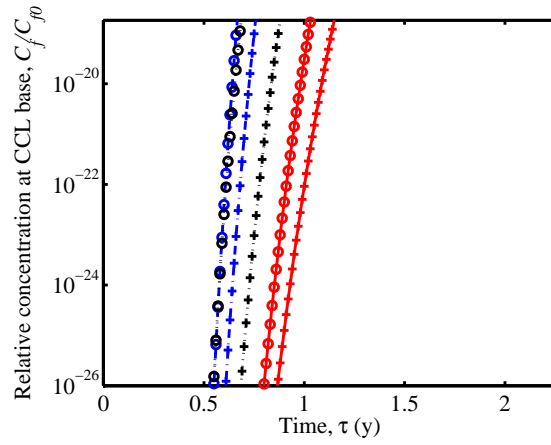


Figure 14: Effect of dispersion on concentration level at CCL base for partially saturated cases ($S_r = 0.8$) with varying D_e and without sorption ($K_d = 0$). Solid lines: $C_c = 0.8$, $k_p = 2 \times 10^{-10}$ m/s; Dashdot lines: $C_c = 0.8$, $k_p = 10^{-9}$ m/s; Dotted lines: $C_c = 0.2$, $k_p = 10^{-9}$ m/s. Cross symbol: $\alpha_L = 0.1$ m; circle symbol: $\alpha_L = 0$ (no dispersion).



(a) $K_d = 0$



(b) $K_d = 1 \text{ ml/g}$

Figure 15: Comparison of the concentration level at CCL base for various variable associative in partially saturation soils ($S_r = 0.8$, $C_c = 0.8$, $k_p = 10^{-9} \text{ m/s}$). Notation: FD: finite deformation model; CD: constant D_e ; NLGD: excluding the dispersion; NCPW: excluding the CPW; ND: no deformation model.